Preheating and inflation in supergravity - the role of flat directions

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Flat direction

- direction in field-space, along which the scalar potential identically vanishes (when all other field VEVs=0)
- general feature of supersymmetric models

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Preheating

very efficient non-perturbative particle production during inflaton oscillations

Preheating and flat directions

Toy model

$$V \supset \frac{1}{2}m^2\varphi^2 + A\varphi^2\chi^2 + Bm\varphi\chi^2 \tag{1}$$

 φ - inflaton field, χ - represents the inflaton decay products

$$\omega_{\chi_k}^2 = k^2 + 2A \langle \varphi \rangle^2 + 2Bm \langle \varphi \rangle \tag{2}$$

$$| au| \equiv \left| \frac{\dot{\omega}}{\omega^2} \right| > 1 \leftrightarrow \textit{preheating}$$
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Toy model with a flat direction

(Allahverdi, Mazumdar '07)

$$V \supset \frac{1}{2}m^2\varphi^2 + A\varphi^2\chi^2 + Bm\varphi\chi^2 + C\alpha^2\chi^2$$
 (4)

 α - parameterizes the flat direction

$$\omega_{\chi_k}^2 = k^2 + 2A \langle \varphi \rangle^2 + 2Bm \langle \varphi \rangle + 2C \langle \alpha \rangle^2$$
 (5)



 the challenge of constructing a consistent model of inflation and particle production in a supersymmetric framework

generate large flat direction VEVs during inflation

- generate large flat direction VEVs during inflation
 - create a potential for the flat direction

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 - consider classical evolution of VEVs during inflation

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- check the impact of large flat direction VEVs on particle production
 - consider excitations around VEVs
 - study the evolution of the mass matrix
 - determine if preheating from the inflaton is possible

Inflaton sector

M. Kawasaki, M. Yamaguchi, T. Yanagida "'Natural Chaotic Inflation in Supergravity"

 Φ - inflaton superfield, X - auxiliary superfield

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 auxiliary field X used to obtain chaotic inflation potential during inflaton domination

$$W \supset mX\Phi$$
 (7)

$$V \stackrel{inflaton \ domination}{\longrightarrow} \frac{1}{2} m^2 \varphi^2$$
 (8)

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Observable sector

MSSM superpotential

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$$W \supset 2hXH_uH_d \tag{10}$$

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representative flat direction udd

$$u_i^{\beta} = d_i^{\gamma} = d_k^{\delta} = \frac{1}{\sqrt{3}}\alpha, \quad \alpha = \rho e^{i\sigma}$$
 (12)

Observable sector

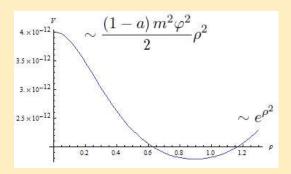
non-minimal Kähler

$$K \supset \left(1 + \frac{a}{M_4^2} X^{\dagger} X\right) \left(H_u^{\dagger} H_u + H_d^{\dagger} H_d + u_i^{\dagger} u_i + d_j^{\dagger} d_j + d_k^{\dagger} d_k\right)$$
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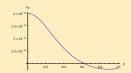
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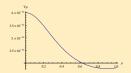
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non-renormalisable terms

$$W \supset \frac{\lambda_{\chi}}{M_{Pl}} (H_u \cdot H_d)^2 + \frac{3\sqrt{3}\lambda_{\alpha}}{M_{Pl}} (u_i d_j d_k \nu_R)$$
 (14)



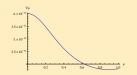
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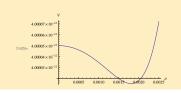
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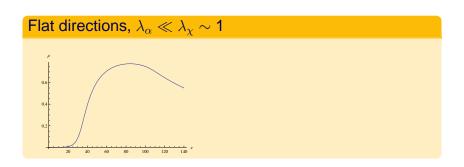
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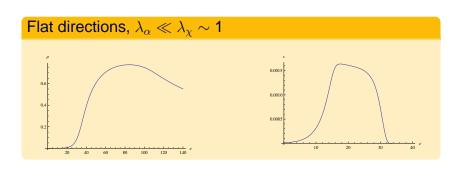
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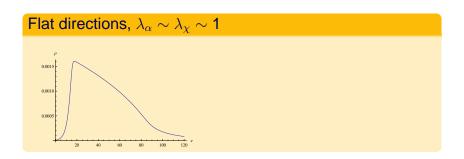


