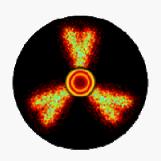
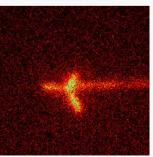
# Two-proton radioactivity and α decay

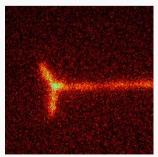


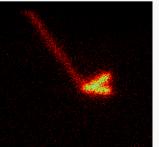
### Marek Pfützner

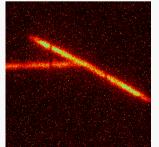
Faculty of Physics, University of Warsaw

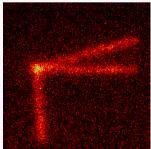












### A disclaimer



#### PHYSICAL REVIEW LETTERS

week ending 31 MAY 2013

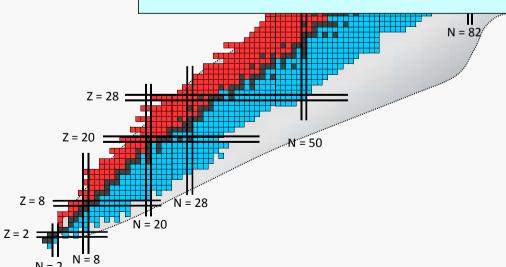
#### Landscape of Two-Proton Radioactivity

E. Olsen, 1,2 M. Pfützner, 3,4 N. Birge, 1,2 M. Brown, 1,5 W. Nazarewicz, 1,2,3 and A. Perhac 1,2 1,2 Department of Physics and Astronomy, University of Tennessee, Knoxville, Tennessee 37996, USA 2,2 Physics Division, Oak Ridge National Laboratory, Oak Ridge, Tennessee 37831, USA 3,4 Faculty of Physics, University of Warsaw, ul. Hoża 69, 00-681 Warsaw, Poland 4,5 Physics Department, 1211 Geneva 23, Switzerland 5,5 Physics Department, Berea College, Berea, Kentucky 40404, USA (Received 4 March 2013; published 29 May 2013)

## The drip lines

➤ The **proton drip-line** is close and almost fully delineated. In most cases, however, it is "invisible" when we cross it. The decay spectroscopy may strech far beyond it.

The questions: how far beyond the proton drip-line we have to go to see the difference? How far is the limit?



➤ The neutron drip-line is far from present experimental reach.
It represents the real limit of decay spectroscopy – the region beyond, if accessible, is a domain of reactions.

N = 126

## Beyond the proton drip-line

#### Competition between two decay modes

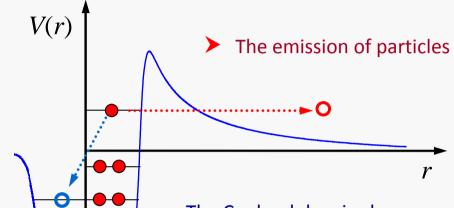
 $\rightarrow$  The  $\beta^+$  decay

Probability of transition:

$$\lambda \sim Q^5$$

Decay energy may be large, but the weak interaction is really weak

→ 
$$T_{1/2} > 1 \text{ ms}$$



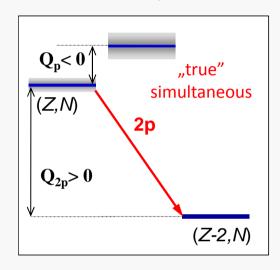
The Coulomb barrier hampers emission of an unbound charged particle ( $\alpha$ , p, 2p,...)

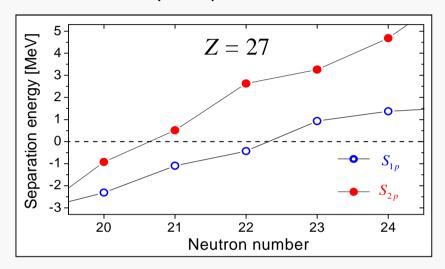
$$\lambda \sim \exp\left\{-\frac{2}{\hbar} \cdot \int_{r_{in}}^{r_{out}} \sqrt{2\mu[V(r)-Q]} \cdot dr\right\}$$

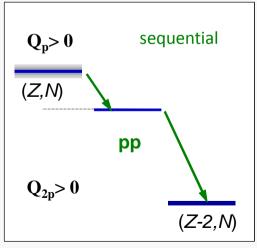
→ To find where the drip-line actually is and to predict which decay will happen, we need: a) atomic masses, b) decay models

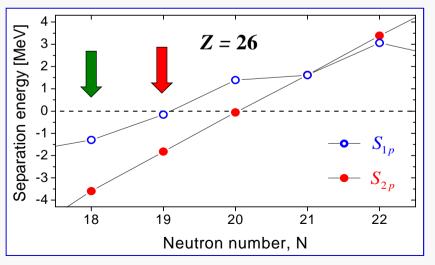
### The answer for even Z

➤ The limit of stability for even-Z elements is determined by two-proton emission









