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*Opracowanie utworu pod tytułem:*

*”Druga kwantyzacja i metoda Bogoljubowa”*

w ramach kursu zaawansowanego, organizowanego w dniach 31.08 - 25.09.09 będącego kontynuacją szkoleń z zakresu eksploatacji i zarządzania dużą infrastrukturą badawczą organizowanego przez Narodowe Laboratorium Technologii Kwantowych

## Second Quantization Formalism

In position representation  $N$ -body Hamiltonian reads

$$\hat{H} = \sum_{i=1}^N \left[ \frac{\hat{\mathbf{p}}_i^2}{2m} + U(\vec{r}_i) \right] + \frac{1}{2} \sum_{i \neq j=1}^N V(\vec{r}_i - \vec{r}_j).$$

Exchange of any two particles does not change the Hamiltonian. Hence eigenstates of  $H$ ,  $|\psi\rangle$ , can be chosen so that

$$P_{12}|\psi\rangle = \lambda|\psi\rangle,$$

where  $\lambda = +1$  (bosons) or  $\lambda = -1$  (fermions).

Let us consider bosons.

For  $V = 0$  the eigenstates are easy to guess

$$|\mathbf{i}_1\rangle |\mathbf{i}_2\rangle \dots |\mathbf{i}_N\rangle ,$$

where  $\mathbf{i}_1 = (i_x, i_y, i_z)_1$  are quantum numbers for the first particle.

In the particle representation the symmetrization is a formidable task.  
Let us try differently. Fock states:

$$|\mathbf{n}\rangle = |n_0, n_1, n_2, \dots\rangle,$$

where we do not label particles – we define only how many particles occupy a given single particle state. Let us define

$$[\hat{a}_i, \hat{a}_j^\dagger] = \delta_{i,j},$$

which imply

$$\begin{aligned} \hat{a}_i |\dots, n_i, \dots\rangle &= \sqrt{n_i} |\dots, n_i - 1, \dots\rangle \\ \hat{a}_i^\dagger |\dots, n_i, \dots\rangle &= \sqrt{n_i + 1} |\dots, n_i + 1, \dots\rangle. \end{aligned}$$

Then

$$\hat{H} = \sum_{i=0}^{\infty} E_i \hat{a}_i^\dagger \hat{a}_i + \frac{1}{2} \sum_{i,j,i',j'} V_{i,j;i',j'} \hat{a}_i^\dagger \hat{a}_j^\dagger \hat{a}_{i'} \hat{a}_{j'},$$

where

$$V_{i,j;i',j'} = \int d^3r \int d^3r' \phi_i^*(\vec{r}) \phi_j^*(\vec{r}') V(\vec{r} - \vec{r}') \phi_{i'}(\vec{r}) \phi_{j'}(\vec{r}').$$

Next step

$$\hat{\psi}(\vec{r}) = \sum_{i=0}^{\infty} \phi_i(\vec{r}) \hat{a}_i,$$

that fulfil

$$\left[ \hat{\psi}(\vec{r}), \hat{\psi}(\vec{r}') \right] = 0, \quad \left[ \hat{\psi}(\vec{r}), \hat{\psi}^\dagger(\vec{r}') \right] = \delta(\vec{r} - \vec{r}').$$

With the help of  $\hat{\psi}(\vec{r})$  the Hamiltonian reads

$$\begin{aligned} \hat{H} = & \int d^3r \hat{\psi}^\dagger(\vec{r}) \left[ -\frac{\hbar^2}{2m} \nabla^2 + U(\vec{r}) \right] \hat{\psi}(\vec{r}) \\ & + \frac{1}{2} \int d^3r \int d^3r' \hat{\psi}^\dagger(\vec{r}) \hat{\psi}^\dagger(\vec{r}') V(\vec{r} - \vec{r}') \hat{\psi}(\vec{r}) \hat{\psi}(\vec{r}'), \end{aligned}$$



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and for  $V(\vec{r} - \vec{r}') = g_0 \delta(\vec{r} - \vec{r}')$  we obtain finally

$$\hat{H} = \int d^3r \hat{\psi}^\dagger(\vec{r}) \left[ -\frac{\hbar^2}{2m} \nabla^2 + U(\vec{r}) + \frac{g_0}{2} \hat{\psi}^\dagger(\vec{r}) \hat{\psi}(\vec{r}) \right] \hat{\psi}(\vec{r}).$$

