



On Schrödinger Operators with Inverse Square Potentials on the Half-Line

Jan Dereziński and Serge Richard

Abstract. The paper is devoted to operators given formally by the expression

$$-\partial_x^2 + \left(\alpha - \frac{1}{4}\right) \frac{1}{x^2}.$$

This expression is homogeneous of degree minus 2. However, when we try to realize it as a self-adjoint operator for real α , or closed operator for complex α , we find that this homogeneity can be broken. This leads to a definition of two holomorphic families of closed operators on $L^2(\mathbb{R}_+)$, which we denote $H_{m,\kappa}$ and H_0^ν , with $m^2 = \alpha$, $-1 < \operatorname{Re}(m) < 1$, and where $\kappa, \nu \in \mathbb{C} \cup \{\infty\}$ specify the boundary condition at 0. We study these operators using their explicit solvability in terms of Bessel-type functions and the Gamma function. In particular, we show that their point spectrum has a curious shape: a string of eigenvalues on a piece of a spiral. Their continuous spectrum is always $[0, \infty[$. Restricted to their continuous spectrum, we diagonalize these operators using a generalization of the Hankel transformation. We also study their scattering theory. These operators are usually non-self-adjoint. Nevertheless, it is possible to use concepts typical for the self-adjoint case to study them. Let us also stress that $-1 < \operatorname{Re}(m) < 1$ is the maximal region of parameters for which the operators $H_{m,\kappa}$ can be defined within the framework of the Hilbert space $L^2(\mathbb{R}_+)$.

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S. Richard: On leave of absence from Univ. Lyon, Université Claude Bernard Lyon 1, CNRS UMR 5208, Institut Camille Jordan, 43 blvd. du 11 November 1918, F-69622 Villeurbanne cedex, France.

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1. Introduction

The family of differential operators

$$-\partial_x^2 + \left(\alpha - \frac{1}{4}\right) \frac{1}{x^2} \quad (1.1)$$

is very special. They are homogeneous of degree -2 . They appear in numerous applications, *e.g.*, as the radial part of the Laplacian in any dimension. Their eigenfunctions can be expressed in terms of Bessel-type functions, and they have a surprisingly long and intricate theory, see [8, 9, 12, 14, 25, 29].

It is natural to try to interpret (1.1) as a closed operator on $L^2(\mathbb{R}_+)$. The most natural interpretation was given and extensively studied in [7]. It involves setting $\alpha = m^2$ and considering a family of closed operators H_m defined for $\operatorname{Re}(m) > -1$, depending on m holomorphically. In fact, the region $\operatorname{Re}(m) > -1$ can be divided into two parts. For $\operatorname{Re}(m) \geq 1$, the operator H_m corresponds to the closure of

$$-\partial_x^2 + \left(m^2 - \frac{1}{4}\right) \frac{1}{x^2} \quad (1.2)$$

restricted to $C_c^\infty(\mathbb{R}_+)$. For $-1 < \operatorname{Re}(m) < 1$, it is necessary to specify a boundary condition: $H_{\pm m}$ are both extensions of (1.2) with elements in their domain behaving like $cx^{\frac{1}{2} \pm m}$ near 0 for some $c \in \mathbb{C}$. Note that all operators H_m are homogeneous of degree -2 .

For $-1 < \operatorname{Re}(m) < 1$, more general boundary conditions can be studied. By considering elements of $L^2(\mathbb{R}_+)$ behaving like $c(\kappa x^{1/2-m} + x^{1/2+m})$ for some $c \in \mathbb{C}$, one naturally obtains a two-parameter family of closed operators $H_{m,\kappa}$, with $\kappa \in \mathbb{C} \cup \{\infty\}$. Note that these operators are no longer homogeneous, except for $H_{m,0} = H_m$ and $H_{m,\infty} = H_{-m}$.

A separate analysis is required for $m = 0$. Possible boundary conditions for this case are $c(x^{1/2} \ln(x) + \nu x^{1/2})$ for some $c \in \mathbb{C}$. These lead to a family of closed operators that we denote by H_0^ν , where $\nu \in \mathbb{C} \cup \{\infty\}$. Note that $H_0^\infty = H_0$, and that this operator is the only homogeneous one amongst the operators H_0^ν .

The study of the two families of operators $H_{m,\kappa}$ and H_0^ν is the object of our paper. We extensively use their explicit solvability, so that for these operators, we can give exact formulas for many constructions of operator theory. Let us mention part of the analysis performed below.

The operators H_m have no point spectrum. However, $H_{m,\kappa}$ and H_0^ν usually have point spectrum, which coincides with the discrete spectrum. All the eigenvalues are simple and depend quite sensitively on the parameters. The number of these eigenvalues can be finite, but also infinite. Their position forms rather interesting patterns on the complex plane: typically, it is a sequence situated along a piece of a spiral.

The continuous spectrum always coincides with the positive half-line $[0, \infty[$. One can express the resolvent of our operators in terms of the *MacDonald* and *modified Bessel functions* K_m and I_m . The resolvent has boundary

values at $[0, \infty[$, which can be expressed in terms of the *Hankel* and *Bessel functions* H_m^\pm and J_m . We make sense out of these boundary values as bounded operators between appropriate weighted Hilbert spaces, a property often called the *Limiting Absorption Principle*. We provide formulas for the spectral projections onto parts of the continuous spectrum, and introduce bounded invertible operators that diagonalize $H_{m,\kappa}$ and H_0^ν , except for a small set of parameters that we call *exceptional*. These operators can be called *generalized Hankel transformations*, see also [33] for related results. Finally, we describe the scattering theory for $H_{m,\kappa}$ and H_0^ν , giving formulas for the *Møller (wave) operators* and for the *scattering operator*.

Let us mention that most of the operators $H_{m,\kappa}$ and H_0^ν are not self-adjoint, but that the subfamilies of self-adjoint ones are also exhibited. More precisely, $H_{m,\kappa}$ are self-adjoint for real m and κ , or for purely imaginary m and $|\kappa| = 1$. Similarly, H_0^ν are self-adjoint for real ν . Thus, our analysis fits into a recent fashion of studying spectral properties of non-self-adjoint operators. Indeed, we observe that in the non-self-adjoint cases, our operators have quite interesting discrete spectrum.

We also stress that the operators that we study are extremely natural and appear in many situations. For example, they describe oscillations of a conical membrane, the Aharonov–Bohm effect [7, 28], and sticky diffusion [16]. They also play an important role in the study of the wave and Klein–Gordon equations on anti-de Sitter spacetime, see [5, 15, 19] and references therein. We are convinced that they possess many more applications we are not aware of. Indeed, since (1.2) is a homogeneous expression and the only non-homogeneity is due to boundary conditions, we expect that the studied operators appear in various scaling limits, or *renormalization group analysis*.

The analysis of our paper can also be considered as a large part of modern theory of the Bessel equation and the Gamma function (which belong to the oldest objects of mathematics, going back at least to Euler in 17th century). Indeed, many identities for Bessel-type functions and the Gamma function have a meaning in the theory of $H_{m,\kappa}$ and H_0^ν . Let us give some simple examples: the identity $z\Gamma(z) = \Gamma(z+1)$ is related to the Møller operators for the pair (H_{m+2}, H_m) , and the identity $\Gamma(z)\Gamma(1-z) = \frac{\pi}{\sin \pi z}$ is related to the Møller operators for the pair (H_{-m}, H_m) . Note that our point of view on Bessel functions is further developed in Sect. 3 and in Appendix A.

Many authors studied various classes of one-dimensional Schrödinger operators on the half-line, also called Sturm–Liouville operators on the half-line. Among classic works on this topic, which included complex potentials and various boundary conditions, let us mention [11, Chap. XX] and [27].

The methods used in our paper are of course adaptations of known approaches. Let us also note that the resolvents of $H_{m,\kappa}$, resp. H_0^ν are rank one perturbations of the resolvents of the operators H_m , resp. H_0 . Therefore, our analysis can be interpreted as an example of the theory of singular rank one perturbation, which is a well-studied subject. Nevertheless, it seems that a large part of our analysis of the operators $H_{m,\kappa}$ and H_0^ν , especially in the non-self-adjoint case, is new. Let us also mention that there are still some topics

about these operators, which are open. In particular, in this paper, we do not analyze fully the *exceptional case*, in which the generalized Hankel transformation are unbounded. In this case, *spectral singularities* appear, first noted in a similar context in [27].

In our paper, we restrict ourselves to a rather narrow, explicit, and solvable family of operators. However, it is not easy to find papers about more general classes of operators that cover the whole family we consider. In many works on one-dimensional Schrödinger operators, the singularity $\frac{1}{x^2}$ is excluded by restrictive assumptions. This is, e.g., the case of [11, 27]. Many papers also assume that the operator is dissipative or accretive (its numerical range is contained in the upper, resp. lower complex half-plane). Our operators are often neither dissipative nor accretive. In fact, it was noted already in [7] that for $\operatorname{Re}(m) < 0$ and $\operatorname{Im}(m) \neq 0$, the numerical range of H_m is the whole \mathbb{C} .

A vast majority of papers about Schrödinger operators considers only self-adjoint cases. One often assumes the essential self-adjointness on $C_c^\infty(\mathbb{R}_+)$. Then, H_m with $|m| < 1$ are not seen at all. If one considers the Friedrichs extension, only the case $m \geq 0$ is covered. If one considers both the Friedrichs and Krein extensions, one throws away the interesting region $m^2 < 0$. Looking at the exactly solvable potential $(m^2 - \frac{1}{4})\frac{1}{x^2}$ and using theory of the Bessel equation, we can check what are the natural assumptions in our case. In particular, we can notice that it is natural to include non-self-adjoint cases, which interpolate in an interesting way between self-adjoint cases, and to study a holomorphic family of closed operators depending on two complex parameters.

With this idea in mind, let us compare our work with some of the recent papers dealing with the operator (1.2). Note that many papers are related to this operator, and therefore, we mention only a few of them (see also the references in these papers). First of all, let us mention [17] in which an extensive study of Schrödinger operators of the form $-\partial_x^2 + V$ with singular V is performed. However, when the special case of $V(x) = (\alpha - \frac{1}{4})\frac{1}{x^2}$ is considered, only the parameters $m \geq 1$ are considered. In [13] the operator (1.2) is also thoroughly analyzed in the range $m \in [0, 1[$, but at the end of the day, only the Friedrichs self-adjoint extension is considered.

In [2] the Friedrichs realization H_m of the operator (1.2) with $m \in [0, 1[$ is considered and the scattering theory is developed for the pair (H_m, H_D) with H_D the Dirichlet Laplacian on the half-line. Again, the set of m considered is rather restrictive. In addition, a sentence like “for $m \geq 1$, no scattering is possible between H_D and H_m ” (see page 85 of that paper) is in contradiction with the scattering theory developed in our Sect. 4 and even further extended in the subsequent sections.

In the paper [24] and in the preprint [3], the Friedrichs realization of the operator (1.2) is also considered for $m \geq 0$. In the former paper, some dispersive estimates are provided for the evolution group $\{e^{-itH_m}\}_{t \in \mathbb{R}}$ with an emphasis in the dependence in m . In [3], a new study of the expression (1.2) on finite intervals or on the half-line is performed with the recently introduced

approach of boundary triplets. The Krein and the Friedrichs extensions are, indeed, considered on the half-line, but the parameter m is always real.

Finally, let us mention the recent paper [23] and the related subsequent preprint [20]. In the former one, dispersive estimates are provided for the operator $-\partial_x^2 + (m^2 - \frac{1}{4}) \frac{1}{x^2} + q(x)$ under some additional conditions on the potential q . Here, only the Friedrichs extension for $m > 0$ is considered. Obviously, the additional potential enlarges the set of operators under investigation, but on the other hand, only the simplest realization of these operators is analyzed. In the preprint [20], only the initial operator (1.2) is considered, (*i.e.*, $q = 0$) for $m \in]0, 1[$, but a rather large family of self-adjoint realizations of this operator is introduced. Dispersive estimates for these operators are obtained, and their dependence on the boundary condition at 0 is emphasized.

The family of operators H_m with $\operatorname{Re}(m) > -1$ was introduced in [7]. Thus, our study of $H_{m,\kappa}$ and H_0^ν can be viewed as a continuation of [7]. It seems, however, that some of the properties of H_m , notably about scattering in the non-self-adjoint case, are described in our present paper for the first time.

One could ask how complete our analysis is. In particular, it is natural to ask whether the operators H_m can be extended holomorphically outside the domain $-1 < \operatorname{Re}(m)$, and $H_{m,\kappa}$ outside $-1 < \operatorname{Re}(m) < 1$. Most probably, if we stick to the framework of the Hilbert space $L^2(\mathbb{R}_+)$, the answer is negative. A question about whether H_m can be extended holomorphically was formulated in [7] and has not been settled rigorously yet. However, one can extend the operators $H_{m,\kappa}$ and H_m to larger domains of parameters, if one goes beyond the framework of the Hilbert space $L^2(\mathbb{R}_+)$. This subject will be considered in a separate paper.

2. Inverse Square Potential

2.1. Notation

\mathbb{C}^\times denotes $\mathbb{C} \setminus \{0\}$. $\bar{\alpha}$ means the complex conjugate of $\alpha \in \mathbb{C}$. If Y is a set, then $\#Y$ denotes the number of elements of Y . If Ξ is a subset of \mathbb{R} , then $\mathbb{1}_\Xi$ represents the characteristic function of Ξ . $C_c^\infty(\mathbb{R}_+)$ denotes the set of smooth functions on $\mathbb{R}_+ :=]0, \infty[$ with compact support.

For an operator A , we denote by $\mathcal{D}(A)$ its domain and by $\sigma_p(A)$ the set of its eigenvalues (its point spectrum). We also use the notation $\sigma(A)$ for its spectrum and $\sigma_{\text{ess}}(A)$ for its essential spectrum. If z is an isolated point of $\sigma(A)$, then $\mathbb{1}_{\{z\}}(A)$ denotes the Riesz projection of A onto z . Similarly, if A is self-adjoint and Ξ is a Borel subset of $\sigma(A)$, then $\mathbb{1}_\Xi(A)$ denotes the spectral projection of A onto Ξ .

A (possibly unbounded) operator A on a Hilbert space \mathcal{H} is invertible (with a bounded inverse) if its null space is $\{0\}$, its range is \mathcal{H} and A^{-1} is bounded. If A is a positive invertible operator on \mathcal{H} and $s \geq 0$, then $A^{-s}\mathcal{H}$ denotes $\mathcal{D}(A^s)$ and $A^s\mathcal{H}$ denotes its (antilinear) dual. Thus, we obtain a nested scale of Hilbert spaces $A^s\mathcal{H}$, $s \in \mathbb{R}$.

In the sequel, we will usually work with the Hilbert space $L^2(\mathbb{R}_+)$ with the generic variable denoted by x or y , and sometimes also by k . It is equipped with the norm denoted by $\|\cdot\|$, the scalar product

$$(f|g) := \int_0^\infty \overline{f(x)}g(x)dx,$$

as well as the bilinear form

$$\langle f|g \rangle := \int_0^\infty f(x)g(x)dx. \tag{2.1}$$

If B is an operator on $L^2(\mathbb{R}_+)$, then B^* denotes the usual Hermitian adjoint of B , whereas $B^\#$ denotes the adjoint (the transpose) of B w.r.t. the (2.1). Clearly, if B is a bounded linear operator on $L^2(\mathbb{R}_+)$ with

$$(Bf)(k) := \int_0^\infty B(k,x)g(x)dx,$$

then

$$(B^*g)(x) = \int_0^\infty \overline{B(k,x)}g(k)dk,$$

while

$$(B^\#g)(x) = \int_0^\infty B(k,x)g(k)dk.$$

We shall use the symbol X to denote the operator of multiplication by the variable in \mathbb{R}_+ , i.e., $(Xf)(x) = xf(x)$ for $f \in \mathcal{D}(X) \subset L^2(\mathbb{R}_+)$ and $x \in \mathbb{R}_+$. Note that if it is clear that the name of the variable is x , we will also write x instead of X . We will often use the scale of Hilbert spaces based on the operator $\langle X \rangle := (1 + X^2)^{1/2}$, denoted by $\langle X \rangle^{-s}L^2(\mathbb{R}_+)$. The Sobolev spaces $\mathcal{H}_0^1(\mathbb{R}_+)$ and $\mathcal{H}^1(\mathbb{R}_+)$ are the subspaces of $L^2(\mathbb{R}_+)$ defined as the form domain of the Dirichlet and Neumann Laplacian, respectively.

We will also consider the unitary group $\{U_\tau\}_{\tau \in \mathbb{R}}$ of dilations acting on $f \in L^2(\mathbb{R}_+)$ as $(U_\tau f)(x) = e^{\tau/2}f(e^\tau x)$. An operator B on $L^2(\mathbb{R}_+)$ is said to be *homogeneous of degree* $\alpha \in \mathbb{R}$ if $U_\tau \mathcal{D}(B) \subset \mathcal{D}(B)$ for any $\tau \in \mathbb{R}$ and if the equality $U_\tau B U_\tau^{-1} = e^{\alpha\tau}B$ holds on $\mathcal{D}(B)$. For instance, X is an operator of degree 1. The generator of dilations is denoted by A , so that $U_\tau = e^{i\tau A}$. On suitable functions f , one has

$$Af(x) = \frac{1}{2i}(\partial_x x + x\partial_x)f(x).$$

The following holomorphic functions are understood as their *principal branches*, that is, their domain is $\mathbb{C} \setminus]-\infty, 0]$ and on $]0, \infty[$, they coincide with their usual definitions from real analysis: $\ln(z)$, \sqrt{z} , z^λ , $\arg(z) := \text{Im}(\ln(z))$. $\text{Ln}(z)$ will denote the multivalued logarithm. This means, if w_0 satisfies $e^{w_0} = z$, then $\text{Ln}(z)$ is the set $w_0 + 2\pi i\mathbb{Z}$.

The Wronskian $W(f, g)$ of two continuously differentiable functions f, g on \mathbb{R}_+ is given by the expression

$$W_x(f, g) \equiv W(f, g)(x) := f(x)g'(x) - f'(x)g(x). \tag{2.2}$$

2.2. Maximal and Minimal Homogenous Schrödinger Operators

For any $\alpha \in \mathbb{C}$, we consider the differential expression

$$L_\alpha := -\partial_x^2 + \left(\alpha - \frac{1}{4}\right) \frac{1}{x^2}$$

acting on distributions on \mathbb{R}_+ , and denote by L_α^{\min} and L_α^{\max} , the corresponding minimal and maximal operators associated with it in $L^2(\mathbb{R}_+)$, see [7, Sec. 4 & App. A] for details. We simply recall from this reference that

$$\mathcal{D}(L_\alpha^{\max}) = \{f \in L^2(\mathbb{R}_+) \mid L_\alpha f \in L^2(\mathbb{R}_+)\}$$

and that $\mathcal{D}(L_\alpha^{\min})$ is the closure of the restriction of L_α to $C_c^\infty(\mathbb{R}_+)$. Clearly, both L_α^{\min} and L_α^{\max} are homogeneous of degree -2 . Note that, from now on, we shall simply say *homogeneous*, without specifying the degree -2 . In addition, the following relation holds:

$$(L_\alpha^{\min})^* = L_{\bar{\alpha}}^{\max}.$$

Let us recall some additional results which have been obtained in [7, Sec. 4]. For that purpose, we say that $f \in \mathcal{D}(L_\alpha^{\min})$ around 0, (or, by an abuse of notation, $f(x) \in \mathcal{D}(L_\alpha^{\min})$ around 0) if there exists $\zeta \in C_c^\infty([0, \infty[)$ with $\zeta = 1$ around 0, such that $f\zeta \in \mathcal{D}(L_\alpha^{\min})$. In addition, it turns out that it is useful to introduce a parameter $m \in \mathbb{C}$, such that $\alpha = m^2$, even though there are two m corresponding to a single $\alpha \neq 0$. In other words, we shall consider, from now on, the operator

$$L_{m^2} := -\partial_x^2 + \left(m^2 - \frac{1}{4}\right) \frac{1}{x^2}.$$

With this notation, if $|\operatorname{Re}(m)| \geq 1$, then $L_{m^2}^{\min} = L_{m^2}^{\max}$, while if $|\operatorname{Re}(m)| < 1$, then $L_{m^2}^{\min} \subsetneq L_{m^2}^{\max}$, and $\mathcal{D}(L_{m^2}^{\min})$ is a closed subspace of codimension 2 of $\mathcal{D}(L_{m^2}^{\max})$. More precisely, if $|\operatorname{Re}(m)| < 1$ and if $f \in \mathcal{D}(L_{m^2}^{\max})$, then there exist $a, b \in \mathbb{C}$, such that:

$$\begin{aligned} f(x) - ax^{1/2-m} - bx^{1/2+m} &\in \mathcal{D}(L_{m^2}^{\min}) \text{ around } 0 \quad \text{if } m \neq 0, \\ f(x) - ax^{1/2} \ln(x) - bx^{1/2} &\in \mathcal{D}(L_0^{\min}) \text{ around } 0. \end{aligned}$$

In addition, the behavior of any function $g \in \mathcal{D}(L_{m^2}^{\min})$ is known, namely $g \in \mathcal{H}_0^1(\mathbb{R}_+)$ and as $x \rightarrow 0$:

$$\begin{aligned} g(x) &= o(x^{3/2}) \quad \text{and} \quad g'(x) = o(x^{1/2}) \quad \text{if } m \neq 0, \\ g(x) &= o(x^{3/2} \ln(x)) \quad \text{and} \quad g'(x) = o(x^{1/2} \ln(x)) \quad \text{if } m = 0. \end{aligned}$$

2.3. Two Families of Schrödinger Operators with Inverse Square Potentials

Let us first recall from [7, Def. 4.1] that for any $m \in \mathbb{C}$ with $\operatorname{Re}(m) > -1$, the operator H_m has been defined as the restriction of $L_{m^2}^{\max}$ to the domain

$$\begin{aligned} \mathcal{D}(H_m) &= \{f \in \mathcal{D}(L_{m^2}^{\max}) \mid \text{for some } c \in \mathbb{C}, \\ &f(x) - cx^{1/2+m} \in \mathcal{D}(L_{m^2}^{\min}) \text{ around } 0\}. \end{aligned}$$

It is then proved in this reference that $\{H_m\}_{\operatorname{Re}(m) > -1}$ is a holomorphic family of closed homogeneous operators in $L^2(\mathbb{R}_+)$. In addition, if $\operatorname{Re}(m) \geq 1$, then

$$H_m = L_{m^2}^{\min} = L_{m^2}^{\max}.$$

For this reason, we shall concentrate on the case $-1 < \operatorname{Re}(m) < 1$, considering a larger family of operators.

For $|\operatorname{Re}(m)| < 1$ and for any $\kappa \in \mathbb{C} \cup \{\infty\}$, we define a family of operators $H_{m,\kappa}$:

$$\begin{aligned} \mathcal{D}(H_{m,\kappa}) = \{ & f \in \mathcal{D}(L_{m^2}^{\max}) \mid \text{for some } c \in \mathbb{C}, \\ & f(x) - c(\kappa x^{1/2-m} + x^{1/2+m}) \in \mathcal{D}(L_{m^2}^{\min}) \text{ around } 0\}, \quad \kappa \neq \infty; \end{aligned} \quad (2.3)$$

$$\begin{aligned} \mathcal{D}(H_{m,\infty}) = \{ & f \in \mathcal{D}(L_{m^2}^{\max}) \mid \text{for some } c \in \mathbb{C}, \\ & f(x) - cx^{1/2-m} \in \mathcal{D}(L_{m^2}^{\min}) \text{ around } 0\}. \end{aligned} \quad (2.4)$$

For $m = 0$, we introduce an additional family of operators H_0^ν with $\nu \in \mathbb{C} \cup \{\infty\}$:

$$\begin{aligned} \mathcal{D}(H_0^\nu) = \{ & f \in \mathcal{D}(L_0^{\max}) \mid \text{for some } c \in \mathbb{C}, \\ & f(x) - c(x^{1/2} \ln(x) + \nu x^{1/2}) \in \mathcal{D}(L_0^{\min}) \text{ around } 0\}, \quad \nu \neq \infty; \end{aligned} \quad (2.5)$$

$$\begin{aligned} \mathcal{D}(H_0^\infty) = \{ & f \in \mathcal{D}(L_0^{\max}) \mid \text{for some } c \in \mathbb{C}, \\ & f(x) - cx^{1/2} \in \mathcal{D}(L_0^{\min}) \text{ around } 0\}. \end{aligned} \quad (2.6)$$

The following properties of these families of operators are immediate:

Lemma 2.1. (i) *For any $|\operatorname{Re}(m)| < 1$ and any $\kappa \in \mathbb{C} \cup \{\infty\}$,*

$$H_{m,\kappa} = H_{-m,\kappa^{-1}}. \quad (2.7)$$

(ii) *The operator $H_{0,\kappa}$ does not depend on κ , and all these operators coincide with H_0^∞ .*

As a consequence of (ii), all the results about the case $m = 0$ will be formulated in terms of the family H_0^ν .

Let us now derive two simple results for this family of operators. The first one is related to the action of the dilation group, while the second is dealing with the Hermitian conjugation.

