

# Information geometric foundations of quantum theory

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Kogóż to z nas tonący nie wiół wrak?  
Któż z nas zaprzeczyć może że ułomny?  
Kogóż nie łudził oślepiiony ptak?  
Kogóż w bezludzie nie wiół pies bezdomny?

A przecież wciąż przyciąga strefa ogrodzona  
I ogrodzona nie bez celu – trzeba wierzyć  
To nie my w Zonie, to nam odebrana Zona  
Nam ją niepewnym, ale własnym krokiem mierzyć  
Póki nadziei gorycz wreszcie nie pokona.

Dlatego, mimo druty, wieże i strażnice  
Tam chcemy dotrzeć, gdzie nam dotrzeć zabroniono  
Bezżyteczne, śmieszne posiąść tajemnice  
Byleby jeszcze raz gorączką tęsknot płonąć  
Nim podmuch jakiś strzepnie chwiejne potylice.

Droga okrężna może być i oszukańcza,  
Może nas wiedzie szalbierz chciwy paru groszy,  
Lecz lepsze to, niż śmierć na wapniejących szańcach  
U progu granic niewidzialnych i aproszy,  
Gdzie ziewa żołnierz tak podobny do skazańca.

Po zatopionych dawno droga to dolinach;  
Pod płytką wodą nieczytelne czasu grypsy:  
Szlak po ikonach, rękopisach, karabinach  
Nad którym wiosło kreśli plusk Apokalipsy;  
Nie po nas płacz – i nie po przodkach – płacz po synach.

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Gdzie nagle dzwonią wyłączone telefony?  
Serdeczna krew snująca w martwym się potoku,  
Bezsilny gniew na obojętność Nieboskłonu  
I magia słów, co chronić ma od złych uroków?

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# Abstract

This thesis is concerned with a new approach to the mathematical foundations of quantum theory. On the basis of noncommutative integration theory and quantum information geometry, a construction of nonlinear quantum kinematics and dynamics is proposed, defined in terms of geometric and algebraic objects that are completely independent of Hilbert spaces.

The starting point is the definition of spaces of quantum states as sets of positive, finite, normal functionals on noncommutative  $W^*$ -algebras. These spaces form a quantum generalisation of sets of finite positive measures on localisable Boole algebras. Geometric structures are introduced by generalised quantum relative entropies, defined as two-point nonsymmetric functions of distance on spaces of quantum states. In particular, under suitable differentiability conditions, generalised relative entropies determine uniquely the corresponding riemannian metrics and pairs of affine connections on these spaces. Two classes of generalised relative entropies are especially important from the point of view of information theory, namely those which are monotone with respect to completely positive linear maps between state spaces, and those that satisfy the generalised pythagorean theorem. A detailed review of known results on these two classes is provided in this work. Next, using the noncommutative  $L_p$  spaces, a canonical family of generalised quantum relative entropies which belong to both classes is constructed. This result generalises earlier results to the form independent of the choice of representation. Moreover, it is shown that for the finite dimensional case, this family satisfies the smooth version of the generalised pythagorean theorem, with respect to the metrics and connections that are induced by these entropies.

The approach to quantum dynamics proposed in this work starts with the notion of an instrument, understood as a map between two state spaces conditioned upon the set of control parameters. Two classes of instruments are considered: active and passive. Their goal is to generalise the roles that are played in the standard formalism of quantum mechanics by quantum measurement and unitary evolution, respectively.

As a nontrivial example of active instruments, maps that maximise the constrained generalised relative entropies are considered. Such instruments are interpreted as quantum analogues of procedures of statistical inference. It is shown that, under some additional conditions, such instruments determine the semi-spectral measures representing the results of measurements. This is an analogue of the same property of completely positive linear instruments considered in the standard quantum information theory. Under some further conditions, these measures are supported by the states corresponding to superselection sectors that are specified by a uniquely determined commutative algebra. This solves the problem of choice of ‘preferred’ Hilbert space basis associated with the description of measurement results. Moreover, it is proved that the standard quantum mechanical description of measurement in terms of the von Neumann–Lüders rule is a special case of the entropic instrument. This result is a quantum analogue of the derivation of Bayes’ rule as a special case of the maximisation of relative entropy with constraints.

The discussion of nontrivial examples of passive instruments starts by considering a generalisation of symplectic structure to the manifold of quantum states, induced by the commutator of noncommutative  $W^*$ -algebra. A specific class of instruments, which are predualisations of such one-parameter weakly- $\star$  continuous, but not necessarily unitary, groups of automorphisms of  $W^*$ -algebras that determine the hamiltonian vector fields on the state spaces, is characterised. Next, it is shown how the geometric structures on the state spaces can be expressed in terms of operators acting on the fibres of a Hilbert space bundle over a state space. In particular, the time-dependent multi-point correla-

tion functions are constructed to include corrections to the unitary evolution which come from the geometric structure of the state space. The class of passive instruments representing this evolution is constructed too. These results are based on the theory of unbounded perturbations of unbounded generators of unitary evolutions that represent one-parameter automorphisms of  $W^*$ -algebras.

In addition, it is analysed in what sense and how the quantum mechanical notion of a time-dependent transition amplitude can be expressed and generalised in terms of the proposed approach. A preliminary discussion is also provided regarding the possibility of a nonperturbative formulation of the renormalisation of quantum dynamical models using the quantum geometry of state spaces.

A procedure for constructing of space-times, understood as ‘emergent’ objects, obtained directly from the geometry of quantum states, is also proposed. This procedure is based on the transformation of a quantum riemannian metric to a lorentzian form by a ‘complex rotation’ with respect to a global vector field of a temporal evolution determined by a passive instrument. As a nontrivial (albeit somewhat restricted) model, a space of quantum states, which recovers the Schwarzschild space-time is constructed.

A separate chapter is dedicated to the discussion of the conceptual meaning of mathematical structures and results that are considered and obtained in this work. The discussion is based on the conceptual framework of information theory, probability theory, and the bayesian approach to statistical inference theory. The relationships of the proposed approach to the standard (Hilbert space based), logico-algebraic (orthocomplemented poset based), algebraic ( $C^*$ -algebra based) and operational (semi-spectral measure based) approaches are extensively studied.

Constructions provided in this work are based on a number of results obtained in recent years in the theories of operator algebras and quantum information geometry. These include: construction of the canonical theory of noncommutative integration and noncommutative  $L_p$  spaces (Falcone–Takesaki), construction of the Lie–Poisson manifold structure on preduals of  $W^*$ -algebras (Odzijewicz–Ratiu), construction of a smooth manifold structure on faithful state spaces over arbitrary  $W^*$ -algebras based on noncommutative Orlicz spaces (Jenčová), construction of the theory of perturbation of liouvilleans by unbounded operators (Dereziński–Jakšić–Pillet), and construction of the bundle of Hilbert spaces over preduals of  $W^*$ -algebras (Odzijewicz–Sliżewska). This work is their first application to the problems of foundations of quantum theory and quantum geometry.

This thesis also includes an extensive appendix, which contains a systematic exposition of notions and results of the theory of noncommutative integration on  $W^*$ -algebras. It is a unique elaboration of this research area, covering many results that are otherwise scattered across the literature of the subject.



# Streszczenie

Tematem tej rozprawy jest nowe podejście do problemu matematycznych podstaw teorii kwantowej. W oparciu o niekomutatywną teorię całkowania oraz kwantową geometrię informacji, zaproponowana została konstrukcja nieliniowej kwantowej kinematyki i dynamiki, w której podstawową rolę odgrywają obiekty geometryczne i algebraiczne całkowicie niezależne od przestrzeni Hilberta.

Punktem wyjścia pracy jest zdefiniowanie przestrzeni stanów kwantowych jako zbiorów dodatnich, skończonych, normalnych funkcjonałów na niekomutatywnych  $W^*$ -algebrach, co stanowi kwantowe uogólnienie zbiorów skończonych dodatnich miar na lokalizowalnych algebrach Boole'a. Struktury geometryczne zostają wprowadzone poprzez dwupunktowe niesymetryczne funkcje odległości na tych przestrzeniach, będące uogólnieniem kwantowej względnej entropii. Różne dodatkowe warunki nałożone na przestrzenie stanów i uogólnione entropie prowadzą do różnych specyficznych geometrii. W szczególności, przy odpowiednich warunkach różniczkowości, uogólnione względne entropie determinują jednoznacznie riemannowską metrykę i parę afinicznych koneksji na przestrzeni stanów kwantowych. Szczególnie istotne, z punktu widzenia teorii informacji, są dwie klasy uogólnionych względnych entropii: entropie monotoniczne ze względu na całkowicie dodatnie liniowe odwzorowania przestrzeni stanów, oraz entropie spełniające uogólnione twierdzenie Pitagorasa. W pracy przeprowadzony jest szczegółowy przegląd znanych wyników dotyczących obydwu klas. Następnie zostaje skonstruowana, w oparciu o niekomutatywne przestrzenie  $L_p$ , kanoniczna rodzina uogólnionych kwantowych względnych entropii należących do przecięcia tych dwóch klas. Wynik ten uogólnia znane wcześniej konstrukcje do postaci niezależnej od wyboru reprezentacji. Prócz tego wykazane zostaje, że dla skończonej wymiarowej sytuacji otrzymana rodzina spełnia różniczkową wersję uogólnionego twierdzenia Pitagorasa, względem metryki i koneksji indukowanych przez te entropie.

Zaproponowane w pracy podejście do kwantowej dynamiki wychodzi od pojęcia instrumentu, jako odwzorowania między dwoma przestrzeniami stanów, zależnego od zbioru parametrów kontrolnych. Rozważane są dwie klasy instrumentów: aktywne i pasywne, mające za zadanie uogólnić role pełnione w standardowym formalizmie mechaniki kwantowej przez, odpowiednio, kwantowy pomiar oraz ewolucję unitarną.

W roli nietrywialnego przykładu aktywnych instrumentów zaproponowane zostają nieliniowe odwzorowania maksymalizujące uogólnione kwantowe entropie z więzami. Instrumenty te są zinterpretowane jako kwantowe odpowiedniki procedur statystycznego wnioskowania. Pokazane zostaje, że, przy pewnych dodatkowych założeniach, instrumenty te determinują miary półspektralne reprezentujące wyniki pomiarów, co stanowi odpowiednik tejże własności posiadanej przez całkowicie dodatnie liniowe instrumenty rozważane w standardowej kwantowej teorii informacji. Ponadto, przy pewnych dodatkowych warunkach, nośnikami tych miar są stany odpowiadające sektorom superselekcji zadanym przez jednoznacznie określoną komutatywną algebrę, co rozwiązuje problem wyboru 'preferowanej' bazy w przestrzeni Hilberta związanej z opisem wyników pomiaru. Ponadto udowodnione zostaje, że standardowy kwantowomechaniczny opis pomiaru przy pomocy reguły von Neumanna–Lüdersa jest szczególnym przypadkiem entropowego instrumentu. Wynik ten stanowi kwantowy odpowiednik wyprowadzenia reguły Bayesa jako szczególnego przypadku maksymalizacji względnej entropii z więzami.

Punktem wyjścia do dyskusji nietrywialnych przykładów nieliniowych pasywnych instrumentów jest rozpatrzenie uogólnienia struktury symplektycznej na różnorodność stanów kwantowych, indukowanej przez komutator niekomutatywnej  $W^*$ -algebry. Scharakteryzowana zostaje klasa instrumentów, które są predualizacją jednoparametrowych słabo- $\star$  ciągłych, ale niekoniecznie unitarnych, grup automor-

fizmów  $W^*$ -algebr, oraz determinują hamiltonowskie pola wektorowe na przestrzeniach stanów. Następnie zostaje pokazane w jaki sposób struktury geometryczne na przestrzeni stanów mogą zostać wyrażone pod postacią operatorów działających na włóknach wiązki przestrzeni Hilberta nad nią. W szczególności, skonstruowane zostają zależne od czasu wielopunktowe funkcje korelacji uwzględniające poprawki do unitarnej ewolucji pochodzące od struktury geometrycznej przestrzeni stanów, a także klasa pasywnych instrumentów reprezentujących tę ewolucję. Konstrukcja ta opiera się na teorii nieograniczonych perturbacji nieograniczonych generatorów unitarnych ewolucji reprezentujących jednoparametrowe automorfizmy  $W^*$ -algebr.

