# $GL(2,\mathbb{R})$ geometry of 5th order ODEs

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Dedicated to Professor Oldřich Kowalski on the occasion of a Conference in his honour

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- This  $\rho(SO(3))$  may be defined as a subgroup of a SO(5) stabilizing rank three *traceless symmetric* tensor  $\Upsilon \in S_0^3 \mathbb{R}^5$ , which is related to the metric gvia:

$$\Upsilon_{jki}\Upsilon_{lmi} + \Upsilon_{lji}\Upsilon_{kmi} + \Upsilon_{kli}\Upsilon_{jmi} = g_{jk}g_{lm} + g_{lj}g_{km} + g_{kl}g_{jm}.$$

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- An irreducible SO(3) structure  $(M^5, g, \Upsilon)$  is called *nearly integrable* if  $\Upsilon$  is a *Killing tensor* for g:

$$\nabla^{LC}_X \Upsilon(X, X, X) = 0, \qquad \forall X \in TM^5.$$

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- Thus, nearly integrable SO(3) structures provide *low-dimensional examples* of *Riemannian* geometries which can be described in terms of a *unique metric* connection ( $\Gamma$ ) with *totally skew symmetric* torsion (T).
- This sort of geometries are studied extensively by the string theorists.

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- We do not know if *nonhomogeneous* examples exist.
- Perhaps these structures are so rigid that they must be homogeneous.

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What are the possible dimensions n in which there exists a tensor  $\Upsilon$  satisfying:

- i)  $\Upsilon_{ijk} = \Upsilon_{(ijk)}$ , (total symmetry)
- ii)  $\Upsilon_{ijj} = 0$ , (no trace)
- $\text{iii)} \Upsilon_{jki}\Upsilon_{lmi} + \Upsilon_{lji}\Upsilon_{kmi} + \Upsilon_{kli}\Upsilon_{jmi} = g_{jk}g_{lm} + g_{lj}g_{km} + g_{kl}g_{jm}?$

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• Such tensor is needed to construct isoparametric hypersurfaces in spheres.

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• A polynomial I, in variables  $a_i$ , is called an *algebraic invariant* of  $w_4(x,y)$  if it changes according to

$$I \to I' = (\det b)^p I, \qquad b \in \mathbf{GL}(2, \mathbb{R})$$

under the action of this 5-dimensional representation on  $a_i$ s.
• The lowest order invariants of  $w_4(x,y)$  are:

$$I_2 = 3a_2^2 - 4a_1a_3 + a_0a_4$$

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• Defining  $\Upsilon_{ijk}$  and  $g_{ij}$  via

$$\Upsilon_{ijk}a_ia_ja_k = 3\sqrt{3I_3}$$

#### $g_{ij}a_ia_j = I_2,$

one can check that the so defined  $g_{ij}$  and  $\Upsilon_{ijk}$  satisfy the desidered relations i)-iii).

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- The stabilizer of the conformal class  $[(g, \Upsilon)]$  is the irreducible  $\operatorname{GL}(2, \mathbb{R})$  in dimension five.

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$$(g,\Upsilon,A) \sim (g',\Upsilon',A') \iff \left(g' = e^{2\phi}g, \ \Upsilon' = e^{3\phi}\Upsilon, \ A' = A - 2d\phi\right),$$

is called an *irreducible*  $GL(2, \mathbb{R})$  structure in dimension five.

# Nearly integrable $\mathbf{GL}(2,\mathbb{R})$ structures in dimension 5

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• An irreducible  $\operatorname{GL}(2,\mathbb{R})$  structure  $(M^5, [(g, \Upsilon, A)])$  is called *nearly integrable* iff tensor  $\Upsilon$  is a *conformal* Killing tensor for  $\stackrel{W}{\nabla}$ :

 $\stackrel{W}{\nabla}_X \Upsilon(X, X, X) + \frac{1}{2}A(X)\Upsilon(X, X, X) = 0, \qquad \forall X \in \mathbf{T}M^5.$ 

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• To achieve the uniqueness one requires the that torsion T of  $\nabla$ , considered as an element of  $\bigotimes^{3} T^{*}M^{5}$ , seats in a 10-dimensional subspace  $\bigwedge^{3} T^{*}M^{5}$ .

• In terms of the connection 1-forms of the Weyl connection  $\Gamma$ , and the characteristic connection  $\Gamma$ , we have

$${\stackrel{W}{\Gamma}} = \Gamma + \frac{1}{2}T,$$

where  $\overset{W}{\Gamma} \in \mathfrak{co}(3,2) \otimes \mathrm{T}^*M^5$ ,  $\Gamma \in \mathfrak{gl}(2,\mathbb{R}) \otimes \mathrm{T}^*M^5$  and  $T \in \bigwedge^3 \mathrm{T}^*M^5$ .