Proposition 2.2. *For any m with $|\operatorname{Re}(m)| < 1$ and any $\kappa, \nu \in \mathbb{C} \cup \{\infty\}$, we have*

$$\begin{aligned} U_\tau H_{m,\kappa} U_{-\tau} &= e^{-2\tau} H_{m, e^{-2\tau} \kappa}, \\ U_\tau H_0^\nu U_{-\tau} &= e^{-2\tau} H_0^{\nu+\tau}, \end{aligned}$$

with the convention that $\alpha\infty = \infty$ for any $\alpha \in \mathbb{C} \setminus \{0\}$ and $\infty + \tau = \infty$ for any $\tau \in \mathbb{C}$. In particular,

(i) Amongst the family of operators $H_{m,\kappa}$ with $m \neq 0$, only

$$H_{m,0} = H_m \quad \text{and} \quad H_{m,\infty} = H_{-m}$$

are homogeneous,

(ii) Amongst the family H_0^ν , only

$$H_0^\infty = H_0$$

is homogeneous.

Proof. If $f \in \mathcal{D}(H_{m,\kappa})$, then $U_\tau f \in \mathcal{D}(H_{m,e^{-2m\tau\kappa}})$. Thus, the only domains which are left invariant are $\mathcal{D}(H_{m,0})$ and $\mathcal{D}(H_{m,\infty})$. Since $L_{m^2}^{\max}$ is homogeneous, the same applies for $H_{m,0}$ and $H_{m,\infty}$.

If $f \in \mathcal{D}(H_0^\nu)$, then $U_\tau f \in \mathcal{D}(H_0^{\nu+\tau})$. Thus, only $\mathcal{D}(H_0^\infty)$ is left invariant, and consequently, only H_0^∞ is homogeneous. \square

Proposition 2.3. For any $m \in \mathbb{C}$ with $|\operatorname{Re}(m)| < 1$ and for any $\kappa, \nu \in \mathbb{C} \cup \{\infty\}$, one has

$$(H_{m,\kappa})^* = H_{\bar{m},\bar{\kappa}} \quad \text{and} \quad (H_0^\nu)^* = H_0^{\bar{\nu}} \tag{2.8}$$

with the convention that $\bar{\infty} = \infty$.

Proof. Recall from [7, App. A] that for any $f \in \mathcal{D}(L_{m^2}^{\max})$ and $g \in \mathcal{D}(L_{\bar{m}^2}^{\max})$, the functions f, f', g, g' are continuous on \mathbb{R}_+ , and that the equality

$$(L_{m^2}^{\max} f|g) - (f|L_{\bar{m}^2}^{\max} g) = -W_0(\bar{f}, g)$$

holds with $W_0(\bar{f}, g) = \lim_{x \rightarrow 0} W_x(\bar{f}, g)$ and W_x defined in (2.2). In particular, if $f \in \mathcal{D}(H_{m,\kappa})$, one infers that

$$(H_{m,\kappa} f|g) = (f|L_{\bar{m}^2}^{\max} g) - W_0(\bar{f}, g).$$

Thus, $g \in \mathcal{D}((H_{m,\kappa})^*)$ if and only if $W_0(\bar{f}, g) = 0$, and then $(H_{m,\kappa})^* g = L_{\bar{m}^2}^{\max} g$. Then, by considering the explicit description of $\mathcal{D}(H_{m,\kappa})$, straightforward computations show that $W_0(\bar{f}, g) = 0$ if and only if $g \in \mathcal{D}(H_{\bar{m},\bar{\kappa}})$. One then deduces that $(H_{m,\kappa})^* = H_{\bar{m},\bar{\kappa}}$.

A similar computation leads to the equality $(H_0^\nu)^* = H_0^{\bar{\nu}}$. \square

Corollary 2.4. (i) The operator $H_{m,\kappa}$ is self-adjoint for $m \in]-1, 1[$ and $\kappa \in \mathbb{R} \cup \{\infty\}$, and for $m \in i\mathbb{R}$ and $|\kappa| = 1$.

(ii) The operator H_0^ν is self-adjoint for $\nu \in \mathbb{R} \cup \{\infty\}$.

Proof. For the operators $H_{m,\kappa}$, one simply has to consider formula (2.8) for the first case, and the same formula together with (2.7) in the second case. Finally, for the operators H_0^ν , considering formula (2.8) leads directly to the result. \square

Remark 2.5. By blowing up around $m = 0$, it is possible to make one single holomorphic function out of $H_{m,\kappa}$ and H_0^ν . In fact, we can extend the function H_0^ν by setting

$$H_m^\nu := H_{m, \frac{\nu m - 1}{\nu m + 1}}, \quad m \neq 0.$$

Then, H_m^ν is holomorphic for $1 < \operatorname{Re}(m) < 1$, $\nu \in \mathbb{C} \cup \{\infty\}$ and covers all values of $H_{m,\kappa}$. To see the holomorphy at $m = 0$ (at least at the level of the boundary conditions), note that for $|m|$ very small, one has

$$\kappa x^{1/2-m} + x^{1/2+m} \approx m(1-\kappa) \left(x^{1/2} \ln(x) + \frac{1+\kappa}{m(1-\kappa)} x^{1/2} \right),$$

which is $c(x^{1/2} \ln(x) + \nu x^{1/2})$ for $\kappa = \frac{\nu m - 1}{\nu m + 1}$.

3. Bessel-Type Equations and Functions

In this section, we establish the link between our initial operator (1.1) and different forms of the Bessel equation. By considering the dependence of the space dimension in these equations, one is naturally led to introduce a new basic family of Bessel-type functions.

Let us recall that the Laplace operator in d dimensions and in spherical coordinates is given by

$$-\Delta_d = -\partial_r^2 - \frac{d-1}{r} \partial_r - \frac{1}{r^2} \Delta_{\mathbb{S}^{d-1}},$$

where r is the radial coordinate and $\Delta_{\mathbb{S}^{d-1}}$ is the *Laplace–Beltrami operator on the sphere* \mathbb{S}^{d-1} . For simplicity, we also use, in this section, the notation r for the corresponding operator of multiplication by the radial coordinate. Eigenvalues of $-\Delta_{\mathbb{S}^{d-1}}$ for $d = 2, 3, \dots$ are

$$l(l+d-2), \quad l \in \{0, 1, 2, \dots\}, \quad (3.1)$$

where l corresponds to the order of spherical harmonics.

Note that in the special case $d = 2$, one has $\Delta_{\mathbb{S}^1} = \partial_\phi^2$ and the *angular momentum operator* $-i\partial_\phi$ has eigenvalues $m \in \mathbb{Z}$. As a consequence, (3.1) can also be written as

$$m^2, \quad m \in \mathbb{Z}.$$

For $d = 1$, the sphere reduces to a pair of points and $\Delta_{\mathbb{S}^0}$ corresponds to the **0**-operator. It has the eigenvalue 0 of multiplicity 2, and this corresponds to the values $l = 0$ and $l = 1$ in (3.1).

Thus, if one sets $m := l + \frac{d}{2} - 1$, then the radial part of the Laplacian takes the following form in any dimension:

$$-\partial_r^2 - \frac{d-1}{r} \partial_r + l(l+d-2) \frac{1}{r^2} = -\partial_r^2 - \frac{d-1}{r} \partial_r + \left(m^2 - \left(\frac{d}{2} - 1 \right)^2 \right) \frac{1}{r^2} \quad (3.2)$$

with an appropriately restricted range of m . In what follows, we will call (3.2) the *Bessel operator of dimension* d , allowing then the parameter m to take arbitrary complex values. In particular, the Bessel operator of dimension 2 is

$$-\partial_r^2 - \frac{1}{r} \partial_r + \frac{m^2}{r^2}, \quad (3.3)$$

while the Bessel operator of dimension 1 is

$$-\partial_r^2 + \left(m^2 - \frac{1}{4}\right) \frac{1}{r^2}. \tag{3.4}$$

Operators (3.2) are related to one another for different d by a simple similarity transformation. Indeed, by a short computation performed on $C_c^\infty(\mathbb{R}_+)$, one easily observes that the following equality holds:

$$\begin{aligned} &-\partial_r^2 - \frac{d-1}{r} \partial_r + \left(m^2 - \left(\frac{d}{2} - 1\right)^2\right) \frac{1}{r^2} \\ &= r^{-\frac{d}{2}+1} \left(-\partial_r^2 - \frac{1}{r} \partial_r + \frac{m^2}{r^2}\right) r^{\frac{d}{2}-1} \\ &= r^{-\frac{d}{2}+\frac{1}{2}} \left(-\partial_r^2 + \left(m^2 - \frac{1}{4}\right) \frac{1}{r^2}\right) r^{\frac{d}{2}-\frac{1}{2}}. \end{aligned}$$

It is then a matter of taste, convenience, and historical circumstances whether the operator (3.3) or (3.4) is taken as the basic one. In the literature, at least since the times of Bessel, it seems that (3.3) has a distinguished status. We prefer (3.4) at least in the context of this paper, since our initial operator (1.1) has a form similar to (3.4).

A simple scaling argument shows that the eigenvalue problem for (3.3) can be reduced to one of the following two equations:

the modified Bessel equation $\left(\partial_r^2 + \frac{1}{r} \partial_r - \frac{m^2}{r^2} - 1\right) v = 0,$ (3.5)

the (standard) Bessel equation $\left(\partial_r^2 + \frac{1}{r} \partial_r - \frac{m^2}{r^2} + 1\right) v = 0.$ (3.6)

Certain distinguished solutions of (3.5) are

- the modified Bessel function $I_m(z),$
- the MacDonald function $K_m(z),$

and of (3.6) are

- the Bessel function $J_m(z),$
- the Hankel function of the 1st kind $H_m^+(z) = H_n^{(1)}(z),$
- the Hankel function of the 2nd kind $H_m^-(z) = H_n^{(2)}(z),$
- the Neumann function $Y_m(z).$

We call them jointly *the Bessel family*. They are probably the best known and the most widely used special functions in mathematics and its applications [1, 4, 18, 34].

Remark 3.1. The notation $H_m^{(1)}, H_m^{(2)}$ for the two kinds of Hankel functions is more common in the literature. We use the notation H_m^+, H_m^- , which is much more convenient.

Instead of (3.5) and (3.6), we will prefer to consider their analogs for dimension 1, namely

$$\left(\partial_r^2 - \left(m^2 - \frac{1}{4}\right)\frac{1}{r^2} - 1\right)v = 0, \quad (3.7)$$

$$\left(\partial_r^2 - \left(m^2 - \frac{1}{4}\right)\frac{1}{r^2} + 1\right)v = 0. \quad (3.8)$$

We will also introduce new special functions that solve (3.7)

the modified Bessel function for dim. 1 $\mathcal{I}_m(r) := \sqrt{\frac{\pi r}{2}} I_m(r),$

the MacDonald function for dim. 1 $\mathcal{K}_m(r) := \sqrt{\frac{2r}{\pi}} K_m(r),$

and new special functions that solve (3.8)

the Bessel function for dim. 1 $\mathcal{J}_m(r) := \sqrt{\frac{\pi r}{2}} J_m(r),$

the Hankel function of the 1st kind for dim. 1 $\mathcal{H}_m^+(r) := \sqrt{\frac{\pi r}{2}} H_m^+(r),$

the Hankel function of the 2nd kind for dim. 1 $\mathcal{H}_m^-(r) := \sqrt{\frac{\pi r}{2}} H_m^-(r),$

the Neumann function for dim. 1 $\mathcal{Y}_m(r) := \sqrt{\frac{\pi r}{2}} Y_m(r).$

Jointly, they will be called *the Bessel family for dimension 1*. Accordingly, the Bessel family should be called *the Bessel family for dimension 2*. As we shall show later on, the Bessel family for dimension 1 contains a number of standard elementary functions: the exponential function, the trigonometric sine and cosine functions, and the hyperbolic sine and cosine functions.

Obviously, properties of the Bessel family for dimension 1 can easily be deduced from the corresponding properties of the Bessel family for dimension 2, and the other way round. Most properties seem to have a simpler form in the case of dimension 1 than in the case of dimension 2. There are some exceptions, mostly involving integer values of m which play a distinguished role in dimension 2 and, more generally, in even dimensions. We collect basic properties of the Bessel family for dimension 1 in Appendix A.

Let us note that in the literature, one sometimes introduces the modified and standard Bessel equations for all dimensions $d \geq 2$, as well as the corresponding functions, see [26]. In particular, there exists a standard notation for the functions of the Bessel family for dimension 3: i_m, k_m, j_m, h_m^\pm , and y_m . In our opinion, however, this introduces too many unnecessary special functions: the main issue is whether the dimension is even or odd. One could argue that the Bessel family for dimension 2, that is I_m, K_m, J_m, H_m^\pm , and Y_m , is better adapted for all even dimensions and integer values of m , whereas the Bessel family for dimension 1, that is $\mathcal{I}_m, \mathcal{K}_m, \mathcal{J}_m, \mathcal{H}_m^\pm$, and \mathcal{Y}_m , is better adapted for odd dimensions, as well as for general values of m .

4. The Homogeneous Case

In this section, we consider the homogeneous operators H_m . Part of the following results was proved in [7], but we add some new material. In particular, some statements and proofs in [7] were restricted to the self-adjoint case. We extend them to all $\operatorname{Re}(m) > -1$. In addition, we shall now express everything in terms of the Bessel family for dimension 1 instead of the Bessel family for dimension 2.

4.1. Resolvent

In [7, Sec. 4.2], the resolvent of H_m is constructed. For completeness, we recall its construction below:

Theorem 4.1. *For any $m \in \mathbb{C}$ with $\operatorname{Re}(m) > -1$, the spectrum of H_m is $[0, \infty[$. In addition, for $k \in \mathbb{C}$ with $\operatorname{Re}(k) > 0$, the resolvent*

$$R_m(-k^2) := (H_m + k^2)^{-1}$$

has the kernel

$$R_m(-k^2; x, y) = \frac{1}{k} \begin{cases} \mathcal{I}_m(kx) \mathcal{K}_m(ky) & \text{if } 0 < x < y, \\ \mathcal{I}_m(ky) \mathcal{K}_m(kx) & \text{if } 0 < y < x. \end{cases} \quad (4.1)$$

Sketch of proof provided in [7]. It is first checked that the kernel provided by (4.1) defines a bounded operator which we denote by $R_m(-k^2)$. Then, it is verified that

$$\left((L_{m^2} + k^2) R_m(-k^2) \right) (x, y) = \delta(x - y).$$

Next, one has $R_m(-k^2)f \in \mathcal{D}(H_m)$ for any $f \in C_c^\infty(\mathbb{R}_+)$. Thus, the previous equality can be reinterpreted as $(H_m + k^2)R_m(-k^2) = \mathbb{1}$. In addition, since $H_m^* = H_{\bar{m}}$ and $R_m(-k^2)^* = R_{\bar{m}}(-\bar{k}^2)$, one then infers that

$$R_m(-k^2)(H_m + k^2) = \left((H_{\bar{m}} + \bar{k}^2) R_{\bar{m}}(-\bar{k}^2) \right)^* = \mathbb{1}.$$

Therefore, $-k^2$ belongs to the resolvent set of H_m and $R_m(-k^2)$ is the resolvent of H_m . \square

For $|\operatorname{Re}(m)| < 1$, we also introduce the operator $P_m(-k^2)$ defined by its kernel

$$P_m(-k^2; x, y) := \frac{\sin(\pi m)}{m} k \mathcal{K}_m(kx) \mathcal{K}_m(ky) \quad \text{if } m \neq 0, \quad (4.2)$$

$$P_0(-k^2; x, y) := \pi k \mathcal{K}_0(kx) \mathcal{K}_0(ky). \quad (4.3)$$

By the bounds (A.7) and (A.8), the function $x \mapsto \mathcal{K}_m(kx)$ is square integrable. Considering (A.16) and (A.17), one easily infers that the operator $P_m(-k^2)$ is a rank one projection. It is orthogonal if k and m^2 are real. By considering the equality $\mathcal{K}_m = \mathcal{K}_{-m}$ and (A.4), one can also deduce the following relation for $|\operatorname{Re}(m)| < 1$:

$$R_{-m}(-k^2) - R_m(-k^2) = \frac{m}{k^2} P_m(-k^2). \quad (4.4)$$

4.2. Boundary Value of the Resolvent and Spectral Density

In this section, we show that a Limiting Absorption Principle holds for the operators H_m . We also compute the kernels of the boundary values of the resolvent and of the spectral density.

Theorem 4.2. *Let $m \in \mathbb{C}$ with $\operatorname{Re}(m) > -1$, and let $k > 0$. Then, the boundary values of the resolvent*

$$R_m(k^2 \pm i0) := \lim_{\epsilon \searrow 0} R_m(k^2 \pm i\epsilon)$$

exist in the sense of operators from $\langle X \rangle^{-s} L^2(\mathbb{R}_+)$ to $\langle X \rangle^s L^2(\mathbb{R}_+)$ for any $s > \frac{1}{2}$, uniformly in k on each compact subset of \mathbb{R}_+ . They have the kernels

$$R_m(k^2 \pm i0; x, y) = \pm \frac{i}{k} \begin{cases} \mathcal{J}_m(kx) \mathcal{H}_m^\pm(ky) & \text{if } 0 < x \leq y, \\ \mathcal{J}_m(ky) \mathcal{H}_m^\pm(kx) & \text{if } 0 < y < x. \end{cases} \quad (4.5)$$

The above theorem describes a property that, at least in the context of self-adjoint Schrödinger operators, is often called the *Limiting Absorption Principle*. Its proof will be based on an explicit estimate of the resolvent kernel:

Proposition 4.3. *Let us consider $\operatorname{Re}(k) > 0$. Then, for $\operatorname{Re}(m) \geq 0$ with $m \neq 0$, one has*

$$|R_m(-k^2; x, y)| \leq \frac{C_m^2}{|k|} e^{-\operatorname{Re}(k)|x-y|} \min(1, |xk|)^{\frac{1}{2}} \min(1, |yk|)^{\frac{1}{2}}, \quad (4.6)$$

for $-1 < \operatorname{Re}(m) \leq 0$ with $m \neq 0$, one has

$$|R_m(-k^2; x, y)| \leq \frac{C_m^2}{|k|} e^{-\operatorname{Re}(k)|x-y|} \min(1, |xk|)^{\operatorname{Re}(m)+\frac{1}{2}} \min(1, |yk|)^{\operatorname{Re}(m)+\frac{1}{2}}, \quad (4.7)$$

while, in the special case $m = 0$, one has

$$|R_0(-k^2; x, y)| \leq \frac{C_0^2}{|k|} e^{-\operatorname{Re}(k)|x-y|} \min(1, |xk|)^{\frac{1}{2}} \min(1, |yk|)^{\frac{1}{2}} \times (1 + |\ln(\min(1, |kx|))|) (1 + |\ln(\min(1, |ky|))|). \quad (4.8)$$

The constants C_m and C_0 are independent of x, y , and k .

Proof. The proof is based on the following estimates on the Bessel and Macdonald functions. For $\epsilon > 0$ and $|\arg z| < \pi - \epsilon$, one has

$$\begin{aligned} |\mathcal{K}_m(z)| &\leq C_m e^{-\operatorname{Re}(z)} \min(1, |z|)^{-|\operatorname{Re}(m)|+\frac{1}{2}}, \quad m \neq 0, \\ |\mathcal{K}_0(z)| &\leq C_0 e^{-\operatorname{Re}(z)} \min(1, |z|)^{\frac{1}{2}} (1 + |\ln(\min(1, |z|))|), \end{aligned} \quad (4.9)$$

and

$$\begin{aligned} |\mathcal{I}_m(z)| &\leq C_m e^{|\operatorname{Re}(z)|} \min(1, |z|)^{\operatorname{Re}(m)+\frac{1}{2}}, \quad m \neq 0, \\ |\mathcal{I}_0(z)| &\leq C_0 e^{|\operatorname{Re}(z)|} \min(1, |z|)^{\frac{1}{2}}. \end{aligned}$$

Using $\operatorname{Re}(k) > 0$ and for $m \neq 0$, observe first that for $0 < x < y$, we obtain

$$|R_m(-k^2; x, y)| \leq \frac{C_m^2}{|k|} e^{x\operatorname{Re}(k)} e^{-y\operatorname{Re}(k)} \min(1, |kx|)^{\operatorname{Re}(m)+\frac{1}{2}} \min(1, |ky|)^{-|\operatorname{Re}(m)|+\frac{1}{2}},$$

while for $0 < y < x$, we have

$$|R_m(-k^2; x, y)| \leq \frac{C_m^2}{|k|} e^{y\operatorname{Re}(k)} e^{-x\operatorname{Re}(k)} \min(1, |kx|)^{-|\operatorname{Re}(m)|+\frac{1}{2}} \min(1, |ky|)^{\operatorname{Re}(m)+\frac{1}{2}}.$$

If $\operatorname{Re}(m) \geq 0$, one observes then that $\frac{|kx|}{|ky|} < 1$ in the first case, while $\frac{|ky|}{|kx|} < 1$ in the second case. This directly leads to (4.6). On the other hand, for $\operatorname{Re}(m) < 0$, one has $-|\operatorname{Re}(m)| = \operatorname{Re}(m)$, from which one infers (4.7). Finally, the special case $m = 0$ is obtained by a straightforward computation. \square

Proof of Theorem 4.2. Define the operator $R_m(k^2 \pm i0)$ by its kernel (4.5). We will show that

$$\langle X \rangle^{-s} R_m(k^2 \pm i\epsilon) \langle X \rangle^{-s}, \tag{4.10}$$

whose kernel is

$$\langle x \rangle^{-s} R_m(k^2 \pm i\epsilon; x, y) \langle y \rangle^{-s}, \tag{4.11}$$

is a Hilbert–Schmidt operator and converges as $\epsilon \searrow 0$ in the Hilbert–Schmidt norm to

$$\langle X \rangle^{-s} R_m(k^2 \pm i0) \langle X \rangle^{-s}. \tag{4.12}$$

Consider first the (slightly more difficult) case $-1 < \operatorname{Re}(m) \leq 0$ with $m \neq 0$. By the estimate (4.7), the expression (4.11) can be bounded by

$$\begin{aligned} & \frac{C}{|k|} e^{-\operatorname{Re}(\sqrt{-k^2 \mp i\epsilon})|x-y|} \langle x \rangle^{-s} \langle y \rangle^{-s} \min(1, |xk|)^{\operatorname{Re}(m)+\frac{1}{2}} \min(1, |yk|)^{\operatorname{Re}(m)+\frac{1}{2}} \\ & \leq \frac{C}{|k|} \langle x \rangle^{-s} \langle y \rangle^{-s} \min(1, |xk|)^{\operatorname{Re}(m)+\frac{1}{2}} \min(1, |yk|)^{\operatorname{Re}(m)+\frac{1}{2}}, \end{aligned} \tag{4.13}$$

where C is a constant independent of x, y , and k . Note that in the computation, the inequality $\operatorname{Re}(\sqrt{-k^2 \mp i\epsilon}) \geq 0$ has been used, and that such an inequality holds by our choice of the principal branch of the square root. One clearly infers that (4.13) belongs to $L^2(\mathbb{R}_+ \times \mathbb{R}_+)$ and dominates (4.11). Since (4.11) converges pointwise to

$$\langle x \rangle^{-s} R_m(k^2 \pm i0; x, y) \langle y \rangle^{-s}, \tag{4.14}$$

one concludes by the Lebesgue Dominated Convergence Theorem that (4.11) converges in $L^2(\mathbb{R}_+ \times \mathbb{R}_+)$ to (4.14). This is equivalent to the convergence of (4.10) to (4.12) in the Hilbert–Schmidt norm. Note finally that the uniform convergence in k on each compact subset of \mathbb{R}_+ can be checked directly on the above expressions.

For $\operatorname{Re}(m) \geq 0$ with $m \neq 0$, the same proof holds with the estimate (4.6) instead of (4.7). Finally, for $m = 0$, the result can be obtained using (4.8), and by observing that the factor with the logarithm is also square integrable near the origin. \square

Based on the previous theorem, one can directly deduce the following statement.

Proposition 4.4. *There exists the spectral density*

$$\begin{aligned} p_m(k^2) &:= \lim_{\epsilon \searrow 0} \frac{1}{2\pi i} \left(R_m(k^2 + i\epsilon) - R_m(k^2 - i\epsilon) \right) \\ &= \frac{1}{2\pi i} \left(R_m(k^2 + i0) - R_m(k^2 - i0) \right), \end{aligned}$$

understood in the sense of operators from $\langle X \rangle^{-s} L^2(\mathbb{R}_+)$ to $\langle X \rangle^s L^2(\mathbb{R}_+)$ for any $s > \frac{1}{2}$. The kernel of this operator is provided for $x, y \in \mathbb{R}_+$ by

$$p_m(k^2; x, y) = \frac{\mathcal{J}_m(kx)\mathcal{J}_m(ky)}{\pi k}.$$

4.3. Hankel Transformation

This section is mostly inspired from Sects. 5 and 6 of [7], from which most of the statements are borrowed. We refer to this reference for more details, or to the subsequent sections of this paper for a more general approach.