V.I. Goldanskii, Nucl. Phys. 19 (60) 482

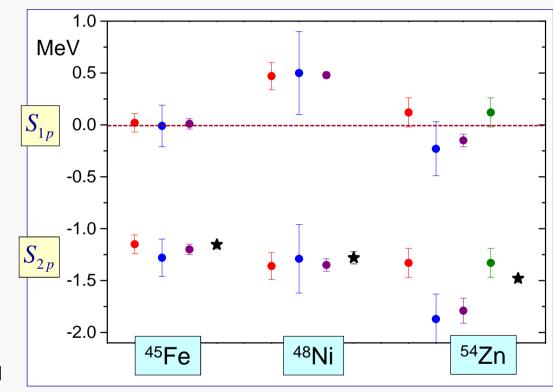
### First 2p candidates

Light and medium masses can be precisely predicted by a trick based on the IMME:

$$BE(T_z = -T) = BE(T_z = T) - 2bT$$

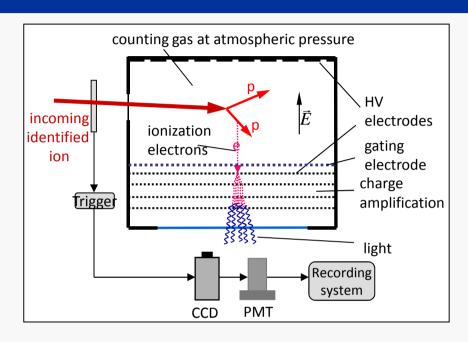
➤ Binding energy of the neutrondeficient nuclide is calculated from the **measured mass** of its neutronrich analogue and from the calculated **coefficient** *b* (shell-model, systematics...)

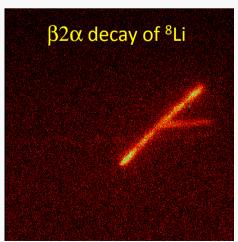
#### Predicted 1p and 2p separation energies

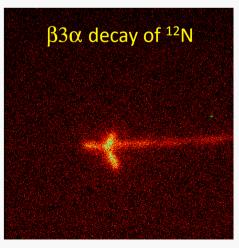


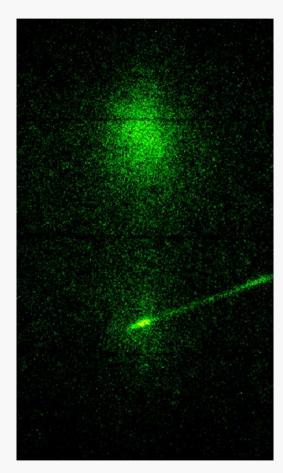
- Brown, PRC 43 (91) R1513
- Ormand, PRC 55 (97) 2407
- **\*** exp
- Cole, PRC 54 (96) 1240
- Brown et al., PRC 65 (02) 045802

## TPC with optical readout



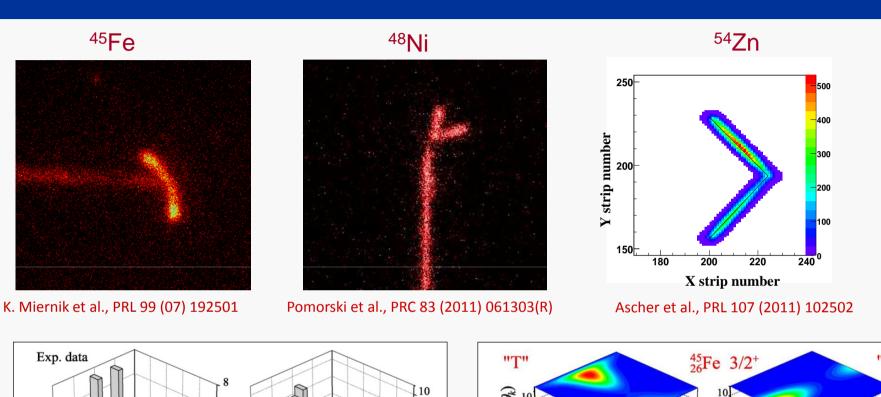


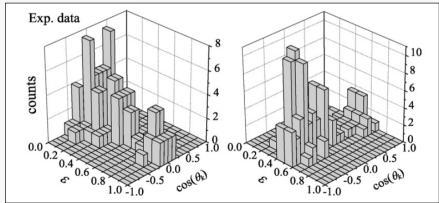


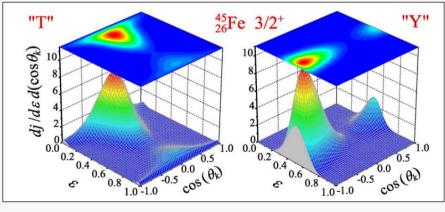


→ Decay event <sup>6</sup>He →  $\alpha$  + dseen on the background of about 10<sup>4</sup> beta rays

