

Notation

$$|\psi\rangle \in \mathcal{H}_d$$

$$\rho \in \mathcal{B}(\mathcal{H}_d), \mathcal{L}(\mathcal{H}_d)$$

qubit:

$$|\psi\rangle \in \mathbb{C}_2$$

$$|\psi\rangle = \alpha|0\rangle + \beta|1\rangle, |\alpha|^2 + |\beta|^2 = 1$$

$$\rho \in \mathcal{B}(\mathbb{C}_2) \quad \rho = \frac{1}{2}(|\underline{1}\rangle + |\underline{0}\rangle)$$

density operator (mixed state)

- $\hat{A} = \hat{A}^\dagger \leftarrow$ observable, e.g. $\{\hat{x}, \hat{p}\}$

- $\{\rho_i, |\psi_i\rangle\} \leftarrow$ quantum ensemble

$$A_{\psi} = \langle \psi | \hat{A} | \psi \rangle = \text{Tr} \{ |\psi\rangle\langle\psi| \hat{A} \}$$

$$\Downarrow A = \sum_i \rho_i A_{\psi_i} = \sum_i \rho_i \langle \psi_i | \hat{A} | \psi_i \rangle$$

$$= \text{Tr} \left\{ \left(\sum_i \rho_i |\psi_i\rangle\langle\psi_i| \right) \hat{A} \right\}$$

$$\rho = \begin{cases} |\psi\rangle\langle\psi|, \text{ pure} \\ \sum_i \rho_i |\psi_i\rangle\langle\psi_i|, \text{ mixed} \end{cases}$$

N.B. $\langle \psi_i | \psi_j \rangle \neq \delta_{ij}$ not necessarily orthogonal

Home: $\rho^\dagger = \rho, \text{tr}(\rho) = 1, \rho \geq 0$

spectral decompⁿ: positive semidefinite matrix

$$\rho = \sum_{i,j=1}^d \rho_{ij} |i\rangle\langle j| = \sum_{k=1}^{\text{rank}} \lambda_k |e_k\rangle\langle e_k|$$

{eigendecomposition} $\langle e_k | e_l \rangle = \delta_{kl}$ rank-1 pure state

Def. "purity"

$$\begin{aligned} \text{Tr}\{\rho^2\} &= \text{Tr}\{\rho\rho\} = \text{Tr}\left\{\sum_k \lambda_k^2 |e_k\rangle\langle e_k|\right\} \\ &= \sum_k \lambda_k^2 \langle e_k | e_k \rangle = \sum_k \lambda_k^2 \end{aligned}$$

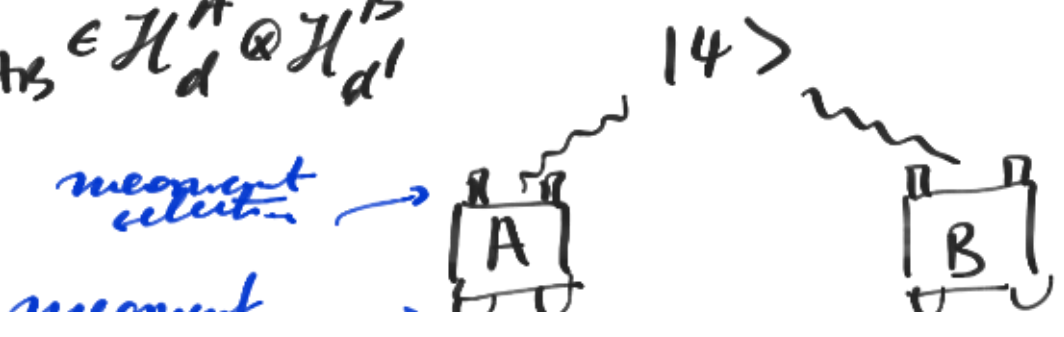
$$\frac{1}{d} \leq \text{Tr}\{\rho^2\} = \sum_k \lambda_k^2 \leq \sum_k \lambda_k = 1$$