For any $m \in \mathbb{C}$ with $\operatorname{Re}(m) > -1$ let us set

$$\mathcal{F}_m : C_c(\mathbb{R}_+) \rightarrow L^2(\mathbb{R}_+)$$

with

$$(\mathcal{F}_m f)(x) := \int_0^\infty \mathcal{F}_m(x, y) f(y) dy$$

and

$$\mathcal{F}_m(x, y) := \sqrt{\frac{2}{\pi}} \mathcal{J}_m(xy). \tag{4.15}$$

We also define the unitary and self-adjoint transformation $J : L^2(\mathbb{R}_+) \rightarrow L^2(\mathbb{R}_+)$ by the formula

$$(Jf)(x) = \frac{1}{x} f\left(\frac{1}{x}\right) \tag{4.16}$$

for any $f \in L^2(\mathbb{R}_+)$ and $x \in \mathbb{R}_+$. Finally, we recall that A denotes the generator of dilations.

Proposition 4.5. *For any $m \in \mathbb{C}$ with $\operatorname{Re}(m) > -1$, the map \mathcal{F}_m continuously extends to a bounded invertible operator on $L^2(\mathbb{R}_+)$ satisfying $\mathcal{F}_m^\# = \mathcal{F}_m = \mathcal{F}_m^{-1}$. In addition, the following equalities hold:*

$$\mathcal{F}_m = J\Xi_m(A) = \Xi_m(-A)J, \tag{4.17}$$

with

$$\Xi_m(t) := e^{i \ln(2)t} \frac{\Gamma\left(\frac{m+1+it}{2}\right)}{\Gamma\left(\frac{m+1-it}{2}\right)}. \tag{4.18}$$

Proof. Let us start by proving the first equality in (4.17). For that purpose, consider the operator $J\mathcal{F}_m : C_c(\mathbb{R}_+) \rightarrow L^2(\mathbb{R}_+)$ whose kernel is given by

$$\sqrt{\frac{2}{\pi}} \frac{1}{x} \mathcal{J}_m\left(\frac{y}{x}\right) = \frac{1}{2\pi} \frac{1}{\sqrt{xy}} \int_{-\infty}^{+\infty} \frac{\Gamma\left(\frac{m+it+1}{2}\right)}{\Gamma\left(\frac{m-it+1}{2}\right)} \left(\frac{1}{2}\right)^{-it} \frac{y^{-it}}{x^{-it}} dt, \quad (4.19)$$

where (A.23) has been used for the second equality. By considering the explicit formula for the kernel of an operator $\psi(A)$, as recalled for example in [7, Lem. 6.4], one infers that the r.h.s. of (4.19) corresponds to the kernel of an operator $\psi(A)$ with ψ provided by the expression (4.18). Then, from the density of $C_c(\mathbb{R}_+)$ in $L^2(\mathbb{R}_+)$ and since the map $t \mapsto \Xi_m(t)$ is bounded, one obtains that $J\mathcal{F}_m$ extends continuously to the bounded operator $\Xi_m(A)$. Since J is unitary and self-adjoint, one directly deduces the first equality in (4.17).

The second equality in (4.17) is a straightforward consequence of the equality

$$J e^{i\tau A} J = e^{-i\tau A}$$

which is easily checked. The equality $\mathcal{F}_m^2 = \mathbb{1}$ can now be deduced from the equalities (4.17). As a consequence, $\mathcal{F}_m^{-1} = \mathcal{F}_m$, and this provides a direct proof of the boundedness of the inverse of \mathcal{F}_m . The equality $\mathcal{F}_m^\# = \mathcal{F}_m$ is finally a direct consequence of the expression (4.15) for the kernel of \mathcal{F}_m . \square

The map \mathcal{F}_m will be called the *Hankel transformation of order m* . It provides a kind of diagonalization of the operator H_m , as shown in the next statement.

Proposition 4.6. *For any $m \in \mathbb{C}$ with $\operatorname{Re}(m) > -1$ and for any $k \in \mathbb{C}$ with $\operatorname{Re}(k) > 0$, the following equality holds:*

$$(H_m + k^2)^{-1} = \mathcal{F}_m (X^2 + k^2)^{-1} \mathcal{F}_m^\#.$$

Proof. The kernel of the operator on the r.h.s. is given by the expression

$$(\mathcal{F}_m (X^2 + k^2)^{-1} \mathcal{F}_m^\#)(x, y) = \frac{2}{\pi} \int_0^\infty \mathcal{J}_m(xp) \mathcal{J}_m(yp) \frac{1}{(p^2 + k^2)} dp.$$

By (A.19), it coincides with the kernel of $R_m(-k^2)$ provided in (4.1). \square

Proposition 4.6 is convenient technically, because it contains only bounded operators. One can rewrite it using unbounded operators as follows:

Theorem 4.7. *The operator H_m is similar to a self-adjoint operator. More precisely, the following equalities hold:*

$$H_m = \mathcal{F}_m X^2 \mathcal{F}_m^\# \quad \text{and} \quad H_m = \Xi_m(-A) X^{-2} \Xi_m(A).$$

4.4. Spectral Projections

For any $0 < a < b$, let us now consider the operator

$$\mathbb{1}_{[a,b]}(H_m) := \int_{\sqrt{a}}^{\sqrt{b}} p_m(k^2) d(k^2) = 2 \int_{\sqrt{a}}^{\sqrt{b}} p_m(k^2) k dk,$$

which is defined as a bounded operator from $\langle X \rangle^{-s} L^2(\mathbb{R}_+)$ to $\langle X \rangle^s L^2(\mathbb{R}_+)$ for any $s > \frac{1}{2}$. The kernel of this operator is given for $x, y \in \mathbb{R}_+$ by the expression

$$\mathbb{1}_{[a,b]}(H_m)(x, y) = \frac{2}{\pi} \int_{\sqrt{a}}^{\sqrt{b}} \mathcal{J}_m(kx) \mathcal{J}_m(ky) dk. \quad (4.20)$$

Clearly, by Stone's formula, for real $m > -1$, the above operator extends to the self-adjoint operator H_m onto the interval $[a, b]$. For complex m , H_m is not self-adjoint, hence strictly speaking Stone's formula is not available. However, H_m is similar to a self-adjoint operator, hence the properties of $\mathbb{1}_{[a,b]}(H_m)$ are almost the same as in the self-adjoint case.

Proposition 4.8. *For any $0 < a < b$ and any $m \in \mathbb{C}$ with $\text{Re}(m) > -1$, one has*

$$\mathbb{1}_{[a,b]}(H_m) = \mathcal{F}_m \mathbb{1}_{[a,b]}(X^2) \mathcal{F}_m^\# \quad (4.21)$$

in $\mathcal{B}(L^2(\mathbb{R}_+))$. In addition, $\mathbb{1}_{[a,b]}(H_m)$ is a projection.

Proof. Let us first compute the r.h.s. of (4.21). For that purpose, recall that $\mathcal{F}_m^\# = \mathcal{F}_m$, and then one gets for any $f \in C_c(\mathbb{R}_+)$ and $x \in \mathbb{R}_+$

$$\begin{aligned} (\mathcal{F}_m \mathbb{1}_{[a,b]}(X^2) \mathcal{F}_m^\# f)(x) &= \sqrt{\frac{2}{\pi}} \int_{\sqrt{a}}^{\sqrt{b}} \mathcal{J}_m(kx) (\mathcal{F}_m^\# f)(k) dk \\ &= \frac{2}{\pi} \int_{\sqrt{a}}^{\sqrt{b}} \mathcal{J}(kx) \left(\int_0^\infty \mathcal{J}(ky) f(y) dy \right) dk \\ &= \int_0^\infty \left(\frac{2}{\pi} \int_{\sqrt{a}}^{\sqrt{b}} \mathcal{J}_m(kx) \mathcal{J}_m(ky) dk \right) f(y) dy, \end{aligned}$$

where Fubini's theorem has been applied for the last equality. By comparing the last expression with (4.20), one directly infers the equality (4.21). Note that since the r.h.s. of (4.21) defines a bounded operator on $L^2(\mathbb{R}_+)$, this equality provides a natural continuous extension of $\mathbb{1}_{[a,b]}(H_m)$ as a bounded operator on $L^2(\mathbb{R}_+)$. Finally, since \mathcal{F}_m satisfies $\mathcal{F}_m^\# \mathcal{F}_m = \mathbb{1}$, one directly infers that $\mathbb{1}_{[a,b]}(H_m)$ is a projection. \square

4.5. Møller Operators and Scattering Operator

In this section, we describe the scattering theory for the operators H_m . Note that we also treat non-self-adjoint operators, and therefore, we cannot always invoke standard results.

Let us start by introducing the *incoming* and *outgoing Hankel transformations of order m* defined by

$$\mathcal{F}_m^\mp := e^{\mp i \frac{\pi}{2} m} \mathcal{F}_m. \quad (4.22)$$

Their kernel is provided by the expressions

$$\mathcal{F}_m^\mp(x, y) = e^{\mp i \frac{\pi}{2} m} \sqrt{\frac{2}{\pi}} \mathcal{J}_m(xy)$$

and the following relations trivially hold:

$$(\mathcal{F}_m^\mp)^{-1} = \mathcal{F}_m^\pm = \mathcal{F}_m^{\pm\#}.$$

Now, for any $m, m' \in \mathbb{C}$ with $\operatorname{Re}(m), \operatorname{Re}(m') > -1$, we define

$$W_{m,m'}^\pm := \mathcal{F}_m^\pm \mathcal{F}_{m'}^{\mp\#}. \tag{4.23}$$

Some easy properties of these operators are gathered in the next statement.

Proposition 4.9. *$W_{m,m'}^\pm$ are bounded invertible operators satisfying*

$$\begin{aligned} W_{m,m'}^{\mp\#} W_{m,m'}^\pm &= \mathbb{1}, \\ W_{m,m'}^\pm W_{m,m'}^{\mp\#} &= \mathbb{1}, \\ W_{m,m'}^{\pm\#} &= W_{m',m}^\mp, \\ W_{m,m'}^\pm H_{m'} &= H_m W_{m,m'}^\pm. \end{aligned}$$

Formally, the kernel of $W_{m,m'}^\pm$ is given by

$$W_{m,m'}^\pm(x, y) = e^{\pm i \frac{\pi}{2}(m-m')} \frac{2}{\pi} \int_0^\infty \mathcal{J}_m(kx) \mathcal{J}_{m'}(ky) dk.$$

On the other hand, using the expression derived in Sect. 4.3, one also gets

$$W_{m,m'}^\pm = e^{\pm i \frac{\pi}{2}(m-m')} \frac{\Gamma(\frac{m+1-iA}{2}) \Gamma(\frac{m'+1+iA}{2})}{\Gamma(\frac{m+1+iA}{2}) \Gamma(\frac{m'+1-iA}{2})}. \tag{4.24}$$

Note also that the equality $(W_{m,m'}^\pm)^{-1} = W_{m,m'}^{\mp\#}$ holds. Therefore, the scattering operator is defined by

$$S_{m,m'} := W_{m,m'}^{-\#} W_{m,m'}^-$$

and for the operators considered above, one simply gets

$$S_{m,m'} = e^{-i\pi(m-m')} \mathbb{1}.$$

Let us now make a link with the traditional approach of scattering theory. By Proposition 4.6 and the boundedness of the Hankel transformations, the operators H_m generate a bounded one-parameter group by the formula:

$$e^{itH_m} = \mathcal{F}_m e^{itX^2} \mathcal{F}_m^\#.$$

Therefore, we can try to apply time-dependent scattering theory, even though the operators H_m may be non-self-adjoint. (We refer to [10, 21, 32] for additional information on scattering theory in the non-self-adjoint setting). In the next statement, we show that $W_{m,m'}^\pm$ coincide with the Møller operators for the pair $(H_m, H_{m'})$.

Theorem 4.10. *For any $m, m' \in \mathbb{C}$ with $\operatorname{Re}(m), \operatorname{Re}(m') > -1$, one has*

$$s\text{-}\lim_{t \rightarrow \pm\infty} e^{itH_m} e^{-itH_{m'}} = W_{m,m'}^\pm.$$

In the following proof, $C([-\infty, \infty])$ denotes the set of continuous functions on \mathbb{R} having a limit at $+\infty$ and a limit at $-\infty$.

Proof. Consider first the equality

$$e^{itH_m} e^{-itH_{m'}} = \mathcal{F}_m e^{itX^2} \mathcal{F}_m^\# \mathcal{F}_{m'} e^{-itX^2} \mathcal{F}_{m'}^\#$$

and observe that

$$\mathcal{F}_m^\# \mathcal{F}_{m'} = \Xi_m(-A) \Xi_{m'}(A) = \frac{\Gamma\left(\frac{m+1-iA}{2}\right) \Gamma\left(\frac{m'+1+iA}{2}\right)}{\Gamma\left(\frac{m+1+iA}{2}\right) \Gamma\left(\frac{m'+1-iA}{2}\right)}.$$

By considering then, the asymptotic behavior of the Γ -function, as presented, for example, in [1, Eq. 6.1.37] or in Lemma A.2, one infers that the map

$$t \mapsto \Xi_m(-t) \Xi_{m'}(t)$$

belongs to $C([-\infty, \infty])$ and that

$$\Xi_m(\mp\infty) \Xi_{m'}(\pm\infty) = e^{\mp i \frac{\pi}{2} (m-m')}. \quad (4.25)$$

One infers then by Lemma B.1 that

$$s\text{-}\lim_{t \rightarrow \pm\infty} e^{itX^2} \mathcal{F}_m^\# \mathcal{F}_{m'} e^{-itX^2} = e^{\pm i \frac{\pi}{2} (m-m')},$$

from which one deduces that

$$\begin{aligned} s\text{-}\lim_{t \rightarrow \pm\infty} e^{itH_m} e^{-itH_{m'}} &= \mathcal{F}_m e^{\pm i \frac{\pi}{2} (m-m')} \mathcal{F}_{m'}^\# \\ &= \mathcal{F}_m^\pm \mathcal{F}_{m'}^\mp. \end{aligned}$$

□

4.6. Some Special Cases

In some special situations, the scattering theory for H_m is very explicit. The Gamma function is not even used in these special cases. Some of them are provided in this section.

Proposition 4.11. *The following identities hold:*

$$\begin{aligned} \mathcal{F}_{-m} &= \mathcal{F}_m \frac{\cos\left(\frac{\pi}{2}(m+iA)\right)}{\cos\left(\frac{\pi}{2}(m-iA)\right)} = \frac{\cos\left(\frac{\pi}{2}(m-iA)\right)}{\cos\left(\frac{\pi}{2}(m+iA)\right)} \mathcal{F}_m, \\ \mathcal{F}_{-m}^\mp &= \mathcal{F}_m^\mp \frac{e^{\pm\pi A} + e^{\pm i\pi m}}{e^{\pm\pi A} + e^{\mp i\pi m}} = \frac{e^{\mp\pi A} + e^{\pm i\pi m}}{e^{\mp\pi A} + e^{\mp i\pi m}} \mathcal{F}_m^\mp, \\ W_{-m,m}^\pm &= \frac{e^{\pm\pi A} + e^{\mp i\pi m}}{e^{\pm\pi A} + e^{\pm i\pi m}}, \\ S_{-m,m} &= e^{i2\pi m} \mathbb{1}. \end{aligned} \quad (4.26)$$

Proof. Using the identity

$$\Gamma\left(z + \frac{1}{2}\right) \Gamma\left(-z + \frac{1}{2}\right) = \frac{\pi}{\cos(\pi z)},$$

one infers that

$$\Xi_{-m}(t) = \Xi_m(t) \frac{\cos\left(\frac{\pi}{2}(m+it)\right)}{\cos\left(\frac{\pi}{2}(m-it)\right)}.$$

All the mentioned equalities can then be easily deduced. □

Proposition 4.12. *The following identities hold:*

$$\begin{aligned} \mathcal{F}_{m+2} &= \mathcal{F}_m \frac{m+1+iA}{m+1-iA} = \frac{m+1-iA}{m+1+iA} \mathcal{F}_m, \\ \mathcal{F}_{m+2}^\pm &= \mathcal{F}_m^\pm \frac{m+1+iA}{m+1-iA} = \frac{m+1-iA}{m+1+iA} \mathcal{F}_m^\pm, \\ W_{m+2,m}^\pm &= \frac{m+1-iA}{m+1+iA}, \\ S_{m+2,m} &= \mathbb{1}. \end{aligned}$$

Proof. Using the identity $\Gamma(z+1) = z\Gamma(z)$, one infers that

$$\Xi_{m+2}(t) = \Xi_m(t) \frac{m+1+it}{m+1-it},$$

from which all the equalities can be easily deduced. □

4.7. Dirichlet and Neumann Laplacians on the Half-Line

The simplest cases of operators H_m are obtained for $m = \pm \frac{1}{2}$. They correspond to the Neumann and the Dirichlet boundary conditions. We denote them with the usual notation, that is

$$H_N := H_{-\frac{1}{2}}, \quad H_D := H_{\frac{1}{2}},$$

and recall some of their properties.

On $L^2(\mathbb{R}_+)$, we define the cosine and sine transformation:

$$\begin{aligned} (\mathcal{F}_N f)(x) &:= \sqrt{\frac{2}{\pi}} \int_0^\infty \cos(xy) f(y) \, dy, \\ (\mathcal{F}_D f)(x) &:= \sqrt{\frac{2}{\pi}} \int_0^\infty \sin(xy) f(y) \, dy. \end{aligned}$$

Note that the transformations \mathcal{F}_N and \mathcal{F}_D are involutive, real, and unitary, and that they correspond to the Hankel transforms, namely

$$\mathcal{F}_{-\frac{1}{2}} = \mathcal{F}_N, \quad \mathcal{F}_{\frac{1}{2}} = \mathcal{F}_D.$$

In addition, it is well known, and also a consequence of our previous computations, that these transformations diagonalize H_N and H_D :

$$\mathcal{F}_N H_N \mathcal{F}_N = X^2, \quad \mathcal{F}_D H_D \mathcal{F}_D = X^2.$$

Let us also recall the resolvents of these operators, their boundary values and their spectral densities:

$$\begin{aligned} R_N(-k^2; x, y) &= \frac{1}{2k} \left(e^{-k|x-y|} + e^{-k(x+y)} \right), \\ R_N(k^2 \pm i0; x, y) &= \pm \frac{i}{k} \begin{cases} \cos(kx) e^{\pm ik y} & \text{if } 0 < x \leq y \\ \cos(ky) e^{\pm ik x} & \text{if } 0 < y < x \end{cases}, \\ p_N(k^2; x, y) &= \frac{1}{\pi k} \cos(kx) \cos(ky), \end{aligned}$$

and

$$R_D(-k^2; x, y) = \frac{1}{2k} \left(e^{-k|x-y|} - e^{-k(x+y)} \right),$$

$$R_D(k^2 \pm i0; x, y) = \frac{1}{k} \begin{cases} \sin(kx) e^{\pm iky} & \text{if } 0 < x \leq y \\ \sin(ky) e^{\pm ikx} & \text{if } 0 < y < x \end{cases},$$

$$p_D(k^2; x, y) = \frac{1}{\pi k} \sin(kx) \sin(ky).$$

The incoming and outgoing Hankel transforms in this case differ from the regular ones by a phase factor closely related to *the Maslov correction*, famous in the late 70's (see the introduction to [6]). According to our definition given in (4.22), one has

$$\mathcal{F}_N^\mp = e^{\pm i\frac{\pi}{4}} \mathcal{F}_N, \quad \mathcal{F}_D^\mp = e^{\mp i\frac{\pi}{4}} \mathcal{F}_D. \quad (4.27)$$

Proposition 4.13. *The wave operators for the pair (H_N, H_D) are given by*

$$W_{ND}^\pm = e^{\mp i\frac{\pi}{2}} \mathcal{F}_N \mathcal{F}_D \quad (4.28)$$

$$= \pm \tanh(\pi A) \mp i \cosh(\pi A)^{-1}. \quad (4.29)$$

Its kernel is

$$W_{ND}^\pm(x, y) = \mp \frac{2i}{\pi} \int_0^\infty \cos(ky) \sin(kx) dk. \quad (4.30)$$

The corresponding scattering operator is simply given by $S_{N,D} = -\mathbb{1}$.

Proof. The equality (4.28) directly follows from Proposition 4.11 together with the explicit formula (4.27). The kernel (4.30) is then a direct consequence of (4.28).

For (4.29), recall that $W_{ND}^\pm = W_{-\frac{1}{2}, \frac{1}{2}}^\pm$, and from (4.26), one infers that

$$W_{-\frac{1}{2}, \frac{1}{2}}^\pm = \pm \frac{e^{\pi A} - i}{e^{\pi A} + i} = \pm \tanh(\pi A) \mp i \cosh(\pi A)^{-1},$$

where some identities involving the hyperbolic functions have been used. \square

Note that with a different approach, the expressions of Proposition 4.13 were already obtained in [31, Sec. 2].

5. Point Spectrum

In this section, we return to the study of the operators $H_{m,\kappa}$ and H_0^ν , and describe their point spectra.

5.1. Eigenvalues

In the first statement, we fix $z \in \mathbb{C}$ and then look for operators from our families which have the eigenvalue z . Obviously, since most of the operators are not self-adjoint, we consider arbitrary complex eigenvalues. Recall also that the notations $z \mapsto \ln(z)$ and $z \mapsto z^m$ are used for the principal branch of both functions, with domain $\mathbb{C} \setminus]-\infty, 0]$. Recall that γ denotes Euler's constant.

Proposition 5.1. *Let $m \in \mathbb{C}^\times$ with $|\operatorname{Re}(m)| < 1$ and let $\kappa, \nu \in \mathbb{C} \cup \{\infty\}$. Then, one has*

(i) $z \in \sigma_p(H_{m,\kappa})$ if and only if $z \in \mathbb{C} \setminus [0, \infty[$ and

$$\kappa = \frac{\Gamma(m)}{\Gamma(-m)} (-z/4)^{-m}. \tag{5.1}$$

(ii) $z \in \sigma_p(H_0^\nu)$ if and only if $z \in \mathbb{C} \setminus [0, \infty[$ and

$$\nu = \gamma + \frac{1}{2} \ln(-z/4).$$

Before providing the proof of this proposition, we deduce from it the main result about the point spectrum for our operators. For that purpose, let us also introduce for $m \neq 0$ the new parameter

$$\varsigma(m, \kappa) = \varsigma := \kappa \frac{\Gamma(-m)}{\Gamma(m)}. \tag{5.2}$$

Theorem 5.2. *Let $m \in \mathbb{C}$ with $|\operatorname{Re}(m)| < 1$.*

(i) *For $m \neq 0$ and $\kappa \in \mathbb{C}^\times$, one has*

$$\sigma_p(H_{m,\kappa}) = \left\{ -4e^{-w} \mid w \in \frac{1}{m} \operatorname{Ln}(\varsigma) \text{ and } -\pi < \operatorname{Im}(w) < \pi \right\}. \tag{5.3}$$

(ii) *For any $\nu \in \mathbb{C}$, $\sigma_p(H_0^\nu)$ is non-empty if and only if $-\frac{\pi}{2} < \operatorname{Im}(\nu) < \frac{\pi}{2}$, and then*

$$\sigma_p(H_0^\nu) = \left\{ -4e^{2(\nu-\gamma)} \right\}.$$

(iii) $\sigma_p(H_{m,0}) = \sigma_p(H_{m,\infty}) = \sigma_p(H_0^\infty) = \emptyset$.

Proof. Only (i) needs a comment, because multivalued functions can be tricky. We can rewrite (5.1) as

$$(-z/4)^{-m} = \varsigma.$$

This is equivalent to

$$\ln(-z/4) \in \frac{-1}{m} \operatorname{Ln}(\varsigma), \tag{5.4}$$

where on the right of (5.4), we have the set of values of the multivalued logarithm. Finally, one deduces from the above inclusion that

$$-z/4 = e^{-w}, \quad w \in \frac{1}{m} \operatorname{Ln}(\varsigma), \quad -\pi < \operatorname{Im}(w) < \pi,$$

which corresponds to (5.3). □

Let us stress that $\sigma_p(H_{m,\kappa})$ depends in a complicated way on the parameters m and κ . There exists a complicated pattern of *phase transitions*, when some eigenvalues “disappear”. This happens if

$$\pi \in \operatorname{Im}\left(\frac{1}{m} \operatorname{Ln}(\varsigma)\right), \quad \text{or} \quad -\pi \in \operatorname{Im}\left(\frac{1}{m} \operatorname{Ln}(\varsigma)\right). \tag{5.5}$$

A pair (m, κ) satisfying (5.5) will be called *exceptional*.

Similarly, for the family of operators H_0^ν , we shall say that ν is exceptional if

$$\operatorname{Im}(\nu) = -\frac{\pi}{2}, \quad \text{or} \quad \operatorname{Im}(\nu) = \frac{\pi}{2}. \quad (5.6)$$

Below, we provide a characterization of $\#\sigma_{\mathbb{P}}(H_{m,\kappa})$, *i.e.*, of the number of eigenvalues of $H_{m,\kappa}$.