### Three cases around Z=28

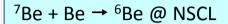


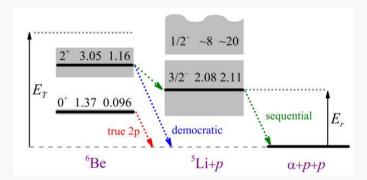


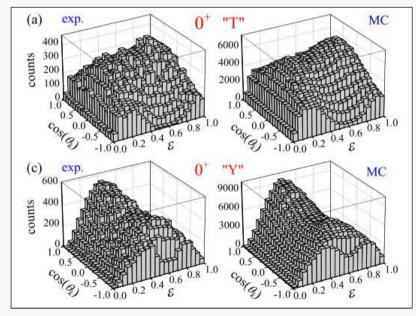


Grigorenko et al., PLB 677 (2009) 30

## <sup>6</sup>Be and <sup>19</sup>Mg

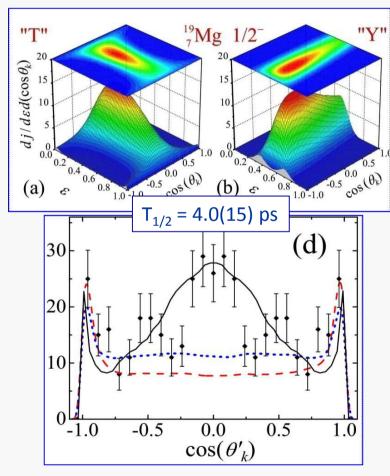






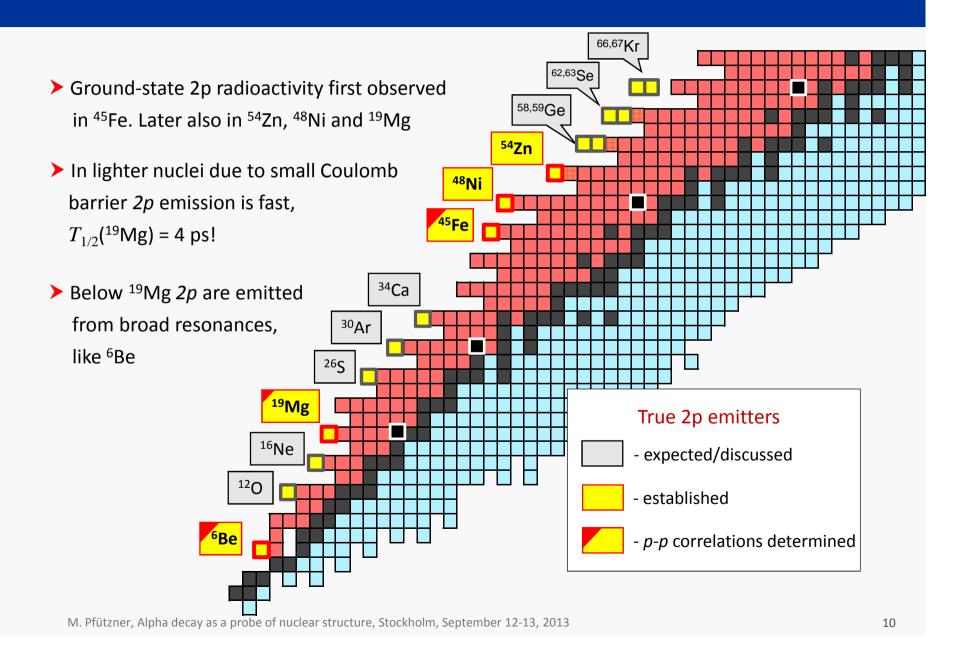
Egorova et al., PRL 109 (2012) 202502

#### $^{20}$ Mg + Be $\rightarrow$ $^{19}$ Mg @ GSI

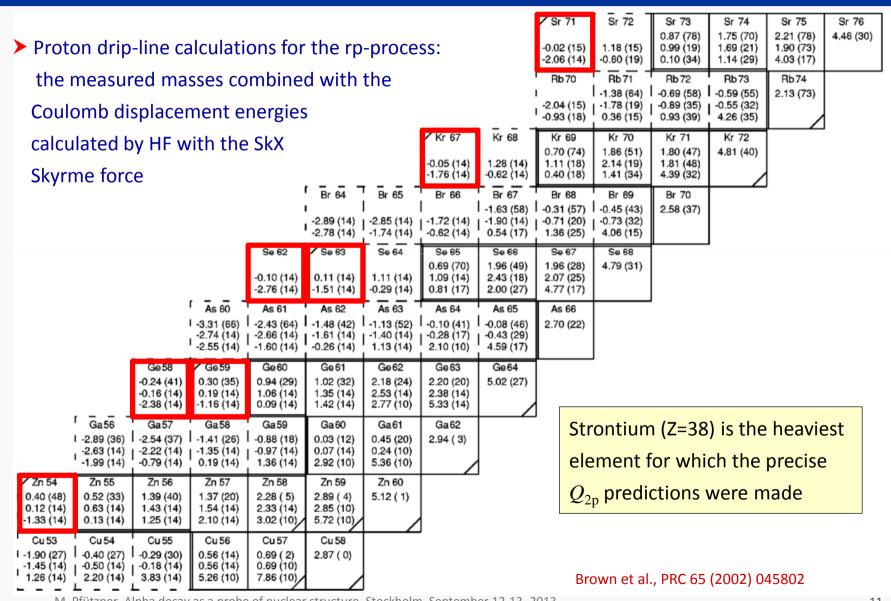


Mukha et al., PRL 99 (2007) 182501 Mukha et al., EPJA 42 (2009) 421

## The current status of 2p emission

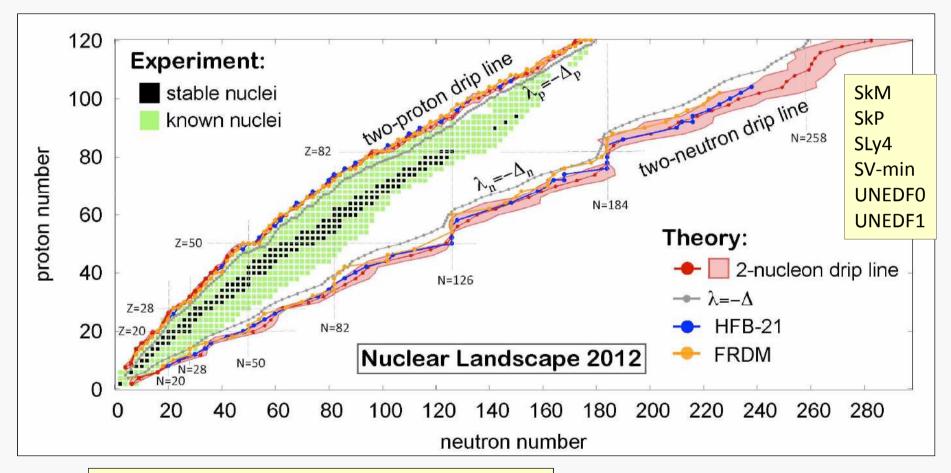


### Heavier 2p candidates



## Nuclear landscape

➤ Global mass predictions using density functional theory with 6 different Skyrme interactions



