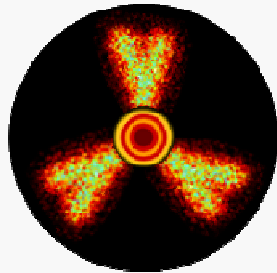


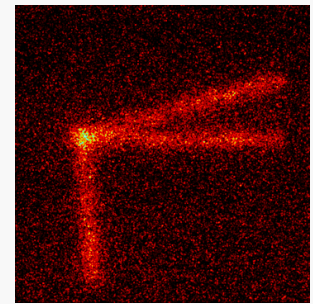
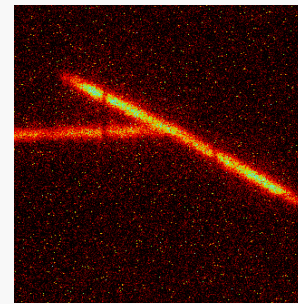
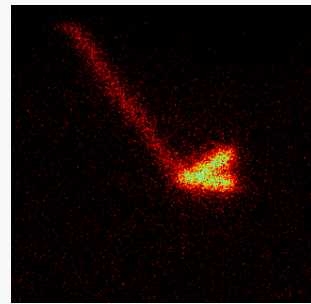
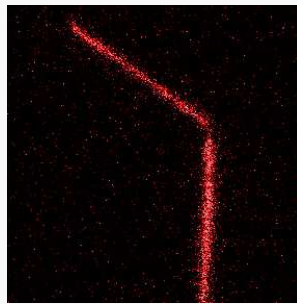
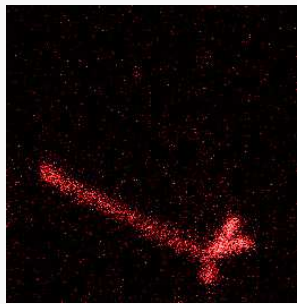
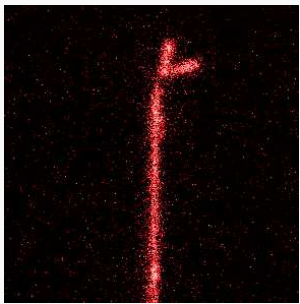
# Particle radioactivity of exotic nuclei