Prócz tego, zanalizowano w jakim sensie i w jaki sposób kwantowomechaniczne pojęcie zależnej od czasu amplitudy przejścia daje się wyrazić i uogólnić w ramach badanego podejścia. Wstępnej dyskusji poddane są także możliwości nieperturbacyjnego sformułowania renormalizacji modeli kwantowej dynamiki opartej o kwantową geometrię przestrzeni stanów.

W pracy zaproponowana została też procedura konstrukcji czasoprzestrzeni jako ‘emergentnych’ obiektów, otrzymywanych bezpośrednio z geometrii przestrzeni stanów kwantowych. Procedura ta opiera się na transformacji kwantowej riemannowskiej metryki do postaci lorentzowskiej poprzez ‘zespolony obrót’ względem globalnego pola wektorowego ewolucji czasowej zadanej pasywnym instrumentem. W roli nietrywialnego (choć posiadającego pewne ograniczenia) modelu, skonstruowana zostaje przestrzeń stanów kwantowych, która odtwarza czasoprzestrzeń Schwarzschilda.

Osobny rozdział pracy poświęcony jest dyskusji znaczenia pojęciowego badanych struktur matematycznych i otrzymanych wyników, w oparciu o aparat pojęciowy teorii informacji, teorii prawdopodobieństwa, oraz teorii wnioskowania statystycznego w ujęciu bayesowskim. Szeroko przedyskutowane zostały związki i różnice zaproponowanego podejścia z podejściami: standardowym (opartym na przestrzeniach Hilberta), logiczno-algebraicznym (opartym na częściowych porządkach z ortokomplementacją), algebraicznym (opartym na  $C^*$ -algebrach), oraz operacyjnym (opartym na miarach półspektralnych).

Przeprowadzone w pracy konstrukcje opierają się o szereg rezultatów otrzymanych w ostatnich kilkunastu latach w teorii algebr operatorów i w kwantowej geometrii informacji. Są to między innymi: konstrukcja kanonicznej teorii niekomutatywnego całkowania i niekomutatywnych przestrzeni  $L_p$  (Falcone–Takesaki), konstrukcja struktury rozmaitości Lie–Poissona na predualach  $W^*$ -algebr (Odziejewicz–Ratiu), konstrukcja struktury gładkiej rozmaitości na stanach wiernych na dowolnych  $W^*$ -algebrach oparta na niekomutatywnych przestrzeniach Orlicza (Jenčová), konstrukcja teorii perturbacji liouvilleanów przez nieograniczone operatory (Dereziński–Jakšić–Pillet), konstrukcja wiązki przestrzeni Hilberta nad predualami  $W^*$ -algebr (Odziejewicz–Sliżewska). Niniejsza praca stanowi pierwsze ich zastosowanie w problemach podstaw teorii kwantowej i kwantowej geometrii.

Rozprawa zawiera także obszerny dodatek zawierający systematyczne opracowanie pojęć i rezultatów teorii niekomutatywnego całkowania na  $W^*$ -algebrach. Jest to jedyne takie opracowanie w literaturze przedmiotu, uwzględniające wiele rezultatów rozproszonych w trudno dostępnych publikacjach.

# List of symbols

(...) żaden symbol graficzny nie wyraża myśli dokładniej, niż słowo lub zdanie, którym opisano znaczenie tego symbolu!

Adam Wiśniewski-Snerg,  
*Jednolita teoria czasoprzestrzeni*<sup>1</sup>

|                               |            |   |          |                                      |             |                                       |             |
|-------------------------------|------------|---|----------|--------------------------------------|-------------|---------------------------------------|-------------|
| $P^x$                         | 38 209     | $\Psi_\gamma$                           | 81       | $C_F^\infty$                         | 135         | $(\cdot)_*$                           | 167 225     |
| E                             | 41 110 114 | $c(\omega, h)$                          | 85       | $\{\cdot, \cdot\}$                   | 135 136 137 | $(\cdot)_0$                           | 167         |
| $\mathcal{M}$                 | 55 56 57   | $\omega^h$                              | 85       | $df$                                 | 135         | $(\cdot)_1$                           | 167         |
| $\mathcal{S}$                 | 56 167     | $\widetilde{\phi}^h$                    | 86       | $\mathfrak{X}_k$                     | 135 136 137 | $\mathcal{N}$                         | 167         |
| $T_*$                         | 58         | $\mathfrak{P}_C^\gamma$                 | 87 88 96 | $\varpi$                             | 135         | $\mathfrak{B}(\mathcal{H})$           | 168         |
| $D$                           | 60         | $\mathcal{S}_0$                         | 89 167   | $w_t^h$                              | 137         | $\mathfrak{m}_\omega$                 | 168         |
| f                             | 61 63      | $\mathbf{T}_p\mathcal{M}$               | 89 247   | Ad                                   | 137         | $\mathcal{W}$                         | 168 228     |
| $f^c$                         | 61         | $\mathbf{T}_p\ell_\gamma$               | 89 92    | $h^\alpha$                           | 139         | $\mathcal{W}_0$                       | 168 228     |
| $D_f$                         | 61 62 63   | $\Upsilon$                              | 90 93    | $\mathcal{HM}$                       | 139         | $\mathfrak{n}_\phi$                   | 169         |
| $\chi^2$                      | 61         | $\Upsilon_1$                            | 90       | $K_\omega^\alpha$                    | 140         | supp                                  | 169 169 178 |
| epi                           | 65         | $\Upsilon_\phi$                         | 93       | $\mathfrak{a}$                       | 140         | $\ll$                                 | 169 236     |
| efd                           | 65 66      | $L_\Upsilon$                            | 90 93    | $\mathbf{A}$                         | 140         | $ \phi $                              | 170         |
| $\partial\Psi$                | 65         | $\mathfrak{g}^{\text{FRJ}}$             | 94       | $K^{\alpha,x}$                       | 141         | $\pi$                                 | 170         |
| $\Psi^L$                      | 65         | $\nabla^\gamma$                         | 95       | $\alpha^x$                           | 141         | $\pi_\omega$                          | 170 171     |
| efc                           | 66         | $\mathfrak{g}^f$                        | 95       | $\mathcal{L}(\omega, t)$             | 145         | $\Omega_\omega$                       | 170         |
| $\mathfrak{D}_+^G$            | 66         | $\mathfrak{c}$                          | 97       | $\mathfrak{I}_\mathcal{L}$           | 145         | $\mathcal{H}_\omega$                  | 170 171     |
| $\mathfrak{D}_+^G$            | 66         | $\mathfrak{h}$                          | 97 98    | $(\otimes^n \mathcal{H})\mathcal{M}$ | 147         | $\langle \cdot, \cdot \rangle_\omega$ | 170         |
| $\mathfrak{D}^F$              | 66         | $\nabla^f$                              | 95       | TP $\mathcal{H}$                     | 149         | $[\cdot]_\omega$                      | 171 171     |
| $\tilde{D}_\Psi$              | 67         | $\mathfrak{g}^h$                        | 97 98    | TP $\mathbb{R}$                      | 149         | $\mathcal{N}^\bullet$                 | 171         |
| $W_\Psi$                      | 68 71      | $\mathfrak{J}_\rho^h$                   | 97       | TP $\text{CU}$                       | 149         | $\mathfrak{Z}_\mathcal{N}$            | 172         |
| $\tilde{\mathfrak{P}}_C^\Psi$ | 69         | $\mathfrak{F}$                          | 98       | TP                                   | 149         | face                                  | 173         |
| $\hat{\mathfrak{P}}_C^\Psi$   | 70         | $\mathfrak{g}^{\text{BKM}}$             | 99       | DeQuant                              | 155         | co                                    | 173         |
| $(\ell, \ell^{\text{a}})$     | 72         | $\mathfrak{g}^\gamma$                   | 101      | $\hat{\mathfrak{g}}$                 | 155 156     | $\mathcal{C}_\infty^\alpha$           | 176         |
| $D_\Psi$                      | 72 74      | $d_{\text{Bures}}$                      | 101 150  | $\mathcal{C}$                        | 164         | Lin                                   | 178         |
| $\mathbf{L}_\Psi$             | 72 73      | $\mathcal{M}_{\text{exp}}$              | 105 107  | $\mathbb{I}$                         | 164 234     | $J_\omega$                            | 178 179 179 |
| $\Theta_\Psi$                 | 73         | $\mathbf{S}_{\text{GS}}$                | 105      | $x^*$                                | 164         | $\Delta_\omega$                       | 178 179 179 |
| add                           | 73         | $\mathbf{S}_{\text{vN}}$                | 106      | $\mathcal{C}^o$                      | 165         | $\sigma^\omega$                       | 178 179 223 |
| $\Xi_\Psi$                    | 73         | $\mathfrak{J}$                          | 109      | Aut                                  | 165         | $\mathcal{N}_{\sigma^\omega}$         | 179         |
| adc                           | 73         | $\mathfrak{M}^{\text{U}}$               | 110      | $\mathcal{C}_\alpha$                 | 165         | $K_\omega$                            | 179         |
| $\bar{D}_\Psi$                | 73 74      | $\Lambda$                               | 112      | $\leq$                               | 165 168 226 | $\mathcal{H}_\Omega^h$                | 180         |
| $\tilde{\mathfrak{P}}_C^\Psi$ | 75         | $\mathfrak{P}_Q^{D,\Lambda}$            | 112      | $(\cdot)^+$                          | 166         | $\mathcal{H}^h$                       | 180         |
| $\mathfrak{P}_K^\Psi$         | 76         | $\mathfrak{D}$                          | 113      | Proj                                 | 166         | $\xi_\pi$                             | 180         |
| $D_\gamma$                    | 76 78 79   | $\mathfrak{P}_\mathfrak{D}^{D,\Lambda}$ | 113      | $(\cdot)^{\text{sa}}$                | 166         | $\xi_\pi^\pi$                         | 180         |
| $f_\gamma$                    | 77         | $\mathfrak{P}_\mathfrak{D}^\gamma$      | 113      | $(\cdot)^{\text{asa}}$               | 166         | $J_{\phi,\omega}$                     | 181         |
| $\ell_\gamma$                 | 78         | $\eta$                                  | 132      | $(\cdot)^{\text{uni}}$               | 166         | $\Delta_{\phi,\omega}$                | 181 182 188 |
| $\tilde{\ell}_\gamma$         | 78         | ad                                      | 135 137  | $(\cdot)^{\text{B}}$                 | 166 225     | $V_{\phi,\omega}$                     | 182         |

<sup>1</sup>«(...) no graphic symbol does express thoughts more accurately than a word or phrase that describes the meaning of this symbol!» Adam Wiśniewski-Snerg, *Uniform theory of spacetime* [1874].