 $\rightarrow$  There are 6900  $\pm$  500 nuclei bound with Z  $\leq$  120

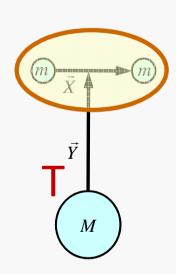
Erler et al., Nature 486 (2012) 509

### Diproton model

 $\blacktriangleright$  By simplifying interactions in the core+p+p system, the three-body decay can be reduced to the <u>combination of two-body processes</u>.

Jacobi T system → diproton model

#### The WKB approximation



$$\Gamma_{2p,dipr} = \theta_{dipr}^2 \mathcal{N} \frac{\hbar^2}{4\mu} \exp\left[-2 \int_{r_{in}}^{r_{out}} k(r) dr\right]$$

$$\mathcal{N}\int_{r_{1}}^{r_{in}}\frac{dr}{2k\left(r\right)}=1 \qquad k\left(r\right)=\sqrt{2\mu\left|Q_{2p}-2V_{p}\left(r\right)\right|}$$

$$\theta_{\text{dipr}}^2 = \frac{(2n)!}{2^{2n} (n!)^2} \left[ \frac{A}{A-2} \right]^{2n} O^2 \qquad n \approx (3Z)^{1/3} - 1$$

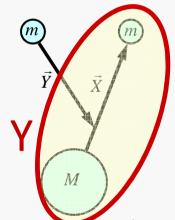
The value of proton overlap function determined from the experimental half-lives of known

2p emitters: <sup>19</sup>Mg, <sup>45</sup>Fe, <sup>48</sup>Ni, and <sup>54</sup>Zn

### Direct model

#### Jacobi Y system → direct model

$$\Gamma_{2p,dir} = \frac{Q_{2p}}{2\pi} \left( Q_{2p} - 2E_p \right)^2 \int_0^1 d\varepsilon \frac{\Gamma_x \left( \varepsilon Q_{2p} \right)}{\left( \varepsilon Q_{2p} - E_p \right)^2 + \Gamma_x \left( \varepsilon Q_{2p} \right)^2 / 4} \times \frac{\Gamma_y \left( (1 - \varepsilon) Q_{2p} \right)}{\left( (1 - \varepsilon) Q_{2p} - E_p \right)^2 + \Gamma_y \left( (1 - \varepsilon) Q_{2p} \right)^2 / 4}$$



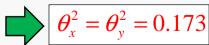
 $\Gamma_i$  is the width of the two-body subsystem:  $\Gamma_i(E) = 2 \gamma_i^2 P_{l_n}(E, R, Z_i)$ 

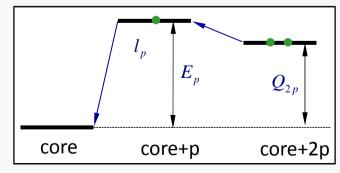
penetrability: 
$$P_{l_p}\left(E,R,Z_i\right) = \frac{kR}{F_{l_p}^2\left(\eta,kR\right) + G_{l_p}^2\left(\eta,kR\right)}$$
 reduced width: 
$$\gamma_i^2 = \frac{\hbar^2}{2\mu_i R^2} \theta_i^2$$

reduced width: 
$$\gamma_i^2 = \frac{\hbar^2}{2\mu_i R^2} \theta_i^2$$

The value of spectroscopic factor determined from the experimental half-lives of known 2p emitters:

$$^{19}$$
Mg,  $^{45}$ Fe,  $^{48}$ Ni, and  $^{54}$ Zn, assuming  $l_p = 0$ 





Grigorenko and Zhukov, PRC 76 (07) 014009 M.P. et al, RMP (2012) 567

### 2p-emission half-lives

**Direct model** 

$$\Gamma_{2p,dir} \cong \frac{8Q_{2p}}{\pi (Q_{2p} - 2E_p)^2} \int_0^1 d\varepsilon \, \Gamma_x (\varepsilon Q_{2p}) \Gamma_y ((1 - \varepsilon) Q_{2p})$$

**Diproton model** 

$$\Gamma_{2p,dipr} = \theta_{dipr}^2 \mathcal{N} \frac{\hbar^2}{4\mu} \exp\left[-2 \int_{r_{in}}^{r_{out}} k(r) dr\right]$$

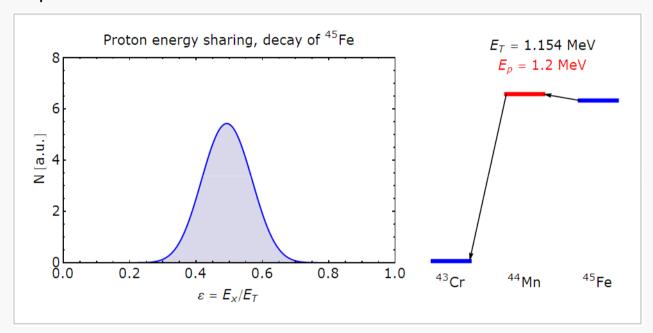
➤ The comparison of predicted half-lives with experiment