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 The converse is also true: if an irreducible GL(2, ℝ) structure in dimension five admits a connection ∇ satisfying

$$\nabla_X g + A(X)g = 0, \qquad \nabla_X \Upsilon + \frac{3}{2}A(X)\Upsilon = 0,$$

and having totally skew symmetric torsion  $T \in \bigwedge^{3} T^{*}M^{5}$  then it is nearly integrable.

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Can we produce examples of the nearly integrable GL(2, ℝ) geometries in dimension five? Can we produce examples with 'pure' torsion in ∧<sub>3</sub> or ∧<sub>7</sub>? Can we produce nonhomogeneous examples?

Ordinary differential equation y<sup>(5)</sup> = 0 has GL(2, ℝ) ×<sub>ρ</sub> ℝ<sup>5</sup> as its group of contact symmetries. Here ρ : GL(2, ℝ) → GL(5, ℝ) is the 5-dimensional irreducible representation of GL(2, ℝ).

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- This, in particular, means that  $y^{(5)} = 0$  may be described in terms of a *flat*  $\mathfrak{gl}(2,\mathbb{R})$ -valued connection on the principal fibre bundle  $\operatorname{GL}(2,\mathbb{R}) \to P \to M^5$  over the solution space  $M^5$  of the ODE.

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## A well known fact

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- What about more complicated 5th order ODEs?

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 $50D^2\overline{F_4} - 75DF_3 + 50F_2 - 60F_4DF_4 + 30F_3F_4 + 8F_4^3 = 0$ 

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 $150F_4DF_3 + 200F_2F_4 - 140F_4^2DF_4 + 130F_3F_4^2 + 14F_4^4 = 0$ 

 $1250D^{2}F_{2} - 6250DF_{1} + 1750DF_{3}DF_{4} - 2750F_{2}DF_{4} - 875F_{3}DF_{3} + 1250F_{2}F_{3} - 500F_{4}DF_{2} + 700(DF_{4})^{2}F_{4} + 1250F_{1}F_{4} - 1050F_{3}F_{4}DF_{4} + 350F_{3}^{2}F_{4} - 350F_{4}^{2}DF_{3} + 550F_{2}F_{4}^{2} - 280F_{4}^{3}DF_{4} + 210F_{3}F_{4}^{3} + 28F_{4}^{5} + 18750F_{y} = 0,$ 

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• Then the 5-dimensional solution space of the equation is naturally equipped with a nearly integrable  $\mathbf{GL}(2,\mathbb{R})$  structure.

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 $550F_2F_4^2 - 280F_4^3DF_4 + 210F_3F_4^3 + 28F_4^5 + 18750F_y = 0,$ 

- Then the 5-dimensional solution space of the equation is naturally equipped with a nearly integrable  $\mathbf{GL}(2,\mathbb{R})$  structure.
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- Then the 5-dimensional solution space of the equation is naturally equipped with a nearly integrable  $\mathbf{GL}(2,\mathbb{R})$  structure.
- Every nearly integrable  $GL(2, \mathbb{R})$  structure obtained in this way has torsion of its characteristic connection of the 'pure' type  $T \in \Lambda_3$ .
- We call the three conditions on F the Wünschmann-like conditions.

# Examples of *F* satisfying the Wünschmann-like conditions

The three differential equations

$$y^{(5)} = c \left( \frac{5y^{(3)3}(5 - 27cy''^2)}{9(1 + cy''^2)^2} + 10 \frac{y''y^{(3)}y^{(4)}}{1 + cy''^2} \right),$$

with c = +1, 0, -1, represent the only three contact nonequivalent classes of Wünschmann-like ODEs having the corresponding nearly integrable  $GL(2, \mathbb{R})$  structures  $(M^5, [g, \Upsilon, A])$  with the characteristic connection with vanishing torsion.

In all three cases the holonomy of the Weyl connection  $\Gamma$  of structures  $(M^5, [g, \Upsilon, A])$  is reduced to the  $\mathbf{GL}(2, \mathbb{R})$ . For all the three cases the Maxwell 2-form  $dA \equiv 0$ . The corresponding Weyl structure is flat for c = 0. If  $c = \pm 1$ , then in the conformal class [g] there is an Einstein metric of positive (c = +1) or negative (c = -1) Ricci scalar. In case c = 1 the manifold  $M^5$  can be identified with the homogeneous space  $\mathbf{SU}(1,2)/\mathbf{SL}(2,\mathbb{R})$  with an Einstein g descending from the Killing form on  $\mathbf{SU}(1,2)$ . Similarly in c = -1 case the manifold  $M^5$  can be identified with the homogeneous space  $\mathbf{SL}(3,\mathbb{R})/\mathbf{SL}(2,\mathbb{R})$  with an Einstein g descending from the Killing form on  $\mathbf{SL}(3,\mathbb{R})$ . In both cases with  $c \neq 0$  the metric g is not conformally flat.