|                                  |                         |  |                 |                         |     |
|----------------------------------|-------------------------|--|-----------------|-------------------------|-----|
| $[\phi : \omega]_t$              | 182 182 184 208 223     | $\mathcal{E}$  | 213             | $\mathfrak{T}^\nabla$   | 248 |
| $\frac{\psi}{\phi^\bullet}$      | 183 188 224             | $\hat{T}$  | 214 215         | $\Gamma^\nabla$         | 249 |
| $J_{\mathcal{N}}$                | 186                     | $\hat{\phi}$   | 214 215 220     | $\kappa^\nabla$         | 249 |
| $\pi_{\mathcal{N}}$              | 186                     | $h_\phi$   | 216             | $\exp^\nabla$           | 249 |
| $\mathfrak{L}$                   | 187 224                 | $\mathcal{M}^p(\hat{\mathcal{N}}, \tau_\psi)$              | 216             | $\nabla^\dagger$        | 249 |
| $\mathfrak{R}$                   | 187 224                 | $L_p(\mathcal{N}, \psi)$                                   | 217 218         | $\tilde{\nabla}^\theta$ | 250 |
| $\mathfrak{d}$                   | 188                     | $L_p(\mathcal{N}, \psi^\bullet)$                           | 218             | $(\ell, \ell^\dagger)$  | 251 |
| $\alpha$                         | 191                     | $\llbracket \cdot, \cdot \rrbracket_\psi$                  | 218             | $\mathbf{g}^\Phi$       | 252 |
| $\tilde{\delta}$                 | 193                     | $L_p(\mathcal{N})$   | 186 219 222 222 | $\mathbf{g}^\Psi$       | 252 |
| $V_\alpha$                       | 195                     | grad   | 221             |                         |     |
| $K^\alpha$                       | 195                     | $\mathcal{M}^p(\hat{\mathcal{N}}, \tilde{\tau})$           | 221             |                         |     |
| $\tilde{\mu}^G$                  | 197                     | sup  | 226             |                         |     |
| $C_c$                            | 197                     | inf  | 226             |                         |     |
| $\pi_\alpha$                     | 198                     | $\mathcal{A}$  | 227             |                         |     |
| $u_G$                            | 198                     | sp <sub>S</sub>  | 228             |                         |     |
| $\times_\alpha$                  | 198                     | $\mu$  | 228             |                         |     |
| $\hat{G}$                        | 198                     | $L_p(\mathcal{A})$   | 229 229 236     |                         |     |
| $\pi_{\sigma\psi}$               | 199                     | $L_p(\mathcal{A}, \mu)$                                    | 229 229         |                         |     |
| $u_{\mathbb{R}}$                 | 199                     | $\chi$   | 230             |                         |     |
| $\phi^\lambda$                   | 186 200 201 219 220 224 | $\mathcal{A}^\mu$  | 230             |                         |     |
| $\hat{\mathcal{N}}$              | 200                     | $\mathfrak{U}$   | 232             |                         |     |
| $\mathcal{N}(t)$                 | 201                     | $\mathfrak{U}^0$   | 232             |                         |     |
| $\mathcal{H}(t)$                 | 202                     | $\tilde{\mu}$  | 233             |                         |     |
| $\tilde{\mathcal{N}}$            | 202                     | Meas <sup>+</sup>  | 233             |                         |     |
| $\tilde{\mathcal{H}}$            | 202                     | null   | 233             |                         |     |
| aff                              | 205 206                 | $\mathfrak{U}^{\tilde{\mu}}$                               | 233             |                         |     |
| $\mathcal{M}(\mathcal{N})$       | 205                     | $L_p(\mathcal{X}, \mathfrak{U}(\mathcal{X}), \tilde{\mu})$ | 233             |                         |     |
| $\tau_h$                         | 206 206                 | $\mathcal{A}_{\tilde{\mu}}$                                | 234             |                         |     |
| $\phi_h$                         | 207                     | $\mathfrak{U}_{\text{Borel}}(\mathcal{X})$                 | 238             |                         |     |
| $L_p(\mathcal{N}, \tau)$         | 209 210                 | Meas <sub>*</sub> <sup>+</sup>                             | 238             |                         |     |
| $\mathcal{M}(\mathcal{N}, \tau)$ | 210                     | Rad( $\cdot$ ) <sup>+</sup>                                | 239             |                         |     |
| $\mathfrak{G}_p$                 | 211 211                 | sp <sub>G</sub>  | 240             |                         |     |
| $\mathcal{N}^{\text{ext}}$       | 211                     | C  | 240             |                         |     |
| $\tilde{\tau}$                   | 212                     | $\mathbf{T}_q^{\otimes} \mathcal{M}$                       | 247             |                         |     |
| $\tilde{\phi}$                   | 212                     | $\mathbf{g}$   | 247 248         |                         |     |
| $\tilde{\tau}_x$                 | 213                     | $d_{\mathbf{g}}$   | 247             |                         |     |
| $T$                              | 59 60 213               | $\nabla$   | 248             |                         |     |
| $\mathfrak{n}_T$                 | 213                     | $\mathfrak{t}^\nabla$                                      | 248             |                         |     |
| $\mathfrak{m}_T$                 | 214                     | $\mathbf{R}^\nabla$  | 248             |                         |     |

Apart from this the following notation is used:  $i, e, \pi, \mathbb{N}, \mathbb{R}, \mathbb{C}, \mathbb{Z}$  are standard, and the convention  $\mathbb{R}^+ = [0, \infty[$  is used;  $\text{id}$  denotes identity morphism in the category of objects of a given kind,  $\text{dom}$  denotes domain,  $\text{cod}$  denotes codomain,  $\text{ran}$  denotes image,  $\text{span}_{\mathbb{K}}$  denotes linear span in the vector space over  $\mathbb{K} \in \{\mathbb{R}, \mathbb{C}\}$ ,  $\text{int}$  denotes interior of a topological space, **Set** denotes the category of sets and functions; given a small category  $\mathcal{C}$ ,  $\text{Ob}(\mathcal{C})$  denotes the set of all objects in  $\mathcal{C}$ , while  $\text{Hom}_{\mathcal{C}}(X, Y)$  denotes the set of all arrows between  $X, Y \in \text{Ob}(\mathcal{C})$ ;  $\wp(\mathcal{X}) := \text{Hom}_{\mathbf{Set}}(\mathcal{X}, \{0, 1\})$  is the power set of a set  $\mathcal{X}$ ;  $\#(\mathcal{X})$  denotes the cardinality of a set  $\mathcal{X}$  if  $\mathcal{X}$  is finite and  $+\infty$  otherwise;  $\oplus$  and  $\otimes$  denote direct product and tensor product, respectively, that are defined dependently of the category of objects under consideration;  $\delta_{ij}$  denotes Kronecker's delta, while  $\delta(x - y)$  denotes Dirac's delta;  $G$  denotes group,  $g$  denotes group element, while  $e$  denotes neutral element of a group;  $\mathcal{H}$  and  $\mathcal{K}$  denote Hilbert spaces, while  $\xi, \zeta$  denote Hilbert space elements;  $\text{sp}$  denotes spectrum of an operator on a Hilbert space, while  $\int^\oplus$  denotes direct integral in von Neumann's [1813] sense.

# Chapter 1

## Introduction

Nie mogę ci powiedzieć, kiedy rojenie zmieniło mi się w hipotezę, bo sam nie wiem. Postawiłem na niezdecydowanie, wiesz przecież. Odkrycie moje jest fizyką, należy do fizyki, ale do takiej, której nikt nie zauważył, bo droga wiodła przez tereny ośmieszane, wyzute z wszelkich praw. Trzeba było przecież zacząć od myśli, że słowo może stać się ciałem, że zakłęcie *m o ż e* się zmaterializować – trzeba było dać nura w ten absurd, wejść w związki zakazane, aby dostać się na drugi brzeg, tam gdzie już oczywiście jest równoważność informacji i masy.

Stanisław Lem,  
*Profesor A. Dońda*<sup>1</sup>

The aim of this thesis is to investigate the possibility of developing new mathematical and conceptual foundations for quantum theory, based on quantum information geometry on preduals of  $W^*$ -algebras instead of spectral theory of operators on Hilbert spaces. Apart from recent tendencies to develop various information theoretic approaches to foundations of quantum *mechanics* (understood mathematically as the theory established by von Neumann [1800]), we are strongly motivated by the problem of lack of nonperturbative foundations of predictive quantum field theory and the problem of mathematical and conceptual unification of general relativity with quantum theory. The impressive resistance of the last two problems, despite many attempts to solve them, make us believe that the general mathematical and conceptual formulation of quantum *theory* still waits to be developed. We consider this problem as requiring essentially different perspective and treatment than the problems of: (i) providing information theoretic reconstruction of quantum mechanics; (ii) constructing some sort of quantisation of general relativity theory; (iii) further development of effective perturbative quantum field theories. The main intention underlying this text is to use the conceptual insights from information theory and bayesian approach to statistics, as well as mathematical tools of quantum information geometry *without* relying on the framework of Hilbert spaces, in order to propose a foundational approach to quantum theory, which would be intrinsically geometric and nonlinear, allowing for a new perspective on the problems of nonperturbative predictive quantum theory and quantum space-time geometry.

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<sup>1</sup>«I cannot tell you when it was that my reverie became a hypothesis, for I do not know this myself. You know that I have chosen indecision. My discovery is physics, belongs to physics, but physics overlooked as the road to it led through the lands derided, deprived of all laws. For it springs from the thought that a word may become flesh, that a spell *m a y* hatch into matter—one had to plunge into this absurdity, and enter the forbidden relationships to emerge on the other side, where information and mass are, obviously, equivalent.» Stanisław Lem, *Profesor A. Dońda* [1056].

In what follows we will present the summary of key mathematical and conceptual ideas of our approach. Conceptual aspects will be further discussed in detail in Section 2, while mathematical aspects will be further discussed in the rest of this work. The detailed discussion of the mathematical contents of this thesis, including the new results and open problems, is given in Section 2.3. The purpose of this Chapter is to explain underlying motivations and insights.

## 1.1 Mathematical framework

The idea to quantify the distance between two probability distributions dates back to Laplace and Gauß. Such quantification was found necessary in order to decide, on the basis of a given evidence, which among a *given set* of abstract probability distributions, provides the ‘best’ fit to this evidence. The notion of ‘best’ was defined as *minimal* in a given *distance*, under constraints that represent the problem under consideration. In particular, Laplace’s principle of minimum variance distance [443] produces median as a ‘best’ unconstrained estimate, while the Gauß principle of least square distance [616] produces expectation as a ‘best’ unconstrained estimate.<sup>2</sup>

Consideration of *distances* on the sets of probability distributions may suggest to investigate geometric properties of these sets by means of some suitable distances. However, despite the use of differential geometry in the study of spaces of thermodynamic parameters (initiated already by Riemann [1441], who calculated the Riemann curvature tensor of some thermodynamic variables), this idea has waited long for its time. The geometric considerations of the spaces of probabilities were initiated by Rao’s [1397] and Jeffreys’ [875] observation that Fisher’s matrix [562]  $\mathbf{g}_{ij}(\theta)$  determines a riemannian metric on the space of probability densities. The systematic development of geometric study of families of probability distributions has begun with the works of Chencov [312, 313, 314, 315, 316, 317], who (among other results) introduced a class of affine connections on the space of probability measures, generalised pythagorean theorem for relative entropies (considered as nonsymmetric distances), and provided a characterisation of the class of riemannian metrics and affine connections that are monotone in the category of probabilistic models with coarse grainings as arrows.

The main mathematical structure of resulting theory of *information geometry* are nonsymmetric distance functions (relative entropies), which can play the role of generalisations of squared euclidean (or Hilbert space) distance because they satisfy nonlinear nonsymmetric generalisations of pythagorean and cosine theorems. The Taylor expansion of relative entropy at a given point allows to define additional local differential geometric structures, such as the above riemannian metric.

Chencov’s revolutionary idea was that the geometric properties of spaces of probability distributions, together with their transformation properties encoded in the structure of a suitable category, should be considered as a basis for foundations of statistical inference theory<sup>3</sup>. The idea that forms the basis of this thesis is to apply the same foundational perspective to quantum theory. However, this immediately leads to two problems. First is to find a suitable generalisation of mathematical framework that would fit this purpose. As discussed in details in the following Sections and Chapters, Hilbert spaces «present a picture of structural monotony» [1166], and cannot serve for this purpose. Second problem is to deal with an ocean of conceptual quandaries associated with variety of different conflicting interpretations of statistical inference theory and quantum theory. These two problems are strongly entangled. In particular, a misguided preference of some conceptual setting might lead to inevitable preference of some mathematical structures, and conversely.