$$T_{1/2} = \frac{\ln 2I}{\Gamma}$$

$$l_{\rm p} = 0$$

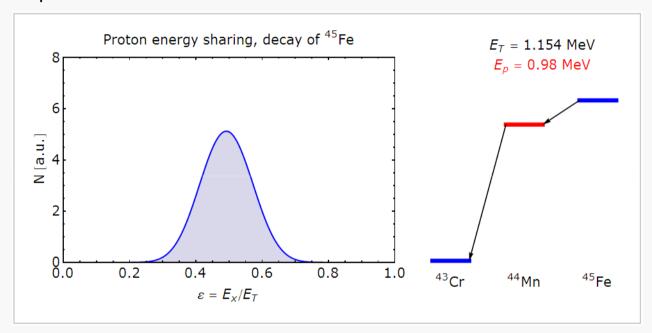
Nucleus	Experiment	Direct	Diproton
<sup>19</sup> Mg [7]	4.0(15) ps	6.2 ps	12.3 ps
<sup>45</sup> Fe [10]	3.7(4) ms	1.1 ms	8.7 ms
<sup>48</sup> Ni [8]	$3.0^{+2.2}_{-1.2}$ ms	6.8 ms	5.3 ms
<sup>54</sup> Zn [9]	$1.98^{+0.73}_{-0.41}$ ms	1.0 ms	0.8 ms

➤ In the direct model we can investigate how the proton's energy spectrum depends on the position of the intermediate state



$$Q_{2p} = 1.15 \text{ MeV}, \ Q_{1p} = -0.05 \text{ MeV}$$
 True 2p decay (simultaneous)

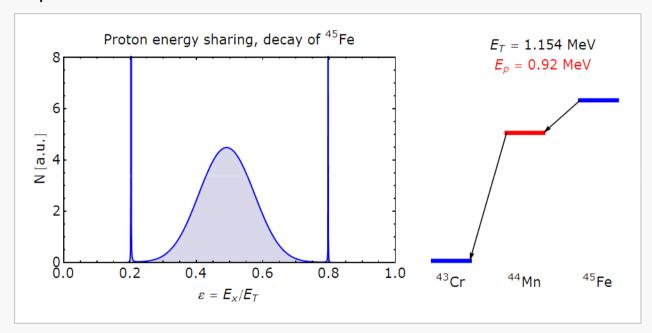
➤ In the direct model we can investigate how the proton's energy spectrum depends on the position of the intermediate state



$$Q_{2p} = 1.15 \text{ MeV}, \ Q_{1p} = 0.17 \text{ MeV}$$

→ Still simultaneous 2p!

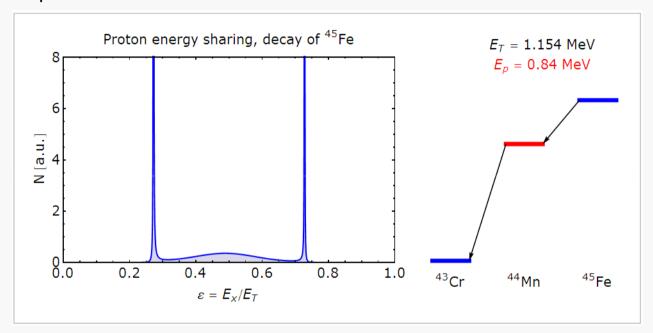
➤ In the direct model we can investigate how the proton's energy spectrum depends on the position of the intermediate state



$$Q_{2p} = 1.15 \text{ MeV}, \ Q_{1p} = 0.23 \text{ MeV}$$

Sequential emission shows up!
Simultaneous component still visible.

➤ In the direct model we can investigate how the proton's energy spectrum depends on the position of the intermediate state



$$Q_{2p} = 1.15 \text{ MeV}, \ Q_{1p} = 0.31 \text{ MeV}$$

Sequential 2p emission dominates

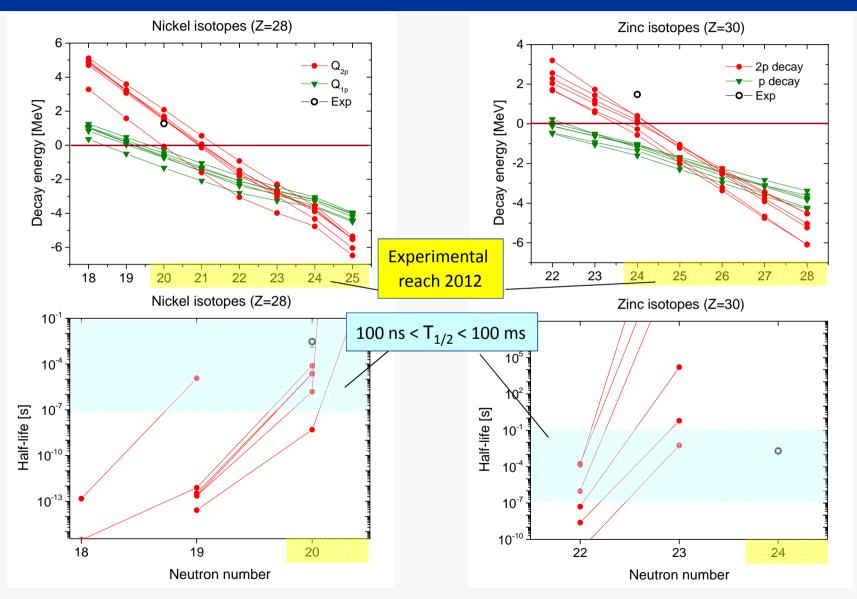
ightharpoonup Rough criterion: for  $Q_p < 0.2 \, Q_{2\,p}$  true, simultaneous 2p decay for  $Q_p > 0.2 \, Q_{2\,p}$  sequential 2p emission

