

The EPRL intertwiners and correct partition function.

Wojciech Kamiński

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- Is the EPRL map isometric, does it preserve the scalar product between the SU(2) intertwiners?
- If not, what is a complete, correct form of the partition function written directly in terms of the SU(2) intertwiners, the preimages of the EPRL map?



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Euclidean 4d QG as spin-foam

- 4d QG is regarded as a BF theory with constraints.
- Spin-networks consists of gauge invariant functions.

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- Constraints are imposed as projections on edges (nodes of spin-network).



• The simplicity constraints are imposed on the elements of \mathcal{H}_{Σ} , locally at each vertex. $\operatorname{Inv}_{\operatorname{Simp}}(\rho_1 \otimes ... \rho_k \otimes \rho_{k+1}^* \otimes ... \otimes \rho_N^*)$



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- In each subsequent spin-network intertwiners *l* should be in this subspace,
- To sum with respect to the spin-network histories with the amplitude as a weight, one fixes an *orthonormal basis* in each space of simple intertwiners.

There are three main proposals for the simple intertwiners:

- 1. that of Barrett-Crane (BC) corresponding to the Palatini action,
- 2. that of Engle-Pereira-Rovelli-Livine (EPRL) corresponding the Holst action with the value of the Barbero-Immirzi parameter $\gamma \neq \pm 1$,
- 3. that of Freidel-Krasnov (FK) also corresponding to the Holst action with the value of the Barbero-Immirzi parameter $\gamma \neq \pm 1$,

Given intertwiner $\eta \in Inv(\rho_{k_1} \otimes ...)$ with the spins k_I

$$j_I^{\pm} = \left| \frac{1 \pm \gamma}{2} \right| k_I$$

 $k_I = j^+ + j^-$, if $|\gamma| < 1$ and $k_I = |j^+ - j^-|$, if $|\gamma| > 1$. This follows from

adjusted/improved constraints.

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The map EPRL

$$\iota_{\mathrm{EPRL}}(\eta) := (P^+ \otimes P^-)c_1 \otimes \ldots \otimes c_n \lrcorner \eta.$$

- P^{\pm} projections onto $SU(2)_{\pm}$ invariants,
- c_i Clebsch-Gordon coefficients $\rho_{k_i} \to \rho_{j_i^+} \otimes \rho_{j_i^-}$.

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- The proof is based on the observation that in suitable basis matrix of the ι_{EPRL} is "upper" triangular and injective on "diagonal" entries.



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then

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We can restrict our attention to $k_{12} = j_{12}^+ + j_{12}^-$



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Imposing additional constraints on j_{12}^{\pm} allows for inductive procedure.



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There are some technical details...

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$$Z[\psi] = \sum_{j_f^{\pm}, \eta_e} \prod_{f \ face} \dim \rho_{j_f^{\pm}, j_f^{\pm}}$$
$$\prod_{v \ vertex} A_v(\{\iota(\eta_e)\})\psi(\{\iota(\eta_{bdr})\}),$$

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$$\prod_{e \ edge} A(\eta_e^{in}, \eta_e^{out}) \prod_{v \ vertex} A_v(\{\iota(\eta_e)\})\psi(\{\iota(\eta_{bdr})\}),$$

 $A(\eta_1, \eta_2)$ is the inverse of the matrix $\langle \iota(\eta_1), \iota(\eta_2) \rangle$ so

$$P_{EPRL} = \sum_{\eta_{in/out}} A_{in,out} |\iota_{EPRL}(\eta_{out})\rangle \langle \iota_{EPRL}(\underline{\eta_{in}})|$$

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- Examples of low j indicate that matrix A is approximately diagonal (conjecture).
- If this is true A does not change (spoil) asymptotic behavior of the spin foam amplitude.



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Summary

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- Simple examples shows that it is not unitary for $|\gamma| \neq 1$.
- The basis labelled by SU(2) intertwiners is not orthonormal and we should introduce an additional factor A in the spin-foam amplitude

$$\langle \eta_1 | A^{-1} \eta_2 \rangle = \langle \iota_{EPRL}(\eta_1) | \iota_{EPRL}(\eta_2) \rangle$$