The crucial mathematical input, allowing us to use a suitable generalisation of information geometry for quantum *foundational* purposes, comes from the noncommutative integration theory on  $W^*$ -algebras. As opposed to the commutative setting, for which there is almost complete equivalence between the approaches based on the notion of measure and the approaches based on the notion of integral, this is no longer the case in the noncommutative setting.

---

<sup>2</sup>Both minimisation problems are stated in terms of the space of estimated functions and not the space of probability measures. However, the modern approach to variational optimisation problems based on duality (see e.g. [1467, 1061, 172]) allows to consider these problems interchangeably (see e.g. [34]).

<sup>3</sup>Later, but independently, the similar idea was proposed by Amari [36, 37, 38, 39, 40].

We begin with postulating that quantum theoretic kinematics should be constructed as a generalisation of mathematical framework of probability theory, with measurable spaces  $(\mathcal{X}, \mathcal{U}(\mathcal{X}))$  replaced by noncommutative  $W^*$ -algebras  $\mathcal{N}$ , probabilistic models  $\mathcal{M}(\mathcal{X}, \mathcal{U}(\mathcal{X}), \tilde{\mu}) \subseteq L_1(\mathcal{X}, \mathcal{U}(\mathcal{X}), \tilde{\mu})_1^+$  replaced by the quantum models<sup>4</sup>  $\mathcal{M}(\mathcal{N}) \subseteq L_1(\mathcal{N})^+ \cong \mathcal{N}_*^+$ , and additional structures associated with probabilistic models (geometry, estimators) replaced by the structures on  $\mathcal{M}(\mathcal{N})$ . This setting can be also thought of as a generalisation and replacement of the Hilbert space based kinematics of quantum mechanics, with elements of  $\mathcal{M}(\mathcal{N})$  generalising the elements of a Hilbert space  $\mathcal{H}$  and the density operators on  $\mathcal{H}$ , and with the quantum relative entropy functional<sup>5</sup>  $D(\cdot, \cdot)$  on  $\mathcal{M}(\mathcal{N})$  replacing the quantitative and geometric role played by the scalar product  $\langle \cdot, \cdot \rangle$  on  $\mathcal{H}$ . From the conceptual point of view, formulation of the new mathematical foundations of quantum theory on the basis of integration theory on  $W^*$ -algebras and geometry of spaces of these integrals is aimed at removing all mathematical reasons for consideration of Hilbert spaces and measure spaces on the foundational level. In such case, the spectral and semi-spectral measures acting from measurable spaces  $(\mathcal{X}, \mathcal{U}(\mathcal{X}))$  to spaces  $\mathfrak{B}(\mathcal{H})^+$  of positive bounded operators on  $\mathcal{H}$  become irrelevant too, and the same applies to probability distributions, arising from traces of spectral and semi-spectral measures with density operators.

The passage from Hilbert spaces to  $W^*$ -algebras is not a simple generalisation, because it changes the roles played by the noncommutative algebra and by the two point functional acting on what is to be considered the space of quantum states. In the former setting the noncommutative algebra consists of operators in a specific representation, while the scalar product is responsible for defining the expectations of elements of an algebra and for defining the geometric structures on the state space. On the contrary, in the latter setting the noncommutative  $W^*$ -algebra  $\mathcal{N}$  is abstract (so it does not consist of operators), the expectations of elements of this algebra are defined by quantum states,  $\phi : \mathcal{N} \ni x \mapsto \phi(x) \in \mathbb{C}$ , without invoking any two point functional, while the quantum relative entropy  $D(\cdot, \cdot)$  is used exclusively to provide the notion of a distance and to define other geometric structures on the spaces  $\mathcal{M}(\mathcal{N})$ .

In particular, under suitable differentiability conditions imposed on  $\mathcal{M}(\mathcal{N})$  and  $D$ , the second and third order Taylor expansions of  $D(\cdot, \cdot)$  determine a riemannian metric  $\mathbf{g}$  on  $\mathcal{M}(\mathcal{N})$  and a pair of torsion-free affine connections  $(\nabla, \nabla^\dagger)$ , respectively. These objects satisfy the characteristic property of the Norden–Sen geometries<sup>6</sup>, namely

$$\mathbf{g}(\mathbf{t}_c^\nabla(u), \mathbf{t}_c^{\nabla^\dagger}(v)) = \mathbf{g}(u, v) \quad \forall u, v \in \mathbf{T}\mathcal{M}(\mathcal{N}), \quad (1.1)$$

and for all smooth curves  $c : \mathbb{R} \rightarrow \mathcal{M}(\mathcal{N})$ , where  $\mathbf{t}_c^\nabla$  denotes a  $\nabla$ -parallel transport along  $c$ . This is a generalisation of the riemannian geometry. On the other hand, for every model  $\mathcal{M}(\mathcal{N})$  one can consider a quantum distance

$$D_{1/2}(\phi, \psi) = \frac{1}{2} \|\xi_{\pi_{\mathcal{N}}}(\phi) - \xi_{\pi_{\mathcal{N}}}(\psi)\|_{L_2(\mathcal{N})}^2, \quad (1.2)$$

where  $\xi_{\pi_{\mathcal{N}}} : \mathcal{N}_*^+ \rightarrow L_2(\mathcal{N})^+$  is a canonical representation in the canonical Hilbert space  $L_2(\mathcal{N})$ .<sup>7</sup> The space  $L_2(\mathcal{N})$  is an operator algebraic generalisation of the Hilbert–Schmidt space  $\mathfrak{S}_2(\mathcal{H})$ , the association  $\mathcal{N} \rightarrow L_2(\mathcal{N})$  is functorial over the category of  $W^*$ -algebras, while the representation  $\xi_{\pi_{\mathcal{N}}}$  is a generalisation of the map  $\mathfrak{S}_1(\mathcal{H})^+ \ni \rho \mapsto \sqrt{\rho} \in \mathfrak{S}_2(\mathcal{H})^+$ , where  $\mathfrak{S}_1(\mathcal{H})$  is the space of all trace class operators on  $\mathcal{H}$ . These properties are quite impressive. However, the main result that requires to consider  $D(\cdot, \cdot)$  as an important *geometric* structure on  $\mathcal{M}(\mathcal{N})$  is the generalised pythagorean equation,

$$D(\phi, \mathfrak{P}_{\mathcal{Q}}^D(\psi)) + D(\mathfrak{P}_{\mathcal{Q}}^D(\psi), \psi) = D(\phi, \psi) \quad \forall (\phi, \psi) \in \mathcal{Q} \times \mathcal{M}(\mathcal{N}), \quad (1.3)$$

<sup>4</sup>Which are the subsets of the positive cones  $\mathcal{N}_*^+$  of Banach preduals  $\mathcal{N}_*$  of  $W^*$ -algebras  $\mathcal{N}$ , which is equivalent to saying the  $\mathcal{M}(\mathcal{N})$  are sets of finite positive normal functionals on  $W^*$ -algebras.

<sup>5</sup>Strictly speaking, the *quantum distance* is a map  $D(\cdot, \cdot) : \mathcal{N}_*^+ \times \mathcal{N}_*^+ \rightarrow [0, \infty]$  such that  $D(\phi, \psi) = 0 \iff \phi = \psi$ , while the *quantum relative entropy* is a map  $(-D)(\cdot, \cdot) : \mathcal{N}_*^+ \times \mathcal{N}_*^+ \rightarrow [-\infty, 0]$ . See Section 3.2 for more discussion. In this Section we will use these notions interchangeably, meaning  $D$  in all cases.

<sup>6</sup>See Section B.2.

<sup>7</sup>See Section A.2.4.

where  $\mathcal{Q} \subseteq \mathcal{M}(\mathcal{N})$  is a closed convex set such that

$$\mathfrak{P}_{\mathcal{Q}}^D(\psi) := \arg \inf_{\omega \in \mathcal{Q}} \{D(\omega, \psi)\} \quad (1.4)$$

exists and is unique. This property is satisfied by a large family of distances  $D(\cdot, \cdot)$ , see Sections 3.2.2 and 3.2.3. Under suitable additional conditions, (1.3) can be restated in local differential geometric terms of (1.1) as a claim that the nonlinear projection of  $\psi \in \mathcal{M}(\mathcal{N})$  along  $\nabla$ -geodesic onto  $\nabla^\dagger$ -autoparallel  $\nabla^\dagger$ -affine submanifold  $\mathcal{Q} \subseteq \mathcal{M}(\mathcal{N})$  coincides with the constrained relative entropic projection (1.4).<sup>8</sup> On the other hand, the generalised pythagorean theorem (1.3) for (1.2) turns to the pythagorean theorem in the Hilbert space  $L_2(\mathcal{N})$  with (1.4) turning to the orthogonal projection onto convex closed subset of this Hilbert space.

These properties show that the quantum information geometry based on quantum distances  $D(\cdot, \cdot)$  over quantum models  $\mathcal{M}(\mathcal{N})$  allows for a sort of “re-unification” of the riemannian geometry with the geometry of Hilbert spaces. This striking feature leads us to consider the geometries of  $\mathcal{M}(\mathcal{N})$  defined by  $D(\cdot, \cdot)$  as the main structural element for the underlying kinematic setting of new foundations of quantum theory. Like riemannian geometry, quantum information geometry naturally carries infinitesimal and local nonlinear structures, but, like in the Hilbert space geometry, individual points have their own internal structure.

Apart from the above structures, quantum models  $\mathcal{M}(\mathcal{N})$  are naturally equipped with one more geometric structure. It is the Poisson bracket  $\{\cdot, \cdot\}$ , which is induced by the commutator of the  $W^*$ -algebra  $\mathcal{N}$ , and which equips  $\mathcal{N}_*^{\text{sa}}$  and some of its subspaces with the Banach–Lie–Poisson manifold structure (see Section 4.2.1). Hence, while  $D(\cdot, \cdot)$  is a generalisation of the scalar product and is able to reproduce riemannian metric and affine connections (which allows to study Killing vector fields, geodesic flows, etc.),  $\{\cdot, \cdot\}$  provides a generalisation of the symplectic form (which allows to study hamiltonian vector fields, Poisson flows, etc.). While one can postulate various quantum distances  $D(\cdot, \cdot)$  on a given quantum model  $\mathcal{M}(\mathcal{N})$ , the Poisson bracket  $\{\cdot, \cdot\}$  on  $\mathcal{M}(\mathcal{N})$  is determined uniquely by  $\mathcal{N}$ . In this sense,  $\{\cdot, \cdot\}$  is a part of quantum geometry of  $\mathcal{M}(\mathcal{N})$ , but is not an independent elements of construction of quantum kinematics. For commutative  $W^*$ -algebras  $\mathcal{N}$ , the BLP structure trivialises, while the quantum information geometry based on  $D(\cdot, \cdot)$  reduces to the corresponding information geometry of statistical models  $\mathcal{M}(\mathcal{X}, \mathcal{U}(\mathcal{X}), \tilde{\mu})$ , where  $L_\infty(\mathcal{X}, \mathcal{U}(\mathcal{X}), \tilde{\mu}) \cong \mathcal{N}$ . On the other hand, the above setting allows to recover the Kähler geometry of pure states (Hilbert space vectors) modulo phase in a special case, for a specific choice of  $D$  which admits an extension of the derived riemannian metric to pure states (see [470, 1672, 1322]). In this sense, the geometry of quantum mechanics and the geometry of probability theory (called «geometrostatistics» by Kolmogorov, cf. [320]) are just the special cases of quantum geometry of  $\mathcal{M}(\mathcal{N})$ .<sup>9</sup>

We also propose to consider quantum theoretic dynamics as a direct generalisation of the statistical inference theory, with various rules of inference, conditioning, and causality replaced by the general setting of *instruments*. An instrument is defined as a function from the set of ‘control parameters’ to the space of mappings  $\mathcal{M}_1(\mathcal{N}) \rightarrow \mathcal{M}_2(\mathcal{N})$  between two quantum models  $\mathcal{M}_1(\mathcal{N}), \mathcal{M}_2(\mathcal{N}) \subseteq \mathcal{N}_*^+$ . These mappings are not assumed to be linear or completely positive. Instruments provide a generalisation of various forms of quantum dynamics and conditioning, including the description of “quantum measurement” and unitary evolution. We investigate various instruments that are dependent on the geometric structure of quantum models  $\mathcal{M}(\mathcal{N})$ . We are especially concerned with instruments that are generated by constrained minimisation (1.4) of the quantum distance  $D(\cdot, \cdot)$  and with instruments that arise from predualisation of automorphisms  $\{\alpha_t \mid t \in \mathbb{R}\} \subseteq \text{Aut}(\mathcal{N})$  of  $W^*$ -algebras  $\mathcal{N}$ . They provide the key examples of ‘active’ and ‘passive’ classes of instruments, respectively.<sup>10</sup>

<sup>8</sup>See Section B.3.

