

EPRL map for Spin 4 and $SL(2, \mathbb{C})$ and integrability of Lorentzian Spin foams.

Spin foam for Spin 4



$V_{j^+ j^-}$ for faces
 i - intertwiners on
edges

Vertex amplitude $\in \text{Inv } V_{j^+ j^-}$
is the evaluation
of this spin-network
on contraction
of the intertwiners
{faces that
meet in the
edge}

Choice of intertwiners:

$$\Upsilon_\gamma: V_k \rightarrow V_{j^+ j^-}$$

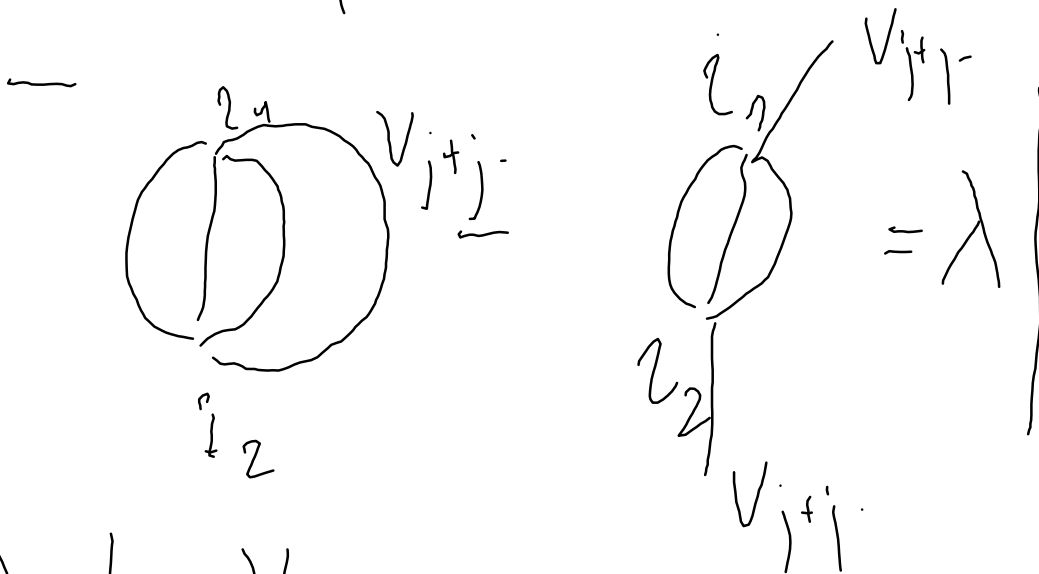
$$j^\pm = \frac{1 \mp \gamma_k}{2}$$

Υ_γ ← Immirzi parameter

$$\text{Res}_{SU(2)} V_{j^+ j^-} = \bigoplus_{j=j^+ - j^-}^{j^+ + j^-} V_j$$

$P_{\text{inv}} : \text{Inv}_{SU(2)} \otimes V_{\lambda} \rightarrow \text{Inv}_{\text{spin4}} \otimes V_{\lambda}$
 ↑
 projection onto invariants spin4.

Problems for $SL(2, \mathbb{C})$



$\lambda \dim V_{j^+ j^-}$

For $SL(2, \mathbb{C})$ contraction of intertwiners ill-defined

$$- \text{Inv}_{SL(2, \mathbb{C})} \otimes V_{n, \xi} \neq \otimes V_{n, \xi}$$

Representations of $SL(2, \mathbb{C})$

$$f(z_1, z_2) - \quad z_1, z_2 \in \mathbb{D}$$

$$f(\lambda z_1, \lambda z_2) = \lambda^k \lambda^{-k} |\lambda|^{2ip-2} f(z_1, z_2)$$

$V_{k, p}$

On the sphere $|z_1|^2 + |z_2|^2 = 1$

$$f(e^{i\varphi} z_1, e^{i\varphi} z_2) = e^{2ik\varphi} f(z_1, z_2)$$

Basis $\left(\bigoplus_{j \in \mathbb{Z}} V_j \right)$ as $SU(2)$ representation

Group averaging

Way of defining $\text{Inv } SL(2, \mathbb{C})$

$$\bigotimes_{k_i, p_i} V_{k_i, p_i} \supset \mathcal{D} \quad [U] \in \mathcal{D}^*$$

$$[\varphi] : [\varphi] \varphi' = \dots$$

$$\varphi \in \mathcal{D} = \int dg \langle \varphi | U(g) \phi \rangle$$

it gives us a natural candidate for scalar product on the space of invariants

Example $L^2(\mathbb{R})$ e^{ikx}

$$\mathcal{D} = C_0^\infty(\mathbb{R})$$

$$\int dk \langle \varphi | e^{ikx} \varphi' \rangle = \bar{\varphi}(0) \varphi'(0)$$

if we take $\mathcal{D} = C_0^\infty(\mathbb{R})$
 $\varphi(0) = 0$

In our case

$$\mathcal{D} = \sum_{d_j: g} V_{j_1} \otimes \dots \otimes V_{j_n}$$

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$$v_{n_i, g_i} \sim \bigcup_{j_i} v_{j_i}$$

$$\int dg \langle \Psi | U(g) | \Psi' \rangle$$

is absolutely convergent
provided that $n \geq 3$

- Contraction:

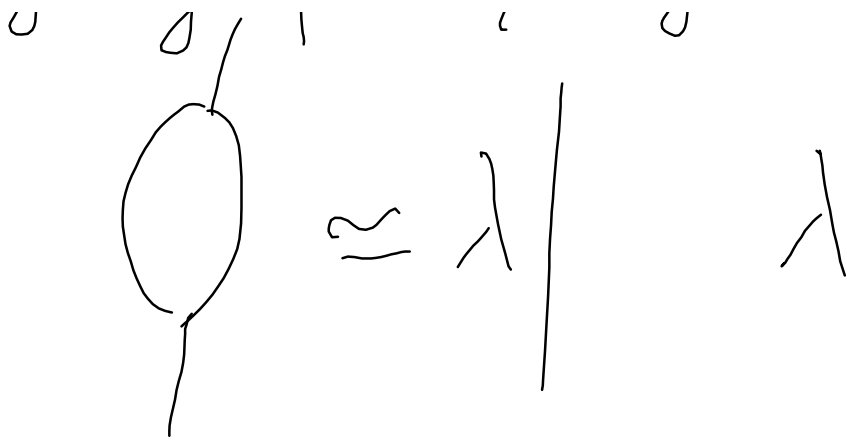
$[\Psi_i]$ for every node
of the graph.

$$\int dg_i \Psi_1 \dots \Psi_n \mathcal{D}(g_1^{-1} g_n)$$

Gauge symmetry

$$g_i \rightarrow h g_i$$

Gauge fixing $g_0 = e$



Spin 4

$$S(g) = \sum_{j^+ j^-} \frac{(l_j^+ + 1)}{(l_j^- + 1)} \Theta_{j^+ j^-}(g)$$

$SL(2, \mathbb{Q})$ Horish
- Chouhox

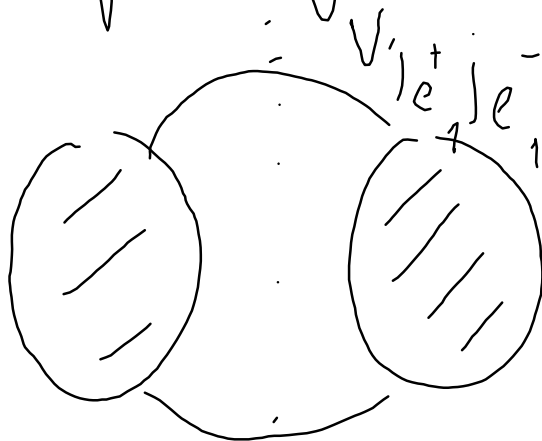
$$\delta(g) = \sum_{k, s} \int dg$$

$$(s^2 + k^2) \Theta_{k, s}(g)$$

$$\Theta_{-k, -s}(g) = \Theta_{k, s}(g)$$

$$\dim V_{k, s} = (k^2 + s^2)$$

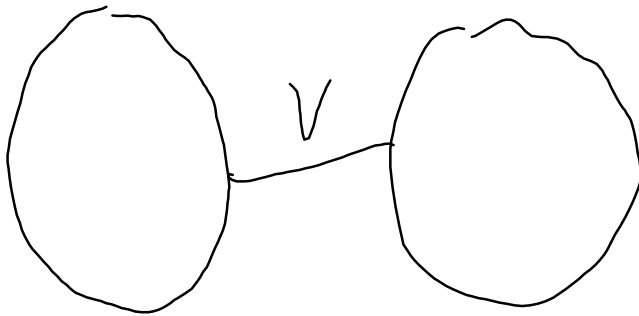
First prescription = $\lambda \dim$



Spin 4.

$$V_i^+ \quad V_i^- \\ |e_2| \quad |e_2|$$

\rightsquigarrow $SLG(\mathbb{Z}, \mathbb{Q})$ 2-edge
 reducible graph
 should not be
 absolutely conv.



What with 3-edge connected graphs

- Every such integral (evaluation) is absolutely convergent.