Proposition 5.3. *Let $m = m_r + im_i \in \mathbb{C}^\times$ with $|m_r| < 1$.*

(i) *Let $m_r = 0$.*

(a) *If $\frac{\ln(|\kappa|)}{m_i} \in]-\pi, \pi[$, then $\#\sigma_{\mathbb{P}}(H_{m,\kappa}) = \infty$.*

(b) *If $\frac{\ln(|\kappa|)}{m_i} \notin]-\pi, \pi[$, then $\#\sigma_{\mathbb{P}}(H_{m,\kappa}) = 0$.*

(ii) *If $m_r \neq 0$ and if $N \in \{0, 1, 2, \dots\}$ satisfies $N < \frac{m_r^2 + m_i^2}{|m_r|} \leq N + 1$, then*

$$\#\sigma_{\mathbb{P}}(H_{m,\kappa}) \in \{N, N + 1\}.$$

Proof. The case (i) can easily be deduced from (5.3). For (ii), let $m = m_r + im_i \in \mathbb{C}^\times$ and $\alpha = \alpha_r + i\alpha_i \in \operatorname{Ln}(\zeta)$. Then, one has

$$\operatorname{Im}\left(\frac{1}{m}(\alpha + 2\pi ij)\right) = \frac{\alpha_i m_r - \alpha_r m_i}{m_r^2 + m_i^2} + \frac{2\pi j m_r}{m_r^2 + m_i^2}. \quad (5.7)$$

If $m_r > 0$, the condition $-\pi < (5.7) < \pi$ is equivalent to

$$0 < j + \frac{\alpha_i}{2\pi} - \frac{\alpha_r m_i}{2\pi m_r} + \frac{m_r^2 + m_i^2}{2m_r} < \frac{m_r^2 + m_i^2}{m_r}.$$

Thus, we can apply Lemma 5.4 below with $\beta := \frac{m_r^2 + m_i^2}{m_r}$ and $\gamma := \frac{\alpha_i}{2\pi} - \frac{\alpha_r m_i}{2\pi m_r} + \frac{m_r^2 + m_i^2}{2m_r}$ and infer the statement (ii). If $m_r < 0$, one obtains

$$0 < j + \frac{\alpha_i}{2\pi} - \frac{\alpha_r m_i}{2\pi m_r} - \frac{m_r^2 + m_i^2}{2m_r} < -\frac{m_r^2 + m_i^2}{m_r}$$

and the same argument leads to the expected result. \square

In the following lemma, $[\beta]$ denotes the integral part of $\beta \in \mathbb{R}$ and $\{\beta\} := \beta - [\beta]$.

Lemma 5.4. *Let $\beta \geq 0$ and $\gamma \in \mathbb{R}$. If $\beta \in \mathbb{Z}$, then*

$$\#\{j \in \mathbb{Z} \mid 0 < \gamma + j < \beta\} = \begin{cases} \beta - 1 & \text{if } \gamma \in \mathbb{Z}, \\ \beta & \text{if } \gamma \notin \mathbb{Z}, \end{cases}$$

while if $\beta \notin \mathbb{Z}$, then

$$\#\{j \in \mathbb{Z} \mid 0 < \gamma + j < \beta\} = \begin{cases} [\beta] & \text{if } \gamma \in \mathbb{Z} \text{ or } \{\beta\} \leq \{\gamma\}, \\ [\beta] + 1 & \text{if } 0 < \{\gamma\} < \{\beta\}. \end{cases}$$

Proof. For $\beta \in \mathbb{Z}$, the following j are suitable:

$$\begin{aligned} j &= -\gamma + 1, \dots, -\gamma + \beta - 1, & \text{if } \gamma \in \mathbb{Z}, \\ j &= -[\gamma], \dots, -[\gamma] + \beta - 1, & \text{if } \gamma \notin \mathbb{Z}. \end{aligned}$$

For $\beta \notin \mathbb{Z}$, the following j are suitable:

$$\begin{aligned} j &= -\gamma + 1, \dots, -\gamma + [\beta], & \text{if } \gamma \in \mathbb{Z}, \\ j &= -[\gamma], \dots, -[\gamma] + [\beta] - 1, & \text{if } \{\beta\} \leq \{\gamma\}, \\ j &= -[\gamma], \dots, -[\gamma] + [\beta], & \text{if } 0 < \{\gamma\} < \{\beta\}. \end{aligned}$$

□

5.2. Proof of Proposition 5.1

In this section, we prove Proposition 5.1 about the location of possible eigenvalues. First of all, instead of looking for solution of the equation $L_{m^2} f = z f$, it will be convenient to write $z = -k^2$ with $k \in \mathbb{C}$ and $\text{Re}(k) \geq 0$. Then, recall that for $k \neq 0$, two linearly independent solutions of the differential equation (acting on distributions)

$$L_{m^2} f = -k^2 f \tag{5.8}$$

are provided by $x \mapsto \mathcal{I}_m(kx)$ and $x \mapsto \mathcal{K}_m(kx)$. On the other hand, in the special case $k = 0$, the equation $L_{m^2} f = 0$ corresponds to Euler’s equation, and its solutions consist in elementary functions. In fact, we shall treat separately the three cases $\text{Re}(k) > 0$, $\text{Re}(k) = 0$ but $k \neq 0$, and the special case $k = 0$.

5.2.1. $\text{Re}(k) > 0$. We first concentrate on the case $k \in \mathbb{C}$ with $\text{Re}(k) > 0$. Since $|\arg(kx)| < \frac{\pi}{2}$, the function $x \mapsto \mathcal{I}_m(kx)$ exponentially increases for large x . One deduces that for any m , it cannot be in $L^2(\mathbb{R}_+)$. On the other hand, by (A.7) and (A.8), the function $x \mapsto \mathcal{K}_m(kx)$ is square integrable. As a consequence, it remains to determine for which pairs (m, κ) it belongs to $\mathcal{D}(H_{m,\kappa})$, or for which ν it belongs to $\mathcal{D}(H_0^\nu)$.

For $|\text{Re}(m)| < 1$ with $m \neq 0$, consider the equality (A.4) and the power expansion of \mathcal{I}_m provided in (A.1). For $x \in]0, 1[$, one obtains

$$\begin{aligned} \mathcal{K}_m(kx) &= \frac{\sqrt{\pi}}{\sin(\pi m)} \frac{1}{\Gamma(1+m)} \left(\frac{\Gamma(1+m)}{\Gamma(1-m)} \left(\frac{kx}{2}\right)^{1/2-m} - \left(\frac{kx}{2}\right)^{1/2+m} \right) + f_k(x) \\ &= \frac{\Gamma(-m)}{\sqrt{\pi}} \left(\frac{k}{2}\right)^{1/2+m} \left(\frac{\Gamma(m)}{\Gamma(-m)} \left(\frac{k}{2}\right)^{-2m} x^{1/2-m} + x^{1/2+m} \right) + f_k(x), \end{aligned}$$

with $f_k \in \mathcal{D}(L_{m^2}^{\min})$ around 0, by [7, Prop. 4.12]. Thus, one infers from this computation that the function $x \mapsto \mathcal{K}_m(kx)$ belongs to $\mathcal{D}(H_{m,\kappa})$ if and only if

$$\kappa = \frac{\Gamma(m)}{\Gamma(-m)} (k/2)^{-2m}. \tag{5.9}$$

Equivalently, $-k^2$ is an eigenvalue of $H_{m,\kappa}$ if and only if the equality (5.9) holds. Note that $H_{m,\infty}$ has no eigenvalue.

In the special case $m = 0$, by (A.6), observe that

$$\mathcal{K}_0(kx) = -\frac{\sqrt{2kx}}{\sqrt{\pi}} \left(\ln(x) + \ln\left(\frac{k}{2}\right) + \gamma \right) + f_k(x)$$

with $f_k \in \mathcal{D}(L_0^{\min})$ around 0, by [7, Prop. 4.12]. One infers from this computation that the function $x \mapsto \mathcal{K}_0(kx)$ does not belong to $\mathcal{D}(H_{0,\kappa})$, for any

$\kappa \in \mathbb{C} \cup \{\infty\}$. On the other hand, this function belongs to $\mathcal{D}(H_0^\nu)$ if and only if

$$\nu = \gamma + \ln(k/2). \quad (5.10)$$

Equivalently, $-k^2$ is never an eigenvalue of $H_{0,\kappa}$, but $-k^2$ is an eigenvalue of H_0^ν if and only if the equality (5.10) holds.

5.2.2. $k \in i\mathbb{R}^\times$. Let us set $k = i\mu$ with $\mu \in \mathbb{R}^\times$. Our aim is to show that $-k^2 > 0$ can never be an eigenvalue of $H_{m,\kappa}$ or of H_0^ν . For that purpose, consider the two linearly independent solutions of $L_{m^2}f = \mu^2 f$ provided by $x \mapsto \mathcal{H}_m^\pm(\mu x)$. By the asymptotics (A.14), it appears that no non-trivial linear combination of these functions is square integrable at infinity. As expected, we have thus shown that if $k \in i\mathbb{R}^\times$, no solution of the Eq. (5.8) is in $L^2(\mathbb{R}_+)$.

5.2.3. $k = 0$. When $k = 0$, the problem consists first in finding solutions to the homogeneous equation $L_{m^2}f = 0$ with $f \in \mathcal{D}(L_{m^2}^{\max})$. Basic solutions for this equations for $m \neq 0$ are the functions f_\pm with $f_\pm(x) = x^{\pm m+1/2}$. In the special case $m = 0$, a second solution for this equation is provided by the function f_0 with $f_0(x) = x^{1/2} \ln(x)$. However, none of these functions belongs to $L^2(\mathbb{R}_+)$, which means that 0 is never an eigenvalue for the operators $H_{m,\kappa}$ or H_0^ν .

Proof of Proposition 5.1. It has been shown above that $-k^2$ is an eigenvalue of some of the operators $H_{m,\kappa}$ or H_0^ν if $\operatorname{Re}(k) > 0$ and if (5.9) or (5.10) holds. The first statement of the proposition corresponds to reformulation of (5.9) with $z = -k^2$, while the second statement corresponds to (5.10) also with $z = -k^2$. \square

5.3. The Self-Adjoint Case

Let us summarize the content of the first part of this section for self-adjoint operators $H_{m,\kappa}$ or H_0^ν . The following statement is a reformulation of Corollary 2.4 and that of Theorem 5.2.

Theorem 5.5. (i) *If $m \in]-1, 1[\setminus \{0\}$, then $H_{m,\kappa}$ is self-adjoint if and only if $\kappa \in \mathbb{R} \cup \{\infty\}$, and then*

$$\begin{aligned} \sigma_{\text{p}}(H_{m,\kappa}) &= \left\{ -4 \left(\kappa \frac{\Gamma(-m)}{\Gamma(m)} \right)^{-1/m} \right\} \quad \text{for } \kappa \in]-\infty, 0[, \\ \sigma_{\text{p}}(H_{m,\kappa}) &= \emptyset \quad \text{for } \kappa \in [0, \infty]. \end{aligned}$$

(ii) *If $m = im_i \in i\mathbb{R} \setminus \{0\}$, then $H_{im_i,\kappa}$ is self-adjoint if and only if $|\kappa| = 1$, and then*

$$\sigma_{\text{p}}(H_{im_i,\kappa}) = \left\{ -4 \exp \left(-\frac{\arg \left(\kappa \frac{\Gamma(-im_i)}{\Gamma(im_i)} \right) + 2\pi j}{m_i} \right) \mid j \in \mathbb{Z} \right\}.$$

(iii) *H_0^ν is self-adjoint if and only if $\nu \in \mathbb{R} \cup \{\infty\}$, and then*

$$\begin{aligned} \sigma_{\text{p}}(H_0^\nu) &= \{ -4e^{2(\nu-\gamma)} \} \quad \text{for } \nu \in \mathbb{R}, \\ \sigma_{\text{p}}(H_0^\infty) &= \emptyset. \end{aligned}$$

Remark 5.6. Let us emphasize that none of the pairs (m, κ) corresponding to a self-adjoint operator $H_{m,\kappa}$ is an exceptional pair. Similarly, the parameter ν corresponding to a self-adjoint operator H'_0 is never exceptional.

6. Continuous Spectrum of $H_{m,\kappa}$

In this section, we extend the results obtained in Sect. 4 to the families of operators $H_{m,\kappa}$.

6.1. Resolvent

In this section, we consider $m \in \mathbb{C}^\times$ with $|\operatorname{Re}(m)| < 1$, $\kappa \in \mathbb{C} \cup \{\infty\}$, and $x \in \mathbb{R}_+$. Let us also fix $k \in \mathbb{C}$ with $-k^2 \notin \sigma(H_{m,\kappa})$, and as before, we impose $\operatorname{Re}(k) > 0$. Note that we have directly imposed that $k \notin i\mathbb{R}$, since later on, we shall show that $[0, \infty[\subset \sigma(H_{m,\kappa})$. Our aim is to compute the integral kernel of the resolvent of $H_{m,\kappa}$

$$R_{m,\kappa}(-k^2) := (H_{m,\kappa} + k^2)^{-1}.$$

First of all, recall from Sect. 5 that the map $x \mapsto \mathcal{K}_m(kx)$ is a solution of the equation

$$(L_{m^2} + k^2)f = 0 \tag{6.1}$$

and belongs to $L^2(\mathbb{R}_+)$. Similarly, both functions $x \mapsto \mathcal{I}_m(kx)$ and $x \mapsto \mathcal{I}_{-m}(kx)$ satisfy the equation (6.1), but only a certain linear combination belongs to $\mathcal{D}(H_{m,\kappa})$ around 0.

Recall now, the parameter that we have introduced in (5.2), namely

$$\varsigma = \kappa \frac{\Gamma(-m)}{\Gamma(m)}.$$

For $(\frac{k}{2})^{2m}\varsigma \neq 1$, by considering (A.2), one infers that the map

$$x \mapsto \frac{\mathcal{I}_m(kx) - \varsigma(\frac{k}{2})^{2m}\mathcal{I}_{-m}(kx)}{1 - \varsigma(\frac{k}{2})^{2m}}, \tag{6.2}$$

satisfies (2.3) or (2.4) around 0, and hence belongs to $\mathcal{D}(H_{m,\kappa})$ around 0. Obviously, it also solves (6.1). Furthermore, using the formulas of Sect. A.3, the Wronskian of (6.2) and $x \mapsto \mathcal{K}_m(kx)$ equals k

From the general theory of Sturm–Liouville operators, as recalled, for example, in [7, Prop. A.1], the kernel of the operator $(H_{m,\kappa} + k^2)^{-1}$ is provided for $\varsigma(\frac{k}{2})^{2m} \neq 1$ by the expression

$$\begin{aligned} & R_{m,\kappa}(-k^2; x, y) \\ &= \frac{1}{k \left(1 - \varsigma(\frac{k}{2})^{2m}\right)} \begin{cases} (\mathcal{I}_m(kx) - \varsigma(\frac{k}{2})^{2m}\mathcal{I}_{-m}(kx)) \mathcal{K}_m(ky) & \text{if } 0 < x < y, \\ (\mathcal{I}_m(ky) - \varsigma(\frac{k}{2})^{2m}\mathcal{I}_{-m}(ky)) \mathcal{K}_m(kx) & \text{if } 0 < y < x. \end{cases} \end{aligned}$$

Let us still provide a relation between $R_{m,\kappa}(-k^2)$ and $R_m(-k^2)$. Indeed, by the definition (6.2), and by considering the equalities $H_{m,0} = H_m$ and $H_{m,\infty} = H_{-m}$, one gets

$$R_{m,\kappa}(-k^2) = \frac{1}{1 - \zeta(\frac{k}{2})^{2m}} \left(R_m(-k^2) - \zeta(\frac{k}{2})^{2m} R_{-m}(-k^2) \right) \quad (6.3)$$

$$= R_m(-k^2) - \frac{\zeta(\frac{k}{2})^{2m}}{1 - \zeta(\frac{k}{2})^{2m}} \frac{m}{k^2} P_m(-k^2), \quad (6.4)$$

where $P_m(-k^2)$ is the projection defined in (4.2), see also (4.4).

Let us finally observe that for fixed κ , the condition $\zeta(\frac{k}{2})^{2m} = 1$ defines a discrete set which corresponds to the eigenvalues of $H_{m,\kappa}$ by Proposition 5.1. On the other hand, since $R_{m,\kappa}(-k^2)$ is a rank one perturbation of $R_m(-k^2)$, one infers that $[0, \infty[$ belongs to the spectrum of $H_{m,\kappa}$, as already mentioned at the beginning of this section.

6.2. Boundary Value of the Resolvent and Spectral Density

Since $[0, \infty[$ belongs to the spectrum of $H_{m,\kappa}$, it is natural to mimic the computations performed in Sect. 4.2 for H_m . Recall that an exceptional pair (m, κ) has been defined in (5.5).

Theorem 6.1. *Let $m \in \mathbb{C}^\times$ with $|\operatorname{Re}(m)| < 1$, let $\kappa \in \mathbb{C} \cup \{\infty\}$, and let $k > 0$.*

- (i) *If (m, κ) is not an exceptional pair, then the boundary values of the resolvent*

$$R_{m,\kappa}(k^2 \pm i0) := \lim_{\epsilon \searrow 0} R_{m,\kappa}(k^2 \pm i\epsilon) \quad (6.5)$$

exist in the sense of operators from $\langle X \rangle^{-s} L^2(\mathbb{R}_+)$ to $\langle X \rangle^s L^2(\mathbb{R}_+)$ for any $s > \frac{1}{2}$, uniformly in k on each compact subset of \mathbb{R}_+ . The kernel of $R_{m,\kappa}(k^2 \pm i0)$ is given for $0 < x \leq y$ by

$$\begin{aligned} & R_{m,\kappa}(k^2 \pm i0; x, y) \\ &= \frac{\pm i}{k \left(1 - \zeta e^{\mp i\pi m} \left(\frac{k}{2} \right)^{2m} \right)} \left(\mathcal{J}_m(kx) - \zeta \left(\frac{k}{2} \right)^{2m} \mathcal{J}_{-m}(kx) \right) \mathcal{H}_m^\pm(ky) \end{aligned}$$

and the same expression with the role of x and y exchanged for $0 < y < x$.

- (ii) *If (m, κ) is an exceptional pair, the same statement holds for k uniformly on each compact subset of*

$$\left\{ k \in \mathbb{R}_+ \mid \left(\frac{k}{2} \right)^{2m} \zeta e^{\mp i\pi m} \neq 1 \right\}. \quad (6.6)$$

Proof. By considering the equality (6.3), one infers that the limiting absorption principle (6.5) can be deduced from Theorem 4.2 for R_m and for R_{-m} . The local uniformity in k is also consequence of these estimates and of the explicit formula for the prefactors appearing in (6.3), as long as the first factor has no singularity.

For the kernel, one directly gets for $0 < x \leq y$ that

$$\begin{aligned} R_{m,\kappa}(k^2 \pm i0; x, y) &= \frac{1}{1 - \varsigma e^{\mp i\pi m} \left(\frac{k}{2}\right)^{2m}} R_m(k^2 \pm i0; x, y) \\ &\quad - \frac{\varsigma e^{\mp i\pi m} \left(\frac{k}{2}\right)^{2m}}{1 - \varsigma e^{\mp i\pi m} \left(\frac{k}{2}\right)^{2m}} R_{-m}(k^2 \pm i0; x, y) \\ &= \frac{\pm i}{k \left(1 - \varsigma e^{\mp i\pi m} \left(\frac{k}{2}\right)^{2m}\right)} \left(\mathcal{J}_m(kx) - \varsigma \left(\frac{k}{2}\right)^{2m} \mathcal{J}_{-m}(kx) \right) \mathcal{H}_m^\pm(ky), \end{aligned}$$

where (A.12) and (4.5) have been considered. Note that for $0 < y < x$, one gets the same expression with the role of x and y exchanged. \square

Based on the previous result, the corresponding *spectral density* can now be defined for any $m \in \mathbb{C}^\times$ with $|\operatorname{Re}(m)| < 1$. More precisely, for any $k > 0$ if (m, κ) is not an exceptional pair, or for k belonging to (6.6) if (m, κ) is an exceptional pair, let us set

$$p_{m,\kappa}(k^2) := \frac{1}{2\pi i} \left(R_{m,\kappa}(k^2 + i0) - R_{m,\kappa}(k^2 - i0) \right),$$

which is bounded from $\langle X \rangle^{-s} L^2(\mathbb{R}_+)$ to $\langle X \rangle^s L^2(\mathbb{R}_+)$ for any $s > \frac{1}{2}$.

Proposition 6.2. *The kernel of the spectral density is given by the following formula:*

$$p_{m,\kappa}(k^2; x, y) = \frac{\left(\mathcal{J}_m(kx) - \varsigma \left(\frac{k}{2}\right)^{2m} \mathcal{J}_{-m}(kx) \right) \left(\mathcal{J}_m(ky) - \varsigma \left(\frac{k}{2}\right)^{2m} \mathcal{J}_{-m}(ky) \right)}{\pi k \left(\sin^2(\pi m) + \left(\cos(\pi m) - \varsigma \left(\frac{k}{2}\right)^{2m} \right)^2 \right)}.$$

Proof. For the kernel of this operator, observe that for $0 < x \leq y$, one has

$$\begin{aligned} 2\pi k p_{m,\kappa}(k^2; x, y) &= \frac{1}{1 - \varsigma e^{-i\pi m} \left(\frac{k}{2}\right)^{2m}} \left(\mathcal{J}_m(kx) - \varsigma \left(\frac{k}{2}\right)^{2m} \mathcal{J}_{-m}(kx) \right) \mathcal{H}_m^+(ky) \\ &\quad + \frac{1}{1 - \varsigma e^{i\pi m} \left(\frac{k}{2}\right)^{2m}} \left(\mathcal{J}_m(kx) - \varsigma \left(\frac{k}{2}\right)^{2m} \mathcal{J}_{-m}(kx) \right) \mathcal{H}_m^-(ky) \\ &= \frac{\mathcal{J}_m(kx) - \varsigma \left(\frac{k}{2}\right)^{2m} \mathcal{J}_{-m}(kx)}{\left(1 - \varsigma e^{-i\pi m} \left(\frac{k}{2}\right)^{2m}\right) \left(1 - \varsigma e^{i\pi m} \left(\frac{k}{2}\right)^{2m}\right)} \\ &\quad \times \left\{ \left(1 - \varsigma e^{i\pi m} \left(\frac{k}{2}\right)^{2m}\right) \mathcal{H}_m^+(ky) + \left(1 - \varsigma e^{-i\pi m} \left(\frac{k}{2}\right)^{2m}\right) \mathcal{H}_m^-(ky) \right\} \\ &= \frac{\left(\mathcal{J}_m(kx) - \varsigma \left(\frac{k}{2}\right)^{2m} \mathcal{J}_{-m}(kx) \right) \left(\mathcal{J}_m(ky) - \varsigma \left(\frac{k}{2}\right)^{2m} \mathcal{J}_{-m}(ky) \right)}{\frac{1}{2} \left(1 - 2\varsigma \cos(\pi m) \left(\frac{k}{2}\right)^{2m} + \varsigma^2 \left(\frac{k}{2}\right)^{4m} \right)}. \end{aligned}$$

Then, using a simple trigonometric equality for the denominator, and since the role of x and y can be exchanged, one directly obtains the desired expression for any $x, y \in \mathbb{R}_+$. □

6.3. Generalized Hankel Transform

We would like to generalize the definition of the Hankel transformations of Sect. 4.3 and to show their relations with the operators $H_{m,\kappa}$. The main idea is to factorize the spectral density, but let us stress that this factorization is not unique.

One possibility, which works well for m and κ real, is as follows. For any $m \in \mathbb{R}^\times$ with $|m| < 1$ and any $\kappa \in \mathbb{R} \cup \{\infty\}$, one could set

$$\mathcal{F}_{m,\kappa} : C_c(\mathbb{R}_+) \rightarrow L^2(\mathbb{R}_+)$$

with

$$(\mathcal{F}_{m,\kappa} f)(x) := \int_0^\infty \mathcal{F}_{m,\kappa}(x, y) f(y) dy,$$

and

$$\mathcal{F}_{m,\kappa}(x, y) := \sqrt{\frac{2}{\pi}} \frac{\mathcal{J}_m(xy) - \varsigma \mathcal{J}_{-m}(xy) \left(\frac{y}{2}\right)^{2m}}{\sqrt{\sin^2(\pi m) + \left(\cos(\pi m) - \varsigma \left(\frac{y}{2}\right)^{2m}\right)^2}}.$$

For real m and κ , the denominator is the square root of a positive number, which has a clear meaning. For general m and κ , we need to choose the branch of the square root, which is ambiguous.

Another option, which we will prefer, since it has always a unique definition, is to define the *incoming* and *outgoing Hankel transformations* $\mathcal{F}_{m,\kappa}^\mp$ given by the kernels

$$\mathcal{F}_{m,\kappa}^\mp(x, y) := e^{\mp i \frac{\pi}{2} m} \sqrt{\frac{2}{\pi}} \frac{\mathcal{J}_m(xy) - \varsigma \mathcal{J}_{-m}(xy) \left(\frac{y}{2}\right)^{2m}}{1 - \varsigma e^{\mp i \pi m} \left(\frac{y}{2}\right)^{2m}}.$$

Note that the denominator of this kernel has a singularity if (m, κ) is an exceptional pair. For simplicity, we shall not consider this situation anymore in the sequel. Thus, if (m, κ) is not an exceptional pair, the following equalities can easily be obtained:

$$\mathcal{F}_{m,\kappa}^\mp = \left(\mathcal{F}_m - \varsigma \mathcal{F}_{-m} \left(\frac{X}{2}\right)^{2m} \right) \frac{e^{\mp i \frac{\pi}{2} m}}{1 - \varsigma e^{\mp i \pi m} \left(\frac{X}{2}\right)^{2m}}. \quad (6.7)$$

Let us now show that these transformations provide a kind of diagonalization of the operator $H_{m,\kappa}$. The statements and the proofs are divided into several shorter statements.