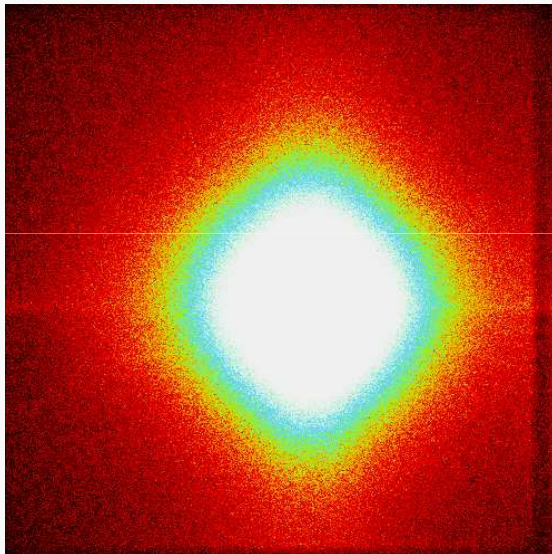
Marek Pfützner

Faculty of Physics, University of Warsaw

GSI Helmholtzentrums

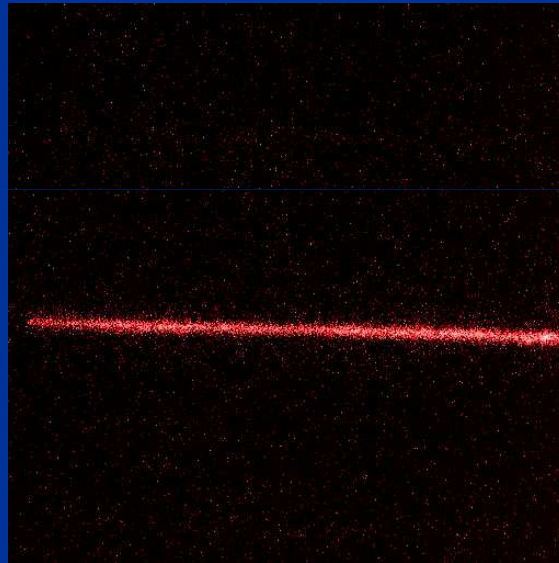


# Outline



- Basic concepts
- In-flight at Coulomb barrier
  - ◊ Proton radioactivity
  - ◊ Alpha emission
- In-flight above Fermi energy
  - ◊ Two-proton radioactivity
  - ◊ Neutron radioactivity?

# Basic concepts



# Gamow picture

➤ Early puzzle of  $\alpha$  decay: half-life is extremely dependent of on energy

**Solution by Gamow (1928) was a triumph of quantum mechanics!**

Geiger-Nuttall (1912)

$$\log T = a + \frac{b}{\sqrt{E_\alpha}}$$

➔ Radioactive process is so slow (decaying state so narrow) that it can be approximated by a **stationary** picture:

$$\hat{H} \Phi(\vec{r}, t) = i\hbar \frac{\partial \Phi(\vec{r}, t)}{\partial t}, \quad \Phi(\vec{r}, t) = \Psi(\vec{r}) e^{-iEt/\hbar}$$

but with the complex energy!

$$\hat{H} \Psi(\vec{r}) = E \Psi(\vec{r}) \quad E = E_0 - i\Gamma/2$$

$$|\Phi(\vec{r}, t)|^2 = |\Psi(\vec{r})|^2 e^{-i\Gamma t/\hbar}$$

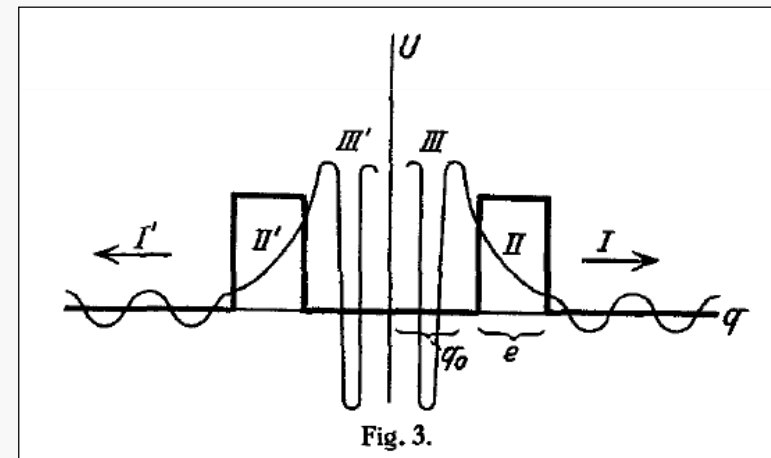


Fig. 3.

Um diese Schwierigkeit zu überwinden, müssen wir annehmen, daß die Schwingungen gedämpft sind, und  $E$  komplex setzen:

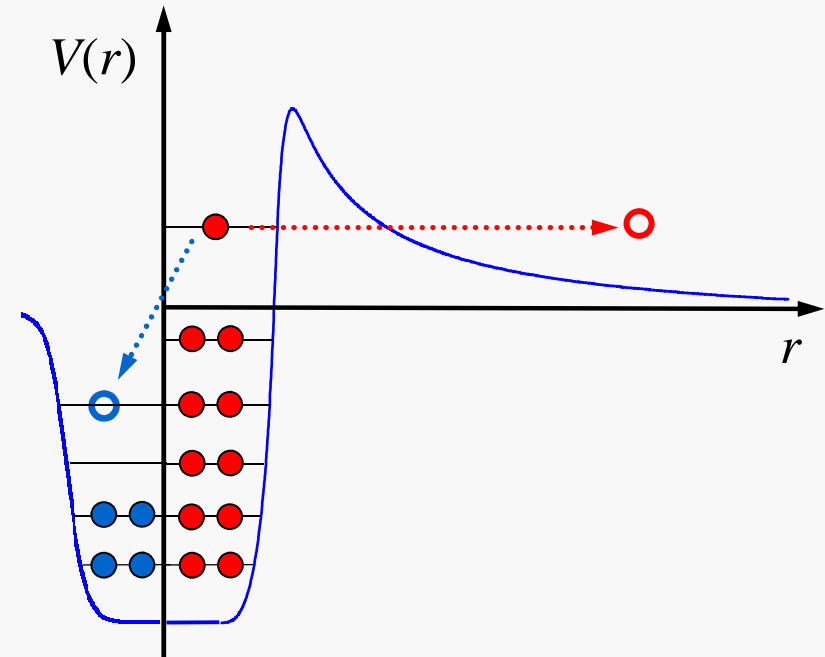
$$E = E_0 + i \frac{\hbar \lambda}{4\pi},$$