"=" max. mixed "=" pure state

Bipartite states

Alice + Bob share a quantum state

$$|\psi\rangle_{AB} \in \mathcal{H}_d^A \otimes \mathcal{H}_{d'}^B$$



maximum outcomes \rightarrow



$$|\psi\rangle_{AB} = \sum_{ij} \alpha_{ij} |i\rangle_A \otimes |j\rangle_B \quad \left\{ \begin{array}{l} = |i,j\rangle = |ij\rangle \\ \text{with } \langle i|j\rangle = \delta_{ij} \end{array} \right.$$

$$\dim(\mathcal{H}^{AB}) = dd'$$

$$d = d' = 2 \quad (2 \text{ qubits}) \Rightarrow \dim(\mathcal{H}_2 \otimes \mathcal{H}_2) = 4$$

basis: (computational) $\{ |00\rangle, |01\rangle, |10\rangle, |11\rangle \}$

Bell basis

$$\begin{aligned} |\psi_+\rangle &= \frac{1}{\sqrt{2}} (|01\rangle + |10\rangle) & |\phi_+\rangle &= \frac{1}{\sqrt{2}} (|00\rangle + |11\rangle) \\ |\psi_-\rangle &= \frac{1}{\sqrt{2}} (|01\rangle - |10\rangle) & &= \frac{1}{\sqrt{2}} (|1\leftrightarrow\rangle - |0\leftrightarrow\rangle) \end{aligned}$$

$\left. \begin{array}{l} \text{SPDC} \\ \text{process} \end{array} \right\}$

$\hat{12}$ form a basis, as

$$\begin{aligned} \langle ij | kl \rangle &= (\langle i | \otimes \langle j |) (|k\rangle_A \otimes |l\rangle_B) = \\ &= \langle i | k \rangle_A \langle j | l \rangle_B = \delta_{ik} \delta_{jl} \end{aligned}$$

density matrices

$$\rho \in \mathcal{B}(\mathcal{H}_d^A \otimes \mathcal{H}_{d'}^B)$$

$$\rho = \sum_i p_i |\psi_i\rangle_{AB} \langle \psi_i| \quad |\psi_i\rangle \in \mathcal{H}_d^A \otimes \mathcal{H}_{d'}^B$$

with $|\psi_i\rangle = \sum_{kl} \alpha_{kl}^{(i)} |kl\rangle$

$$\rho = \sum_i p_i \sum_{kl} \alpha_{kl}^{(i)} |kl\rangle (\alpha_{k'l'}^{(i)})^* \langle k'l'|$$

$$= \sum_{\substack{kl \\ k'l'}} \underbrace{\left(\sum_i \rho_i \alpha_{kl}^{(i)} \alpha_{k'l'}^{(i)*} \right)}_{\rho_{kl, k'l'}} |kl\rangle \langle k'l'|$$

2 qubits: $\{ |00\rangle, |01\rangle, |10\rangle, |11\rangle \}$

$$\begin{pmatrix} 1 \\ 0 \\ 0 \\ 0 \end{pmatrix} \quad \begin{pmatrix} 0 \\ 1 \\ 0 \\ 0 \end{pmatrix} \quad \begin{pmatrix} 0 \\ 0 \\ 1 \\ 0 \end{pmatrix} \quad \begin{pmatrix} 0 \\ 0 \\ 0 \\ 1 \end{pmatrix}$$

e.g. $\int_{11}^{01} |01\rangle \langle 11| = \int_{11}^{01} \begin{pmatrix} 0 \\ 1 \\ 0 \\ 0 \end{pmatrix} (0001) = \begin{pmatrix} 0000 \\ 0001 \\ 0000 \\ 0000 \end{pmatrix}$

\Rightarrow

$$\rho = \begin{pmatrix} \rho_{00}^{00} & \rho_{01}^{00} & \rho_{10}^{00} & \rho_{11}^{00} \\ \rho_{00}^{01} & \rho_{01}^{01} & \rho_{10}^{01} & \rho_{11}^{01} \\ \rho_{00}^{10} & \rho_{01}^{10} & \rho_{10}^{10} & \rho_{11}^{10} \\ \rho_{00}^{11} & \rho_{01}^{11} & \rho_{10}^{11} & \rho_{11}^{11} \end{pmatrix} \begin{matrix} |00\rangle \\ |01\rangle \\ |10\rangle \\ |11\rangle \end{matrix}$$

Definition

1) Pure states: $|\psi\rangle \in \mathcal{H}_A \otimes \mathcal{H}_B$ notation!
 is product iff $|\psi\rangle = |\psi_A\rangle \otimes |\psi_B\rangle$

2) Mixed states: $\rho \in \mathcal{B}(\mathcal{H}_A \otimes \mathcal{H}_B)$

is separable iff "mixture of product states"

$$\rho = \sum_i p_i |\phi_i\rangle\langle\phi_i| \otimes |\chi_i\rangle\langle\chi_i|$$