Lemma 6.3. *For any $m \in \mathbb{C}^\times$ with $|\operatorname{Re}(m)| < 1$ and any $\kappa \in \mathbb{C} \cup \{\infty\}$ with (m, κ) not an exceptional pair, the operators $\mathcal{F}_{m,\kappa}^\mp$ extend continuously to the*

following operators:

$$\begin{aligned} & \left(\Xi_m(-A) - \varsigma \Xi_{-m}(-A)(2X)^{-2m} \right) \frac{e^{\mp i \frac{\pi}{2} m}}{1 - \varsigma e^{\mp i \pi m} (2X)^{-2m}} J \\ &= J \left(\Xi_m(A) - \varsigma \Xi_{-m}(A) \left(\frac{X}{2} \right)^{2m} \right) \frac{e^{\mp i \frac{\pi}{2} m}}{1 - \varsigma e^{\mp i \pi m} \left(\frac{X}{2} \right)^{2m}} \end{aligned} \quad (6.8)$$

where Ξ_m and Ξ_{-m} have been defined in (4.18).

Proof. The proof simply consists in using the equalities obtained in Proposition 4.5 and in a short computation based on the definition of J , see equation (4.16). \square

Let us still provide the expressions for the adjoint of these operators, namely $\mathcal{F}_{m,\kappa}^{\mp\#}$ are equal to the operators

$$\begin{aligned} & J \frac{e^{\mp i \frac{\pi}{2} m}}{1 - \varsigma e^{\mp i \pi m} (2X)^{-2m}} \left(\Xi_m(A) - \varsigma (2X)^{-2m} \Xi_{-m}(A) \right) \\ &= \frac{e^{\mp i \frac{\pi}{2} m}}{1 - \varsigma e^{\mp i \pi m} \left(\frac{X}{2} \right)^{2m}} \left(\Xi_m(-A) - \varsigma \left(\frac{X}{2} \right)^{2m} \Xi_{-m}(-A) \right) J. \end{aligned} \quad (6.9)$$

Additional information on the operators $\mathcal{F}_{m,\kappa}^{\pm\#}$ is provided in the next statement.

Lemma 6.4. *For any $m \in \mathbb{C}^\times$ with $|\operatorname{Re}(m)| < 1$ and $\kappa \in \mathbb{C} \cup \{\infty\}$ with (m, κ) not an exceptional pair, the following equalities hold:*

$$\mathcal{F}_{m,\kappa}^{\pm\#} \mathcal{F}_{m,\kappa}^{\mp} = \mathbb{1}.$$

Proof. The proof consists simply in considering the terms (6.8) and (6.9) and in checking that their product (for the correct sign) is equal to $\mathbb{1}$. Indeed, observe first that on $C_c(\mathbb{R}_+)$, the following equalities hold:

$$\begin{aligned} & \left(\Xi_m(A) - \varsigma (2X)^{-2m} \Xi_{-m}(A) \right) \left(\Xi_m(-A) - \varsigma \Xi_{-m}(-A)(2X)^{-2m} \right) \\ &= 1 + \varsigma^2 (2X)^{-4m} \\ & \quad - \varsigma \left((2X)^{-2m} \Xi_{-m}(A) \Xi_m(-A) + \Xi_m(A) \Xi_{-m}(-A)(2X)^{-2m} \right). \end{aligned}$$

If one then shows that

$$(2X)^{-2m} \Xi_{-m}(A) \Xi_m(-A) + \Xi_m(A) \Xi_{-m}(-A)(2X)^{-2m} = 2 \cos(\pi m) (2X)^{-2m}, \quad (6.10)$$

the statement of the lemma directly follows using the full expressions provided in (6.8) and (6.9).

For the proof of (6.10), recall $\{U_\tau\}_{\tau \in \mathbb{R}}$ corresponds to the dilation group, namely $U_\tau = e^{i\tau A}$ as introduced in Sect. 2.1. Then, for any $f \in C_c(\mathbb{R}_+)$ and $x \in \mathbb{R}_+$, one has

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$$\begin{aligned}
 ((2X)^{-2m} U_\tau (2X)^{2m} f)(x) &= e^{-2m \ln(x)} (U_\tau X^{2m} f)(x) \\
 &= e^{-2m \ln(x)} e^{\tau/2} (X^{2m} f)(e^\tau x) \\
 &= e^{-2m \ln(x)} e^{\tau/2} e^{2m \ln(e^\tau x)} f(e^\tau x) \\
 &= e^{2m\tau} (e^{i\tau A} f)(x) \\
 &= \left(e^{i\tau(A-2m)} f \right)(x).
 \end{aligned}$$

One infers from this computation that the l.h.s. of (6.10) is equal to

$$\left(\Xi_{-m}(A - i2m) \Xi_m(-A + i2m) + \Xi_m(A) \Xi_{-m}(-A) \right) (2X)^{-2m}.$$

Finally, by considering the explicit formula (4.18) for Ξ_m and the equality $\Gamma(z)\Gamma(1-z) = \frac{\pi}{\sin(\pi z)}$, one gets (with t instead of A)

$$\begin{aligned}
 &\Xi_{-m}(t - i2m) \Xi_m(-t + i2m) + \Xi_m(t) \Xi_{-m}(-t) \\
 &= \Gamma\left(\frac{m+1+it}{2}\right) \Gamma\left(\frac{-m+1-it}{2}\right) \left(\frac{1}{\Gamma\left(\frac{-3m+1-it}{2}\right) \Gamma\left(\frac{3m+1+it}{2}\right)} + \frac{1}{\Gamma\left(\frac{m+1-it}{2}\right) \Gamma\left(\frac{-m+1+it}{2}\right)} \right) \\
 &= \frac{1}{\sin(\alpha)} (\sin(\alpha - 2\beta) + \sin(\alpha + 2\beta))
 \end{aligned}$$

with $\alpha := \pi\left(\frac{-m+1-it}{2}\right)$ and $\beta := \pi\frac{m}{2}$. Some trigonometric identities lead then directly to the equality

$$\frac{1}{\sin(\alpha)} (\sin(\alpha - 2\beta) + \sin(\alpha + 2\beta)) = 2 \cos(2\beta) = 2 \cos(\pi m),$$

as expected. □

Let us now set

$$\mathbb{1}_{\mathbb{R}_+}(H_{m,\kappa}) := \mathcal{F}_{m,\kappa}^\pm \mathcal{F}_{m,\kappa}^{\mp\#} \tag{6.11}$$

and observe that this is again a projection. In the self-adjoint case, this operator corresponds to the usual projection onto the continuous spectrum of $H_{m,\kappa}$. One can then prove the analog of Proposition 4.6.

Proposition 6.5. *Let $m \in \mathbb{C}^\times$ with $|\operatorname{Re}(m)| < 1$ and $\kappa \in \mathbb{C} \cup \{\infty\}$ with (m, κ) not an exceptional pair. Then, for any $k \in \mathbb{C}$ with $\operatorname{Re}(k) > 0$ and $-k^2 \notin \sigma_p(H_{m,\kappa})$, the following equalities hold:*

$$R_{m,\kappa}(-k^2) \mathbb{1}_{\mathbb{R}_+}(H_{m,\kappa}) = \mathcal{F}_{m,\kappa}^\pm (X^2 + k^2)^{-1} \mathcal{F}_{m,\kappa}^{\mp\#} = \mathbb{1}_{\mathbb{R}_+}(H_{m,\kappa}) R_{m,\kappa}(-k^2).$$

Proof. We will use the following convenient expression for $\mathcal{F}_{m,\kappa}^{\mp\#}$ and two formulas for the resolvent:

$$\mathcal{F}_{m,\kappa}^{\mp\#} = \frac{e^{\mp i \frac{\pi}{2} m}}{1 - \zeta e^{\mp i \pi m} \left(\frac{X}{2}\right)^{2m}} \left(\mathcal{F}_m - \zeta \left(\frac{X}{2}\right)^{2m} \mathcal{F}_{-m} \right), \tag{6.12}$$

$$R_{m,\kappa}(-k^2) = R_m(-k^2) - \frac{\zeta \left(\frac{k}{2}\right)^{2m}}{1 - \zeta \left(\frac{k}{2}\right)^{2m}} \frac{m}{k^2} P_m(-k^2), \tag{6.13}$$

$$R_{m,\kappa}(-k^2) = R_{-m}(-k^2) - \frac{1}{1 - \zeta \left(\frac{k}{2}\right)^{2m}} \frac{m}{k^2} P_m(-k^2). \tag{6.14}$$

By multiplying (6.13) and (6.14) from the left by \mathcal{F}_m and \mathcal{F}_{-m} , respectively, we obtain

$$\begin{aligned} \mathcal{F}_m R_{m,\kappa}(-k^2) &= (X^2 + k^2)^{-1} \mathcal{F}_m - \frac{\varsigma(\frac{k}{2})^{2m}}{1 - \varsigma(\frac{k}{2})^{2m}} \frac{m}{k^2} \mathcal{F}_m P_m(-k^2), \\ \mathcal{F}_{-m} R_{m,\kappa}(-k^2) &= (X^2 + k^2)^{-1} \mathcal{F}_{-m} - \frac{1}{1 - \varsigma(\frac{k}{2})^{2m}} \frac{m}{k^2} \mathcal{F}_{-m} P_m(-k^2). \end{aligned}$$

By combining then the above two identities, and by considering (6.12), we get

$$\begin{aligned} \mathcal{F}_{m,\kappa}^{\mp\#} R_{m,\kappa}(-k^2) &= (X^2 + k^2)^{-1} \mathcal{F}_{m,\kappa}^{\mp\#} \\ &\quad - \frac{e^{\mp i \frac{\pi}{2} m}}{\left(1 - \varsigma e^{\mp i \pi m} \left(\frac{X}{2}\right)^{2m}\right) \left(1 - \varsigma \left(\frac{k}{2}\right)^{2m}\right)} \\ &\quad \times \frac{m\varsigma}{k^2} \left(\left(\frac{k}{2}\right)^{2m} \mathcal{F}_m - \left(\frac{X}{2}\right)^{2m} \mathcal{F}_{-m} \right) P_m(-k^2). \end{aligned}$$

By considering the equality (A.18), one can deduce that for any $z > 0$ and $m \in \mathbb{C}$ with $|\operatorname{Re}(m)| < 1$, one has

$$\int_0^\infty \mathcal{K}_m(z^{-1}x) \mathcal{J}_m(x) dx = z^{2m} \int_0^\infty \mathcal{K}_m(z^{-1}x) \mathcal{J}_{-m}(x) dx. \quad (6.15)$$

We infer then from this equality that

$$\left(\frac{k}{2}\right)^{2m} \mathcal{F}_m P_m(-k^2) = \left(\frac{X}{2}\right)^{2m} \mathcal{F}_{-m} P_m(-k^2), \quad (6.16)$$

and as a direct consequence,

$$\mathcal{F}_{m,\kappa}^{\mp\#} R_{m,\kappa}(-k^2) = (X^2 + k^2)^{-1} \mathcal{F}_{m,\kappa}^{\mp\#}. \quad (6.17)$$

This equality corresponds to one of the identities of our theorem, the proof of the other identity being analogous. \square

6.4. Spectral Projections

We first describe the spectral projections corresponding to eigenvalues of $H_{m,\kappa}$.

Proposition 6.6. *For any $-k^2 \in \sigma_p(H_{m,\kappa})$, one has*

$$\mathbb{1}_{\{-k^2\}}(H_{m,\kappa}) = P_m(-k^2).$$

Proof. Let γ be a contour encircling $-k^2$ in the complex plane, with no other eigenvalue inside γ and with no intersection with $[0, \infty[$. We use (6.4), and then compute the residue of the resolvent at $z = -k^2$:

$$\begin{aligned}
 \mathbb{1}_{\{-k^2\}}(H_{m,\kappa}) &= -\frac{1}{2\pi i} \int_{\gamma} R_{m,\kappa}(z) dz \\
 &= -\frac{1}{2\pi i} \int_{\gamma} \frac{\varsigma \frac{(-z)^m}{2^{2m}}}{1 - \varsigma \frac{(-z)^m}{2^{2m}}} \frac{m}{z} P_m(z) \\
 &= -\frac{\varsigma \frac{(-z)^m}{2^{2m}}}{\frac{d}{dz} \left(1 - \varsigma \frac{(-z)^m}{2^{2m}} \right)} \frac{m}{z} P_m(z) \Big|_{z=-k^2} \\
 &= P_m(-k^2).
 \end{aligned}$$

□

Let us now assume that (m, κ) is not an exceptional pair. As in Sect. 4.4, we can consider for any $0 < a < b < \infty$, the operator

$$\mathbb{1}_{[a,b]}(H_{m,\kappa}) := 2 \int_{\sqrt{a}}^{\sqrt{b}} p_{m,\kappa}(k^2) k dk, \quad (6.18)$$

which is bounded from $\langle X \rangle^{-s} L^2(\mathbb{R}_+)$ to $\langle X \rangle^s L^2(\mathbb{R}_+)$ for any $s > \frac{1}{2}$. Its kernel is given for $x, y \in \mathbb{R}_+$ by the expression

$$\begin{aligned}
 &\mathbb{1}_{[a,b]}(H_{m,\kappa})(x, y) \\
 &= 2 \int_{\sqrt{a}}^{\sqrt{b}} \frac{\left(\mathcal{J}_m(kx) - \varsigma \left(\frac{k}{2}\right)^{2m} \mathcal{J}_{-m}(kx) \right) \left(\mathcal{J}_m(ky) - \varsigma \left(\frac{k}{2}\right)^{2m} \mathcal{J}_{-m}(ky) \right)}{\pi \left(\sin^2(\pi m) + \left(\cos(\pi m) - \varsigma \left(\frac{k}{2}\right)^{2m} \right)^2 \right)} dk.
 \end{aligned} \quad (6.19)$$

Proposition 6.7. *For any $0 < a < b$, any $m \in \mathbb{C}^\times$ with $|\operatorname{Re}(m)| < 1$, and any $\kappa \in \mathbb{C} \cup \{\infty\}$ with (m, κ) not an exceptional pair, one has*

$$\mathbb{1}_{[a,b]}(H_{m,\kappa}) = \mathcal{F}_{m,\kappa}^\pm \mathbb{1}_{[a,b]}(X^2) \mathcal{F}_{m,\kappa}^{\mp\#} \quad (6.20)$$

in $\mathcal{B}(L^2(\mathbb{R}_+))$. In addition, $\mathbb{1}_{[a,b]}(H_{m,\kappa})$ is a projection.

Proof. Let us recall that the l.h.s. has been defined in (6.18), and check that the r.h.s. corresponds to this expression. Indeed, for any $f \in C_c(\mathbb{R}_+)$ and $x \in \mathbb{R}_+$, one has

$$\begin{aligned}
 &(\mathcal{F}_{m,\kappa}^\pm \mathbb{1}_{[a,b]}(X^2) \mathcal{F}_{m,\kappa}^{\mp\#} f)(x) \\
 &= e^{\pm i \frac{\pi}{2} m} \sqrt{\frac{2}{\pi}} \int_{\sqrt{a}}^{\sqrt{b}} \frac{\mathcal{J}_m(xk) - \varsigma \left(\frac{k}{2}\right)^{2m} \mathcal{J}_{-m}(xk)}{1 - \varsigma e^{\pm i \pi m} \left(\frac{k}{2}\right)^{2m}} (\mathcal{F}_{m,\kappa}^{\mp\#} f)(k) dk \\
 &= \frac{2}{\pi} \int_{\sqrt{a}}^{\sqrt{b}} \left(\frac{\mathcal{J}_m(xk) - \varsigma \left(\frac{k}{2}\right)^{2m} \mathcal{J}_{-m}(xk)}{1 - \varsigma e^{\pm i \pi m} \left(\frac{k}{2}\right)^{2m}} \right)
 \end{aligned}$$

$$\begin{aligned} & \times \int_0^\infty \frac{\mathcal{J}_m(ky) - \varsigma \left(\frac{k}{2}\right)^{2m} \mathcal{J}_{-m}(ky)}{1 - \varsigma e^{\mp i\pi m} \left(\frac{k}{2}\right)^{2m}} f(y) dy \Big) dk \\ & = \int_0^\infty \mathbb{1}_{[a,b]}(H_{m,\kappa})(x, y) f(y) dy, \end{aligned}$$

where the kernel $\mathbb{1}_{[a,b]}(H_{m,\kappa})(x, y)$ has been defined in (6.19). Note that since the r.h.s. of (6.20) defines a bounded operator on $L^2(\mathbb{R}_+)$, this equality provides a natural continuous extension of $\mathbb{1}_{[a,b]}(H_{m,\kappa})$ as a bounded operator on $L^2(\mathbb{R}_+)$. Finally, by Lemma 6.4 one readily infers that $\mathbb{1}_{[a,b]}(H_{m,\kappa})$ is a projection. \square

Note that the equality

$$\mathbb{1}_\Xi(H_{m,\kappa}) = \mathcal{F}_{m,\kappa}^\pm \mathbb{1}_\Xi(X^2) \mathcal{F}_{m,\kappa}^{\mp\#} \tag{6.21}$$

extends (6.20) to any Borel subset Ξ of \mathbb{R}_+ . In particular, $\mathbb{1}_{\mathbb{R}_+}(H_{m,\kappa})$ obtained with this definition corresponds to the one already introduced in (6.11).

6.5. Møller Operators and Scattering Operator

In this section, we extend the results obtained for the Møller and the scattering operators to the larger family of operators $H_{m,\kappa}$. For that purpose, we consider pairs (m, κ) and (m', κ') which are not exceptional. The easiest way to introduce the wave operators for the pair $(H_{m,\kappa}, H_{m',\kappa'})$ is to define them using the Hankel transformations:

$$W_{m,\kappa;m',\kappa'}^\pm = \mathcal{F}_{m,\kappa}^\pm \mathcal{F}_{m',\kappa'}^{\mp\#}. \tag{6.22}$$

These definitions immediately imply the following identities:

Proposition 6.8. *The following identities hold:*

$$\begin{aligned} W_{m,\kappa;m',\kappa'}^{\mp\#} W_{m,\kappa;m',\kappa'}^\pm &= \mathbb{1}_{\mathbb{R}_+}(H_{m',\kappa'}), \\ W_{m,\kappa;m',\kappa'}^\pm W_{m,\kappa;m',\kappa'}^{\mp\#} &= \mathbb{1}_{\mathbb{R}_+}(H_{m,\kappa}), \\ W_{m,\kappa;m',\kappa'}^{\pm\#} &= W_{m',\kappa';m,\kappa}^\mp, \\ W_{m,\kappa;m',\kappa'}^\pm H_{m',\kappa'} &= H_{m,\kappa} W_{m,\kappa;m',\kappa'}^\pm. \end{aligned} \tag{6.23}$$

By (6.23), $W_{m,\kappa;m',\kappa'}^{-\#}$ can be treated as an inverse of $W_{m,\kappa;m',\kappa'}^+$. Therefore, we define the scattering operator as

$$S_{m,\kappa;m',\kappa'} := W_{m,\kappa;m',\kappa'}^{-\#} W_{m,\kappa;m',\kappa'}^-.$$

Clearly, the scattering operator can be expressed in terms of the Hankel transform:

$$S_{m,\kappa;m',\kappa'} = \mathcal{F}_{m',\kappa'}^+ \mathcal{F}_{m,\kappa}^{-\#} \mathcal{F}_{m,\kappa}^- \mathcal{F}_{m',\kappa'}^{+\#}.$$

To analyze the scattering operator, it is convenient to introduce the operators

$$\mathcal{G}_{m,\kappa}^\mp := \mathcal{F}_{m,\kappa}^{\mp\#} \mathcal{F}_{m,\kappa}^\mp. \tag{6.24}$$

Proposition 6.9. *The following equalities hold:*

$$\mathcal{G}_{m,\kappa}^\mp = e^{\mp i\pi m} \frac{\mathbb{1} - \zeta e^{\pm i\pi m} \left(\frac{X}{2}\right)^{2m}}{\mathbb{1} - \zeta e^{\mp i\pi m} \left(\frac{X}{2}\right)^{2m}}.$$

Proof. The proof consists in a simple computation, starting from the expressions (6.8) and (6.9) and considering the equality (6.10). \square

Let us stress that $\mathcal{G}_{m,\kappa}^\mp$ are simply functions of X . Finally, using the Hankel transformations, one can obtain a diagonal representation of scattering operators. These operators are expressed in terms of the operators (6.24), namely

$$\mathcal{F}_{m',\kappa'}^{-\#} S_{m,\kappa;m',\kappa'} \mathcal{F}_{m',\kappa'}^+ = \mathcal{G}_{m,\kappa}^- \mathcal{G}_{m',\kappa'}^+ = \mathcal{F}_{m',\kappa'}^{+\#} S_{m,\kappa;m',\kappa'} \mathcal{F}_{m',\kappa'}^-.$$

In the non-exceptional case, the operator $H_{m,\kappa}$ generates a bounded one-parameter group, at least on the range of the projection $\mathbb{1}_{\mathbb{R}_+}(H_{m,\kappa})$, by the formula

$$e^{itH_{m,\kappa}} \mathbb{1}_{\mathbb{R}_+}(H_{m,\kappa}) = \mathbb{1}_{\mathbb{R}_+}(H_{m,\kappa}) e^{itH_{m,\kappa}} = \mathcal{F}_{m,\kappa}^\pm e^{itX^2} \mathcal{F}_{m,\kappa}^{\mp\#}. \quad (6.25)$$

In this context, we can then show that $W_{m,\kappa;m',\kappa'}^\pm$ correspond to the Møller operators as usually defined in the time-dependent scattering theory.

Proposition 6.10. *For any $m, m' \in \mathbb{C}$ with $|\operatorname{Re}(m)| < 1$ and $|\operatorname{Re}(m')| < 1$, and for any $\kappa, \kappa' \in \mathbb{C} \cup \{\infty\}$, such that (m, κ) and (m', κ') are not exceptional pairs, the Møller operators exist and coincide with the operators defined in (6.22):*

$$s\text{-}\lim_{t \rightarrow \pm\infty} \mathbb{1}_{\mathbb{R}_+}(H_{m,\kappa}) e^{itH_{m,\kappa}} e^{-itH_{m',\kappa'}} \mathbb{1}_{\mathbb{R}_+}(H_{m',\kappa'}) = W_{m,\kappa;m',\kappa'}^\pm.$$

Proof. By (6.25), we have

$$\mathbb{1}_{\mathbb{R}_+}(H_{m,\kappa}) e^{itH_{m,\kappa}} e^{-itH_{m',\kappa'}} \mathbb{1}_{\mathbb{R}_+}(H_{m',\kappa'}) = \mathcal{F}_{m,\kappa}^\pm e^{itX^2} \mathcal{F}_{m,\kappa}^{\mp\#} \mathcal{F}_{m',\kappa'}^\pm e^{-itX^2} \mathcal{F}_{m',\kappa'}^{\mp\#}.$$

Let us then observe that

$$\begin{aligned} \mathcal{F}_{m,\kappa}^{\mp\#} \mathcal{F}_{m',\kappa'}^\pm &= \frac{e^{\mp i\frac{\pi}{2}m}}{1 - \zeta e^{\mp i\pi m} \left(\frac{X}{2}\right)^{2m}} \left(\Xi_m(-A) - \zeta \left(\frac{X}{2}\right)^{2m} \Xi_{-m}(-A) \right) \\ &\quad \times \left(\Xi_{m'}(A) - \zeta' \Xi_{-m'}(A) \left(\frac{X}{2}\right)^{2m'} \right) \frac{e^{\pm i\frac{\pi}{2}m'}}{1 - \zeta' e^{\pm i\pi m'} \left(\frac{X}{2}\right)^{2m'}}. \end{aligned} \quad (6.26)$$

Using (4.25) and Lemma B.1, one infers that

$$\begin{aligned} &s\text{-}\lim_{t \rightarrow \pm\infty} e^{itX^2} \left(\Xi_m(-A) - \zeta \left(\frac{X}{2}\right)^{2m} \Xi_{-m}(-A) \right) \\ &\quad \times \left(\Xi_{m'}(A) - \zeta' \Xi_{-m'}(A) \left(\frac{X}{2}\right)^{2m'} \right) e^{-itX^2} \\ &= \left(e^{\pm i\frac{\pi}{2}m} - \zeta \left(\frac{X}{2}\right)^{2m} e^{\mp i\frac{\pi}{2}m} \right) \left(e^{\mp i\frac{\pi}{2}m'} - \zeta' e^{\pm i\frac{\pi}{2}m'} \left(\frac{X}{2}\right)^{2m'} \right) \\ &= e^{\pm i\frac{\pi}{2}(m-m')} \left(1 - \zeta e^{\mp i\pi m} \left(\frac{X}{2}\right)^{2m} \right) \left(1 - \zeta' e^{\pm i\pi m'} \left(\frac{X}{2}\right)^{2m'} \right). \end{aligned}$$

This together with (6.26) yields

$$s\text{-}\lim_{t \rightarrow \pm\infty} e^{itX^2} \mathcal{F}_{m,\kappa}^{\mp\#} \mathcal{F}_{m',\kappa'}^{\pm} e^{-itX^2} = \mathbb{1} \tag{6.27}$$

which directly implies the statement. □

7. Continuous Spectrum of H'_0

In this section, we mimic the computations and results of the previous section analyzing the family of operators H'_0 .