Gamow, Z. Phys. 51 (1928) 204

# Particle radioactivity

- The fundamental concept: **potential barrier**
- ➔ The (Coulomb) barrier stops an unbound object ( $\alpha$ , p, 2p,  $^{14}\text{C}$ ,...) from flying out immediately.
- ➔ Neutrons can still be hampered by the centrifugal barrier.
- ➔ Beyond proton drip-line, there is always competition with  $\beta$  decay!

Particle observable if  $T_{1/2}^p \leq T_{1/2}^\beta$



- In spherical case, WKB-like method:

$$\Gamma/\hbar = S \nu \exp \left\{ -\frac{2}{\hbar} \cdot \int_{r_{in}}^{r_{out}} \sqrt{2\mu[V(r) - Q_p]} \cdot dr \right\}$$

$\nu$  – frequency of assaults  
 $S$  – spectroscopic factor (p)  
 preformation factor ( $\alpha$ )

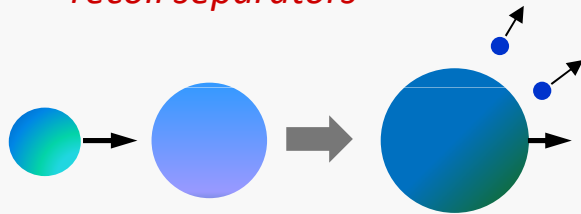
- ➔ This simple approach works surprisingly good and is still frequently used in the analysis of proton and  $\alpha$  radioactivity.

Gurvitz and Kalbermann, PRL 59 (1987) 262

# Production methods

➤ To produce short-lived and very proton-rich radioactive nuclei in-flight techniques proved advantageous.

- Fusion-evaporation  
reactions between heavy-ions  
*GSI, Argonne, Oak Ridge, Jyväskylä,...*  
*recoil separators*



Low energy:  $\approx$  Coulomb barrier

- large beam intensity
- thin target  
identification by decays

*p* and  $\alpha$  radioactivity

- Fragmentation  
of relativistic heavy-ions  
*GSI, NSCL, GANIL, RIKEN,...*  
*fragment separators*

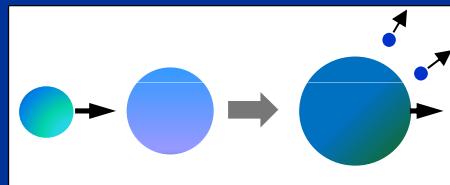


High energy:  $\approx$  above Fermi energy

- lower beam intensity
- thick target  
identification in-flight  
single ion sensitivity

