Hartree-Fock-Bogoliubov

A. Two-body (density-dependent) interaction:

$$\hat{H} = \sum_{ij} t_{ij} c_i^+ c_j + \frac{1}{4} \sum_{ijkl} \overline{v}_{ijkl} c_i^+ c_j^+ c_l c_k$$

B. Variational principle:
$$\delta E[\Psi] = 0; \quad E[\Psi] = \frac{\langle \Psi | \hat{H} | \Psi \rangle}{\langle \Psi | \Psi \rangle}; \quad E[\Psi] > E_0$$

C. Trial wave functions: product states $\beta_{K}^{+} = \sum_{i} \left(u_{iK} c_{i}^{+} + v_{iK} c_{i} \right)$ general Bogoliubov transformation $|\Psi\rangle = \prod_{K} \beta_{K} |-\rangle; \quad \beta_{K} |\Psi\rangle = 0$

HFB - quasiparticles incorporated into the HF formalism HFB wave function is the quasuparticle vacuum

- HFB wave function the most general product wave function consisting of independently moving quasiparticles (in HF: particles)
- Selfconsistent description of coupling between p-h an p-p channels

HFB - density matrix and pairing tensor

HFB density matrix

HFB pairing tensor

$$\rho_{ij} = \langle \Psi | c_j^{\dagger} c_i | \Psi \rangle, \quad \kappa_{ij} = \langle \Psi | c_j c_i | \Psi \rangle$$
$$\hat{\rho} = v^* v^T, \quad \hat{\kappa} = v^* u^T$$
$$\hat{\rho}^2 - \hat{\rho} = -\hat{\kappa}\hat{\kappa}^+, \quad \hat{\rho}\hat{\kappa} = \hat{\kappa}\hat{\rho}$$

$$\hat{\Re} = \begin{pmatrix} \hat{\rho} & \hat{\kappa} \\ -\hat{\kappa}^* & 1 - \hat{\rho}^* \end{pmatrix}, \quad \hat{\Re}^2 = \hat{\Re}$$
$$\hat{\Re} \begin{pmatrix} u_K \\ v_K \end{pmatrix} = 0, \quad \hat{\Re} \begin{pmatrix} v_K^* \\ u_K^* \end{pmatrix} = \begin{pmatrix} v_K^* \\ u_K \end{pmatrix}$$

Generalized density matrix

Quasiparticle

interaction

Eigenvalues are 0 or 1 (thus defining occupations of quasiparticle states)

$$\hat{H}(c^+,c) \Rightarrow \hat{H}(\beta^+,\beta) = E_0 + \hat{H}_{20} + \hat{H}_{11} + \hat{H}_{int}$$

Independent quasiparticle Hamiltonian (H_{HFB})

$$\left[\hat{H},\hat{N}\right] = 0, \text{ but } \left[\hat{H}_{HFB},\hat{N}\right] \neq 0$$

However, we require that $\langle \hat{N} \rangle = N \Rightarrow \hat{H}' = \hat{H} - \lambda \hat{N}$

HFB equations

$$\begin{split} & \mathcal{K} = \begin{pmatrix} \hat{h} - \lambda & \hat{\Delta} \\ -\hat{\Delta}^* & -\hat{h}^* + \lambda \end{pmatrix} \\ & h_{ij} = t_{ij} + \Gamma_{ij} \\ & \Gamma_{ij} = \sum_{kl} \nabla_{iljk} \rho_{kl} \\ & \Delta_{ij} = \frac{1}{2} \sum_{kl} \nabla_{ijkl} \kappa_{kl} \\ \end{split}$$
HFB equations
$$\begin{aligned} & \left(\hat{h} - \lambda & \hat{\Delta} \\ -\hat{\Delta}^* & -\hat{h}^* + \lambda \right) \begin{pmatrix} u_K \\ v_K \end{pmatrix} = E_K \begin{pmatrix} u_K \\ v_K \end{pmatrix} \\ & \text{[K, $\hat{S}t$] = 0$ \\ \end{aligned} \\ \end{split}$$
HFB equations treat p-h and p-p on the same footing Δ is a state-dependent field. In general, it depends on density and has a kinetic term The generalized density matrix and the HFB Hamiltonian can be diagonalized simultaneously \\ \end{split}

Fermi level determined from the particle number equation

Often it is convenien to express the HFB equations in the canonical basis

Independent quasiparticle approach