### **Predictions**

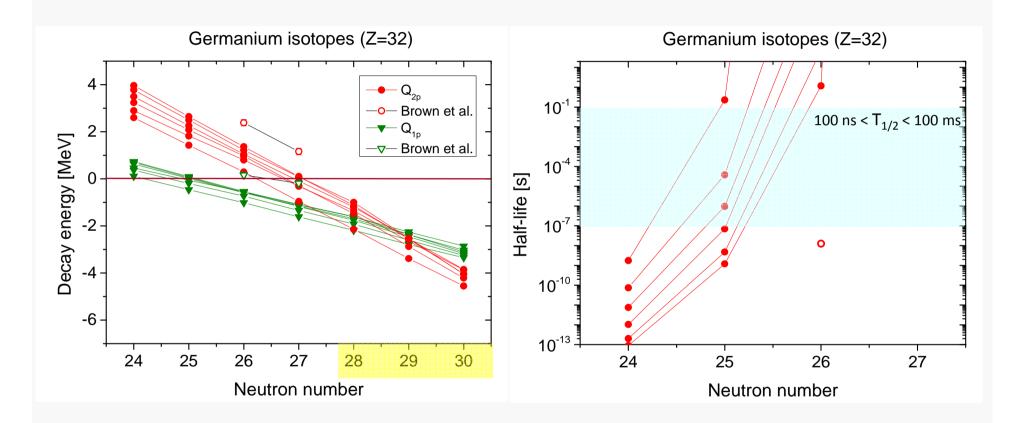
- ➤ Nuclear binding energies: deformed DFT with six effective Skyrme interaction plus density-dependent zero-range pairing term (Erler et al., Nature 486 (2012) 509)
- The half-lives for 2p emission: estimated with the direct and diproton models. The α decay half-lives calculated using global, fenomenological formula by Koura, J. Nucl. Science and Tech. 49 (2012) 816
- The adopted decay-time criterion (arbitrary): we consider a nucleus to be a 2p decay candidate predicted by a given mass (and decay) model when 100 ns < T<sub>1/2</sub> < 100 ms.</p>
  Longer half-life will loose competition with β decay.
  Shorter will be difficult to detect using in-flight separation and implantation technique.
- **Counting:**

a candidate has the model multiplicity m(Z,N) = k when it is predicted by k mass models.

### Nickel and zinc in the direct model



### Germanium

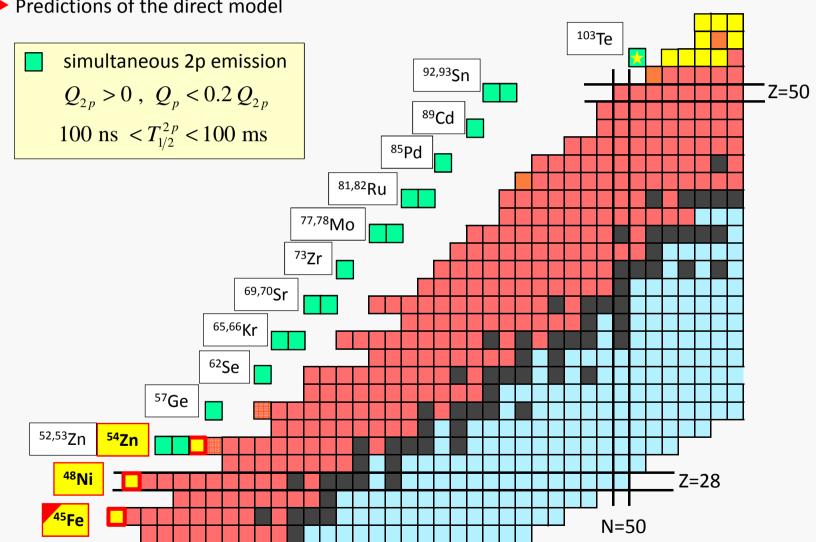


- $\rightarrow$  We predict <sup>57</sup>Ge to be 2p radioactive (m=2)
- ➤ Taking decay energies from Brown, the 2p half-life of <sup>58</sup>Ge comes shorter than 100 ns and that of <sup>59</sup>Ge longer than 100 ms