<sup>9</sup>We will use term ‘quantum information geometry’ to speak of geometries on  $\mathcal{M}(\mathcal{N})$  arising from  $D(\cdot, \cdot)$ , and we will speak of ‘quantum geometry’ to refer to all geometric structures on  $\mathcal{M}(\mathcal{N})$ , including those that arise from  $D(\cdot, \cdot)$  and those that arise from  $\{\cdot, \cdot\}$ .

<sup>10</sup>The idea of two different and independent temporal evolutions, *external*—governed by maximisation of entropy, and *internal*—governed by the energy, each one associated with its own notion of time, is due to Kępiński [966, 967, 968]. Our concept of active and passive instruments is intended as a suitable generalisation of this idea (we were also inspired



An instrument is said to be ‘active’ if it is used for the purposes of inductive inference conditioned on given set of epistemic constraints. Thus, it is a ‘rule of conditioning’ of information states. Bayes’ rule and the von Neumann–Lüders rules are the most standard examples of conditioning rules. (Recently it has been shown that Bayes’ rule is an entropic instrument in the above sense [962, 1609, 1822, 292, 633]. In Sections 4.1.2 and 4.2.3 we prove that the same is true for the von Neumann–Lüders rules. This justifies consideration of entropic instruments as nonlinear active instruments.) Moreover, as we show in Section 4.1.1, under some additional conditions, active instruments determine the semi-spectral measures representing the results of measurements. This is an analogue of the same property of completely positive linear instruments considered in the standard quantum information theory. Under some further conditions, these measures are supported by the states corresponding to superselection sectors that are specified by a uniquely determined commutative algebra. This provides a specific solution to the problem of choice of ‘preferred’ Hilbert space basis associated with the description of measurement results.

An instrument is called ‘passive’ if it is used for the purpose of defining the referential system of causality, understood as a transformation of the information states that is independent of any action of active instruments. This provides a generalisation of the role played by the unitary evolution of density operators in quantum mechanics and the predualised action  $\{\alpha_\star^t : \mathcal{N}_\star^+ \rightarrow \mathcal{N}_\star^+ \mid t \in \mathbb{R}\}$  of the weakly- $\star$  continuous group of automorphisms of a  $W^*$ -algebra  $\mathcal{N}$  in the algebraic approach. This predualisation works also in the commutative case, and one can apply it to automorphisms of an algebra  $L_\infty(\mathcal{X}, \mathcal{U}(\mathcal{X}), \tilde{\mu})$  induced by automorphisms of an underlying boolean algebra  $\mathcal{U}(\mathcal{X})/\mathcal{U}^{\tilde{\mu}}(\mathcal{X})$  (see e.g. [591]) or its sample-space representation. So, the passive instruments can be also considered as a specific method of modelling causality in the sense of statistical inference theory (see e.g. [1286] for a review). The presence of the BLP structure, which is a generalisation of the symplectic structure, leads us to analyse its role in the construction of passive instruments. We characterise the class of predualised automorphisms  $\alpha_\star$  that are also Poisson flows of  $\{\cdot, \cdot\}$ , and show that under certain conditions hamiltonian vector fields of these flows can be expressed by generators of unitary evolutions in the Hilbert space bundle over  $\mathcal{M}(\mathcal{N})$ . Using this correspondence, we construct a new family of nonlinear passive instruments, which arise from the additional perturbations by one forms on the BLP manifold. Trivialisation of the BLP structure for commutative  $W^*$ -algebras shows an important difference between quantum theory and statistical inference theory: in former causality may be analysed in terms of hamiltonians and their perturbations, while in the latter this is impossible. Yet, both have the same conceptual meaning.

Quantum distances  $D(\cdot, \cdot)$  play a crucial role in construction of active instruments that implement the nonlinear inductive inference procedures, generalising widely the setting of Bayes’ and von Neumann–Lüders’ rules. In infinitesimal approximation, they determine riemannian metrics and affine connections acting on the tangent bundle of the model  $\mathcal{M}(\mathcal{N})$ . On the other hand, the BLP structure  $\{\cdot, \cdot\}$  plays a crucial role in construction of passive instruments that implement the nonlinear causality systems, generalising the unitary evolutions. It is genuinely related to the cotangent bundle of the model  $\mathcal{M}(\mathcal{N})$ . In this sense, two independent geometric structures on the quantum models  $\mathcal{M}(\mathcal{N})$  play the key role in the construction of nonlinear examples of two independent forms of dynamics on models  $\mathcal{M}(\mathcal{N})$ .

The question regarding the relationship between those two structures leads us to the final step of our investigations, which amounts to proposing that the *space-times* should be thought of as objects that emerge from information geometric kinematics and causal (passive) dynamics, while the space-time trajectories (not necessarily causal) should arise from the active dynamics. We propose to implement this general postulate by changing the signature of the riemannian metric induced by  $D(\cdot, \cdot)$  to lorentzian form in a way that makes the global vector field of a given passive (causal) evolution timelike.

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by Savvidou’s work [1515, 1516] on two different time evolutions in the Isham–Linden approach to consistent quantum histories). This terminology stems from observation that entropic evolution can be regarded as *active*, because it describes the change of knowledge and represents a *decision*, while energy-governed evolution (such as unitary evolution) can be considered as passive, because it describes invariance properties of knowledge and represents the criteria of judging the quality of decision (e.g. whether a given active instrument is causal or not).

## 1.2 Conceptual motivations

Let us move to the discussion of conceptual justification of this specific choice of mathematical structure and its relationship to quantum mechanics. Our starting point is an opinion that the predictive content of quantum theory should be independent of any other theory, so it should rely neither on any ‘classical theory’ nor on any additional (and chosen *ad hoc*) methods of statistical inference used for probabilistic estimation. This leads us to part with two foundational principles of the orthodox approach to quantum mechanics: quantisation<sup>11</sup> and spectral representation.

The idea of quantisation was originally proposed [464, 467] as a method of construction of quantum theoretic models that relies on some additional (‘classical’) theory, which is a subject of quantisation procedure. But there are many inequivalent variants of this procedure,<sup>12</sup> and most of them are in some sense mathematically ambiguous,<sup>13</sup> while those that are not ambiguous have a seriously restricted applicability. In addition, the ‘classical’ theories subjected to quantisation are quite often devoid of any predictive content (e.g. most of nonabelian Yang–Mills theories), so they are also *ad hoc*. Finally, quantisation seems awkward from the conceptual point of view, if one considers quantum theory as more fundamental than ‘classical’ theory (or if one considers quantum theory as an analogue of probability theory)—why more fundamental theory should rely on the less fundamental one?

On the other hand, the purpose of spectral representation [1794, 1797, 1813] is to construct a probability measure supported over the subsets of sample space that are indexed by eigenvalues of a self-adjoint operator or a commutative algebra of self-adjoint operators. Hence, its use is based on the assumption, or leads to a conclusion, that the predictive validity of quantum theory has to rely on other additional theories, namely probability theory and statistical estimation theory. Moreover, it also reflects a belief (or a postulate) that eigenvalues of self-adjoint operators are mutually independent «in principle observable» [767], which (from the frequentist statistical point of view) justifies using them as labels for subsets of sample space. However, the eigenvalues of self-adjoint operators *are not* directly observable or measurable: they are abstract symbols, while this what is measured consists of values of various experimental variables.<sup>14</sup> The relationship between them is *not* direct, as it relies on the choice of *ad hoc* conventions of statistical inference. The idealistic character of the assumption of “in principle observability” was criticised in length in the literature on operational approach to quantum theory (see Section 2.2 for an extended discussion), while the idealistic character of the assumption of mutual independence of eigenvalues comes under criticism in the literature on renormalisation theory<sup>15</sup>. Eigenvalues, as well as any other abstract mathematical concept, such as expectation values or higher order correlation functions, can be used only to *estimate* the values of some functions that are, more or less directly, related to some selective choice of what is considered as experimental ‘data’. Finally, if the mathematical formalism of kinematics of quantum theory can be shown to be a strict

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<sup>11</sup>More precisely, quantisation is not a *principle* of von Neumann formalisation of quantum mechanics. However, it became the main method of practical quantum theoretic model construction, and is widely considered to be a constitutive component of an orthodox approach to quantum mechanics.

<sup>12</sup>Such as Dirac’s and other canonical quantisations, the Feynman–Kac and other path integral quantisations, the Kirillov–Konstant–Souriau and other geometric quantisations, the Weyl–Wigner–Groenewold–Moyal quantisation, etc. See e.g. [626] for a partial review.

<sup>13</sup>E.g. from the van Hove theorem [1776] it follows that it is impossible to preserve the whole ‘classical’ Poisson bracket algebra under ‘canonical quantisation procedure’, defined as replacing the Poisson brackets between smooth functions on the phase space by commutators between the corresponding operators on the Hilbert space. Moreover, the choice of a particular subalgebra is arbitrary. In general, this means that the equivalent ‘classical’ Poisson algebras of smooth functions on the phase (= sample) space lead to unitarily inequivalent ‘canonically quantised’ noncommutative algebras of operators on a Hilbert space. See also [1589] and references therein.

<sup>14</sup>In our opinion, the role of eigenvalues in foundations of quantum mechanics is quite similar to the role of ‘virtual particles’ in the perturbative approach to quantum field theory: they are abstract symbols devoid of any *direct* relationship to epistemic evidence, but are kept due to their applicability in calculations within available mathematical formalism (in fact of lack of alternative to it) *and* due to the ontological conceptual content they carry. In the case of perturbative quantum field theory the corresponding formalism is given by the Fok–Cook Hilbert space representation, while the conceptual content is the ontological interpretation in terms of particles.

<sup>15</sup>In principle, introduction of an epistemically motivated cut-off on unbounded ‘ideal’ (and predictively unreliable) self-adjoint operators leads to a mutual dependence of eigenvalues of a predictive, renormalised, operator. (A similar problem, but encountered for expectation values, will be analysed in details in Section 4.1.3.)

generalisation of formalism of probability theory, then the predictive applicability of former should require no further translation to the latter.

In this sense, denying the principles of quantisation and spectral representation we aim at removing dependence of quantum theory on auxiliary and, in our opinion, foundationally irrelevant theories: classical mechanics and probability theory.

In order to explain better the reasons behind, and consequences of, the conceptual change we propose, let us take a short comparative look at the mathematical and conceptual foundations of quantum mechanics and of frequentist statistics. The origins of both theories seem to be heavily influenced by the postulate that the full content of a predictive theory can be *reduced* to manipulation with tables (matrices) of experimental data (numbers). Most clearly these ideas were expressed by the founders of the modern formulations of these theories: Heisenberg [767] postulated<sup>16</sup> that all experimental data should be in principle reducible in terms of real eigenvalues of the matrices with complex valued entries, while Fisher [561, 562] postulated that all experimental data should be in principle reducible in terms of sufficient statistics over sample spaces. These theories have been elevated to their nowadays orthodox mathematical formalisms in the monographs of von Neumann [1800] and Kolmogorov [972], respectively. The latter text established the abstract measure theory as a framework for probability theory (following earlier ideas of [1623]), while the former one established the abstract Hilbert space theory as a framework for quantum mechanics (following earlier ideas of [1794]). However, none of these books provided solid *conceptual* grounds for frameworks provided there.<sup>17</sup> This has fruitfully contributed to long debates on the interpretations of probability theory and quantum theory. More importantly, none of these books have offered general and unambiguous methods of quantitative model construction and conditioning<sup>18</sup> that would be based on the presented formalisms, and would reflect the quantitative descriptions and constraints of the practical (experimental) situations under consideration in a unique way. This has led to long debates on *the correct* methods of statistical inference and *the correct* methods of construction of ‘interacting’ quantum theories (and quantum field theories in particular).