{ N.B. product state is separable with TRIVIAL

$$p_i = \begin{cases} 1, & i=1 \\ 0, & i \neq 1 \end{cases}$$

3) Any state that is not separable is entangled

Motivation: classical correlation



observable $\hat{A} \otimes \hat{B}$ {e.g. $\hat{X} \otimes \hat{P}$ }

correlator:

$$\langle AB \rangle = \text{Tr} \left\{ \rho_{AB} \hat{A} \otimes \hat{B} \right\}$$

$$\rho_{AB} \in \text{SEP} \Rightarrow \rho_{AB} = \sum_i p_i |\phi_i\rangle\langle\phi_i|_A \otimes |\chi_i\rangle\langle\chi_i|_B$$

$$\langle AB \rangle_{\text{sep}} = \sum_i p_i \langle \hat{A} \rangle_{\phi_i} \langle \hat{B} \rangle_{\chi_i}$$

classical correlations

• convex mixture of separable states

is separable ρ

$$\rho_1, \rho_2 \in \text{SEP} \Rightarrow \lambda \rho_1 + (1-\lambda) \rho_2 \in \text{SEP}$$

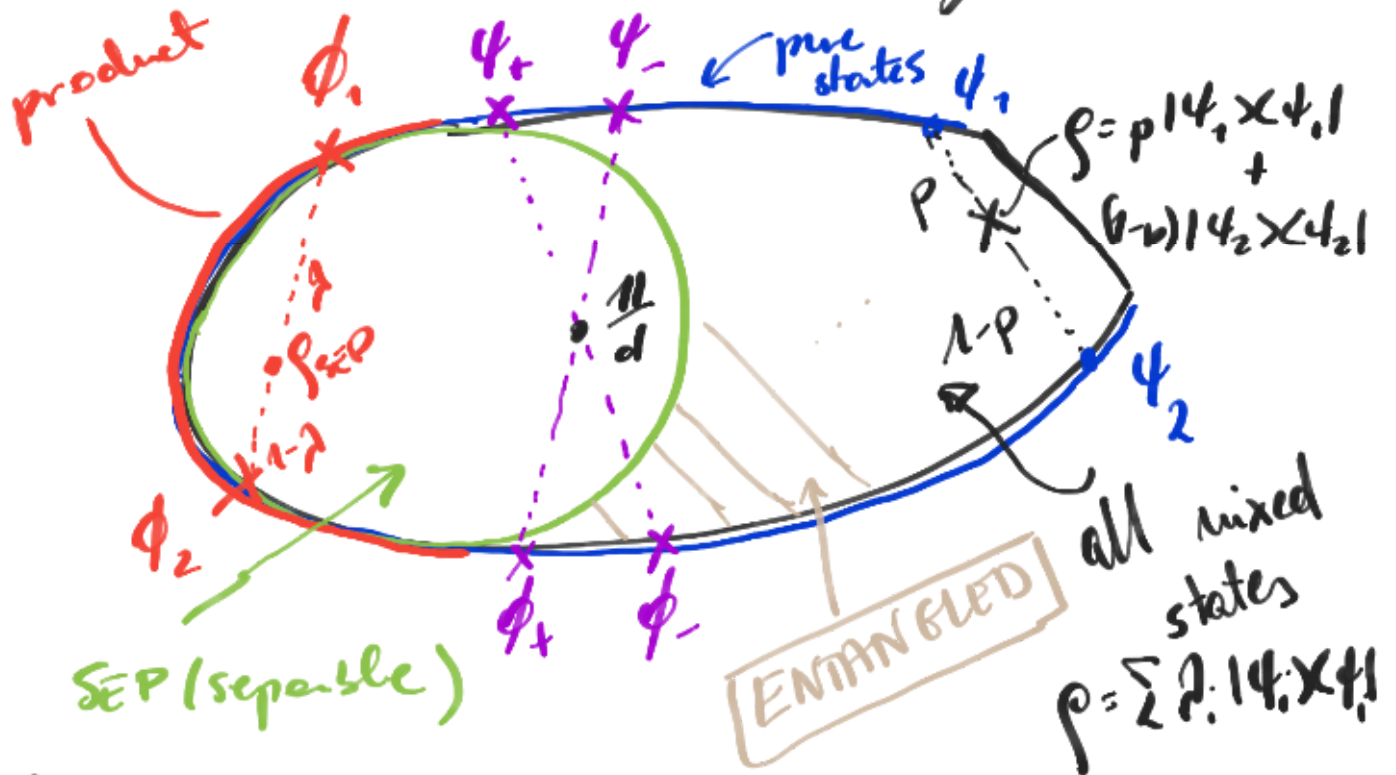
👉 (Proof)