## Bogoliubov theory

Hamiltonian of  $N$  interacting particles supplemented with a constant term

$$-\mu \int d^3r \hat{\psi}^\dagger(\vec{r}) \hat{\psi}(\vec{r}) = -\mu \hat{N},$$

reads

$$\hat{H} = \int d^3r \left[ \hat{\psi}^\dagger(\vec{r}) H_0 \hat{\psi}(\vec{r}) + \frac{g_0}{2} \hat{\psi}^\dagger(\vec{r}) \hat{\psi}^\dagger(\vec{r}) \hat{\psi}(\vec{r}) \hat{\psi}(\vec{r}) - \mu \hat{\psi}^\dagger(\vec{r}) \hat{\psi}(\vec{r}) \right],$$

where  $H_0 = -\frac{\hbar^2}{2m} \nabla^2 + U(\vec{r})$ .

Suppose the ground state is, at the first approximation, a perfect BEC

$$|N\rangle = \frac{(\hat{a}_0^\dagger)^N}{\sqrt{N!}} |0\rangle,$$

where  $\hat{a}_0^\dagger$  creates atom in a certain (unknown) mode  $\phi_0$ .

Let us decide

$$\hat{\psi}(\vec{r}) = \phi_0(\vec{r}) \hat{a}_0 + \delta\hat{\psi}(\vec{r}),$$

where

$$\delta\hat{\psi}(\vec{r}) = \hat{Q}\hat{\psi}(\vec{r}) = [1 - |\phi_0\rangle\langle\phi_0|]\hat{\psi}(\vec{r}).$$

Assume

$$\langle \hat{a}_0^\dagger \hat{a}_0 \rangle = N_0 \simeq N.$$

that implies

$$dN = \int d^3r \langle \delta\hat{\psi}^\dagger(\vec{r}) \delta\hat{\psi}(\vec{r}) \rangle \ll N,$$

and we can define a *small parameter*, i.e.  $dN/N$ .

Hence, we get estimates

$$\hat{a}_0 \sim \sqrt{N}, \quad \delta\hat{\psi} \sim \sqrt{dN}.$$

► Zero order

In zero order

$$\hat{H}^{(0)} = \int d^3r \left[ \phi_0^*(\vec{r}) H_0 \phi_0(\vec{r}) \hat{a}_0^\dagger \hat{a}_0 + \frac{g_0}{2} |\phi_0(\vec{r})|^4 \hat{a}_0^\dagger \hat{a}_0^\dagger \hat{a}_0 \hat{a}_0 - \mu |\phi_0(\vec{r})|^2 \hat{a}_0^\dagger \hat{a}_0 \right],$$

and we deal with a perfect BEC. One obtains an extremal value of energy

$$\langle N | \hat{H}^{(0)} | N \rangle = \int d^3r \left[ \phi_0^*(\vec{r}) H_0 \phi_0(\vec{r}) N + \frac{g_0}{2} |\phi_0(\vec{r})|^4 N(N-1) - \mu |\phi_0(\vec{r})|^2 N \right],$$

if the Gross-Pitaevskii equation is fulfilled

$$H_0 \phi_0 + g_0(N-1) |\phi_0|^2 \phi_0 = \mu \phi_0.$$



► Second order

Operators

$$\hat{\Lambda}(\vec{r}) = \frac{\hat{a}_0^\dagger}{\sqrt{N}} \delta\psi(\vec{r}),$$

do not change particle number and they fulfil

$$\left[ \hat{\Lambda}(\vec{r}_1), \hat{\Lambda}^\dagger(\vec{r}_2) \right] = \delta(\vec{r}_1 - \vec{r}_2) - \phi_0(\vec{r}_1)\phi_0^*(\vec{r}_2) + \mathcal{O}\left(\frac{dN}{N}\right).$$

Employing  $\hat{\Lambda}$  one obtains

$$\hat{H}^{(2)} = \frac{1}{2} \int d^3r (\hat{\Lambda}^\dagger, -\hat{\Lambda}) \mathcal{L} \left( \begin{array}{c} \hat{\Lambda} \\ \hat{\Lambda}^\dagger \end{array} \right),$$

where

$$\mathcal{L} = \begin{pmatrix} H_{GP} + g_0 N \hat{Q} |\phi_0|^2 \hat{Q} & g_0 N \hat{Q} \phi_0^2 \hat{Q}^* \\ -g_0 N \hat{Q}^* \phi_0^{*2} \hat{Q} & -H_{GP}^* - g_0 N \hat{Q}^* |\phi_0|^2 \hat{Q}^* \end{pmatrix}. \quad (1)$$

