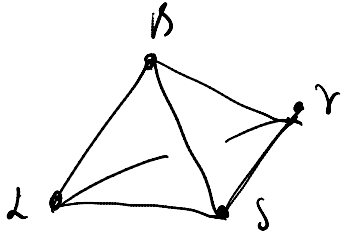


associated to the triangles of M

Amplitudes: $A_t : \bigotimes_{\Delta \in t} H_\Delta \longrightarrow \mathbb{C}$

$$A_c(2 \otimes 1 \otimes 1 \otimes 1) = \{ \delta_{\mathcal{I}} \}_t$$



The model $PR(M) = \sum_{\mathcal{I}} \prod_{e \in \mathcal{I}} \dim \mathcal{I}(e) \prod_t \{ \delta_{\mathcal{I}} \}_t$

$\dim \mathcal{I} = 2\mathcal{I} + 1$

Sum over all possible assignments

[Bonzano-Regge '68]

$PR(M)$ regularization of $Z(M) = \int Dg e^{i \int d^4x \sqrt{g} R}$

Problem: $PR(M)$ diverges for a large class of M

→ need a regularisation

Idea: quantum group $U_q(\mathfrak{so}(2))$ with $q^r = 1$ admits only a finite number of irreducible representations

$$I \in \left\{ 0, \frac{1}{2}, 1, \dots, \frac{r-2}{2} \right\}$$

↳ Replace $SU(2)$ in PR model by $U_q(SU(2))$

$$TV_r(M) = K \sum_I \prod_c [\dim I(c)]_q \prod_t \left\{ \frac{1}{q} \right\}_t$$

finite sum \rightarrow q -dimension $[n]_q = \frac{q^n - q^{-n}}{q - q^{-1}}$ quantum q -symbol \rightarrow

↳ $TV_r(M)$ converges for all manifolds M

$TV_r(M)$ is a regularisation of

$$Z_\Lambda(M) = \int \rho_g e^{i \int_M \text{tr}(\sigma_g) (\mathbb{R} - 2\Lambda)}, \quad \Lambda > 0$$

$$r = k+2, \quad k = \frac{\pi}{\epsilon \rho \Lambda}$$

Idea: apply the same procedure to regularise the Lorentzian EPRL model

Outline:
 I The quantum Lorentz group
 II The quantum EPRL intertwiner
 III The 1 -simplex amplitude

IV The 4-simplex amplitude

IV The model

Hopf algebras

G Lie group, $\Pi: G \rightarrow \text{Aut } V$
representation on V (field k)

Consider:

i) Tensor product of representations

$$g \in G, \quad \Pi_{V \otimes W}(g) = \Pi_V(g) \otimes \Pi_W(g)$$

ii) Dual representation

$$\alpha \in V^* \quad \Pi_V^*(g) \alpha = \alpha \circ \Pi_V(g^{-1})$$

iii) Trivial representation

$$\Pi_k(g) = \text{id}_k$$

Question: Suppose I replace G with an (associative, unital) H . What structure do I need to put on H for representations to mimic those of G ?

i) Co-product $\Delta: H \rightarrow H \otimes H$

$$a \in H \quad \Pi_{V \otimes W}(a) = \Pi_V \otimes \Pi_W(\Delta(a))$$

$$a \in H \quad \Pi_{\text{view}}(a) = \widehat{\Pi}_V \otimes \widehat{\Pi}_W (\Delta(a))$$

$$(\Delta g = g \otimes g)$$

ii) Antipode : $S : H \rightarrow H$

$$\Pi_V^+(a) \Delta = \Delta \circ \Pi(S(a))$$

$$(S(g) = \bar{g}^{\wedge})$$

iii) Co-unit $\varepsilon : H \rightarrow k$

$$\Pi_k(a) = \varepsilon(a) \text{id}_k$$

↳ These structures must satisfy compatibility and consistency conditions

↳ H becomes a Hopf algebra

I The quantum Lorentz group (\mathcal{QLG})

I.1 Hopf algebra structures

\mathcal{QLG} : quantum double of $U_q(\mathfrak{su}(2))$
with $q = e^{-\hbar} \in]0, 1[$ real