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$$F = \frac{5(8y_3^3 - 12y_2y_3y_4 + 3y_1y_4^2)}{6(2y_1y_3 - 3y_2^2)},$$
$$F = \frac{5y_4^2}{3y_3} \pm y_3^{5/3},$$

represent four nonequivalent nearly integrable  $GL(2,\mathbb{R})$  structures corresponding to the different signs in the second expression and to the different signs of the denominator in the first expression. These structures have 6-dimensional symmety group.

$$F = \frac{1}{9(y_1^2 + y_2)^2} \times$$

 $\left( 5w \left( y_1^6 + 3y_1^4 y_2 + 9y_1^2 y_2^2 - 9y_2^3 - 4y_1^3 y_3 + 12y_1 y_2 y_3 + 4y_3^2 - 3y_4 (y_1^2 + y_2) \right) + 45y_4 (y_1^2 + y_2) (2y_1 y_2 + y_3) - 4y_1^9 - 18y_1^7 y_2 - 54y_1^5 y_2^2 - 90y_1^3 y_2^3 + 270y_1 y_2^4 + 15y_1^6 y_3 + 45y_1^4 y_2 y_3 - 405y_1^2 y_2^2 y_3 + 45y_2^3 y_3 + 60y_1^3 y_3^2 - 180y_1 y_2 y_3^2 - 40y_3^3 \right),$  where

 $w^{2} = y_{1}^{6} + 3y_{1}^{4}y_{2} + 9y_{1}^{2}y_{2}^{2} - 9y_{2}^{3} - 4y_{1}^{3}y_{3} + 12y_{1}y_{2}y_{3} + 4y_{3}^{2} - 3y_{1}^{2}y_{4} - 3y_{2}y_{4}.$ 

This again has 6-dimensional symmetry group.

An ansatz

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This equation may be solved explicitly giving example of ODEs having its nearly integrable structure being nonhomogeneous.

• If a 3rd order ODE y''' = F(x, y, y', y'') satisfies the Wünschmann condition  $9D^2F_2 - 18F_2DF_2 - 27DF_1 + 4F_2^3 - 18F_1F_2 + 54F_y = 0,$ 

 $D = \partial_x + y_1 \partial_y + y_2 \partial_{y_1} + F \partial_{y_2},$ 

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then it defines a *Lorentzian* conformal structure on the **3**-dimensional space of its solutions.

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$$D = \partial_x + y_1 \partial_y + y_2 \partial_{y_1} + F \partial_{y_2}$$

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• This conformal structure in dimension *three* is related to the quadratic  $\mathbf{GL}(2,\mathbb{R})$  invariant  $\Delta = a_0a_2 - a_1^2$  of  $w_2(x,y) = a_0x^2 + 2a_1xy + a_2y^2$ .

 $4D^2F_3 - 8DF_2 + 8F_1 - 6DF_3F_3 + 4F_2F_3 + F_3^3 = 0,$ 

$$4D^2F_3 - 8DF_2 + 8F_1 - 6DF_3F_3 + 4F_2F_3 + F_3^3 = 0,$$

 $160D^2F_2 - 640DF_1 + 144(DF_3)^2 - 352DF_3F_2 + 144F_2^2 - 640DF_1 + 64$ 

 $80DF_2F_3 + 160F_1F_3 - 72DF_3F_3^2 + 88F_2F_3^2 + 9F_3^4 + 16000F_y = 0,$ 

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 $80DF_2F_3 + 160F_1F_3 - 72DF_3F_3^2 + 88F_2F_3^2 + 9F_3^4 + 16000F_y = 0,$ 

$$D = \partial_x + y_1 \partial_y + y_2 \partial_{y_1} + y_3 \partial_{y_2} + F \partial_{y_3}$$

then it defines an irreducible  $\mathbf{GL}(2,\mathbb{R})$  structure on the 4-dimensional space  $M^4$  of its solutions.

• This  $GL(2,\mathbb{R})$  structure in dimension *four* may be understood in terms of a *conformal* Weyl-like structure associated with the *quartic*  $GL(2,\mathbb{R})$  invariant

$$I_4 = -3a_1^2a_2^2 + 4_0a_2^3 + 4a_1^3a_3 - 6a_0a_1a_2a_3 + a_0^2a_3^2,$$

of

$$w_3(x,y) = a_0 x^3 + 3a_1 x^2 y + 3a_2 x y^2 + a_3 y^3$$

and a certain 1-form A on  $M^4$ 

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This is a report on a *joint* work with my student Michał Godliński.