7.1. Resolvent

From now on, we consider $\nu \in \mathbb{C} \cup \{\infty\}$ and $x \in \mathbb{R}_+$. Let us also fix $k \in \mathbb{C}$ with $\text{Re}(k) > 0$, and $-k^2 \notin \sigma(H'_0)$. Our first aim is to compute the integral kernel of the resolvent of H'_0

$$R'_0(-k^2) := (H'_0 + k^2)^{-1}.$$

For that purpose, recall first that the map $x \mapsto \mathcal{K}_0(kx)$ satisfies (6.1) and belongs to $L^2(\mathbb{R}_+)$. We then consider the map

$$x \mapsto \mathcal{I}_0(kx) + \frac{\pi}{2(\gamma + \ln(\frac{k}{2}) - \nu)} \mathcal{K}_0(kx),$$

and infer from (A.2) and (A.8) that this map and the map $x \mapsto \mathcal{I}_0(kx)$ satisfy the equation (6.1) as well as the equation (2.5) or (2.6) around 0. In addition, their Wronskian is equal to k .

From the general theory of Sturm–Liouville operators, as recalled for example in [7, Prop. A.1], one deduces that the kernel of $R'_0(-k^2)$ is given by the expression

$$R'_0(-k^2; x, y) = \frac{1}{k} \begin{cases} \left(\mathcal{I}_0(kx) + \frac{\pi}{2(\gamma + \ln(\frac{k}{2}) - \nu)} \mathcal{K}_0(kx) \right) \mathcal{K}_0(ky) & \text{if } 0 < x < y, \\ \left(\mathcal{I}_0(ky) + \frac{\pi}{2(\gamma + \ln(\frac{k}{2}) - \nu)} \mathcal{K}_0(ky) \right) \mathcal{K}_0(kx) & \text{if } 0 < y < x. \end{cases}$$

Let us also observe that the following relation holds:

$$R'_0(-k^2) = R_0(-k^2) + \frac{1}{2k^2(\gamma + \ln(\frac{k}{2}) - \nu)} P_0(-k^2), \tag{7.1}$$

where $P_0(-k^2)$ is the projection defined in (4.3). Hence, $R'_0(-k^2)$ is well defined except for $\nu = \gamma + \ln(\frac{k}{2})$. This restriction corresponds to the eigenvalue of H'_0 , as already mentioned in Proposition 5.1. Note also that since $R'_0(-k^2)$ is a rank one perturbation of $R_0(-k^2)$, one again infers that $[0, \infty[$ belongs to the spectrum of H'_0 . This justifies *a posteriori* our choice for $\text{Re}(k) > 0$.

7.2. Boundary Value of the Resolvent and Spectral Density

Since $[0, \infty[$ belongs to the spectrum of H_0^ν , it is natural to mimic the computations performed in Sect. 4.2. Note that it will be convenient to use the function \mathcal{Y}_0 , as introduced in Sect. A.6. Note also that since the special case $\nu = \infty$ has already been treated in Sect. 4, when considering the operator H_0 , we shall not consider it again.

Proposition 7.1. *Let $\nu \in \mathbb{C}$, and let $k > 0$.*

(i) *If ν is not an exceptional value, then the boundary values of the resolvent*

$$R_0^\nu(k^2 \pm i0) := \lim_{\epsilon \searrow 0} R_0^\nu(k^2 \pm i\epsilon)$$

exist in the sense of operators from $\langle X \rangle^{-s} L^2(\mathbb{R}_+)$ to $\langle X \rangle^s L^2(\mathbb{R}_+)$ for any $s > \frac{1}{2}$, uniformly in k on each compact subset of \mathbb{R}_+ . The kernel of $R_0^\nu(k^2 \pm i0)$ is given for $0 < x \leq y$ by

$$\begin{aligned} & R_0^\nu(k^2 \pm i0; x, y) \\ &= \frac{\pm i}{k \left(\gamma + \ln \left(\frac{k}{2} \right) - \nu \mp i \frac{\pi}{2} \right)} \left(\left(\gamma + \ln \left(\frac{k}{2} \right) - \nu \right) \mathcal{J}_0(kx) - \frac{\pi}{2} \mathcal{Y}_0(kx) \right) \mathcal{H}_0^\pm(ky), \end{aligned}$$

and the same expression with the role of x and y exchanged for $0 < y < x$.

(ii) *If ν is an exceptional value, then the same statement holds for k uniformly on each compact subset of $\mathbb{R}_+ \setminus \{2e^{\operatorname{Re}(\nu) - \gamma}\}$.*

Proof. The starting point for the proof of this statement is formula (7.1). In addition, since the first term on the r.h.s. of this equality has already been treated in Sect. 4.2, we shall concentrate on the second term only. For that purpose and as in the proof of Theorem 4.2, we consider for $s > \frac{1}{2}$ and $x, y > 0$, the expression

$$\langle x \rangle^{-s} \mathcal{K}_0 \left(\sqrt{-k^2 \mp i\epsilon} x \right) \mathcal{K}_0 \left(\sqrt{-k^2 \mp i\epsilon} y \right) \langle y \rangle^{-s}. \quad (7.2)$$

By the estimate provided in (4.9), one easily infers that this kernel belongs to $L^2(\mathbb{R}_+ \times \mathbb{R}_+)$. In addition, since this kernel converges pointwise to

$$\pm i \langle x \rangle^{-s} \mathcal{H}_0^\pm(kx) \mathcal{H}_0^\pm(ky) \langle y \rangle^{-s}$$

one concludes by the Lebesgue Dominated Convergence Theorem that this convergence also holds in $L^2(\mathbb{R}_+ \times \mathbb{R}_+)$, which is equivalent to a convergence in the Hilbert–Schmidt norm. Note that the uniform convergence in k on compact subsets of \mathbb{R}_+ can be directly checked, as well as the convergence of the prefactors, as long as this factor has no singularity.

For the computation of the kernel of $R_0^\nu(k^2 \pm i0)$, one has for $0 < x \leq y$,

$$\begin{aligned} & R_0^\nu(k^2 \pm i0; x, y) \\ &= R_0(k^2 \pm i0; x, y) - \frac{\pi}{2k(\gamma + \ln(\frac{\mp ik}{2}) - \nu)} \mathcal{H}_0^\pm(kx) \mathcal{H}_0^\pm(ky) \\ &= \pm \frac{i}{k} \left(\mathcal{J}_0(kx) \pm i \frac{\frac{\pi}{2}}{(\gamma + \ln(\frac{k}{2}) - \nu \mp i\frac{\pi}{2})} \mathcal{H}_0^\pm(kx) \right) \mathcal{H}_0^\pm(ky) \\ &= \frac{\pm i}{k(\gamma + \ln(\frac{k}{2}) - \nu \mp i\frac{\pi}{2})} \left((\gamma + \ln(\frac{k}{2}) - \nu) \mathcal{J}_0(kx) - \frac{\pi}{2} \mathcal{Y}_0(kx) \right) \mathcal{H}_0^\pm(ky), \end{aligned}$$

as expected. For $0 < y \leq x$, the same expression can be obtained, with the role of x and y exchanged. \square

Based on the previous result, the corresponding spectral density can now be computed, namely if ν is not exceptional for any $k > 0$ and for any $s > \frac{1}{2}$, one has

$$p_0^\nu(k^2) := \frac{1}{2\pi i} \left(R_0^\nu(k^2 + i0) - R_0^\nu(k^2 - i0) \right) \in \mathcal{B}(\langle X \rangle^{-s} L^2(\mathbb{R}_+), \langle X \rangle^s L^2(\mathbb{R}_+)).$$

If ν is exceptional, the same formulas hold once a suitable restriction on k has been imposed. In the sequel, this restriction will be made tacitly.

Proposition 7.2. *The kernel of the spectral density is given by the following formula:*

$$\begin{aligned} & p_0^\nu(k^2; x, y) \\ &= \frac{((\gamma + \ln(\frac{k}{2}) - \nu) \mathcal{J}_0(kx) - \frac{\pi}{2} \mathcal{Y}_0(kx)) ((\gamma + \ln(\frac{k}{2}) - \nu) \mathcal{J}_0(ky) - \frac{\pi}{2} \mathcal{Y}_0(ky))}{\pi k ((\gamma + \ln(\frac{k}{2}) - \nu)^2 + (\frac{\pi}{2})^2)}. \end{aligned}$$

Proof. For $0 < x \leq y$, one has

$$\begin{aligned} & 2\pi k p_0^\nu(k^2; x, y) \\ &= \frac{1}{\gamma + \ln(\frac{k}{2}) - \nu - i\frac{\pi}{2}} \left((\gamma + \ln(\frac{k}{2}) - \nu) \mathcal{J}_0(kx) - \frac{\pi}{2} \mathcal{Y}_0(kx) \right) \mathcal{H}_0^+(ky) \\ &+ \frac{1}{\gamma + \ln(\frac{k}{2}) - \nu + i\frac{\pi}{2}} \left((\gamma + \ln(\frac{k}{2}) - \nu) \mathcal{J}_0(kx) - \frac{\pi}{2} \mathcal{Y}_0(kx) \right) \mathcal{H}_0^-(ky) \\ &= \frac{((\gamma + \ln(\frac{k}{2}) - \nu) \mathcal{J}_0(kx) - \frac{\pi}{2} \mathcal{Y}_0(kx)) ((\gamma + \ln(\frac{k}{2}) - \nu) \mathcal{J}_0(ky) - \frac{\pi}{2} \mathcal{Y}_0(ky))}{\frac{1}{2} (\gamma + \ln(\frac{k}{2}) - \nu - i\frac{\pi}{2}) (\gamma + \ln(\frac{k}{2}) - \nu + i\frac{\pi}{2})}. \end{aligned}$$

Since the role of x and y can be exchanged, one directly gets the statement. \square

7.3. Generalized Hankel Transform

Let us now define the *incoming* and *outgoing Hankel transformations* $\mathcal{F}_0^{\nu\mp}$ given by the kernels

$$\begin{aligned}\mathcal{F}_0^{\nu\mp}(x, y) &:= \sqrt{\frac{2}{\pi}} \left(\mathcal{J}_0(xy) \pm \frac{i\frac{\pi}{2}}{\gamma + \ln\left(\frac{y}{2}\right) - \nu \mp i\frac{\pi}{2}} \mathcal{H}_0^\pm(xy) \right) \\ &= \sqrt{\frac{2}{\pi}} \left(\frac{(\gamma + \ln\left(\frac{y}{2}\right) - \nu) \mathcal{J}_0(xy) - \frac{\pi}{2} \mathcal{Y}_0(xy)}{\gamma + \ln\left(\frac{y}{2}\right) - \nu \mp i\frac{\pi}{2}} \right).\end{aligned}\quad (7.3)$$

Note that we have written two expressions, since they will be both useful later on. Note also since the denominator has a singularity if ν is exceptional, we shall ignore this special case in the sequel.

To have a better picture of the maps $\mathcal{F}_0^{\nu\mp}$, we recall that Mellin–Barnes representation of \mathcal{H}_0^\pm has been provided in (A.25). Based on this formula, the following statement can be proved. Note that $C_b(\mathbb{R})$ denotes the set of continuous and bounded functions on \mathbb{R} .

Lemma 7.3. *The map $\mathcal{H}_0^\pm : C_c(\mathbb{R}_+) \rightarrow L^2(\mathbb{R}_+)$ with kernel*

$$\mathcal{H}_0^\pm(x, y) := \sqrt{\frac{2}{\pi}} \mathcal{H}_0^\pm(xy)$$

continuously extends to a bounded operator of the form $J\Xi_0^\pm(A) = \Xi_0^\pm(-A)J$ with

$$\Xi_0^\pm(t) := \frac{1}{\pi} \Gamma\left(\frac{1+it}{2}\right)^2 e^{i\ln(2)t} e^{\mp\frac{\pi}{2}t} \quad (7.4)$$

and $\Xi_0^\pm \in C_b(\mathbb{R})$.

Proof. The operator $J\mathcal{H}_0^\pm : C_c(\mathbb{R}_+) \rightarrow L^2(\mathbb{R}_+)$ has kernel

$$\sqrt{\frac{2}{\pi}} \frac{1}{x} \mathcal{H}_0^\pm\left(\frac{y}{x}\right) = \frac{1}{2\pi^2} \frac{1}{\sqrt{xy}} \int_{-\infty}^{+\infty} \Gamma\left(\frac{1+it}{2}\right)^2 e^{i\ln(2)t} e^{\mp\frac{\pi}{2}t} \frac{y^{-it}}{x^{-it}} dt.$$

Using [7, Lem. 6.4], one infers that this kernel corresponds to the kernel of the operator defined by $\Xi_0^\pm(A)$, see also the proof of Proposition 4.5.

Clearly, the map $t \mapsto \frac{1}{\pi} \Gamma\left(\frac{1+it}{2}\right)^2 e^{i\ln(2)t} e^{\mp\frac{\pi}{2}t}$ is continuous and locally bounded. To show that it is bounded, let us estimate its asymptotic values as $t \rightarrow \pm\infty$. By a consequence of Lemma A.2, one gets

$$\left| \frac{1}{\pi} \Gamma\left(\frac{1+it}{2}\right)^2 e^{i\ln(2)t} e^{\mp\frac{\pi}{2}t} \right| = 2e^{-\frac{\pi}{2}|t|} e^{\mp\frac{\pi}{2}t} \left(1 + O(t^{-1})\right).$$

One then infers that $\Xi_0^\pm \in C_b(\mathbb{R})$, as expected. It also means that $\Xi_0^\pm(A)$ extends continuously to a bounded operator in $L^2(\mathbb{R}_+)$. \square

Based on the previous result, one directly infers the following statement:

Lemma 7.4. *For any $\nu \in \mathbb{C}$ with ν not exceptional, the operator $\mathcal{F}_0^{\nu\mp}$ extends continuously to the following operator:*

$$\begin{aligned} & J\left(\Xi_0(A) \pm \Xi_0^\pm(A) \frac{i\frac{\pi}{2}}{\gamma + \ln(\frac{X}{2}) - \nu \mp i\frac{\pi}{2}}\right) \\ &= \left(\Xi_0(-A) \pm \Xi_0^\pm(-A) \frac{i\frac{\pi}{2}}{\gamma - \ln(2X) - \nu \mp i\frac{\pi}{2}}\right) J \end{aligned} \quad (7.5)$$

with Ξ_0 defined in (4.18) and Ξ_0^\pm defined in (7.4).

Let us still provide the expression for the adjoint of this operator, namely

$$\begin{aligned} (\mathcal{F}_0^{\nu\mp})^\# &= \left(\Xi_0(-A) \pm \frac{i\frac{\pi}{2}}{\gamma + \ln(\frac{X}{2}) - \nu \mp i\frac{\pi}{2}} \Xi_0^\pm(-A)\right) J \\ &= J \left(\Xi_0(A) \pm \frac{i\frac{\pi}{2}}{\gamma - \ln(2X) - \nu \mp i\frac{\pi}{2}} \Xi_0^\pm(A)\right). \end{aligned} \quad (7.6)$$

To derive alternative formulas for these operators, let us first recall the equality

$$\Gamma\left(\frac{1+it}{2}\right) \Gamma\left(\frac{1-it}{2}\right) = \frac{\pi}{\cosh\left(\frac{\pi}{2}t\right)}, \quad \forall t \in \mathbb{R}, \quad (7.7)$$

and prove the following statement.

Lemma 7.5. *The map $\mathcal{Y}_0 : C_c(\mathbb{R}_+) \rightarrow L^2(\mathbb{R}_+)$ with kernel*

$$\mathcal{Y}_0(x, y) := \sqrt{\frac{2}{\pi}} \mathcal{Y}_0(xy)$$

continuously extends to the bounded operator $iJ\Xi_0(A) \tanh\left(\frac{\pi}{2}A\right)$.

Proof. From the equality $\mathcal{Y}_0 = \mp i(\mathcal{H}_0^\pm - \mathcal{J}_0)$, one infers that

$$\mathcal{Y}_0 = \mp i(\mathcal{H}_0^\pm - \mathcal{F}_0) = \mp iJ(\Xi_0^\pm(A) - \Xi_0(A))$$

with

$$\begin{aligned} \Xi_0^\pm(t) - \Xi_0(t) &= e^{i\ln(2)t} \frac{\Gamma\left(\frac{1+it}{2}\right)}{\Gamma\left(\frac{1-it}{2}\right)} \left(\frac{1}{\pi} \Gamma\left(\frac{1+it}{2}\right) \Gamma\left(\frac{1-it}{2}\right) e^{\mp \frac{\pi}{2}t} - 1\right) \\ &= \Xi_0(t) \left(\frac{1}{\cosh\left(\frac{\pi}{2}t\right)} e^{\mp \frac{\pi}{2}t} - 1\right) = \mp \Xi_0(t) \frac{\sinh\left(\frac{\pi}{2}t\right)}{\cosh\left(\frac{\pi}{2}t\right)}, \end{aligned}$$

which directly leads to the statement. □

Corollary 7.6. *For any $\nu \in \mathbb{C}$ not exceptional, the following alternative description of $\mathcal{F}_0^{\nu\mp}$ and $(\mathcal{F}_0^{\nu\mp})^\#$ holds:*

$$\mathcal{F}_0^{\nu\mp} = \mathcal{F}_0\left(\gamma + \ln\left(\frac{X}{2}\right) - \nu - i\frac{\pi}{2} \tanh\left(\frac{\pi}{2}A\right)\right) \frac{1}{\gamma + \ln\left(\frac{X}{2}\right) - \nu \mp i\frac{\pi}{2}} \quad (7.8)$$

and

$$(\mathcal{F}_0^{\nu\mp})^\# = \frac{1}{\gamma + \ln\left(\frac{X}{2}\right) - \nu \mp i\frac{\pi}{2}} \left(\gamma + \ln\left(\frac{X}{2}\right) - \nu + i\frac{\pi}{2} \tanh\left(\frac{\pi}{2}A\right)\right) \mathcal{F}_0. \quad (7.9)$$

Proof. The statement follows directly from the Definition (7.3), the previous result, and the equality $\tanh(-t) = -\tanh(t)$ for any $t \in \mathbb{R}$. \square

We can now derive additional properties of the operator $\mathcal{F}_0^{\nu\pm}$.

Lemma 7.7. *For any $\nu \in \mathbb{C}$ not exceptional, the equalities $(\mathcal{F}_0^{\nu\pm})^\# \mathcal{F}_0^{\nu\mp} = \mathbb{1}$ hold.*

Proof. The proof consists in computing the product of the terms (7.5) and (7.6) (with the correct sign) and in checking that this product is equal to $\mathbb{1}$. For that purpose, one first observes that for any $t \in \mathbb{R}$, one has

$$\Xi_0(-t)\Xi_0(t) = 1. \quad (7.10)$$

By considering (7.7), one also observes that

$$\Xi_0(-t)\Xi_0^\pm(t) = \Xi_0^\mp(-t)\Xi_0(t) = \frac{e^{\mp\frac{\pi}{2}t}}{\cosh\left(\frac{\pi}{2}t\right)} \quad (7.11)$$

and that

$$\Xi_0^\mp(-t)\Xi_0^\pm(t) = \left(\frac{e^{\mp\frac{\pi}{2}t}}{\cosh\left(\frac{\pi}{2}t\right)} \right)^2. \quad (7.12)$$

Then, by a few algebraic manipulations, one easily reduces the statement to the following equality in the form sense on $C_c(\mathbb{R}_+)$:

$$\mp \frac{1}{\pi} \left[i \frac{e^{\mp\frac{\pi}{2}A}}{\cosh\left(\frac{\pi}{2}A\right)}, \ln(X) \right] - \frac{e^{\mp\frac{\pi}{2}A}}{\cosh\left(\frac{\pi}{2}A\right)} + \frac{1}{2} \left(\frac{e^{\mp\frac{\pi}{2}A}}{\cosh\left(\frac{\pi}{2}A\right)} \right)^2 = 0. \quad (7.13)$$

To check this equality, let us recall that the equality $[iA, \ln(X)] = \mathbb{1}$ holds, once suitably defined. One then infers that

$$\begin{aligned} \mp \frac{1}{\pi} \left[i \frac{e^{\mp\frac{\pi}{2}A}}{\cosh\left(\frac{\pi}{2}A\right)}, \ln(X) \right] &= \frac{1}{\cosh\left(\frac{\pi}{2}A\right)^2} \left(\frac{1}{2} e^{\mp\frac{\pi}{2}A} \cosh\left(\frac{\pi}{2}A\right) \pm \frac{1}{2} e^{\mp\frac{\pi}{2}A} \sinh\left(\frac{\pi}{2}A\right) \right) \\ &= \frac{e^{\mp\frac{\pi}{2}A}}{\cosh\left(\frac{\pi}{2}A\right)} - \frac{1}{2} \frac{(e^{\mp\frac{\pi}{2}A})^2}{\cosh\left(\frac{\pi}{2}A\right)^2}, \end{aligned} \quad (7.14)$$

where the equalities $\cosh(y) = \frac{e^y + e^{-y}}{2}$ and $\sinh(y) = \frac{e^y - e^{-y}}{2}$ have been used for the last equality. This final expression leads directly to the equality (7.13). \square

Once again, it is natural to set

$$\mathbb{1}_{\mathbb{R}_+}(H_0^\nu) := \mathcal{F}_0^{\nu\pm} (\mathcal{F}_0^{\nu\mp})^\#$$

which is again a projection.

Proposition 7.8. *Let $\nu \in \mathbb{C}$ be not exceptional. Then, for any $k \in \mathbb{C}$ with $\operatorname{Re}(k) > 0$ and $-k^2 \notin \sigma_p(H_0^\nu)$, the following equalities hold:*

$$R_0^\nu(-k^2) \mathbb{1}_{\mathbb{R}_+}(H_0^\nu) = \mathcal{F}_0^{\nu\pm} (X^2 + k^2)^{-1} (\mathcal{F}_0^{\nu\mp})^\# = \mathbb{1}_{\mathbb{R}_+}(H_0^\nu) R_0^\nu(-k^2).$$

Proof. Let us first prove that

$$(\mathcal{F}_0^{\nu\mp})^\# R_0^\nu(-k^2) - (X^2 + k^2)^{-1} (\mathcal{F}_0^{\nu\mp})^\# = 0 \tag{7.15}$$

which implies the second equality of the statement. By considering the expression for $(\mathcal{F}_0^{\nu\mp})^\#$ provided in (7.9) together with the equality (7.1), one observes that the l.h.s. of (7.15) is equal to

$$\begin{aligned} & \frac{1}{\gamma + \ln\left(\frac{X}{2}\right) - \nu \mp i\frac{\pi}{2}} \left\{ \left(\gamma + \ln\left(\frac{X}{2}\right) - \nu + i\frac{\pi}{2} \tanh\left(\frac{\pi}{2}A\right) \right) \mathcal{F}_0 R_0^\nu(-k^2) \right. \\ & \quad \left. - (X^2 + k^2)^{-1} \left(\gamma + \ln\left(\frac{X}{2}\right) - \nu + i\frac{\pi}{2} \tanh\left(\frac{\pi}{2}A\right) \right) \mathcal{F}_0 \right\} \\ &= \frac{1}{\gamma + \ln\left(\frac{X}{2}\right) - \nu \mp i\frac{\pi}{2}} \left\{ i\frac{\pi}{2} [\tanh\left(\frac{\pi}{2}A\right), (X^2 + k^2)^{-1}] \mathcal{F}_0 \right. \\ & \quad \left. + \frac{1}{2k^2(\gamma + \ln\left(\frac{k}{2}\right) - \nu)} \left(\gamma + \ln\left(\frac{X}{2}\right) - \nu + i\frac{\pi}{2} \tanh\left(\frac{\pi}{2}A\right) \right) \mathcal{F}_0 P_0(-k^2) \right\}. \end{aligned}$$

By a few algebraic computations, one then observes that the term inside the curly bracket would be equal to 0 if

$$\pi k^2 [i \tanh\left(\frac{\pi}{2}A\right), (X^2 + k^2)^{-1}] = -\mathcal{F}_0 P_0(-k^2) \mathcal{F}_0 \tag{7.16}$$

and

$$\left(\ln\left(\frac{X}{k}\right) + i\frac{\pi}{2} \tanh\left(\frac{\pi}{2}A\right) \right) \mathcal{F}_0 P_0(-k^2) \mathcal{F}_0 = 0. \tag{7.17}$$

From now on, let us compute some kernels, by always considering $x, y \in \mathbb{R}_+$. Since by (A.18) $\int_0^\infty \mathcal{J}_0(xy) \mathcal{K}_0(ky) dy = \frac{\sqrt{kx}}{x^2 + k^2}$, one first infers that the kernel of the operator $\mathcal{F}_0 P_0(-k^2) \mathcal{F}_0$ is given by

$$\mathcal{F}_0 P_0(-k^2) \mathcal{F}_0(x, y) = 2k^2(xy)^{1/2} \frac{1}{x^2 + k^2} \frac{1}{y^2 + k^2}. \tag{7.18}$$

On the other hand, the kernel of $\tanh\left(\frac{\pi}{2}A\right)$ is given by (see the proof of [31, Lem. 9.2])

$$i \tanh\left(\frac{\pi}{2}A\right)(x, y) = -\frac{2}{\pi} \text{Pv} \left(\frac{1}{\frac{x}{y} - \frac{y}{x}} \right) (xy)^{-1/2},$$

where Pv denotes the principal value distribution. One then infers the following kernel:

$$\begin{aligned} & \pi k^2 [i \tanh\left(\frac{\pi}{2}A\right), (X^2 + k^2)^{-1}](x, y) \\ &= -2k^2 \text{Pv} \left(\frac{1}{\frac{x}{y} - \frac{y}{x}} \right) (xy)^{-1/2} \left(\frac{1}{y^2 + k^2} - \frac{1}{x^2 + k^2} \right) \\ &= -2k^2(xy)^{1/2} \frac{1}{y^2 + k^2} \frac{1}{x^2 + k^2}. \end{aligned} \tag{7.19}$$

By comparing (7.18) with (7.19), one directly gets (7.16).