Ingredients : i) Hopf algebra $U_q(\mathfrak{su}(2))$

Ingredients: i) Hopf algebra $U_q(\mathfrak{su}(2))$

⊗ Algebra structure: generators $q^{\pm J_z}, \bar{J}_{\pm}$

+ relations: $q^{\pm J_z} q^{\mp J_z} = 1$
 $q^{\pm J_z} \bar{J}_{\pm} q^{\mp J_z} = q^{\pm 1} \bar{J}_{\pm}$
 $[\bar{J}_+, \bar{J}_-] = [2 \bar{J}_z]$

⊗ Co-algebra structure:

$$\Delta(q^{\pm J_z}) = q^{\pm J_z} \otimes q^{\pm J_z},$$

$$\Delta(\bar{J}_{\pm}) = q^{\mp J_z} \otimes \bar{J}_{\pm} + \bar{J}_{\pm} \otimes q^{J_z}$$

$$\Sigma(q^{\pm J_z}) = 1, \quad \Sigma(\bar{J}_{\pm}) = 0$$

⊗ Antipode $S(q^{\pm J_z}) = q^{\mp J_z}$
 $S(\bar{J}_{\pm}) = -q^{\pm 1} \bar{J}_{\pm}$

⊗ UIR behave as in the classical case

→ $\Pi_{\mathbb{Z}}: U_q(\mathfrak{su}(2)) \rightarrow \text{End } V_{\mathbb{Z}}, \mathbb{Z} \in \mathbb{N}/2,$
 spin \mathbb{Z} UIR of $U_q(\mathfrak{su}(2)), V_{\mathbb{Z}} = \mathbb{C} \left\{ e_{k\mathbb{Z}} \right\}_{k \in \mathbb{Z}}$

→ Fusion rules

$$V_{\mathbb{Z}} \otimes V_{\mathbb{J}} \cong \bigoplus_{k=|\mathbb{Z}-\mathbb{J}|}^{\mathbb{Z}+\mathbb{J}} V_k \dots \mathbb{Z}$$

→ Clebsch-Gordan maps (CG)

$$C_{IJ}^K : V_I \otimes V_J \rightarrow V_K, \quad C^{IJ}_K : V_K \rightarrow V_I \otimes V_J,$$

defined as:

$$C_{IJ}^K \left(\begin{matrix} I & J \\ e_b \otimes e_c \end{matrix} \right) = \sum_a \left(\begin{matrix} a & I & J \\ K & b & c \end{matrix} \right) e_a^K,$$

$$C^{IJ}_K (e_a^K) = \sum_{b,c} \left(\begin{matrix} b & c & K \\ I & J & a \end{matrix} \right) e_b^I \otimes e_c^J$$

CG coefficients ($\in \mathbb{R}$)

ii) Dual Hopf algebra

$$U_q(\mathfrak{su}(2))^* = F_q(\mathfrak{su}(2))$$

\hookrightarrow q -analogue of the algebra of functions on $SU(2)$

QLG: Hopf algebra

$$\textcircled{1} U_q(\mathfrak{su}(2)) = U_q(\mathfrak{su}(2)) \hat{\otimes} F_q(\mathfrak{su}(2))^{\text{op}}$$

2 copies do not commute

Rem: This is a quantum deformation of $\mathfrak{sl}(2, \mathbb{C})_{\mathbb{R}} \cong \mathfrak{so}(3, 1)$

Turaeva decomposition: $\mathfrak{sl}(2, \mathbb{C})_{\mathbb{R}} \cong \mathfrak{su}(2) \oplus \mathfrak{so}(3)$

Zwasara decomposition : $sl(2, \mathbb{C})_{\mathbb{R}} \cong su(2) \oplus an(2)$

$an(2)$ Lie algebra $AN(2) = \left\{ \begin{pmatrix} d & n \\ 0 & d^{-1} \end{pmatrix} \mid d \in \mathbb{R}^{\pm}, n \in \mathbb{C} \right\}$

($SL(2, \mathbb{C})_{\mathbb{R}} \cong SU(2) \ltimes AN(2)$)

$$\text{D } U_q(su(2)) = U_q(su(2)) \hat{\otimes} U_q(an(2))$$

$$\text{L}_s \left| \text{D } U_q(su(2)) = U_q(sl(2, \mathbb{C})_{\mathbb{R}}) \right.$$