  Brown et al., PRC 65 (2002) 045802

## Heavy 2p landscape

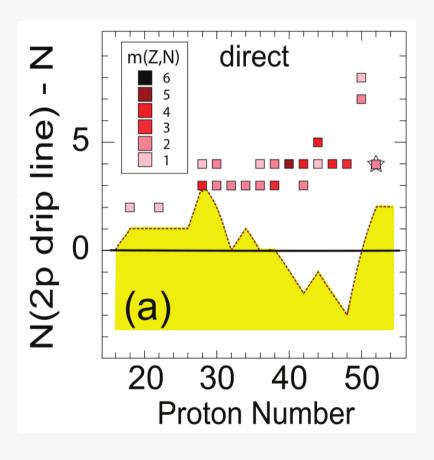
> Predictions of the direct model

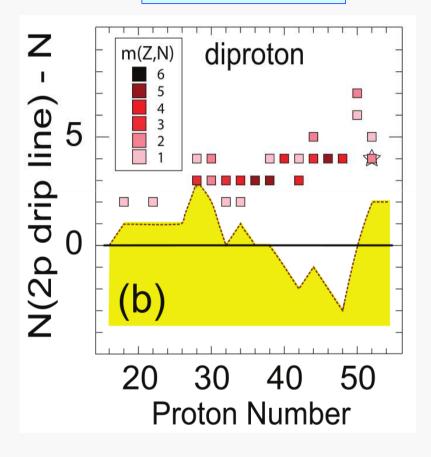


### True 2p landscape

Predicted candidates relative to the 2p dripline

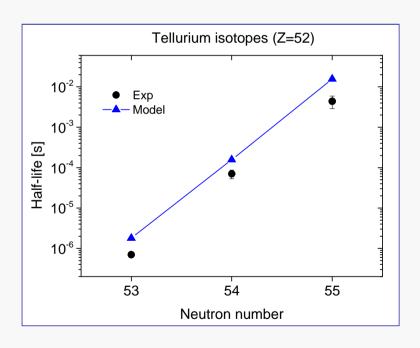
 $100 \text{ ns} < T_{2p} < 100 \text{ ms}$ 

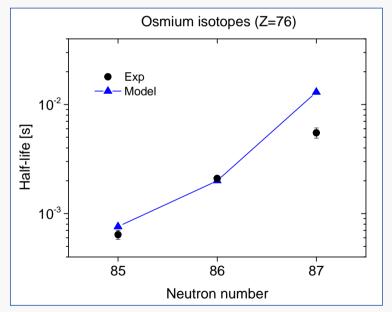




### $\alpha$ -emission

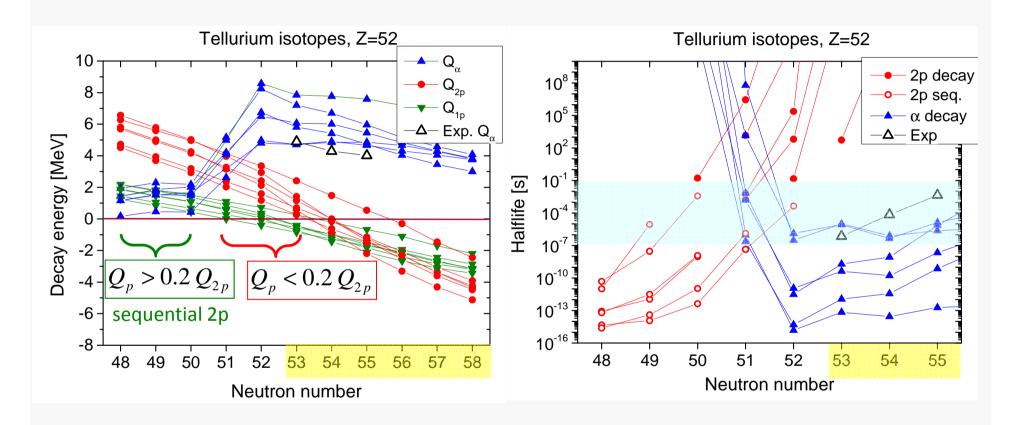
 $\blacktriangleright$  Global, fenomenological formula for  $\alpha$  decay half-lives: H. Koura 2012





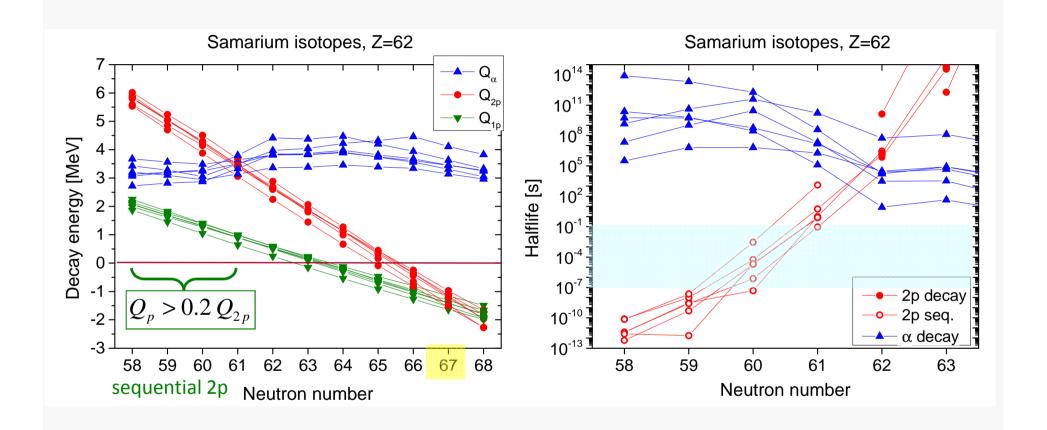
Koura, J. Nucl. Science and Tech. 49 (2012) 816

### **Tellurium**

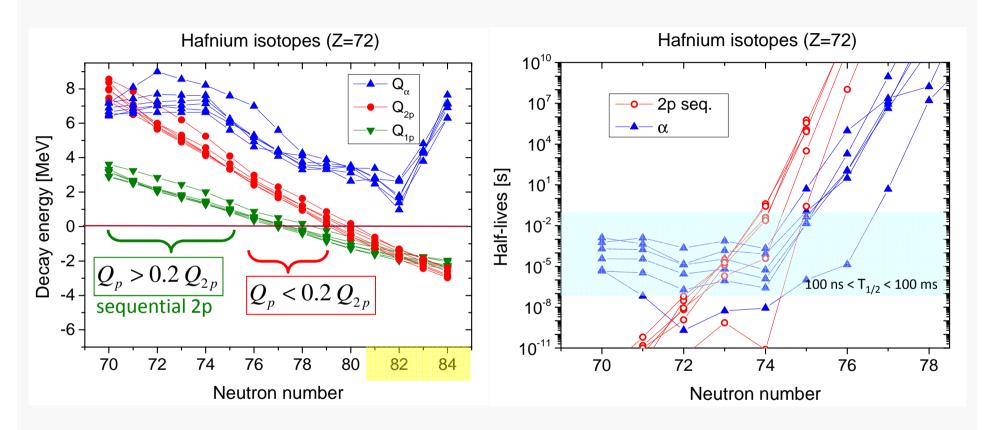


- ➤ At <sup>103</sup>Te a transition from the simultaneous 2p to the sequential emission occurs
- $\blacktriangleright$  In addition, in <sup>103</sup>Te both decays,  $\alpha$  and 2p may be observable!