$$\text{for } 0 \leq \lambda \leq 1$$

• ok, but ^{all} mixed states $\rho \in \mathcal{B}(\mathcal{H}_A \otimes \mathcal{H}_B) \geq 0$
form also a convex set

$$\forall \rho, \sigma \geq 0 : \lambda \rho + (1-\lambda) \sigma \geq 0$$

$$\text{for } 0 \leq \lambda \leq 1$$



$$1) \quad |\phi_1\rangle = |\chi_1\rangle_A |\psi_1\rangle_B$$

$$|\phi_2\rangle = |\chi_2\rangle_A |\psi_2\rangle_B$$

$$\rho_{\text{sep}} = \lambda |\phi_1\rangle\langle\phi_1| + (1-\lambda) |\phi_2\rangle\langle\phi_2|$$

2) Max. mixed state

$$\underline{I} = \frac{1}{d} \sum |ij\rangle\langle ij| = \sum \frac{1}{d} |i\rangle\langle i| \otimes |j\rangle\langle j|$$

$d \quad d \quad ij \quad \dots \quad ij \quad d$
 mixture of product states

but! for 2 qubits

$$\frac{\mathbb{1}}{4} = \frac{1}{4} (|\psi_+\rangle\langle\psi_+| + |\psi_-\rangle\langle\psi_-| + |\phi_+\rangle\langle\phi_+| + |\phi_-\rangle\langle\phi_-|)$$

Bell states (entangled!)
 form an orthonormal basis!

\Rightarrow by mixing entangled states we can obtain separable states
 (noise)

Entanglement verification

1) Pure states

Schmidt decomposition

$$|\psi\rangle_{AB} = \sum_{i=1}^{d_A} \sum_{j=1}^{d_B} d_{ij} |i\rangle_A |j\rangle_B$$

Singular value decomposition (SVD)
 of a complex matrix

$$d_{ij} = \sum_{k=1}^r U_{ik} \sigma_k V_{kj}^*$$

$\alpha: d_A \neq d_B$

$$d_A \left[\begin{array}{c} \alpha \\ \vdots \\ \alpha \end{array} \right]_{d_B} = d_A \left[\begin{array}{c} U^* \\ \vdots \\ U \end{array} \right]_r \left[\begin{array}{c} \lambda_1 \\ \vdots \\ \lambda_r \\ 0 \\ \vdots \\ 0 \end{array} \right] \underbrace{\left[\begin{array}{c} v \\ \vdots \\ v \end{array} \right]}_{d_B} \Bigg\}^r$$

$r \rightarrow$ "Schmidt rank"

$$\begin{aligned} |\psi\rangle_{AB} &= \sum_{ij} \sum_k U_{ik}^* \lambda_k V_{kj} |i\rangle |j\rangle = \\ &= \sum_{k=1}^r \lambda_k \left(\sum_{i=1}^{d_A} U_{ki}^* |i\rangle \right) \left(\sum_{j=1}^{d_B} V_{kj} |j\rangle \right) \\ &= \sum_{k=1}^r \lambda_k |e_k\rangle_A \otimes |f_k\rangle_B \end{aligned}$$

"Schmidt coefficients"

SVD : $\Rightarrow 1 \leq r(\text{Schmidt rank}) \leq \min\{d_A, d_B\}$

$\bullet r=1 \Rightarrow |\psi\rangle_{AB} = |e\rangle_A \otimes |f\rangle_B$

product state (separable)

$\bullet r > 1 \Rightarrow$ entangled state

$\bullet r = \min\{d_A, d_B\} \wedge \forall_k: |\lambda_k| = \frac{1}{\sqrt{r}}$

(equal weights)

$$|\psi\rangle_{AB} = \frac{1}{\sqrt{r}} \sum_k e^{i\phi_k} |e_k\rangle_A \otimes |f_k\rangle_B$$

\Rightarrow maximally entangled state

2 qubits:

• Maximally entangled

$$|\psi_{ME}\rangle = \frac{1}{\sqrt{2}} (|\phi\rangle|\chi\rangle + e^{i\varphi} |\phi_{\perp}\rangle|\chi_{\perp}\rangle)$$

e.g.: $|\psi_{-}\rangle = \frac{1}{\sqrt{2}} (|01\rangle - |10\rangle)$

and so $|\psi_{+}\rangle, |\phi_{\pm}\rangle \Rightarrow$ Bell states are max. entangled

• Partially entangled

e.g. $|\psi_{+}(\lambda)\rangle = \frac{1}{\sqrt{2}} (\sqrt{\lambda} |01\rangle + \sqrt{1-\lambda} |10\rangle)$

with $0 < \lambda < \frac{1}{2} < \lambda < 1$
product \nearrow $\frac{1}{2}$ \nwarrow product
max. ent.