II.2 Irreducible representations

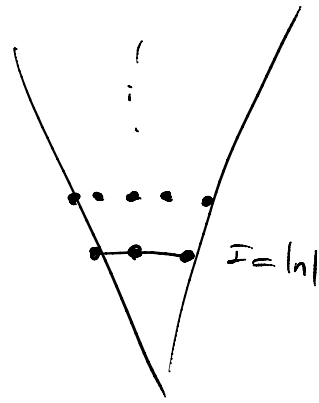
Unitary irreducible representations (of the principal series) of the GL₂ are labelled

$$\text{by } \lambda = (n, \rho) \quad , \quad n \in \mathbb{Z}/2 \quad , \quad \rho \in [0, \frac{4\pi}{\kappa} [$$
$$\left(\gamma = e^{-\kappa} \right)$$

Decomposition in terms of representations

of $U_q(\mathfrak{su}(L))$

$$V_\lambda = \bigoplus_{I=|\lambda|}^{\infty} V_I$$



II The quantum EPRC intertwiner

$$\text{Irrep}(U_q(\mathfrak{sl}(2, \mathbb{R}))) = \left\{ (n, p) \mid n \in \mathbb{Z}/2, p \in [0, \frac{4\pi}{r}] \right\}$$

$$\begin{aligned} \phi_r : \text{Irrep}(U_q(\mathfrak{su}(L))) &\longrightarrow \text{Irrep}(U_q(\mathfrak{sl}(2, \mathbb{R}))) \\ \kappa &\longmapsto (n(\kappa), p(\kappa)) := \underbrace{(\kappa, r\kappa)} \end{aligned}$$

$v \in \mathbb{R}^+$ fixed parameter q-EPRC representation

Rem: pre-image of ϕ_r is restricted

$$\mathcal{L} = \left\{ \kappa \in \mathbb{N}/2 \mid \kappa \leq \frac{4\pi}{r\kappa} \right\}$$

II.2 Quantum EPRC intertwiner

Need to generalise expressions of the form

$$\int_{\mathbb{R} \setminus \{0\}} \delta(x) \bigotimes_{i=1}^n \Pi_{\lambda_i}(x) : \bigotimes_{i=1}^n V_{\lambda_i} \rightarrow \bigotimes_{i=1}^n V_{\lambda_i} \quad (\text{I})$$

Introduce a Haar measure on the Hopf algebra $F_q(SL(2, \mathbb{C})_{\mathbb{R}})$ dual to $U_q(\mathfrak{sl}(2, \mathbb{C})_{\mathbb{R}})$

$$h: F_q(SL(2, \mathbb{C})_{\mathbb{R}}) \rightarrow \mathbb{C}$$

[Muffler, Roche '99]

Satisfies : $(h \otimes \text{id}) \Delta(x) = (\text{id} \otimes h) \Delta(x) = h(x) 1$

Ex: $\mathbb{C}(G)$, $h(f) = \int_G f(g) dg$

(I) is generalised as

$$T_{d_1, \dots, d_n} = \sum_A \left(\bigotimes_{i=1}^n \Pi_{\lambda_i}(\Delta^{(n)}(x^A)) \right) h(x^A)$$

dual basis of $U_q(\mathfrak{sl}(2, \mathbb{C})_{\mathbb{R}})$
basis of $F_q(SL(2, \mathbb{C})_{\mathbb{R}})$

[Noui, Roche '02]

The EPRL representation $\mathcal{L}(K) = (K, \mathbb{R}K)$

factors as $V_{\mathcal{L}(K)} = \bigoplus_{\mathbb{Z} \geq K} V_{\mathbb{Z}}$

Let $f_{\mathbb{Z}}^K: V_{\mathbb{Z}} \rightarrow V_K$ be the projection on lowest weight factor

on lowest weight factor

The dual map induces an embedding

$$F^* : \text{Hom}_{U_q(\mathfrak{sl}(n))} \left(\bigotimes_{i=1}^n V_{\lambda_i}, \mathbb{C} \right) \rightarrow \text{Hom}_{U_q(\mathfrak{sl}(n))} \left(\bigotimes_{i=1}^n V_{\lambda_i}, \mathbb{C} \right)$$

The EPRL intertwiner is defined as

$$\Lambda_{\lambda}(z) = F^*(\Lambda_{\lambda}) = \Lambda_{\lambda} \circ \bigotimes_{i=1}^n f_{z_i}^{k_i} = T_{z_1, \dots, z_n}$$

where $\Lambda_{\lambda} : \bigotimes_{i=1}^n V_{\lambda_i} \rightarrow \mathbb{C}$ is an

$U_q(\mathfrak{sl}(n))$ -intertwiner.

infinite sum

Theorem: The evaluation of the q -EPRL intertwiner $\Lambda_{\lambda}(z) = F^*(\Lambda_{\lambda})$ on the basis element $\bigotimes_{i=1}^n \zeta_{c_i}^{k_i}$ of $\bigotimes_{i=1}^n V_{\lambda_i}$ is a multiple series that converges absolutely.