It became soon clear that the idea underlying the original postulates of Fisher and Heisenberg is idealistic (metaphysical) and, more importantly, it is invalid in most of situations of *practical predictive concern*. In quantum theory this has been shown ultimately by Wintner [1873] and Wielandt [1848], who proved that the Born–Jordan–Dirac–Heisenberg commutation relations cannot be represented by bounded operators in Hilbert space. In statistical theory this has been shown by Darmais [428], Koopman [974] and Pitman [1345], who proved that exponential families are unique probabilistic models that admit sufficient statistics whose dimension remains bounded with the increase of the sample size. Thus, the experimental data alone are insufficient to determine the structure of, as well as quantitative predictions drawn from quantum and statistical models. In order to construct the model and to draw inferences from it, some additional principles are necessary. This leads to the problems of justification of the choice of specific principles<sup>19</sup> and reevaluation of the conceptual and mathematical foundations of the respective theories in the light of the key role played by these principles for predictive verifiability of those theories.

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<sup>16</sup>As it was made precise by Born and Jordan [208, 207].

<sup>17</sup>[972] explicitly states that measure theoretic foundations of probability theory are not justified by frequencies, but also does not give any preference to a subjective, or other, point of view. [1800] contains allusions to several different interpretations, including ontological idealism (while speaking about “quantum systems” as ontological objects), subjectivism (while speaking of “quantum measurements” as provided by consciousness), and operationalism (while associating projections with experimental filters).

<sup>18</sup>[972] introduced conditional expectations as a method of conditioning, while [1800] introduced projection postulates (as shown by Nakamura and Umegaki [1226], a version of projection postulate, given by the weak Lüders’ rule, is a noncommutative analogue of a conditional expectation). However, the practical applicability of both methods is strongly restricted: most of epistemic situations under consideration *do not* provide a type of information that can be handled by these tools. For example, the information specified by means of average values of some experimental quantities, estimated theoretically by means of expectation, cannot be recast using a projection on a closed convex subset of a Hilbert space.

<sup>19</sup>For example, in quantum information theory one deals with the problem of construction of a specific semi-spectral measure corresponding to a given epistemic situation (see Section 2.2), while in effective quantum field theory one deals with the problem of choice of the renormalisation method encoding known epistemic constraints.

However, the pragmatic success of several models based on adaptation of *ad hoc* additional principles has overshadowed these problems. As a result, the failure of initial postulates has not led to downfall of the theories built upon them, despite clear further indications of the foundational crisis. In quantum theory the problem of multitude of different unitarily inequivalent Hilbert space representations of a given abstract algebra generated by commutation relations was discovered<sup>20</sup>, together with the closely related problem of ambiguity of construction of ‘interacting dynamics’. The former discovery implied that the eigenvalues of operators, which are quantities dependent on the choice of a specific representation, cannot be given absolute (universal, objective, ontological) status, while the latter was indicated by the widespread appearance of infinities in quantum field theory. In statistical theory the problems of lack of general criteria of model selection, optimal estimation, and the choice among different inequivalent statistical tests and methods of introducing causality, arose. For the wide range of cases, the solution that have been adapted to *pragmatically* cure this issue was: in quantum theory—to stay within one Hilbert space representation (provided by Schrödinger [1534] and Heisenberg–Born–Jordan [767, 208, 207] in the finite dimensional case and by Fok [567] and Cook [372] in the infinite dimensional case) and to construct predictive dynamical models by means of perturbative expansions; in statistical theory—to stay within the range of asymptotic applicability of exponential and normal models (with sample size assumed to be infinite), and to use such tools as Fisher’s maximum likelihood principle [561], or, more generally, Le Cam’s local asymptotic normality theory [1043, 1044]. The unsolved critical issue of lack of general methods of control over the quality of these approximations were swept under the carpet by conceptual downgrading the problem of construction of quantum dynamics and the problem of choice of statistical inference procedure that corresponds to a *given* epistemic situation to the status of *technical* (and no longer *foundational*) issues, governed by the set of *pragmatic* rules.

As a result, these two fields of research are plagued by a plethora of *ad hoc* techniques which lack strict mathematical derivation from the single set of underlying principles. At the level of practical applications, asymptotic techniques<sup>21</sup> (in frequentist statistics) and perturbative techniques<sup>22</sup> (in quantum field theory) still play the main role. This is a dark *predictive dynamical* side of the coin,

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<sup>20</sup>See Footnote 3 in Section 2.1 for details.

<sup>21</sup>According to Weil, «although we know that statistical mathematics have been of considerable importance to science (and in particular to biological sciences), it needs to be mentioned that statistical books amount in fact to a collection of recipes and precepts we would like to believe to be well chosen. Being written in a highly algebraic form, sometimes using logarithms, exponentials and integrals, they all have the prestige of mathematical exactitude to the untrained eye while the so-called demonstrations statistics are wrapped in, even highly sophisticated, most often make no sense for the mathematician and are simply made of more or less convincing heuristic considerations» [1831] (engl. transl. in: [1879]). The average level of mathematical rigour in statistics has greatly improved since Weil wrote this judgement (see [1821, 1044, 1045, 317, 821, 673, 1327, 1642, 40, 1046, 186, 1722, 1597]), yet the choice of specific mathematical structures and methods corresponding to a given epistemic situation is still based on «more or less convincing heuristic considerations», and there is also a large gap between this what can be strictly proved and this what is applied to a given epistemic situation. In fact, most of mathematical results apply only to the asymptotic case of infinite number of data points, which is devoid of practical meaning from the *strict* point of view. For example, Le Cam notes that «limits theorems ‘as  $n$  tends to infinity’ are logically devoid of content about what happens at any particular  $n$ . All they can do is to suggest certain approaches whose performance must then be checked on the case at hand. Unfortunately the approximation bounds we could get were too often too crude and cumbersome to be of any practical use» [1046]. The routinely used *vague* point of view is based upon using *ad hoc* conventions that relate the asymptotic results with nonasymptotic situations: «in statistics no ‘general’ theory has yet been formulated that would investigate from a single point of view the properties of estimates, tests and other inference procedures and would be accepted by the majority specialists. A fundamental concept such as optimality is usually defined *ad hoc* in each special case, depending on the nature of the problem» [1596].

<sup>22</sup>As Borchers notes, the lagrangean approach to quantum field theory «has the disadvantage that the expressions which appear in this theory have only formal meaning. Up to now there is no convincing scheme which puts the formal expressions onto a solid and consistent mathematical basis. The existing perturbation and renormalization theory does not, in most cases, indicate anything about the quality of the approximation. Therefore, only comparison with experiment can indicate the quality of the Lagrange function and the approximation» [199]. The comparison of this remark with quotes in footnote 21 leaves a noticeable impression. It is quite striking to observe that two seemingly distinct fields of scientific inquiry share a strong structural similarity of the large variety of *ad hoc* methods mixed with a language of specialised, but rigorously undefined, concepts and semi-mathematical results, which hide the lack of conceptually clear *and* mathematically strict principles. In fact, this observation was a motivation for a conceptual reverse engineering of relationships of these two theories carried in Chapter 1.

bright side of which is a mathematical clarity of the *abstract kinematical* frameworks for probability theory and quantum mechanics that were provided and turned into orthodoxy by Kolmogorov and von Neumann, respectively. Persistence of this division despite all its negative consequences indicates that this division is necessary in order to satisfy some implicit but very important postulate. For the purpose of identification of it, let us summarise the main classes of problems that are encountered at the predictive side of the above two theories:

- (1) the arbitrariness of choice of a particular statistical inference method that is applied to a particular experimental situation. This choice might be suggested by intuition, tradition, or some *convention*, but is not derived from any general principle that would explicitly take into consideration the available knowledge that defines a given experimental situation. This situation is based on the frequentist interpretation, which equips probability distributions with ontological meaning, which forbids from considering them as carriers of epistemic knowledge, encoded and processed directly by means of some conventions. As a result, the division of inductive inference theory to probability theory and statistical inference theory is enforced, with probability theory keeping the ideal ‘ontological’ look and statistical inference handling all epistemic contents by means of *ad hoc* methods. See [870] for an extensive critical discussion concerning these issues, and [960, 1043, 1358, 959, 862] for various examples of spectacular failures of standard *ad hoc* methods;
- (2) the arbitrariness of construction of *predictive* (‘effective’, ‘interacting’) quantum dynamics. Quantum mechanics does not provide any prescription how to construct hamiltonians that would directly encode the quantitative constraints imposed by experimentally verifiable temporal relationships of experimental control-and-response parameters. There are various, inequivalent methods of ‘quantisation’ of ‘classical’ hamiltonians or lagrangeans, which lack common conceptual or mathematical foundation and are, in general, ambiguous. The results of use of these techniques often do not correspond directly to particular quantitative results and descriptions of experimental procedures under consideration (they do not define *predictive* dynamics), hence they have to be modified using additional *ad hoc* procedures (with more or less *ad hoc* choice of parameters), like cut-offs, parameter fitting, regularisation, or renormalisation. These procedures lead to serious problems of conceptual and mathematical character. In particular, they might move the operators between different unitarily inequivalent representations. In such case, the originally postulated lagrangean function or hamiltonian operator, far from being unique, has also no experimental meaning, because it cannot be justified by the reference to quantitative results of experimental procedures neither in terms of spectral representation, nor in terms of the expectation values or higher correlation moments. Thus, the choice of a hamiltonian might be suggested by intuition or tradition, but is not derived from any general principle (or explicitly specified class of principles) which in *some* clearly defined sense would optimally encode the available information about given experimental situation (or a class of them);
- (3) the arbitrariness of relationship between the choice of specific method of model construction, statistical inference procedure, and the assumptions regarding the causal structure that is considered as responsible for the changes of experimental data (see [1618, 1285, 1182, 1286] for a detailed discussion);
- (4) the arbitrariness of relationship between unitary (Born–Jordan–Dirac–Schrödinger) temporal evolution and nonunitary (von Neumann–Lüders) temporal evolution. Quantum mechanics does not give any definite prescription how to relate these two forms of temporal evolution of density operators. This conflict is sometimes rephrased by saying that «the border between the field of application of Schrödinger’s equation and the one of the projection postulate is not well defined» [100] (see also [183]), or that «no formulation of a projection postulate tells us exactly at which point [of time parameter of the unitary evolution] to apply it» [1700]. This sometimes leads to denial of the structure associated with the nonunitary evolution as an attempt to solve this conflict, see e.g. [252]. However, one should note that the above elements of mathematical

structure of quantum mechanics do not give any argument which favours unitary continuous temporal structure over nonunitary discrete temporal structure. Such preference may come only from some additional assumption (interpretation), which can be usually identified as a belief that quantum mechanics is a theory in some sense analogous to classical mechanics, as opposed to probability theory.