To check (7.17), let us first deduce from (7.18) that

$$\left(\ln \left(\frac{X}{k} \right) \mathcal{F}_0 P_0(-k^2) \mathcal{F}_0 \right) (x, y) = 2k^2 \ln \left(\frac{x}{k} \right) (xy)^{1/2} \frac{1}{x^2 + k^2} \frac{1}{y^2 + k^2}. \quad (7.20)$$

On the other hand, since

$$\int \left(\frac{1}{x^2 - y^2} \right) \frac{y}{y^2 + 1} dy = \frac{1}{2(x^2 + 1)} \ln \left(\frac{x^2 + 1}{|x^2 - y^2|} \right),$$

one can easily compute the kernel $-i \frac{\pi}{2} \tanh \left(\frac{\pi}{2} A \right) \mathcal{F}_0 P_0(-k^2) \mathcal{F}_0$ and observe that it corresponds to the expression obtained in (7.20). This finishes the proof of the equality (7.17). \square

7.4. Spectral Projections

Let us start again with a result about the spectral projection corresponding to the eigenvalues of H_0^ν . The proof of the following statement can be mimicked from the proof of Proposition 6.6.

Proposition 7.9. *For any $-k^2 \in \sigma_p(H_0^\nu)$, one has*

$$\mathbb{1}_{\{-k^2\}}(H_0^\nu) = P_0(-k^2).$$

As in Sects. 4.4 or 6.4, we can also consider for any $0 < a < b$ the operator

$$\mathbb{1}_{[a,b]}(H_0^\nu) := 2 \int_{\sqrt{a}}^{\sqrt{b}} p_0^\nu(k^2) k dk,$$

which is bounded from $\langle X \rangle^{-s} L^2(\mathbb{R}_+)$ to $\langle X \rangle^s L^2(\mathbb{R}_+)$ for any $s > \frac{1}{2}$. Its kernel is given for $x, y \in \mathbb{R}_+$ by the expression

$$\begin{aligned} & \mathbb{1}_{[a,b]}(H_0^\nu)(x, y) \\ &= 2 \int_{\sqrt{a}}^{\sqrt{b}} \frac{((\gamma + \ln(\frac{k}{2}) - \nu) \mathcal{J}_0(kx) - \frac{\pi}{2} \mathcal{Y}_0(kx)) ((\gamma + \ln(\frac{k}{2}) - \nu) \mathcal{J}_0(ky) - \frac{\pi}{2} \mathcal{Y}_0(ky))}{\pi((\gamma + \ln(\frac{k}{2}) - \nu)^2 + (\frac{\pi}{2})^2)} dk. \end{aligned} \quad (7.21)$$

One can now obtain a result similar to the one contained in Proposition 6.7.

Proposition 7.10. *For any $0 < a < b$ and any $\nu \in \mathbb{C} \cup \{\infty\}$ not exceptional, one has*

$$\mathbb{1}_{[a,b]}(H_0^\nu) = \mathcal{F}_0^{\nu \pm} \mathbb{1}_{[a,b]}(X^2) (\mathcal{F}_0^{\nu \mp})^\# \quad (7.22)$$

in $\mathcal{B}(L^2(\mathbb{R}_+))$. In addition, $\mathbb{1}_{[a,b]}(H_0^\nu)$ is a projection.

Proof. The proof can be mimicked from the one of Proposition 6.7. The new necessary information is the kernel of $\mathbb{1}_{[a,b]}(H_0^\nu)(x, y)$, which is provided in (7.21), and the equality recalled in (7.3). \square

Finally, observe that the equality

$$\mathbb{1}_\Xi(H_0^\nu) = \mathcal{F}_0^{\nu \pm} \mathbb{1}_\Xi(X^2) (\mathcal{F}_0^{\nu \mp})^\#$$

extends (7.22) to any Borel subset Ξ of \mathbb{R}_+ .

7.5. Møller Operators and Scattering Operator

In this section, we extend the results obtained for the Møller and the scattering operators to the family of operators H_0^ν . As before, we shall consider any $\nu \in \mathbb{C}$ which is not an exceptional value.

For the pair $(H_0^\nu, H_0^{\nu'})$ using the Hankel transformations, we define

$$W_{0;0}^{\nu;\nu';\pm} = \mathcal{F}_0^{\nu\pm} (\mathcal{F}_0^{\nu'\mp})^\# . \tag{7.23}$$

Note that the Møller operators for any pairs $(H_0^\nu, H_{m,\kappa})$ or $(H_{m,\kappa}, H_0^\nu)$ could be defined similarly, but the corresponding expressions can also be directly obtained using the chain rule. For that reason, we shall not analyze them separately.

Based on these definitions and on the results obtained so far, one easily infers some identities:

Proposition 7.11. *The following identities hold:*

$$\begin{aligned} (W_{0;0}^{\nu;\nu';\mp})^\# W_{0;0}^{\nu;\nu';\pm} &= \mathbb{1}_{\mathbb{R}_+}(H_0^{\nu'}), \\ W_{0;0}^{\nu;\nu';\pm} (W_{0;0}^{\nu;\nu';\mp})^\# &= \mathbb{1}_{\mathbb{R}_+}(H_0^\nu), \\ (W_{0;0}^{\nu;\nu';\pm})^\# &= W_{0;0}^{\nu';\nu;\mp}, \\ W_{0;0}^{\nu;\nu';\pm} H_0^{\nu'} &= H_0^\nu W_{0;0}^{\nu;\nu';\pm}. \end{aligned} \tag{7.24}$$

By (7.24), $(W_{0;0}^{\nu;\nu';-})^\#$ can be treated as an inverse of $W_{0;0}^{\nu;\nu';+}$. Therefore, we define the scattering operator as

$$S_{0;0}^{\nu;\nu'} := (W_{0;0}^{\nu;\nu';-})^\# W_{0;0}^{\nu;\nu';+} .$$

Clearly, the scattering operator can be expressed in terms of the Hankel transform:

$$S_{0;0}^{\nu;\nu'} = \mathcal{F}_0^{\nu'+} (\mathcal{F}_0^{\nu-})^\# \mathcal{F}_0^{\nu-} (\mathcal{F}_0^{\nu'+})^\# .$$

To analyze the scattering operator, it is convenient to introduce the operators

$$\mathcal{G}_0^{\nu\mp} := (\mathcal{F}_0^{\nu\mp})^\# \mathcal{F}_0^{\nu\mp} . \tag{7.25}$$

Proposition 7.12. *The following equalities hold:*

$$\mathcal{G}_0^{\nu\mp} = \frac{\gamma + \ln\left(\frac{X}{2}\right) - \nu \pm i\frac{\pi}{2}}{\gamma + \ln\left(\frac{X}{2}\right) - \nu \mp i\frac{\pi}{2}} .$$

Proof. Starting with (7.5) and (7.6), and by considering successively the equalities (7.10), (7.11), and (7.7), one infers that

$$\begin{aligned}
 (\mathcal{F}_0^{\nu\mp})^\# \mathcal{F}_0^{\nu\mp} &= J \left(\mathbb{1} \pm \frac{i\frac{\pi}{2}}{\gamma - \ln(2X) - \nu \mp i\frac{\pi}{2}} \frac{e^{\mp\frac{\pi}{2}A}}{\cosh\left(\frac{\pi}{2}A\right)} \right. \\
 &\quad \pm \frac{e^{\pm\frac{\pi}{2}A}}{\cosh\left(\frac{\pi}{2}A\right)} \frac{i\frac{\pi}{2}}{\gamma - \ln(2X) - \nu \mp i\frac{\pi}{2}} \\
 &\quad \left. + \frac{i\frac{\pi}{2}}{\gamma - \ln(2X) - \nu \mp i\frac{\pi}{2}} \frac{1}{\cosh^2\left(\frac{\pi}{2}A\right)} \frac{i\frac{\pi}{2}}{\gamma - \ln(2X) - \nu \mp i\frac{\pi}{2}} \right) J.
 \end{aligned}$$

By observing then that

$$\begin{aligned}
 &\frac{i\frac{\pi}{2}}{\gamma - \ln(2X) - \nu \mp i\frac{\pi}{2}} \frac{e^{\mp\frac{\pi}{2}A}}{\cosh\left(\frac{\pi}{2}A\right)} \\
 &= \frac{e^{\mp\frac{\pi}{2}A}}{\cosh\left(\frac{\pi}{2}A\right)} \frac{i\frac{\pi}{2}}{\gamma - \ln(2X) - \nu \mp i\frac{\pi}{2}} \\
 &\quad + \frac{\pi}{2} \frac{1}{\gamma - \ln(2X) - \nu \mp i\frac{\pi}{2}} \left[i \frac{e^{\mp\frac{\pi}{2}A}}{\cosh\left(\frac{\pi}{2}A\right)}, \ln(X) \right] \frac{1}{\gamma - \ln(2X) - \nu \mp i\frac{\pi}{2}}
 \end{aligned}$$

and using (7.14) together with some algebraic manipulations, one directly infers that

$$(\mathcal{F}_0^{\nu\mp})^\# \mathcal{F}_0^{\nu\mp} = J \left(\mathbb{1} \pm 2 \frac{i\frac{\pi}{2}}{\gamma - \ln(2X) - \nu \mp i\frac{\pi}{2}} \right) J = J \frac{\gamma - \ln(2X) - \nu \pm i\frac{\pi}{2}}{\gamma - \ln(2X) - \nu \mp i\frac{\pi}{2}} J.$$

The statement follows then by the definition of J . \square

Let us stress once again that $\mathcal{G}_0^{\nu\mp}$ are simply functions of X . Finally, using the Hankel transformations, one can obtain a diagonal representation of scattering operators. These operators are expressed in terms of the operators (7.25), namely

$$(\mathcal{F}_0^{\nu'-})^\# S_{0;0}^{\nu';\nu'} \mathcal{F}_0^{\nu'+} = \mathcal{G}_0^{\nu-} \mathcal{G}_0^{\nu'+} = (\mathcal{F}_0^{\nu'+})^\# S_{0;0}^{\nu';\nu'} \mathcal{F}_0^{\nu'-}.$$

In the non-exceptional case, the operator H_0^ν generates a bounded one-parameter group, at least on the range of the projection $\mathbb{1}_{\mathbb{R}_+}(H_0^\nu)$, by the formula

$$e^{itH_0^\nu} \mathbb{1}_{\mathbb{R}_+}(H_0^\nu) = \mathbb{1}_{\mathbb{R}_+}(H_0^\nu) e^{itH_0^\nu} := \mathcal{F}_0^{\nu\pm} e^{itX^2} (\mathcal{F}_0^{\nu\mp})^\#.$$

We finally show that $W_{0;0}^{\nu';\nu';\pm}$ correspond to the Møller operators as usually defined by the time-dependent scattering theory.

Proposition 7.13. *For any $\nu, \nu' \in \mathbb{C}$ not exceptional, the Møller operators exist and coincide with the operators defined in (7.23):*

$$s\text{-}\lim_{t \rightarrow \pm\infty} \mathbb{1}_{\mathbb{R}_+}(H_0^\nu) e^{itH_0^\nu} e^{-itH_0^{\nu'}} \mathbb{1}_{\mathbb{R}_+}(H_0^{\nu'}) = W_{0;0}^{\nu';\nu';\pm}.$$

Proof. The proof is parallel to the proof of Proposition 6.10. One first observes that

$$\begin{aligned}
 (\mathcal{F}_0^{\nu\mp})\# \mathcal{F}_0^{\nu'\pm} &= \left(\Xi_0(-A) \pm \frac{i\frac{\pi}{2}}{\gamma + \ln\left(\frac{X}{2}\right) - \nu \mp i\frac{\pi}{2}} \Xi_0^\pm(-A) \right) \\
 &\quad \times \left(\Xi_0(A) \mp \Xi_0^\mp(A) \frac{i\frac{\pi}{2}}{\gamma + \ln\left(\frac{X}{2}\right) - \nu' \pm i\frac{\pi}{2}} \right).
 \end{aligned}$$

Using the explicit expression for the products of Ξ_0 and Ξ_0^\pm as given in (7.10), (7.11), and (7.12), one then infers that for the four maps $t \mapsto \Xi_0(-t)\Xi_0(t)$, $t \mapsto \Xi_0^\pm(-t)\Xi_0(t)$, $t \mapsto \Xi_0(-t)\Xi_0^\mp(t)$, and $t \mapsto \Xi_0^\pm(-t)\Xi_0^\mp(t)$ belong to $C([-\infty, \infty])$ with

$$\begin{aligned}
 \Xi_0(-\infty)\Xi_0(\infty) &= 1 \quad \text{and} \quad \Xi_0(\infty)\Xi_0(-\infty) = 1, \\
 \Xi_0^+(-\infty)\Xi_0(\infty) &= 2 \quad \text{and} \quad \Xi_0^+(\infty)\Xi_0(-\infty) = 0, \\
 \Xi_0^-(-\infty)\Xi_0(\infty) &= 0 \quad \text{and} \quad \Xi_0^-(\infty)\Xi_0(-\infty) = 2, \\
 \Xi_0^+(-\infty)\Xi_0^-(\infty) &= 4 \quad \text{and} \quad \Xi_0^+(\infty)\Xi_0^-(-\infty) = 0.
 \end{aligned}$$

Based on these observations, one directly deduces from Lemma B.1 that

$$\begin{aligned}
 & \lim_{t \rightarrow \pm\infty} e^{itX^2} (\mathcal{F}_0^{\nu\mp})\# \mathcal{F}_0^{\nu'\pm} e^{-itX^2} \\
 &= \lim_{t \rightarrow \pm\infty} e^{itX^2} \left(\Xi_0(-A) \pm \frac{i\frac{\pi}{2}}{\gamma + \ln\left(\frac{X}{2}\right) - \nu \mp i\frac{\pi}{2}} \Xi_0^\pm(-A) \right) \\
 &\quad \times \left(\Xi_0(A) \mp \Xi_0^\mp(A) \frac{i\frac{\pi}{2}}{\gamma + \ln\left(\frac{X}{2}\right) - \nu' \pm i\frac{\pi}{2}} \right) e^{-itX^2} \\
 &= \mathbb{1}.
 \end{aligned}$$

The remaining argument is similar to the one of the proof of Proposition 6.10. □

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Appendix A: Bessel Family for Dimension 1

In this section, we gather several properties of the Bessel family for dimension 1. If no additional restriction is imposed, the parameter m is an arbitrary element of \mathbb{C} , but a special attention is often required when $m \in \{\dots, -2, -1\}$. All the described properties follow from the theory of the usual Bessel family of dimension 2 described in literature. However, in general, we outline an independent derivation.

A.1: The Function \mathcal{I}_m

The *modified Bessel function for dimension 1*, denoted \mathcal{I}_m , is defined by an everywhere convergent power series or by the so-called Schläfli integral representation:

$$\begin{aligned} \mathcal{I}_m(z) &= \sum_{n=0}^{\infty} \frac{\sqrt{\pi} \left(\frac{z}{2}\right)^{2n+m+\frac{1}{2}}}{n! \Gamma(m+n+1)} \\ &= \frac{\sqrt{z}}{\sqrt{2\pi}} \int_0^\pi e^{z \cos(\phi)} \cos(m\phi) d\phi \\ &\quad - \frac{\sqrt{z}}{\sqrt{2\pi}} \sin(\pi m) \int_0^\infty e^{-z \cosh(\beta) - m\beta} d\beta \quad \text{for } \operatorname{Re}(z) > 0. \end{aligned} \quad (\text{A.1})$$

Note that since $z \mapsto 1/\Gamma(z)$ is an entire function, the first expression is meaningful even for $m \in \{\dots, -2, -1\}$. Clearly, the function \mathcal{I}_m is analytic on $\mathbb{C} \setminus]-\infty, 0]$, and it is analytic on \mathbb{C} for $m \in \mathbb{Z} + \frac{1}{2}$.

The function \mathcal{I}_m has the following asymptotics for z near 0:

$$\mathcal{I}_m(z) = \frac{\sqrt{\pi}}{\Gamma(m+1)} \left(\frac{z}{2}\right)^{m+\frac{1}{2}} + O(|z|^{\operatorname{Re}(m)+\frac{5}{2}}). \quad (\text{A.2})$$

On the other hand, for $m \in \mathbb{Z}$, we have

$$\mathcal{I}_{-m}(z) = \mathcal{I}_m(z) = \frac{\sqrt{z}}{\sqrt{2\pi}} \int_0^\pi e^{z \cos(\phi)} \cos(m\phi) d\phi.$$

For any m , the analytic continuation around 0 by the angle $\pm\pi$ multiplies \mathcal{I}_m by a phase factor, namely

$$\mathcal{I}_m(e^{\pm i\pi} z) = e^{\pm i\pi(m+\frac{1}{2})} \mathcal{I}_m(z).$$

The value $\mathcal{I}_m(z)$ is real for $z > 0$ and for $m \in \mathbb{R}$, and more generally, one has

$$\overline{\mathcal{I}_m(z)} = \mathcal{I}_m(\bar{z}).$$

There also exists a generating function, namely for any $t \in \mathbb{C}^\times$ and any $z \in \mathbb{C} \setminus]-\infty, 0]$:

$$\frac{\sqrt{\pi z}}{\sqrt{2}} \exp\left(\frac{z}{2}(t+t^{-1})\right) = \sum_{m=-\infty}^{\infty} t^m \mathcal{I}_m(z).$$

A.2: The Function \mathcal{K}_m

The *MacDonald function for dimension 1*, denoted \mathcal{K}_m , can be defined for $\operatorname{Re}(z) > 0$ by the integral representation

$$\mathcal{K}_m(z) := \frac{\sqrt{z}}{\sqrt{2\pi}} \int_0^\infty \exp\left(-\frac{z}{2}(s+s^{-1})\right) s^{-m-1} ds. \quad (\text{A.3})$$

It extends by analytic continuation onto a larger domain with a possible branch point at 0. Outside of the basic region $\operatorname{Re}(z) > 0$, one can obtain other integral formulas by an appropriate deformation of the contour of integration in (A.3), see [34, Sec. 6.22].

For $m \notin \mathbb{Z}$, the functions \mathcal{I}_m and \mathcal{K}_m are connected by the equality

$$\mathcal{K}_m(z) = \frac{1}{\sin(\pi m)} (\mathcal{I}_{-m}(z) - \mathcal{I}_m(z)). \tag{A.4}$$

For $m \in \mathbb{Z}$, the relation (A.4) can be extended by l'Hôpital's rule:

$$\mathcal{K}_m(z) = \frac{(-1)^{m+1}}{\pi} \left(\frac{d}{dn} \mathcal{I}_n(z) \Big|_{n=-m} + \frac{d}{dn} \mathcal{I}_n(z) \Big|_{n=m} \right). \tag{A.5}$$

By setting $h(n) := \sum_{k=1}^{n-1} \frac{1}{k}$ and $h(1) = 0$, the equality (A.5) leads for $m \in \mathbb{Z}$ to

$$\begin{aligned} \mathcal{K}_m(z) &= (-1)^{m+1} \frac{2}{\pi} \left(\ln \left(\frac{z}{2} \right) + \gamma \right) \mathcal{I}_m(z) \\ &+ \frac{1}{\sqrt{\pi}} \sum_{k=0}^{m-1} (-1)^k \frac{(m-k-1)!}{k!} \left(\frac{z}{2} \right)^{2k-m+\frac{1}{2}} \\ &+ \frac{(-1)^m}{\sqrt{\pi}} \sum_{k=0}^{\infty} \frac{h(k+1) + h(m+k+1)}{k!(m+k)!} \left(\frac{z}{2} \right)^{2k+m+\frac{1}{2}}, \end{aligned} \tag{A.6}$$

where γ is the Euler's constant. Note that only the first and the third terms are present in the special case $m = 0$.

The function \mathcal{K}_m is analytic for $\mathbb{C} \setminus]-\infty, 0]$ and satisfies $\mathcal{K}_{-m}(z) = \mathcal{K}_m(z)$. We also have

$$\overline{\mathcal{K}_m(z)} = \mathcal{K}_{\overline{m}}(\overline{z}),$$

and the value $\mathcal{K}_m(z)$ is real for $z > 0$ and for $m \in \mathbb{R}$ or $m \in i\mathbb{R}$.

The function \mathcal{K}_m has a well-defined asymptotic at infinity, namely for any $\epsilon > 0$ and $|\arg(z)| < \pi - \epsilon$:

$$\mathcal{K}_m(z) = e^{-z} (1 + O(|z|^{-1})). \tag{A.7}$$

The asymptotics near 0 can be obtained from (A.4) and (A.6). We present the asymptotics in the strip $|\operatorname{Re}(m)| < 1$:

$$\mathcal{K}_m(z) = \begin{cases} -\frac{\sqrt{2z}}{\sqrt{\pi}} \left(\ln \left(\frac{z}{2} \right) + \gamma \right) + O\left(|z|^{\frac{5}{2}} \ln(|z|)\right) & \text{if } m = 0, \\ \frac{\Gamma(m)}{\sqrt{\pi}} \left(\frac{z}{2} \right)^{-m+\frac{1}{2}} + O(|z|^{\operatorname{Re}(m)+\frac{1}{2}}) & \text{if } 1 > \operatorname{Re}(m) > 0, \\ \frac{\Gamma(-m)}{\sqrt{\pi}} \left(\frac{z}{2} \right)^{m+\frac{1}{2}} + O(|z|^{-\operatorname{Re}(m)+\frac{1}{2}}) & \text{if } -1 < \operatorname{Re}(m) < 0, \\ \frac{\sqrt{\pi}}{\sin(\pi m)} \left(\frac{z}{2} \right)^{\frac{1}{2}} \left(\frac{(z/2)^{-m}}{\Gamma(1-m)} - \frac{(z/2)^m}{\Gamma(1+m)} \right) + O(|z|^{\frac{5}{2}}) & \text{if } m \in i\mathbb{R}^\times. \end{cases} \tag{A.8}$$

Actually, we will only need the following estimates, valid for all m : For $|z| < 1$, we have

$$|\mathcal{K}_m(z)| \leq \begin{cases} C_0 |z|^{\frac{1}{2}} (1 + |\ln(z)|) & \text{if } m = 0, \\ C_m |z|^{-|\operatorname{Re}(m)|+\frac{1}{2}} & \text{if } m \neq 0, \end{cases}$$

for some constants C_0 and C_m independent of z .

Let us also mention another relation between the functions \mathcal{I}_m and \mathcal{K}_m , namely

$$\mathcal{I}_m(z) = \frac{1}{2} \left(\mathcal{K}_m(e^{-i\pi} z) - e^{i\pi(m-\frac{1}{2})} \mathcal{K}_m(z) \right).$$

A.3: Additional Properties of \mathcal{I}_m and \mathcal{K}_m

As already mentioned in Sect. 3, the functions \mathcal{I}_m and \mathcal{K}_m are solutions of the modified Bessel equation for dimension 1:

$$\left(\partial_z^2 - \left(m^2 - \frac{1}{4} \right) \frac{1}{z^2} - 1 \right) v(z) = 0. \quad (\text{A.9})$$

They form a basis of solutions of this equation.