Properties of  $\mathcal{L}$

$$\sigma_x \mathcal{L} \sigma_x = -\mathcal{L}^*,$$

$$\sigma_z \mathcal{L} \sigma_z = \mathcal{L}^\dagger,$$

where  $\sigma_x$  and  $\sigma_z$  are the Pauli matrices. These properties imply that if

$$|\psi_k^R\rangle = \begin{pmatrix} u_k \\ v_k \end{pmatrix}, \quad \mathcal{L} |\psi_k^R\rangle = E_k |\psi_k^R\rangle,$$

and  $E_n$  is real-valued, then the left eigenvector is the following

$$|\psi_k^L\rangle \propto \sigma_z |\psi_k^R\rangle = \begin{pmatrix} u_k \\ -v_k \end{pmatrix},$$

and

$$\sigma_x |\psi_k^R\rangle^* = \begin{pmatrix} v_k^* \\ u_k^* \end{pmatrix},$$

is the right eigenvector corresponding to eigenvalue  $-E_n$ .

The eigenvectors of  $\mathcal{L}$  can be divided into a family " + " and " - " according to  $\langle \psi_k^R | \sigma_z | \psi_k^R \rangle = \langle u_k | u_k \rangle - \langle v_k | v_k \rangle = \pm 1$ .

We can now expand the operators

$$\begin{pmatrix} \hat{\Lambda}(\vec{r}) \\ \hat{\Lambda}^\dagger(\vec{r}) \end{pmatrix} = \sum_{k \in \text{"+"}} \hat{b}_k \begin{pmatrix} u_k(\vec{r}) \\ v_k(\vec{r}) \end{pmatrix} + \hat{b}_k^\dagger \begin{pmatrix} v_k^*(\vec{r}) \\ u_k^*(\vec{r}) \end{pmatrix},$$

where the operators

$$\hat{b}_k = \int d^3r [u_k^*(\vec{r}), -v_k^*(\vec{r})] \begin{bmatrix} \hat{\Lambda}(\vec{r}) \\ \hat{\Lambda}^\dagger(\vec{r}) \end{bmatrix} = \frac{1}{\sqrt{N}} [\langle u_k | \delta \hat{\psi} \rangle \hat{a}_0^\dagger - \langle v_k | \delta \hat{\psi}^\dagger \rangle \hat{a}_0],$$

fulfil bosonic commutation relations  $[\hat{b}_k, \hat{b}_{k'}^\dagger] = \delta_{k,k'}$ .

Finally we obtain

$$\hat{H} \simeq E_0(N) + \sum_{k \in "+"} E_k \hat{b}_k^\dagger \hat{b}_k.$$

Eigenstates of  $\hat{H}$ :

The ground state (Bogoliubov vacuum)

$$\hat{b}_k |0_b\rangle = 0.$$

Excited states:

$$|m_1, m_2, \dots\rangle_b = \prod_{k=1}^{\infty} \frac{(\hat{b}_k^\dagger)^{m_k}}{\sqrt{m_k!}} |0_b\rangle.$$

Condensate depletion

$$dN = \int d^3r {}_b \langle m_1, \dots | \hat{\Lambda}^\dagger \hat{\Lambda} | m_1, \dots \rangle_b = \sum_{k=1}^{\infty} [m_k \langle u_k | u_k \rangle + (m_k + 1) \langle v_k | v_k \rangle].$$

The Bogoliubov vacuum in the particle representation:

Suppose  $|0_b\rangle \sim (\hat{d}^\dagger)^M |0\rangle$  where  $[\hat{b}_k, \hat{d}^\dagger] = 0$  for all  $k$ . Then

$$\hat{b}_k (\hat{d}^\dagger)^M |0\rangle = (\hat{d}^\dagger)^M \hat{b}_k |0\rangle = 0.$$