Moreover, when it comes to assertion of predictive verifiability of dynamical quantum theoretic models, all above problems combine into a new quality, due to:

- (5) the arbitrariness (and *a priori* postulated necessity) of division of a knowledge describing (*defining*) a given experimental situation into a part that is subjected to probabilistic statistical inference and a part that is subjected to quantum theoretical dynamics.<sup>23</sup> This perspective is based on *a priori* division of actual scientific inquiry into two different parts: “ideal/ontic theories” and “statistical/epistemic estimation”. As a result, the predictive application of quantum theoretic formalism becomes split into: (i) the *ad hoc* methods of probabilistic statistical estimation (inductive inference), which aims at selecting the unique semi-spectral measure and unique density operator that fit the specified quantitative results (see e.g. [777, 807, 809, 269, 427] and discussion in Section 2.2), (ii) the abstract Hilbert space based dynamical model construction, based on *ad hoc* procedures. The orthodox conceptual perspective considers dynamical model of some theory as something that can be arbitrarily postulated (“let us consider the following langrangean/hamiltonian of an interacting system...” and “let us postulate the following initial state of an interacting system...”), and later, in face of experimental evidence, it can be either dismissed or arbitrarily changed. As a result, “ideal/ontic theories” do not fit *optimally* the experimental data, and require some fine tuning, which is more or less arbitrary (and is provided either on the level of statistical estimation, or on the level of perturbative construction of ‘interacting’ hamiltonian). This shows strong contrast between the idealistic principles of a theory and the pragmatic procedures of its application (and is also an illustration of Duhem’s thesis [507, 508]). From the pragmatic (or renormalisation-minded) point of view, the *ad hoc* hamiltonians/langrangeans serve only as a starting reference point for some other procedures whose aim is to encode the particular experimental evidence.

In this context, let us note that the idea of “true probability distribution”, whose best incomplete description is provided by means of frequentist statistic, is an analogue of the idea of “classical theory”, whose incomplete description is provided by means of quantum operators. The former idea is based on consideration of the notion of an ‘elementary event’ (modelled by an element of a sample space) as a fundamental ontic entity, which is quantified by means of a statistic, while the latter idea is grounded in consideration of the notion of an ‘elementary property’ (modelled by an element of a phase space) as a fundamental ontic entity, which is quantified by means of a function on a phase space that is further subjected to quantisation. In this sense the notion of a statistic in frequentist statistics and the notion of an observable in quantum mechanics carry strong ontological undertones.<sup>24</sup> Note also that, while quantum mechanics does not rely formally on “classical theories” and quantisation on the level of foundations, the spectral representation reintroduces the ontological interpretation by relying on frequentist statistics over the sample space associated with eigenvalues. As a specific but crucial consequence, this leads to interpretation of the procedure of perturbation of ‘quantised’ hamiltonians as construction of the models of ‘interaction’ of some ontic entities, and not as providing the methods of inference conditioned upon epistemic control parameters.

Thus, while it is disputable whether the quantum theory and statistical inference theory are currently the prisoners of ontological postulates or of a lack of alternative mathematical framework (probably both), it seems for us that the above mathematical and conceptual problems require more than to postulate some new mathematical setting. It is essential to relate the structure of this setting with

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<sup>23</sup>One can consider this sentence as nonontological reading (or the practical meaning) of the so-called “Bohr–Heisenberg cut”.

<sup>24</sup>It is quite ironic, as compared with the postulates of Fisher and Heisenberg. Yet, it is a common problem of all reductionisms.

some specific postulates regarding the dynamical model construction and verifiability of predictions in face of epistemic (experimental) evidence.

### 1.3 Postulates

The conceptual layer of our approach (as summarised in the main postulate below) provides an interpretation of the underlying mathematical structure that is aimed at two goals: eliminating the ontological assumptions imposed on the formalisms of quantum mechanics and statistical inference, and clarifying the relationship between the predictive verifiability of dynamical models and their dependence on the conventions used for the purpose of their construction.

Our main postulate is to regard quantum theoretic kinematics and dynamics as generalisations of, respectively, probability theory *and* statistical inference to a regime where the underlying object of inquiry is not a set of measurable sets, but an integrable noncommutative algebra, with nonlinear quantum kinematics determined by the quantum relative entropy, and nonlinear quantum dynamics determined by two independent notions of evolution, provided by active and passive instruments. The corresponding conceptual postulate is to regard quantum theory as a strictly epistemic theory of information and intersubjective inference<sup>25</sup>, as opposed to ontic theories of substance its action (such as classical mechanics and general relativity).

The specific interpretation of the mathematical elements of our framework is:

1. **Algebras.** The  $W^*$ -algebras  $\mathcal{N}$  are viewed as the collections of abstract ‘elementary qualities’, which are *intersubjectively regarded* as basic entities of a theoretical inquiry, and are subjected to quantification by integration. The integrals on  $\mathcal{N}$  are given by quantum states  $\phi \in \mathcal{N}_*^+$ , and the corresponding quantifications are given by (not necessarily normalised) expectations  $\phi : \mathcal{N}^{\text{sa}} \mapsto \phi(x) \in \mathbb{R}$ . This is in strong analogy to viewing abstract boolean algebras as collections of abstract ‘elementary events’, which are subjected to quantification in terms of measures. (If these measures are normalised, then the corresponding quantifications are probabilities.) See Sections 2.1 and 2.2 for a thorough discussion of these issues. Elements of  $\mathcal{N}$  *are not* regarded as “observables”.
2. **Models.** Quantum models  $\mathcal{M}(\mathcal{N})$  are viewed as sets of quantitative descriptions of epistemic<sup>26</sup> knowledge about ‘abstract elementary qualities’ in  $\mathcal{N}$ . The passage from abstract notion of ‘quantification’ to more specific term, given in the previous sentence, is conditional upon additional requirement. It states that an application of  $\mathcal{M}(\mathcal{N})$  for the purpose of description of some epistemic situation has to be based on explicit specification of a set theoretic bijection between  $\mathcal{M}(\mathcal{N})$  and the set of all those possible quantitative descriptions of this situation which are *intersubjectively regarded* as relevant. In other words, our semantic requirement of the quantum model construction is: everything, that is considered to be an intersubjectively relevant quantitative knowledge that *defines* a given epistemic situation, should be bijectively translatable (= encodable and decodable) in terms of  $\mathcal{M}(\mathcal{N})$ .<sup>27</sup> This is independent of the language

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<sup>25</sup>In a sense similar, to some extent, to statistical mechanics and thermodynamics, when considered from the perspective of Jaynes [856, 858, 871, 859, 860, 863, 867], Ingarden [832, 834, 825, 826, 831], and others [1725, 1517, 1726, 658, 659]. From another point of view, one could say that our approach regards quantum theory as a framework for epistemic information processing. This resembles some similarity with the viewpoints of approaches of Srinivas [1619, 1620] and Abramsky–Coecke [3]. However, apart from other differences, we allow a large variety of inference rules, which are not required to be linear or local.

<sup>26</sup>Our preference to the adjective ‘epistemic’ instead of ‘experimental’ stems from information theoretic and inference theoretic perspective we adopt. According to it, quantum theory provides mathematical and conceptual framework for processing *any* quantifiable knowledge. To some extent, this can be traced back to Hartley [756], who postulated that information should be measured regardless of its semantic content. However, we impose an additional *semantic* constraint, requiring that this knowledge has to be intersubjective (the meaning of this term is discussed onwards).

<sup>27</sup>The idea that the probability density function should be constructed as a result of maximisation of an absolute entropy functional subjected to constraints that are expressed exclusively in terms of experimentally defined parameters was introduced by Gibbs [627] and was further developed by Jaynes [856, 863, 864, 870] (see also [651, 1844]). As opposed to Jaynes, we allow arbitrary intersubjective conventions of model construction.

used for the description of epistemic situation, an intentional subject of concern of experiments, experimental setting handled, etc.

3. **Geometry.** The meaning we assign to defining quantum theoretic kinematics in terms of quantum information geometry stems from the general purpose of any kinematics as a framework for quantitative analysis of the changes of states provided by dynamics, with a special emphasis on initial and final states. In the case of information dynamics, the initial states are called ‘priors’, while the final states are called ‘posteriors’.<sup>28</sup> The specific choice of geometric structure on a quantum model is an explicit declaration of *intersubjective criteria of relative relevance (distinguishability)*<sup>29</sup> that are used for the purpose of selection of priors on the basis of available epistemic evidence<sup>30</sup> and for the purpose of judging the quality of inference on  $\mathcal{M}(\mathcal{N})$  by comparison of posteriors with some reference (e.g. some state or a subset of states). The main issue at stake here is that the prior selection and posterior quality judgement, which are pre- and post-inferential procedures, respectively, require explicit specification of the criteria of ‘best’/‘optimal’ fit of the quantum state to some standard of reference. This is implemented by the quantum distance  $D(\cdot, \cdot)$ , as well as by the derived notions, such as the riemannian metric  $\mathbf{g}$ . Note that quantum distance  $D(\cdot, \cdot)$  on  $\mathcal{M}(\mathcal{N})$  is irrelevant for construction of a model  $\mathcal{M}(\mathcal{N})$  as a whole (this is provided already by the Item 2 above). Putting it into a slogan: quantum models can be *constructed*; quantum states can be *distinguished*.
4. **Dynamics.** The active instrument defines the *intersubjective convention* regarding inference procedure. It can be considered as an ‘information theoretic force’<sup>31</sup>. On the other hand, the passive instrument defines the *intersubjective convention* regarding causality. It encodes the assumptions about transformations of the epistemic evidence used for the purpose of model construction. By itself, it carries no predictive content, but it serves as a reference for judgement about inferences provided in terms of active instruments. In this sense, it is an ‘information theoretic reference frame of motion’.<sup>32</sup> For this reason, we will use the phrases ‘causality system’ and ‘referential evolution’ interchangeably with the notion of ‘passive instrument’.

The post-inferential judgement of quality of inference may be provided with respect to a quantum state selected ( $:=$  distinguished) on the base of some epistemic evidence. But, as discussed in Section 4.2.3, it can be also provided with respect to a passive instrument. Thus, instead of postulating a hamiltonian/lagrangean of a “free evolution” of a “noninteracting system” and further perturbing it by means of an “interaction” (with *indefinite* methods and criteria of quantification of quality of this approximation with respect to the available epistemic constraints), we postulate that one should specify the referential evolution (passive instrument)  $\mathfrak{I}_{\text{pass}}$ , the specific choice of  $D(\cdot, \cdot)$  used for comparison, and then impose constraints on the allowed range of active instruments (quantum inference procedures) by means of  $D$  and  $\mathfrak{I}_{\text{pass}}$ . The concept of transition distance, introduced in Section 4.2.3, is intended to serve this purpose. In other words, we postulate that the individual acts of inference (represented by active instruments) should be considered as relative to an intersubjective choice of the causality system, which describes the motion of information states when no specific evidence and no active instrument is given.

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<sup>28</sup>The entropic instrument that uses a prior *measure* as an input and produces an effect *measure* as an output is a generalisation of this idea. (See Item 4 below and Section 4.1.1 for definitions of these notions.) Actually, this suggests to construct a category of such generalised mappings, but we leave it as a topic for future research. See also Footnote 5 in Chapter 4.

<sup>29</sup>«this geometry does not describe the dynamics, or evolution, of a quantum system; rather it places limits on our ability to distinguish one state from another through measurements» [241].

<sup>30</sup>We use term ‘epistemic evidence’ an ‘quantitative epistemic knowledge’ interchangeably.

<sup>31</sup>This interpretation is especially suggestive in the case of entropic instruments  $\mathfrak{P}_{\mathfrak{D}}^{D, \Lambda}$ . In this case the prior measure  $\Lambda$  can be interpreted as an ‘information theoretic mass’ (see Section 4.1.1).

<sup>32</sup>This interpretation is especially suggestive for any instrument defined as a predualisation of automorphism of  $\mathcal{N}$  (thus, a causality system defined strictly at the algebraic level as a group of self-transformations of a set  $\mathcal{N}$  of ‘abstract qualities’) which coincides with the Poisson flow of a globally defined hamiltonian vector field (see Section 4.2.2). All unitary evolutions of the finite dimensional parametric quantum models belong to this class of instruments.