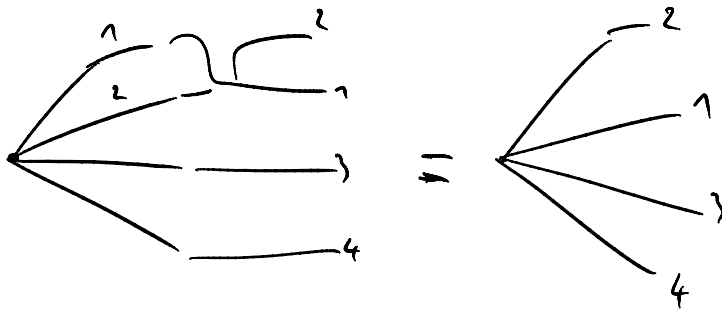
$\{ \zeta_c \mid c \geq |l|, c = -l, \dots, l \}$ basis of V_{λ}

Property: The q -EPRL intertwiner is not invariant under the braiding

$$C_{n, k_1, k_2} = \tilde{C}_{n, k_1, k_2} \cdot \Pi_{2n} \otimes \Pi_{k_1} \otimes \Pi_{k_2} (R) : V_{n_1} \otimes V_{k_1} \rightarrow V_{k_2} \otimes V_{n_2}$$

↑ population ↑ R-matrix of QLF

Prop: The q-BC intertwiner $[N_{n_1}, R_{k_1}, k_2]$ is invariant

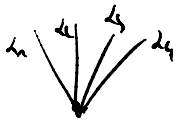
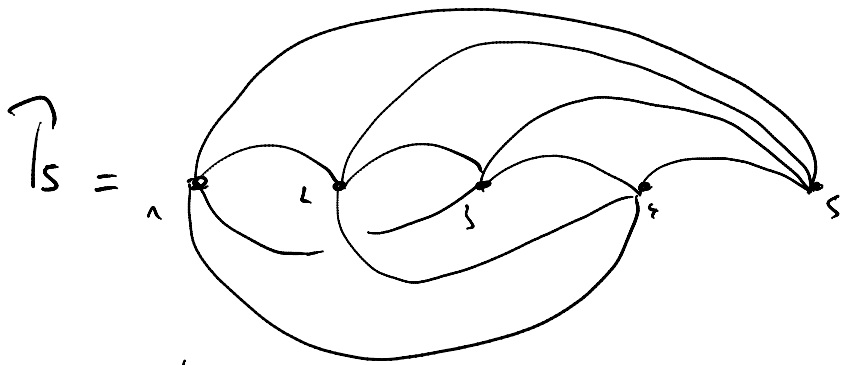


III The 4-simplex amplitude

We define a procedure to associated a number (possibly infinite) $ev(\hat{\Gamma}_n)$ to all closed diagrams $\hat{\Gamma}_n$ with n vertices

III.2 Amplitude for the 4-simplex

$$A_4 = ev(\hat{\Gamma}_4) \quad \text{4-simplex amplitude}$$



q -EMC interchain



bilinear form $V_{\alpha\beta} \rightarrow \mathbb{C}$



$c_{2,1}$

Theorem: $ev(T_S)$ is a multiple series that converges absolutely

IV The model

M triangulated 4-manifold

$$Z_K(M) = \sum_{K, \mathcal{T}} \prod_{\Delta} [2K(\Delta) + 1]_q \prod_{\sigma} A_{\sigma}(L(\Delta), \gamma_{\sigma})$$

sum ranges over all $K \in \mathcal{L}$

Property: $Z_K(M)$ is finite

Physics: Setting $q = e^{-k_1/k_2}$, $l_p = \sqrt{6}k_1$, $l_c = \frac{1}{\sqrt{2}}$

we obtain a bound :

$$\text{Area}(\Delta) \leq 32\pi^2 l_c^2 \quad (l_p \ll l_c)$$

for the area of Δ .