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Proceedings of the XXI International School of Semiconducting Compounds, Jaszowiec 1992

HYDROSTATIC PRESSURE SPECTROSCOPY OF THE VANADIUM LUMINESCENCE IN GaAs

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We report luminescence measurements of the intracenter transition ${}^{3}T_{2} \rightarrow {}^{3}A_{2}$ of the V³⁺(3d²) charge state in semi-insulating GaAs under hydrostatic pressure up to 0.8 GPa at liquid helium temperature. The hydrostatic pressure coefficient of the zero-phonon line is found to be equal to $6.9 \pm 0.2 \text{ meV/GPa}$. This result enables us to determine the Huang-Rhys parameter, which characterizes the coupling to the symmetric mode of vibration, as $S_{A} = 1.4 \pm 0.1$. Using this parameter, computer simulation leading to a reconstruction of the shape of both luminescence and corresponding absorption spectra were performed.

PACS numbers: 71.55.Eq, 78.55.Cr

1. Introduction

Optical absorption and luminescence bands in GaAs at 0.74 eV corresponding to ${}^{3}A_{2} \leftrightarrow {}^{3}T_{2}$ transitions of the V³⁺(3d²) charge state crystals have been investigated for many years (see for example review articles [1, 2]). The results of several experiments were analysed in terms of a dynamical Jahn-Teller effect with the ε -mode existing in the excited ${}^{3}T_{2}$ state [3-7]. The obtained values of the Huang-Rhys parameter $S_{E} = E_{\rm JT}/\hbar\omega_{E}$ (where $E_{\rm JT}$ is the Jahn-Teller energy and $\hbar\omega_{E}$ — the phonon energy) vary between 3.4 and 4. It was also shown [8] that in order to explain the total shape of luminescence and absorption bands, an interaction with the symmetric α -mode of vibrations should also be taken into account.

The lack of any effect of hydrostatic pressure up to 1 GPa on the transport properties of p-type GaAs:V [9] suggests a minor influence of the α -mode on the ground ${}^{3}A_{2}$ state (degenerate with the GaAs valence band). The parameter S_{A} describing the interaction of the α -mode with the excited ${}^{3}T_{2}$ state can be thus obtained by a moment analysis of the luminescence or absorption bands [8] or from the hydrostatic pressure coefficient of one of these bands. This pressure coefficient is known from a uniaxial stress experiment [3], but with a significant experimental error. This paper presents hydrostatic pressure measurements of the vanadium luminescence band which gave a much more accurate value of the pressure coefficient and, subsequently the S_A parameter. This result and the ε -mode Jahn-Teller data from literature make it possible to reproduce the shape of both luminescence and absorption spectra

2. Experimental

Vanadium doped, semi-insulating GaAs crystals were prepared by the liquid encapsulated Czochralski (LEC) technique with a vanadium concentration of $7 \div 8 \times 10^{16}$ cm⁻³ [10]. Zero pressure luminescence was measured in a CF1204 (Oxford) cryostat. For pressure investigations samples were placed in a high pressure optical cell with benzine as the transmitting medium. The cell with a sample was mounted in the exchange helium gas cryostat and cooled to low temperatures. Hydrostatic pressure up to 0.8 GPa was measured with a calibrated InSb manometer. Photoluminescence (PL) was excited by the 488 nm line of an Ar ion laser with power density around 1 W/cm². PL was analysed with a SPEX 500M monochromator equipped with a North Coast Optics Ge photodiode or cooled PbS detector (Hamamatsu).