5. **Observables.** The downfall of the principle of “in principle observability” of elements of spectra of self-adjoint operators on a Hilbert space, as well as regarding the elements of an abstract  $W^*$ -algebra as ‘elementary qualities’ leads us to dissolve the quantum mechanical notion of “observable” into a kinematic notion of ‘estimable’<sup>33</sup> and dynamic notion of an ‘effect’. The ‘estimable’ is defined as an arbitrary function on an information model that can be used in order to provide estimation ( $:=$  prior selection or posterior quality judgement) using the available epistemic evidence. In particular, every element of an abstract  $W^*$ -algebra  $\mathcal{N}$  determines a *linear* estimable by means of  $\hat{x} : \mathcal{M}(\mathcal{N}) \ni \phi \mapsto \phi(x) \in \mathbb{C}$ . Yet, one can consider arbitrary nonlinear estimables on  $\mathcal{M}(\mathcal{N})$  as well. Whenever  $\mathcal{M}(\mathcal{N})$  is equipped with a smooth manifold structure, then the algebra of smooth functions on  $\mathcal{M}(\mathcal{N})$  might be of particular concern. A detailed analysis of this class of nonlinear estimables and their relationships with the Banach–Lie–Poisson structure  $\{\cdot, \cdot\}$  on  $\mathcal{M}(\mathcal{N})$ , as well as an extensive discussion of the relationship of these objects with an orthodox setting of quantum mechanical “observables” is due to Bóna [188, 190] (he uses the term “generalised observable”, cf. also [1167]). On the other hand, the dynamic notion of an ‘effect’ is an “estimable-like” quantification of the posterior of an inference procedure that is uniquely determined by this procedure. This should be opposed to the *relative* quantification of this posterior by means of  $D(\cdot, \cdot)$ . More precisely, an ‘effect’ is defined as a function on  $\mathcal{M}(\mathcal{N})$  such that it is uniquely determined by a given active instrument for any choice of constraints for this instrument. In Section 4.1.1 we introduce and analyse the nontrivial class of effects defined by the entropic instruments. We regain the relationship with the quantum mechanical setting by showing that, under some conditions, entropic effects determine the corresponding unique semi-spectral measures, and this procedure is structurally analogous to determination of semi-spectral measures by completely positive linear instruments.
6. **Space-times.** The notion of space-time does not belong to foundations of quantum mechanics. However, one of our main motivations for developing new foundations of quantum theory is to recover space-times from quantum theory as emergent objects (which is opposite to the idea of “quantisation of space-time”). The mathematical aspects of this problem are discussed in Sections 4.2.4 and 2.4. From the conceptual point of view, construction of emergent space-times amounts to regarding the lorentzian metric as an object encoding the conventions about the *local* criteria of optimal estimation (infinitesimal measure of distinguishability of information states) and the conventions about the *global* system of causality. This can be summarised by saying that: (i) two space-time points are distinguishable to a given extent *because* they represent the states of information that are distinguishable to a given extent; (ii) the space-time notion of causality, provided by space-time metric, is determined by the information theoretic notion of causality, provided by the passive instrument. Hence, the space-time causality encodes *intersubjective convention* regarding the ‘causal factuality’ and ‘causal counterfactuality’ in the same way as quantum information geometry encodes intersubjective conventions regarding ‘verifiability’ and ‘falsifiability’. As a result, each epistemic situation determines in principle its own space-time. Different space-times can be related to each other whenever their corresponding epistemic subjects, as well as intersubjective conventions used for the purpose of their construction, can be related.

## 1.4 Conceptual framework

Note that we do not aim at avoiding a convention-dependent character of quantum dynamics/inductive inference and its predictive verification. Our intention is to get rid of nonoptimality and intersubjectively incoherent ad hocery. Following Spengler [1617] and Fleck [565], we do not expect that anything stronger than intersubjective coherence in the above sense can be provided (the ontological postulates serve the same purpose, after all, just without acclaiming it). In this sense, intersubjective bayesianism,

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<sup>33</sup>We do not call it an ‘estimator’ in order to stress that its interpretation differs from, and is independent of, frequentist statistics.

as we define it here, can be understood as placed inside an intersection of two epistemic nominalisms: moderate conventionalism and moderate operationalism, where ‘moderate’ means ‘lacking reductionist overtones’.

This perspective is strongly influenced by the ideas of Jaynes [858, 860, 863, 870]. Like earlier founders of modern subjective bayesian approach to statistical theory, he denied identification of probability with frequency. However, he argued that the statistical inference needs to be constrained more strongly than by individual subjective judgements. The point of view proposed by Jaynes was a specific information theoretic interpretation of probability distributions, according to which they are theoretical objects that serve for the purpose of encoding and decoding information regarding *intersubjectively* communicable and verifiable knowledge. Jaynes attempted to justify the uniqueness of some methods of statistical inference on the base of various arguments that were shown to be indefinite (see [1741, 1742]). Our own point of view is that, while scientific inference is definitively intersubjective and not personal, there is no absolute basis allowing to justify any unique rules of inductive inference, so one necessarily deals with intersubjective conventions. The reconsideration of information geometric approach to foundations of quantum theory from this perspective, leads to an interesting return of Poincaré’s ideas [1350, 1351]: the choice of quantum geometry is conventional, so is the emergent space-time geometry. However, this is not a subjective convention, because the criteria of predictive verifiability have intersubjective character.

So far we have not defined the semantic term ‘intersubjective’. This would require us to coin some specific metaphysics (not necessarily of ontological type), and argue in favour of it. But we consider such attempts as arbitrary, hence irrelevant for the foundations (as opposed to applications). Instead of this, we will try to specify the minimal conditions required for regarding some semantic setting as providing necessary and sufficient meaning to ‘intersubjectivity’ in the context of *scientific inquiry*. The starting point is an observation, which follows Heraclitus, that any particular individual instance of an experimental situation is never reproducible. Only abstract *schematic* [339] descriptions of experimental situations can be reproduced, so only these abstract descriptions can be used as a basis for *scientific* inquiry. Hence, the postulate of reproducibility forces to use some abstract language of description of experimental situations. A *scientific fact* is an abstract statement of the above language. Next, the scientific requirement of an *intersubjective* reproducibility implies that the criteria of providing intersubjective meaning to sentences of this abstract language *are already given*. Finally, the last necessary condition for the scientific inference is a strict separation between abstract language used for the purpose of schematic description and communication of experimental situations and an abstract language used for the purpose of theoretic inference. As discussed in details by Fleck [565], scientific facts (and more generally, whole epistemic layer of scientific inquiry) can evolve similarly to scientific theories, and these two evolutions are mutually dependent. The above postulate of strict separation warrants that there exists an intersubjective class of reference (community of users) for which the division between experiment and theory is kept clear and fixed. Our interpretation amounts to saying that: (1) an “information” and “information dynamics” is a theoretic quantification of knowledge and its changes<sup>34</sup>; (2) this knowledge and its changes refer to the scientific facts expressed in terms of an abstract language used for description of an epistemic layer of inquiry; (3) the quantification in (1) and the relationship in (2) are provided accordingly to intersubjectively shared conventions, and under strict separation of scientific inquiry into epistemic and theoretic layers.

In this context, let us now consider the conceptual aspect of our main postulate, as stated at the beginning of Section 1.3. By the phrase ‘strictly epistemic’ we mean a requirement that a given theory can be subjected to predictive quantitative verification in terms of the available epistemic evidence without invoking any external concepts, notions, and theories, and without assigning ontological status to any of its elements.<sup>35</sup> By ‘experimental evidence’ we understand sentences of some abstract lan-

<sup>34</sup>Quantification amounts to turning knowledge and its changes into *evidence*.

<sup>35</sup>«no verification is possible unless the relevant inference method is an integral part of the theory» [1677]. Note that we postulate a judgement of the quality of inference to be a part of a theory (in line with the methods of model construction, prior selection, and inference). In this sense, every theory constructed within our approach should carry its own criteria of predictive verifiability.

guage which is suitable for intersubjective communication of quantitative descriptions of experimental designs, actions, and results, but does not involve any model construction, inference, or verification.

Interpreting inductive inference procedures as referring to something that is of abstract and intersubjective character draws a difference between the conceptual setting of our approach and the personalistic (‘subjective’) bayesianism [440, 441, 1514, 442, 534, 603, 297, 844, 605, 604, 298, 1707, 606] on one side and rationalistic (‘objective’) bayesianism [961, 874, 376, 858, 1726, 148, 868, 870] on the other. Similarly to subjective bayesianism, we consider the choice of particular method of inductive inference as *in principle* arbitrary. Similarly to objective bayesianism, we consider the *scientific inference* as a specific case of inductive inference theories, in which the constraint of *intersubjective reproducibility* of the subject of inference should play main role.

Our implementation of both insights (which we tentatively call ‘intersubjective bayesianism’) stems from observations of Spengler [1617], Granet [660], Fleck [564, 565, 566], and Chwistek [339]. According to them, the subject of concern of scientific inquiry at *experimental* level is neither a content of an individual subjective experience, nor an ontic “material reality”, but a structuralised system of abstract ideas, which is shared with other people educated and cognitively prepared in the similar way, and which frames the range of individual perceptions and thought style that can be subjected to intersubjective communication and predictive verification.

## 1.5 Published and joint work

The new perspective and main definitions contained in the above Introduction, and in Sections 3.1.1 and 4.1.1 were partially published in [992, 995]. The discussion of some of the conceptual issues connected with this new mathematical approach was provided in [994]. The new results contained in Sections 3.2.3 and 3.2.4 were published in [993]. The content of Section 4.1.2 is a result of the joint work with Wojciech Kamiński [924]<sup>36</sup>. The quantum Schwarzschild space-time model of Section 4.2.4 is a joint work with Paweł Duch [506].

## 1.6 Conventions

The following conventions are used: 1) the symbols denote the same mathematical entities throughout all work; 2) the special syntactic forms used in text are: *definitions*, «citations», ‘notions subjectable to strict definition’, “vague notions”, and *attention markers*; 3) the mathematical style of text formatting (definition/theorem/proof) is used only for stating essentially new mathematical results; 4) whenever possible, we refer to original works containing results that are discussed or used; 5) the folk attributions of surnames to mathematical concepts are changed to historically correct ones whenever there is a definite evidence for the latter, and the concepts with attribution of three or four surnames to it are turned to an acronym after first use, while in the case of more authors we use only the descriptive naming of objects;<sup>37</sup> 6) for the Latin transliteration of the Cyrillic script (in references and surnames) we use the following modification of the system GOST 7.79-2000 B: ц = c, ч = ch, x = kh, ж = zh, ш = sh, щ = shh, ю = yu, я = ya, ё = ë, ъ = ‘, ь = ’, ы = y, э = e’, й = ĭ, with an exception that surnames beginning with X are transliterated to H.<sup>38</sup>

<sup>36</sup>The idea to use the generalised pythagorean equation as a method of proof, as well as some calculations suggesting the correctness of this idea, emerged in our joint discussion with Frank Hellmann and Carlos Guedes.

<sup>37</sup>The only definite exception that we have made consciously is the abstract notion of a ‘Hilbert space’. It was invented (and named) by von Neumann. Hilbert invented only its special case, an  $\ell_2$  space, while an  $L_2(\mathbb{R}, d\lambda)$  space was invented by Riesz. Yet, the standard naming convention is too influential to stand up against it. More generally, we have adopted the sequential adaptation of rules: (1) strong relevance of temporal priority, (2) weak relevance of folk popularity, (3) weak avoidance of acronyms. As a borderline example, we speak of the “Morse–Transue–Nakano–Luxemburg norm”, abbreviated to the “MTNL norm”, despite 5 years of difference between [1192, 1228] and [1108] due to wide popularity of the term “Luxemburg norm” (enforced by [998, 1400]), but we speak of the “Csiszár–Morimoto f-distance” despite 3 years of difference between [380, 1184] and the Ali–Silvey paper [31]. On the contrary to this partially restrictive *naming* system, the *references* are made as complete as it is reasonably possible.

<sup>38</sup>This is required for agreement with the widespread practice to transliterate Холево as Holevo, etc.



# References

Я мысленно вхожу в ваш кабинет;  
Здесь те, кто был, и те, кого уж нет,  
Но чья для нас не умерла химера;  
И бьётся сердце, взятое в их плен...

Максимилиан А. Волошин

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