For the three functions \mathcal{I}_m , \mathcal{I}_{-m} , and \mathcal{K}_m , their respective Wronskians (2.2) can be computed and are independent of z , namely:

$$\begin{aligned} W_z(\mathcal{K}_m, \mathcal{I}_m) &= 1, \\ W_z(\mathcal{K}_m, \mathcal{I}_{-m}) &= 1, \\ W_z(\mathcal{I}_m, \mathcal{I}_{-m}) &= -\sin(\pi m). \end{aligned}$$

Let \mathcal{L}_m denote either the functions \mathcal{I}_m or the function $e^{i\pi m} \mathcal{K}_m$. Then, the following *contiguous relations* are satisfied:

$$\begin{aligned} 2m\mathcal{L}_m(z) &= z\mathcal{L}_{m-1}(z) - z\mathcal{L}_{m+1}(z), \\ 2m\partial_z\mathcal{L}_m(z) &= \left(m + \frac{1}{2} \right) \mathcal{L}_{m-1}(z) + \left(m - \frac{1}{2} \right) \mathcal{L}_{m+1}(z). \end{aligned}$$

They imply the following *recurrence relations*:

$$\begin{aligned} \partial_z \left(z^{m-\frac{1}{2}} \mathcal{L}_m(z) \right) &= z^{m-\frac{1}{2}} \mathcal{L}_{m-1}(z), \\ \left(\partial_z + \left(m - \frac{1}{2} \right) \frac{1}{z} \right) \mathcal{L}_m(z) &= \mathcal{L}_{m-1}(z), \\ \partial_z \left(z^{-m-\frac{1}{2}} \mathcal{L}_m(z) \right) &= z^{-m-\frac{1}{2}} \mathcal{L}_{m+1}(z), \\ \left(\partial_z + \left(-m - \frac{1}{2} \right) \frac{1}{z} \right) \mathcal{L}_m(z) &= \mathcal{L}_{m+1}(z). \end{aligned}$$

Let us finally observe that for $m = \pm\frac{1}{2}$, the functions \mathcal{I}_m and \mathcal{K}_m coincide with well-known elementary functions:

$$\begin{aligned} \mathcal{K}_{\pm\frac{1}{2}}(z) &= e^{-z}, \\ \mathcal{I}_{-\frac{1}{2}}(z) &= \cosh(z), \\ \mathcal{I}_{\frac{1}{2}}(z) &= \sinh(z). \end{aligned}$$

The simplicity of these relations is one of the motivations for introducing the Bessel family for dimension 1.

A.4: The Function \mathcal{J}_m

The *Bessel function for dimension 1*, denoted \mathcal{J}_m , is defined by the following relations:

$$\begin{aligned} \mathcal{J}_m(z) &= e^{\pm i\frac{\pi}{2}(m+\frac{1}{2})} \mathcal{I}_m(e^{\mp i\frac{\pi}{2}} z), \\ &= \sum_{n=0}^{\infty} \frac{(-1)^n \sqrt{\pi} \left(\frac{z}{2}\right)^{2n+m+\frac{1}{2}}}{n! \Gamma(m+n+1)} \end{aligned} \quad (\text{A.10})$$

$$= \frac{1}{2} \left(e^{-i\frac{\pi}{2}(m+\frac{1}{2})} \mathcal{K}_m(e^{-i\frac{\pi}{2}} z) + e^{i\frac{\pi}{2}(m+\frac{1}{2})} \mathcal{K}_m(e^{i\frac{\pi}{2}} z) \right). \quad (\text{A.11})$$

This function is clearly analytic on $\mathbb{C} \setminus]-\infty, 0]$, and it is analytic on \mathbb{C} for $m \in \mathbb{Z} + \frac{1}{2}$.

Some additional properties of this function are

$$\mathcal{J}_m(e^{\pm i\pi} z) = e^{\pm i\pi(m+\frac{1}{2})} \mathcal{J}_m(z)$$

and

$$\overline{\mathcal{J}_m(z)} = \mathcal{J}_{\overline{m}}(\bar{z}).$$

From the Taylor expansion, one infers that near 0, one has

$$\mathcal{J}_m(z) = \frac{\sqrt{\pi}}{\Gamma(m+1)} \left(\frac{z}{2}\right)^{m+\frac{1}{2}} + O(|z|^{\operatorname{Re}(m)+\frac{5}{2}}).$$

For large z with $|\arg(z)| < \frac{\pi}{2} - \epsilon$ for some $\epsilon > 0$, one also has

$$\mathcal{J}_m(z) = \cos\left(z - \frac{1}{2}\pi m - \frac{1}{4}\pi\right) + e^{|\operatorname{Im}(z)|} O(|z|^{-1}).$$

A.5: The Functions \mathcal{H}_m^{\pm}

The *Hankel functions for dimension 1*, denoted \mathcal{H}_m^{\pm} , are essentially analytic continuations of the function \mathcal{K}_m , one for the lower part and the other one for the upper part of the complex plane. Indeed, the following relations are satisfied:

$$\begin{aligned} \mathcal{H}_m^{\pm}(z) &= e^{\mp i\frac{\pi}{2}(m+\frac{1}{2})} \mathcal{K}_m(e^{\mp i\frac{\pi}{2}} z) \\ &= \frac{e^{\mp i\frac{\pi}{2}(m+\frac{1}{2})}}{\sin(\pi m)} \left(\mathcal{I}_{-m}(e^{\mp i\frac{\pi}{2}} z) - \mathcal{I}_m(e^{\mp i\frac{\pi}{2}} z) \right), \end{aligned}$$

from which one also infers that

$$\mathcal{H}_{-m}^{\pm}(z) = e^{\pm i\pi m} \mathcal{H}_m^{\pm}(z). \quad (\text{A.12})$$

Some additional relations between \mathcal{J}_m and \mathcal{H}_m^{\pm} are:

$$\begin{aligned} \mathcal{J}_m(z) &= \frac{1}{2} (\mathcal{H}_m^+(z) + \mathcal{H}_m^-(z)), \\ \mathcal{J}_{-m}(z) &= \frac{1}{2} (e^{i\pi m} \mathcal{H}_m^+(z) + e^{-i\pi m} \mathcal{H}_m^-(z)), \\ \mathcal{H}_m^{\pm}(z) &= \frac{-e^{\mp i\pi(m+\frac{1}{2})} \mathcal{J}_m(z) + e^{\mp i\frac{\pi}{2}} \mathcal{J}_{-m}(z)}{\sin(\pi m)}. \end{aligned}$$

Inverse Square Potentials

The following asymptotic expansions will also be necessary: from (A.8), one infers that for any θ , $|\arg(z)| < \theta$, as $z \rightarrow 0$,

$$\mathcal{H}_m^\pm(z) = \begin{cases} \pm i e^{\mp i \frac{\pi}{2} m} \frac{\sqrt{2}}{\sqrt{\pi}} z^{\frac{1}{2}} (\ln(z) + \gamma \mp i \frac{\pi}{2}) + O(|z|^{\frac{5}{2}} \ln(|z|)) & \text{if } m = 0, \\ \mp i \frac{\Gamma(m)}{\sqrt{\pi}} \left(\frac{z}{2}\right)^{-m+\frac{1}{2}} + O(|z|^{\operatorname{Re}(m)+\frac{1}{2}}) & \text{if } 1 > \operatorname{Re}(m) > 0, \\ \mp i e^{\mp i \pi m} \frac{\Gamma(-m)}{\sqrt{\pi}} \left(\frac{z}{2}\right)^{m+\frac{1}{2}} + O(|z|^{-\operatorname{Re}(m)+\frac{1}{2}}) & \text{if } -1 < \operatorname{Re}(m) < 0, \\ \mp i \frac{\sqrt{\pi}}{\sin(\pi m)} \left(\frac{z}{2}\right)^{\frac{1}{2}} \left(\frac{(z/2)^{-m}}{\Gamma(1-m)} - \frac{e^{\mp i \pi m} (z/2)^m}{\Gamma(1+m)}\right) + O(|z|^{\frac{5}{2}}) & \text{if } m \in i\mathbb{R}^\times. \end{cases} \quad (\text{A.13})$$

On the other hand, for any $\epsilon > 0$ the following asymptotic formulas are true for $|z| \rightarrow \infty$ with $\arg(z) \notin [\mp \frac{\pi}{2} - \epsilon, \mp \frac{\pi}{2} + \epsilon]$:

$$\mathcal{H}_m^\pm(z) = e^{\pm i(z - \frac{1}{2}\pi m - \frac{1}{4}\pi)} (1 + O(|z|^{-1})). \quad (\text{A.14})$$

A.6: Function \mathcal{Y}_m

The *Neumann function for dimension 1*, denoted \mathcal{Y}_m , is defined by

$$\mathcal{Y}_m(z) = \frac{1}{2i} (\mathcal{H}_m^+(z) - \mathcal{H}_m^-(z)), \quad \mathcal{H}_m^\pm = \mathcal{J}_m \pm i\mathcal{Y}_m.$$

The function \mathcal{Y}_m is especially useful for $m \in \{0, 1, 2, \dots\}$, when we have

$$\begin{aligned} \mathcal{Y}_m(z) &= \frac{2}{\pi} \left(\log\left(\frac{z}{2}\right) + \gamma \right) \mathcal{J}_m(z) \\ &\quad - \frac{1}{\sqrt{\pi}} \sum_{k=0}^{m-1} \frac{(m-k-1)!}{k!} \left(\frac{z}{2}\right)^{2k-m+\frac{1}{2}} \\ &\quad - \frac{1}{\sqrt{\pi}} \sum_{k=0}^{\infty} (-1)^k \frac{h(k+1) + h(m+k+1)}{k!(m+k)!} \left(\frac{z}{2}\right)^{2k+m+\frac{1}{2}}, \end{aligned}$$

with the function h introduced in Sect. A.2.

A.7: Additional Properties of \mathcal{J}_m , \mathcal{H}_m^\pm and \mathcal{Y}_m

As already mentioned in Sect. 3, the functions \mathcal{J}_m , \mathcal{H}_m^\pm , and \mathcal{Y}_m are solutions of the Bessel equation for dimension 1:

$$\left(\partial_z^2 - \left(m^2 - \frac{1}{4}\right) \frac{1}{z^2} + 1 \right) v(z) = 0. \quad (\text{A.15})$$

In addition, their respective Wronskian can be computed and are independent of z , for example

$$\begin{aligned} W_z(\mathcal{J}_m, \mathcal{J}_{-m}) &= -\sin(\pi m), \\ W_z(\mathcal{H}_m^-, \mathcal{H}_m^+) &= 2i. \end{aligned}$$

Let us still observe that for $m = \pm\frac{1}{2}$, the resulting functions coincide with well-known elementary functions:

$$\begin{aligned} \mathcal{J}_{\frac{1}{2}}(z) &= \sin(z), \\ \mathcal{J}_{-\frac{1}{2}}(z) &= \cos(z), \\ \mathcal{H}_{\frac{1}{2}}^{\pm}(z) &= e^{\pm i(z - \frac{\pi}{2})}, \\ \mathcal{H}_{-\frac{1}{2}}^{\pm}(z) &= e^{\pm iz}. \end{aligned}$$

A.8: Integral Identities

The following indefinite integrals follow from the recurrence relations of Sect. A.3:

$$\begin{aligned} &\int_y^\infty \mathcal{K}_m(ax)\mathcal{K}_m(bx)dx \\ &= \frac{1}{a/b - b/a} \left(\frac{1}{b}\mathcal{K}_{m-1}(ay)\mathcal{K}_m(by) - \frac{1}{a}\mathcal{K}_m(ay)\mathcal{K}_{m-1}(by) \right), \quad \operatorname{Re}(a+b) > 0; \\ &\int_y^\infty \mathcal{K}_m(ax)^2 dx \\ &= -\frac{y}{2}\mathcal{K}_m(ay)^2 + \frac{m}{a}\mathcal{K}_m(ay)\mathcal{K}_{m-1}(ay) + \frac{y}{2}\mathcal{K}_{m-1}(ay)^2, \quad \operatorname{Re}(a) > 0. \end{aligned}$$

They imply the following definite integrals:

$$\begin{aligned} \int_0^\infty \mathcal{K}_m(ax)\mathcal{K}_m(bx)dx &= \frac{1}{\sin(\pi m)} \frac{(a/b)^m - (b/a)^m}{\sqrt{ab}(a/b - b/a)}, \\ &m \neq 0, \quad |\operatorname{Re}(m)| < 1, \quad \operatorname{Re}(a+b) > 0, \\ \int_0^\infty \mathcal{K}_0(ax)\mathcal{K}_0(bx)dx &= \frac{2}{\pi} \frac{\ln(a) - \ln(b)}{\sqrt{ab}(a/b - b/a)}, \quad \operatorname{Re}(a+b) > 0, \\ \int_0^\infty \mathcal{K}_m(ax)^2 dx &= \frac{m}{\sin(\pi m)a}, \tag{A.16} \\ &m \neq 0, \quad |\operatorname{Re}(m)| < 1, \quad \operatorname{Re}(a) > 0, \end{aligned}$$

$$\int_0^\infty \mathcal{K}_0(ax)^2 dx = \frac{1}{\pi a}, \quad \operatorname{Re}(a) > 0. \tag{A.17}$$

In the same vein, let us also mention the definite integral

$$\int_0^\infty \mathcal{K}_m(ax)\mathcal{J}_m(bx)dx = \frac{(a/b)^m}{\sqrt{ab}(a/b + b/a)}, \tag{A.18}$$

see [18, Eq. 6.521]. We also derive an additional relation which will be useful later on. For Bessel functions for dimension 2, this result corresponds to [18, Eq. 6.541].

Proposition A.1. *For any $m \in \mathbb{C}$ with $\operatorname{Re}(m) > -1$, one has*

$$\frac{2}{\pi} \int_0^\infty \mathcal{J}_m(xp)\mathcal{J}_m(yp) \frac{1}{(p^2 + k^2)} dp = \frac{1}{k} \begin{cases} \mathcal{I}_m(kx)\mathcal{K}_m(ky) & \text{if } 0 < x < y, \\ \mathcal{I}_m(ky)\mathcal{K}_m(kx) & \text{if } 0 < y < x. \end{cases} \tag{A.19}$$

Proof. For $0 < x < y$, one has by (A.10) and (A.11)

$$\begin{aligned}
 & \frac{2}{\pi} \int_0^\infty \mathcal{J}_m(xp) \mathcal{J}_m(yp) \frac{1}{(p^2 + k^2)} dp \\
 &= \frac{1}{\pi} \int_0^\infty \left(\mathcal{I}_m(e^{-i\frac{\pi}{2}} xp) \mathcal{K}_m(e^{-i\frac{\pi}{2}} yp) + \mathcal{I}_m(e^{i\frac{\pi}{2}} xp) \mathcal{K}_m(e^{i\frac{\pi}{2}} yp) \right) \frac{dp}{(p^2 + k^2)} \\
 &= \frac{1}{\pi} \int_{-\infty}^\infty \mathcal{I}_m(-ixp) \mathcal{K}_m(-iyp) \frac{dp}{(p^2 + k^2)} \\
 &= \frac{2\pi i}{\pi} \frac{\mathcal{I}_m(xk) \mathcal{K}_m(yk)}{2ik} \\
 &= \frac{1}{k} \mathcal{I}_m(kx) \mathcal{K}_m(ky).
 \end{aligned}$$

In the last step, we used the fact that the integral over the semicircle on the upper half-plane $p = Re^{i\phi}$, with $\phi \in [0, \pi]$, goes to zero as $R \rightarrow \infty$. Besides, we have a single singularity of the integrand inside the contour of integration at $p = ik$, which is a simple pole, whose residue has been evaluated.

A similar proof holds for $0 < y < x$. □

A.9: Barnes Identities

Integral representations of Bessel-type functions in terms of the Gamma function are sometimes called *Barnes identities* from the name of their discoverer. Before we present them, let us quote a useful result about asymptotics of the Gamma function given in [4, Cor. 1.4.4], which is a consequence of the Stirling formula.

Lemma A.2. *Let $a, b \in \mathbb{R}$ with $a_1 \leq a \leq a_2$ for two constants a_1 and a_2 . Then, one has*

$$|\Gamma(a + ib)| = \sqrt{2\pi} |b|^{a-\frac{1}{2}} e^{-\frac{\pi}{2}|b|} (1 + O(|b|^{-1})),$$

where the constant implied by the term $O(|\cdot|)$ depends only on a_1 and a_2 .

Let $m \in \mathbb{C}$ and $c \in \mathbb{R}$ with

$$c < \frac{\operatorname{Re}(m)}{2}. \tag{A.20}$$

The following representation is a reformulation of an identity found in [34, Ch. VI.5]:

$$\mathcal{J}_m(x) = \frac{1}{4i\sqrt{\pi}} \int_\gamma \frac{\Gamma(c + \frac{s}{2})}{\Gamma(m + 1 - c - \frac{s}{2})} \left(\frac{x}{2}\right)^{m+\frac{1}{2}-2c-s} ds, \tag{A.21}$$

where γ is a contour which asymptotically coincides with the vertical line $]-\infty, +\infty[$ and passes on the right of $-2c$. Note that by Lemma A.2

$$\left| \frac{\Gamma(c + i\frac{t}{2})}{\Gamma(m + 1 - c - i\frac{t}{2})} \right| \leq C(1 + |t|)^{2c - \operatorname{Re}(m) - 1},$$

hence, the condition (A.20) implies the integrability of the r.h.s. of (A.21).

If we want that the contour is a straight vertical line, we need to assume that $c \in]0, \frac{\operatorname{Re}(m)}{2}[$ (which implies $c > 0$ and $\operatorname{Re}(m) > 0$), and then, we can rewrite (A.21) as

$$\mathcal{J}_m(x) = \frac{1}{4\sqrt{\pi}} \int_{-\infty}^{+\infty} \frac{\Gamma(c + i\frac{t}{2})}{\Gamma(m + 1 - c - i\frac{t}{2})} \left(\frac{x}{2}\right)^{m+\frac{1}{2}-2c-it} dt. \quad (\text{A.22})$$

As shown in the proof of [7, Lem. 6.3], the validity of (A.22) can then be extended in the sense of distribution up to $\operatorname{Re}(m) > -1$ and $0 < c < \operatorname{Re}(m) + 1$. In particular, by choosing $c = \frac{m+1}{2}$, one infers that in the sense of distributions for $\operatorname{Re}(m) > -1$, we have

$$\mathcal{J}_m(x) = \frac{1}{4\sqrt{\pi}} \int_{-\infty}^{+\infty} \frac{\Gamma(\frac{m+it+1}{2})}{\Gamma(\frac{m-it+1}{2})} \left(\frac{x}{2}\right)^{-it-\frac{1}{2}} dt. \quad (\text{A.23})$$

Let us also consider a representation of Hankel functions similar to (A.21) and valid under the condition (A.20). The next formula follows from [34, Sec. 6.5]:

$$\mathcal{H}_m^\pm(x) = \frac{e^{\mp i\pi(m+\frac{1}{2})\pm i\pi c}}{4i\pi^{\frac{3}{2}}} \int_{\gamma'} \Gamma\left(c + \frac{s}{2}\right)\Gamma\left(c - m + \frac{s}{2}\right)e^{\pm i\frac{\pi}{2}s} \left(\frac{x}{2}\right)^{-2c-s+m+\frac{1}{2}} ds, \quad (\text{A.24})$$

where γ' is a contour which asymptotically coincides with the vertical line $] -i\infty, +i\infty[$ and passes on the right of $-2c$ and $-2c + 2m$. As a consequence of Lemma A.2, one infers that

$$\left| \Gamma\left(c + \frac{it}{2}\right)\Gamma\left(c - m + \frac{it}{2}\right)e^{\mp \frac{\pi}{2}t} \right| \leq C(1 + |t|)^{2c-\operatorname{Re}(m)-1} e^{-\frac{\pi}{2}|t|} e^{\mp \frac{\pi}{2}t},$$

which guarantees the integrability of (A.24). Here, we cannot choose γ' to be a straight line. However, if we are satisfied with the interpretation of the integral (A.24) in the sense of distributions, then under conditions $0 < c$ and $\operatorname{Re}(m) < c$, a straight vertical line will work and we can rewrite (A.24) as

$$\mathcal{H}_m^\pm(x) = \frac{e^{\mp i\pi(m+\frac{1}{2})\pm i\pi c}}{4\pi^{\frac{3}{2}}} \int_{-\infty}^{+\infty} \Gamma\left(c + \frac{it}{2}\right)\Gamma\left(c - m + \frac{it}{2}\right)e^{\mp \frac{\pi}{2}t} \left(\frac{x}{2}\right)^{-2c-it+m+\frac{1}{2}} dt.$$

In particular, for $-1 < \operatorname{Re}(m) < 1$ and by setting $c = \frac{\operatorname{Re}(m)+1}{2}$, we obtain after a few manipulations and in the sense of distributions

$$\mathcal{H}_m^\pm(x) = \frac{e^{\mp i\frac{\pi}{2}m}}{4\pi^{\frac{3}{2}}} \int_{-\infty}^{+\infty} \Gamma\left(\frac{-m+1+it}{2}\right)\Gamma\left(\frac{m+1+it}{2}\right)e^{\mp \frac{\pi}{2}t} \left(\frac{x}{2}\right)^{-it-\frac{1}{2}} dt. \quad (\text{A.25})$$

Appendix B: Propagation of the Generator of Dilations

We derive some relations between the generator of dilations A and the multiplication operator X^2 . Note first that

$$e^{it \ln(X)} e^{i\tau A} e^{-it \ln(X)} = e^{i\tau(A-t)}.$$

Therefore, $\ln(X)$ and A satisfy the canonical commutation relations, which determines their properties up to unitary equivalence. It is easy to see that for $\psi \in C([-\infty, \infty])$, one has

$$e^{it \ln(X)} \psi(-A) e^{-it \ln(X)} = \psi(-A + t),$$

and consequently

$$s\text{-}\lim_{t \rightarrow \pm\infty} e^{it \ln(X)} \psi(-A) e^{-it \ln(X)} = \psi(\pm\infty). \quad (\text{B.1})$$

The next lemma contains a similar result with the operator $\ln(X)$ replaced by X^2 . It can be obtained by an abstract argument based on Mourre theory, see [30] for the details. We will give an alternative elementary proof. Note that since $X^2 = \varphi(\ln(X))$ with $\varphi(u) = e^{2u}$ for any $u \in \mathbb{R}$ and since $\varphi' > 0$, the following statement has a flavor similar to the invariance principle in scattering theory.

Lemma B.1. *For any $\psi \in C([-\infty, \infty])$, the following equalities hold:*

$$s\text{-}\lim_{t \rightarrow \pm\infty} e^{itX^2} \psi(-A) e^{-itX^2} = \psi(\pm\infty). \quad (\text{B.2})$$

Proof. Let us first note that it is enough to show that

$$w\text{-}\lim_{t \rightarrow \pm\infty} e^{itX^2} \psi(-A) e^{-itX^2} = \psi(\pm\infty). \quad (\text{B.3})$$

Indeed, this easily follows from the equality:

$$\begin{aligned} & \left\| \left(e^{itX^2} \psi(-A) e^{-itX^2} - \psi(\pm\infty) \right) f \right\|^2 \\ &= \left(f | e^{itX^2} |\psi|^2 (-A) e^{-itX^2} f \right) + \|\psi(\pm\infty) f\|^2 \\ & \quad - \left(\psi(\pm\infty) f | e^{itX^2} \psi(-A) e^{-itX^2} f \right) - \left(e^{itX^2} \psi(-A) e^{-itX^2} f | \psi(\pm\infty) f \right). \end{aligned}$$

We now introduce the unitary transformation $W : L^2(\mathbb{R}) \rightarrow L^2(\mathbb{R}_+)$ by

$$\begin{aligned} (Wf)(x) &= x^{-\frac{1}{2}} f(\ln(x)), \quad \forall x \in \mathbb{R}_+, \\ (W^{-1}g)(t) &= e^{\frac{t}{2}} g(e^t), \quad \forall t \in \mathbb{R}, \end{aligned}$$

and check that

$$W^{-1}X^2W = e^{2Q} \quad \text{and} \quad W^{-1}AW = P$$

with Q and P , the usual self-adjoint operators of position and momentum in $L^2(\mathbb{R})$. Therefore, one infers that

$$W^{-1} e^{itX^2} \psi(-A) e^{-itX^2} W = e^{ite^{2Q}} \psi(-P) e^{-ite^{2Q}}.$$

For $f_1, f_2 \in L^2(\mathbb{R})$ with compact support, one then observes that

$$\begin{aligned}
 & \left(f_1 | e^{ite^{2Q}} \psi(-P) e^{-ite^{2Q}} f_2 \right) \\
 &= \frac{1}{2\pi} \int dx \int d\xi \int dy \overline{f_1(x)} e^{ite^{2x}} \psi(-\xi) e^{i(x-y)\xi} e^{-ite^{2y}} f_2(y) \\
 &= \frac{1}{2\pi} \int dx \int d\xi \int dy \overline{f_1(x)} \psi(-\xi) \exp\left(i(x-y)\left(\xi + t \frac{e^{2x} - e^{2y}}{x-y}\right)\right) f_2(y) \\
 &= \frac{1}{2\pi} \int dx \int d\xi \int dy \overline{f_1(x)} \psi\left(-\xi + t \frac{e^{2x} - e^{2y}}{x-y}\right) e^{i(x-y)\xi} f_2(y). \tag{B.4}
 \end{aligned}$$

Clearly, for any x, y , one has

$$\frac{e^{2x} - e^{2y}}{x - y} > 0,$$

and thus for $y \in \text{supp} f_2$ and $x \in \text{supp} f_1$, there exists a strictly positive c_0 , such that

$$\frac{e^{2x} - e^{2y}}{x - y} \geq c_0 > 0.$$

Finally, one easily obtains that (B.4) converges as $t \rightarrow \pm\infty$ to

$$(f_1 | \psi(\pm\infty) f_2),$$

which shows (B.3) by a density argument. \square

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Jan Dereziński

Department of Mathematical Methods in Physics, Faculty of Physics

University of Warsaw

Pasteura 5

02-093 Warszawa

Poland

e-mail: jan.derezinski@fuw.edu.pl

Serge Richard

Graduate school of mathematics

Nagoya University

Chikusa-ku

Nagoya 464-8602

Japan

e-mail: richard@math.nagoya-u.ac.jp

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