Proof:

$$\int \prod_{i \in E_0 \setminus \{0\}} dg_i \prod_i \varphi_i^{\mu_{e_1} \dots \mu_{e_n}}$$

$$\prod_{e \in E^1} D(g_i^{-1} g_j)$$

$(i \rightarrow j)$
 (e)

$$\langle \sum_{i \in e} D(g_i^{-1} g_j) \sum_{j \in e} \rangle$$

$$\varphi = \bigotimes_{e \text{ coming to } i} \sum_{j \in e} \sum_{j \in e} \in V_{j \in e}$$

$$\int \prod_i dg_i \prod_e \langle \quad \rangle$$

$$K > 1 \in (r_{j \in e}, j \in e) \frac{1}{(\cosh r)^{1-r}}$$

$$\cosh r = \frac{1}{2} \text{tr} (g_i^{-1} g_j)^{\dagger} (g_i^{-1} g_j)$$

\uparrow

$$dg = dp du$$

↑
of angle

$$\ll \int \prod dp_i \prod \frac{1}{e (\cosh r_{ij})^{1-\alpha}}$$

measure on the
unit hyperboloid

$$(n, \bar{\eta}) \quad n_0^2 - \bar{\eta}^2 = 1$$

$$(E, \bar{\xi})$$

↑
3 dim vector
with norm 1

$$\int \prod \frac{dE_i}{E_i^3} \prod d\xi_i$$

$$\prod \left(\frac{E_i E_j}{\Theta_{ij}^2 + E_i^2 + E_j^2} \right)^{1-\alpha}$$

Θ_{ij} - distance on 3-dim sphere between ξ_i and ξ_j .

Subdivide region of integration into sub-region of different scales,

Θ_{ij}, ϵ_i

$$\int \epsilon_1 \leq \Theta_{12} \leq \dots$$

$$\int_0^{\Theta_{12}} d\epsilon_1 \dots \pi(\dots)^{1-\gamma}$$

this allows us to estimate

Example

$$\int_0^{\Theta_{12}} d\epsilon \left(\frac{\epsilon_1 \epsilon_2}{\Theta_{12}^2 + \epsilon_1^2 + \epsilon_2^2} \right)^{1-\gamma} \approx$$

$$\left(\frac{\epsilon_1 \epsilon_2}{\theta_{12}^2 + \epsilon_1^2 + \epsilon_2^2} \right)^{1-\alpha}$$

→ 1

if we integrate
over ϵ_1, ϵ_2
and θ_{12}

$$d\bar{\xi}_i d\bar{\xi}_j \approx \theta_{12} d\theta_{12} d\phi d\bar{\xi}_{ij}$$

After every integration we have
additional χ coming from $d\epsilon_i$
 χ^2 coming from $\theta_{ij} \text{ and } \theta_{ij}$

Power counting \geq
external legs of the following
graph:

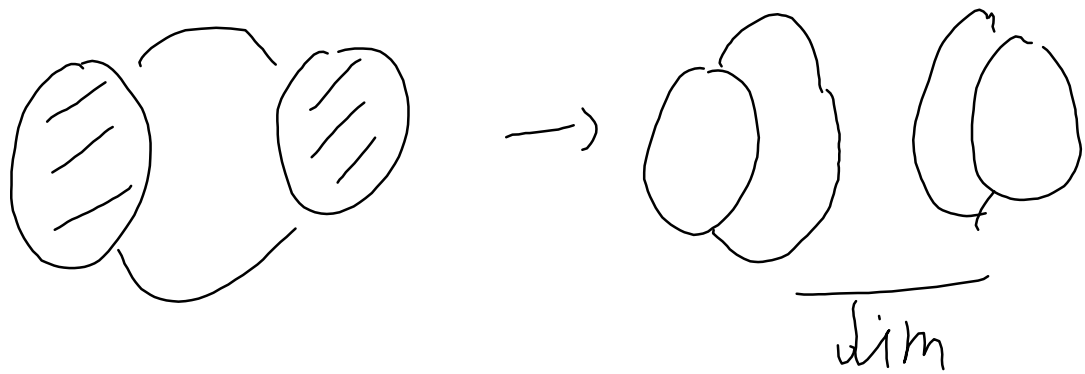
In our spin-network we take all vertices that correspond to E_i integrated out \Rightarrow edges we take all edges with Θ_{ij} integrated out E_i, E_j int.

$$\text{power} \Rightarrow \sum_{\text{conn.}} -2 + E_{\mu} (1 - \gamma)$$

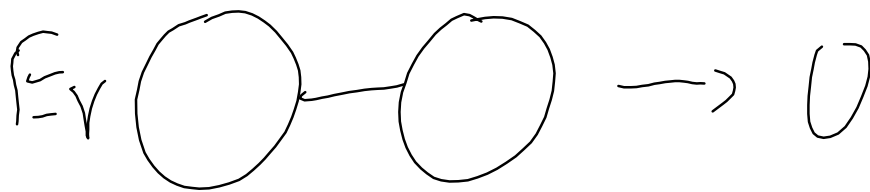
\uparrow
 number of external legs

technicalities

2-edge reducible.



There are ways to obtain
 reasonable evaluation for
 every non-edge-reducible graph



Remark: There is dependence
 on the h'w'le of Haar measure