### Samarium

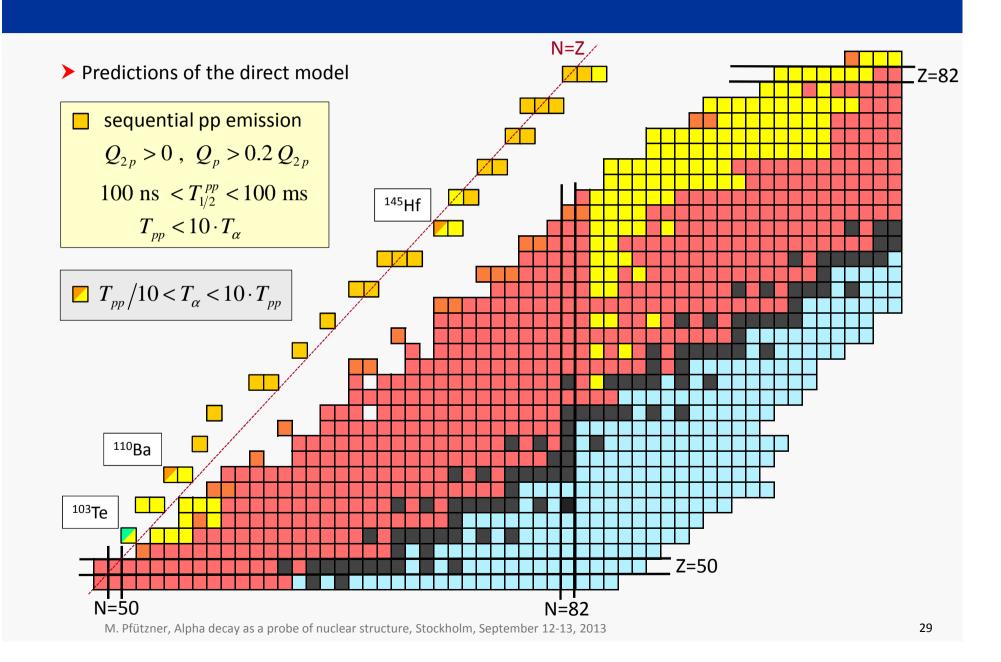


### Hafnium

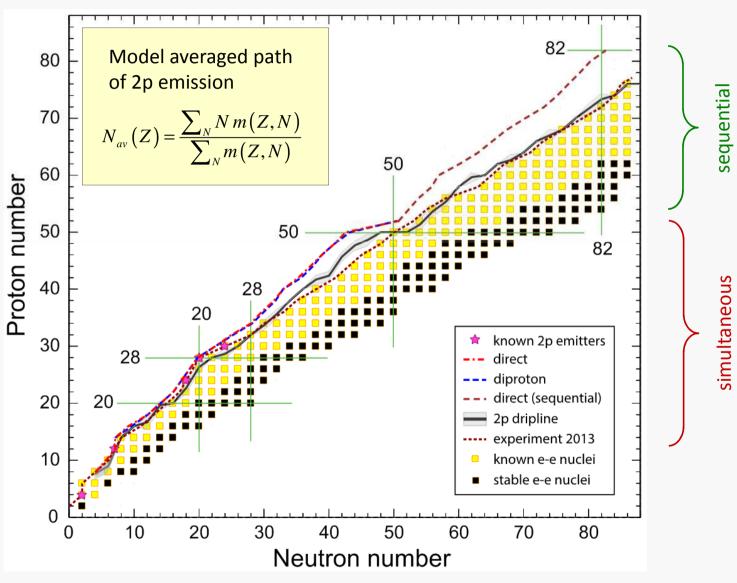


- When the energy condition for the true 2p decay is fulfilled, the predicted half-life is extremely long
- > When the fast proton emission becomes possible, it proceeds as the sequential 2p decay

### Between tellurium and lead



## Full 2p landscape



### Summary

- The direct (simultaneous) ground-state 2p emission established for <sup>6</sup>Be, <sup>19</sup>Mg, <sup>45</sup>Fe, <sup>48</sup>Ni, and <sup>54</sup>Zn.
   The hunt for other cases continues: <sup>30</sup>Ar, <sup>59</sup>Ge,....
- For every even-Z element between zinc and tellurium (Z=52) the isotopes decaying by 2p radioactivity in the time window  $100 \text{ ns} < T_{1/2} < 100 \text{ ms}$  are predicted.
- In  $^{103}$ Te the competition between simultaneous 2p, sequential pp, and  $\alpha$  emission may occur. For  $^{145}$ Hf the competition between  $\alpha$  and sequential pp is predicted.
- Above tellurium the limit of decay spectroscopy is represented by sequential pp emission, except for xenon (Z=54) where  $\alpha$  decay dominates.
- Above lead (Z=82)  $\alpha$  decay dominates, no 2p emission is expected to be observed.

# Thank you!