3. Results and discussion

A typical low-temperature (4 K) atmospheric pressure photoluminescence spectrum related to the ${}^{3}T_{2} \rightarrow {}^{3}A_{2}$ transition is shown in Fig. 1. It consists of a strong zero-phonon line (ZPL) A with accompanying hot line B (see inset)

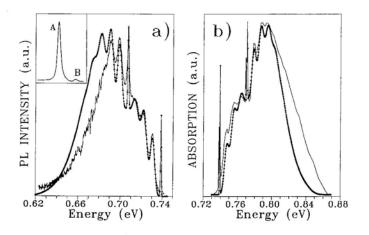


Fig. 1. Luminescence (a) and absorption (b) bands corresponding to ${}^{3}A_{2} \leftrightarrow {}^{3}T_{2}$ transitions of the V³⁺(3d²) charge state in GaAs at about 4 K (solid lines). Theoretical Poisson shapes (star lines) were calculated for the parameters $S_{E} = 3.9$, $\hbar\omega_{E} = 8.1$ meV; $S_{A} = 1.4$, $\hbar\omega_{E} = 31.5$ meV. Inset — luminescence zero-phonon lines A and B.

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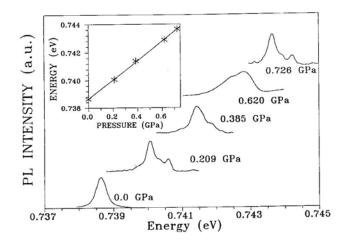


Fig. 2. Shift of the luminescence zero-phonon line A under hydrostatic pressure at about 4 K (10 K for 0.62 GPa). Inset — a linear fit to the experimental data with a pressure coefficient equal to 6.9 ± 0.2 meV/GPa.

followed by phonon replicas. The energy shift of the A line with increasing pressure is presented in Fig. 2. With the exception of the measurements at 0.62 GPa, all photoluminescence data were collected at about 4 K. The low intensity shoulders of the main line are due to axial stresses induced on the GaAs surface by frozen benzene. The measured positions of line A were plotted versus the stress value. A straight line fit to the data gave the hydrostatic pressure coefficient of the ZPL being equal to $6.9 \pm 0.2 \text{ meV/GPa}$. From this value it is easy to calculate the symmetric vibration energy $E_A = 43 \pm 3 \text{ meV}$. Taking the phonon energy $\hbar\omega_A = 31.5 \text{ meV}$ from the experimental spectra one obtains the Huang-Rhys parameter being equal to $S_A = E_A/\hbar\omega_A = 1.4 \pm 0.1$.

Computer simulation taking into account optical transitions involving two phonons ($\hbar\omega_E$ and $\hbar\omega_A$) with the ε -mode Jahn-Teller effect at the 3T_2 state known from other experiments [3-7] and the obtained value of the S_A parameter lead to a reconstruction of the shape of both luminescence and absorption bands (see Fig. 1).

Acknowledgment

This work was supported by the Committee for Scientific Research grant No. 2 0179 91 01.

References

- [1] B. Clerjaud, J. Phys. C Solid State Phys. 18, 3615 (1985).
- [2] A.M. Hennel, in: Imperfections in III-V Materials, Ed. E. Weber; in series Semiconductors and Semimetals, Eds. R.K. Willardson, A.C. Beer, Vol. 38, Academic Press, in press.

[3] G. Armelles, J. Barrau, D. Thébaut, J. Phys. C, Solid State Phys. 17, 6883 (1984).

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- [4] G. Armelles, J. Barrau, D. Thébaut, M. Brousseau, J. de Phys. 45, 1795 (1984).
- [5] G. Aszodi, U. Kaufmann, Phys. Rev. B 32, 7108 (1985).
- [6] K.J. Nash, M.S. Skolnik, B. Cockayne, W.R. MacEwan, J. Phys. C, Solid State Phys. 17, 6199 (1984).
- [7] F.G. Anderson, F.S. Ham, in: Defects in Semiconductors, Ed. G. Ferenczi, Material Science Forum 38-41, 305 (1989).
- [8] A. Wysmołek, Z. Liro, A.M. Hennel, in: Defects in Semiconductors, Ed.G. Ferenczi, Material Science Forum 38-41, 827 (1989).
- [9] D. Wasik, A.M. Hennel, unpublished.
- [10] A.M. Hennel, C.D. Brandt, K.Y. Ko, J. Lagowski, H.C. Gatos, J. Appl. Phys. 62, 163 (1987).