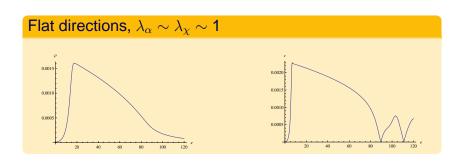
Flat directions, $\lambda_{\alpha} \ll \lambda_{\chi} \sim 1$





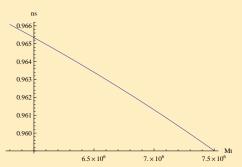
Flat directions, $\lambda_{\alpha} \sim \lambda_{\chi} \sim 1$





Spectral index

values of the spectral index 50-60 e-folds before the end of inflation in the slow-roll approximation



WMAP5: $n_s = 0.960^{+0.014}_{-0.013}$

Parameterization of excitations

• consider excitations around fields belonging to H_u , H_d , u_i , d_j and d_k multiplets

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$$VEV = 0 \longrightarrow field \sim \delta_a + i\delta_b$$
 (16)

Analyzing the mass matrix evolution, $\lambda_{\alpha} \ll \lambda_{\chi} \sim 1$

heavy eigenvalues

$$m_{udd}^2 \approx \frac{g^2}{3}\rho^2 + \underbrace{-\frac{m^2\varphi^2}{2}(a-1)}_{SUGRA} + \dots$$
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Toy model analogy

$$m_{\chi}^{2} = 2A \langle \varphi \rangle^{2} + 2Bm \langle \varphi \rangle + 2C \langle \alpha \rangle^{2}$$
 (19)



Analyzing the mass matrix evolution, $\lambda_{\alpha} \ll \lambda_{\chi} \sim 1$

naturally light eigenvalues corresponding to

$$\left(\xi_{u_i} + \xi_{d_j} + \xi_{d_k}\right)/\sqrt{3}$$

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and the excitation around the phase of the flat direction VEV

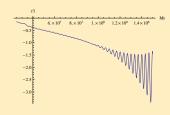
$$m_{phase}^2 \approx \underbrace{(1-a)\frac{m^2\varphi^2}{2} + g(a)\frac{m^2\varphi^2}{2}\rho^2}_{SUGRA} + \dots$$
 (21)

Analyzing the mass matrix evolution, $\lambda_{\alpha} \ll \lambda_{\chi} \sim$ 1

the time evolution of both m_{abs}^2 and m_{phase}^2 leads to non-perturbative particle production

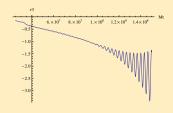
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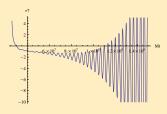
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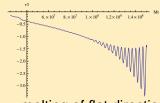
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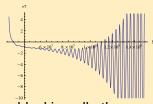




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— melting of flat direction VEV and unblocking all other channels of preheating

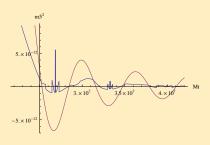
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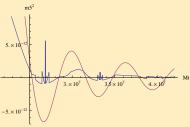
an example of a naturally light eigenvalue corresponding to a combination of excitations around VEVs of complex fields α and χ parameterizing the (quasi) flat directions



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 $SU(3)\times SU(2)\times U(1)\to U(1)$

an example of a naturally light eigenvalue corresponding to a combination of excitations around VEVs of complex fields α and χ parameterizing the (quasi) flat directions



— very efficient preheating into Higgs particles allowed from the beginning of inflaton oscillations

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- Non-perturbative particle production from the inflaton is likely to remain the source of preheating even in the initial presence of large flat direction VEVs.