It easy to show that

$$\hat{d}^\dagger = \hat{a}_0^\dagger \hat{a}_0^\dagger + \sum_{k,l=1}^{\infty} Z_{kl} \hat{a}_k^\dagger \hat{a}_l^\dagger,$$

where  $Z = U^{-1} V$  with  $V_{kl} = \langle v_k | \phi_l \rangle$  and  $U_{kl} = \langle u_k | \phi_l \rangle$ . Thus the Bogoliubov vacuum in the particle representation reads

$$|0_b\rangle \sim (\hat{d}^\dagger)^{N/2} |0\rangle = \left( \hat{a}_0^\dagger \hat{a}_0^\dagger + \sum_{k=1}^{\infty} \lambda_k \hat{a}_k^\dagger \hat{a}_k^\dagger \right)^{N/2} |0\rangle.$$

Time evolution:

$$i\hbar \frac{d}{dt} \begin{pmatrix} \hat{\Lambda} \\ \hat{\Lambda}^\dagger \end{pmatrix} = \mathcal{L} \begin{pmatrix} \hat{\Lambda} \\ \hat{\Lambda}^\dagger \end{pmatrix}.$$

With

$$\begin{pmatrix} \hat{\Lambda}(\vec{r}, t) \\ \hat{\Lambda}^\dagger(\vec{r}, t) \end{pmatrix} = \sum_{k \in "+"} \hat{b}_k \begin{pmatrix} u_k(\vec{r}, t) \\ v_k(\vec{r}, t) \end{pmatrix} + \hat{b}_k^\dagger \begin{pmatrix} v_k^*(\vec{r}, t) \\ u_k^*(\vec{r}, t) \end{pmatrix},$$

we get

$$i\hbar \frac{d}{dt} \begin{pmatrix} u_k(\vec{r}, t) \\ v_k(\vec{r}, t) \end{pmatrix} = \mathcal{L}(t) \begin{pmatrix} u_k(\vec{r}, t) \\ v_k(\vec{r}, t) \end{pmatrix},$$

and

$$\frac{d}{dt} \hat{b}_k = 0.$$



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Example:

Gross-Pitaevskii equation

$$-\frac{1}{2} \frac{d^2 \phi_0}{dx^2} + \frac{1}{2} x^2 \phi_0 + g_0 N |\phi_0|^2 \phi_0 = \mu \phi_0,$$

possesses dark *soliton* solution

$$\phi_0 \simeq \sqrt{\frac{\mu_0 - x^2/2}{g_0 N}} \tanh\left(\frac{x}{\xi}\right).$$

Anomalous mode:  $E_1 \simeq -1/\sqrt{2}$ ,

$$f_+ \equiv u_1 + v_1 \simeq \frac{\sqrt{3g_0 N}}{2\sqrt{2} \cosh^2(x/\xi)},$$

$$f_- \equiv u_1 - v_1 \simeq \sqrt{\frac{2}{3}} \sqrt{\frac{\mu_0 - x^2/2}{g_0 N}}.$$

For  $N \simeq 1.5 \cdot 10^5$  rubidium atoms and  $g_0 N = 7500$  we get

$$dN = \int dx \langle 0_b | \hat{\Lambda}^\dagger(x) \hat{\Lambda}(x) | 0_b \rangle = \int dx |v_1(x)|^2 \simeq 60,$$

$$\rho(x) = \langle 0_b | \hat{\psi}^\dagger(x) \hat{\psi}(x) | 0_b \rangle = N |\phi_0(x)|^2 + |v_1(x)|^2,$$

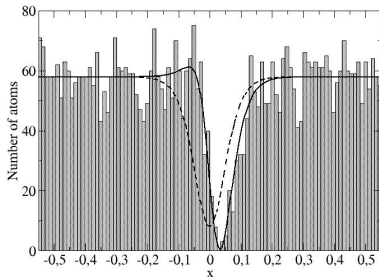
and

$$|0_b\rangle \sim \left( \hat{a}_0^\dagger \hat{a}_0^\dagger + \lambda \hat{a}_1^\dagger \hat{a}_1^\dagger \right)^{N/2} |0\rangle,$$

where

$$\phi_1(x) \simeq \frac{f_+(x)}{\sqrt{\langle f_+ | f_+ \rangle}}, \quad \lambda \simeq \frac{\langle f_+ | f_+ \rangle - \langle f_- | f_- \rangle}{\langle f_+ | f_+ \rangle + \langle f_- | f_- \rangle}.$$





**Figure:** Dashed line: atomic density averaged over many experimental realizations. Histogram: example of a single measurement. Solid line: curve fitted to the histogram.