Mixed states

• no universal way!

• sufficient condition:

... partial transposition

negativity under partial transposition
(PT)

suff. and necessary
only for $d_A \times d_B$
 2×2
 2×3

Transposition

$$\begin{aligned} \rho^T &= \sum_{ij, kl} \rho_{kl}^{ij} (|ij\rangle\langle kl|) = \sum_{ij, kl} \rho_{kl}^{ij} (|i\rangle\langle k|)_A \otimes (|j\rangle\langle l|)_B \\ &= \sum_{ij, kl} \rho_{kl}^{ij} |k\rangle\langle i| \otimes |l\rangle\langle j| = \sum_{ij, kl} \rho_{kl}^{ij} |kl\rangle\langle ij| \\ &= \sum_{ij, kl} \rho_{ij}^{kl} |ij\rangle\langle kl| \end{aligned}$$

Partial transposition

$$\begin{aligned} \rho^{T_A} &= \sum_{ij, kl} \rho_{kl}^{ij} (|i\rangle\langle k|)_A \otimes |j\rangle\langle l|_B \\ &= \sum_{ij, kl} \rho_{kl}^{ij} |ij\rangle\langle kl| \end{aligned}$$

similarly $\rho^{T_B} = \sum_{ij, kl} \rho_{kl}^{ij} |ij\rangle\langle kl|$

For separable states

Proof:

$$\rho \in \text{SEP} \Rightarrow \rho = \sum_i p_i |4_i\rangle\langle 4_i| \otimes |\phi_i\rangle\langle \phi_i|$$

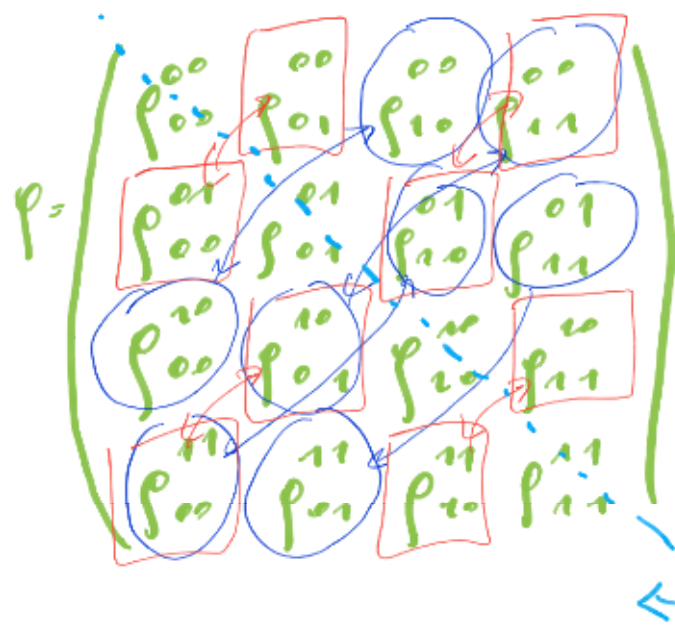
$$\Rightarrow \rho^{T_A} = \sum_i p_i (|4_i\rangle\langle 4_i|)^T \otimes |\phi_i\rangle\langle \phi_i|$$

similarly: $\rho^{T_B} \geq 0$ also a state
 { N.B. $(\rho^{T_A})^T = \rho^{T_B}$ }

Hence: $\rho \in \text{SEP} \Rightarrow (\rho^{T_A} \geq 0 \Leftrightarrow \rho^{T_B} \geq 0)$

Hence: $(\rho^{T_A} < 0 \Leftrightarrow \rho^{T_B} < 0) \Rightarrow \rho$ is entangled

2 qubits ($d_A = d_B = 2$)



$\rho^T, \rho^{T_A}, \rho^{T_B}$

State is entangled if and only if for 2×2 (2×3)
 $\rho^{T_A} < 0, \rho^{T_B} < 0$