*2p* radioactivity

# In-flight technique at Coulomb barrier

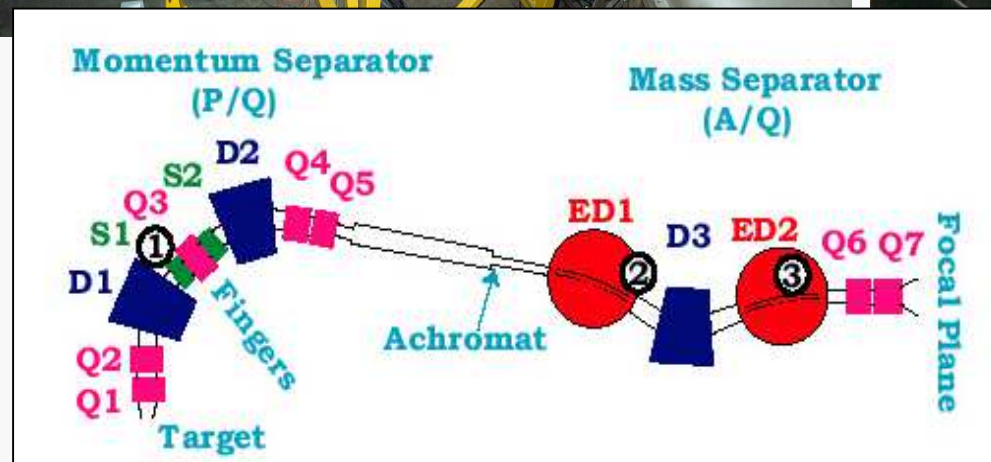


# Recoil separators

Recoil Mass Separator @ ORNL

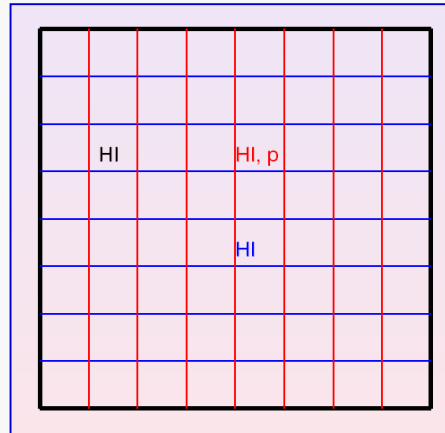
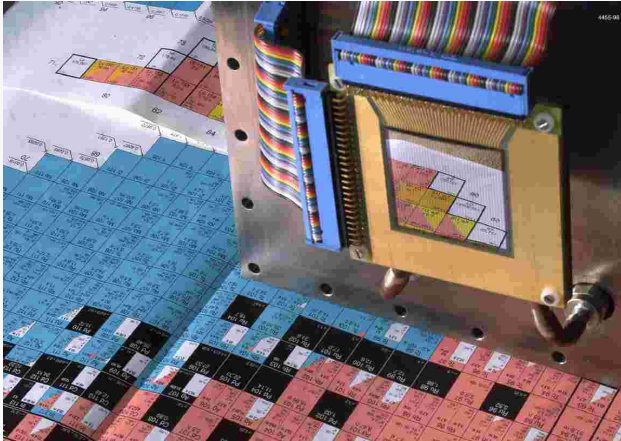


Fragment Mass Analyser @ ANL



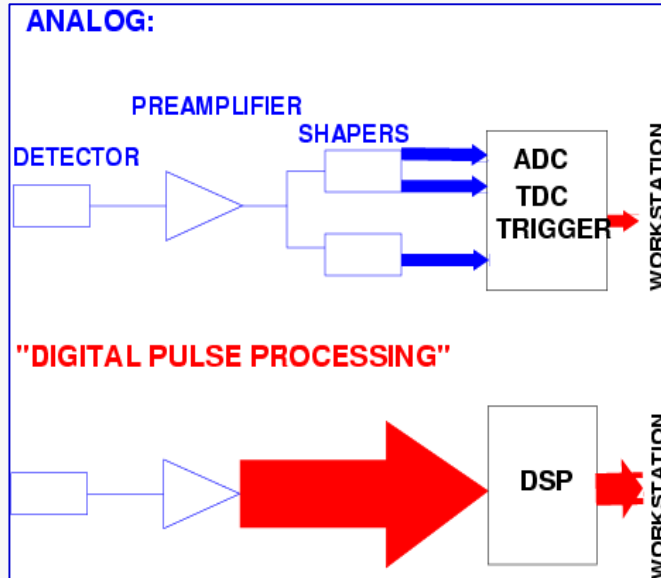


# DSSSD and Digital

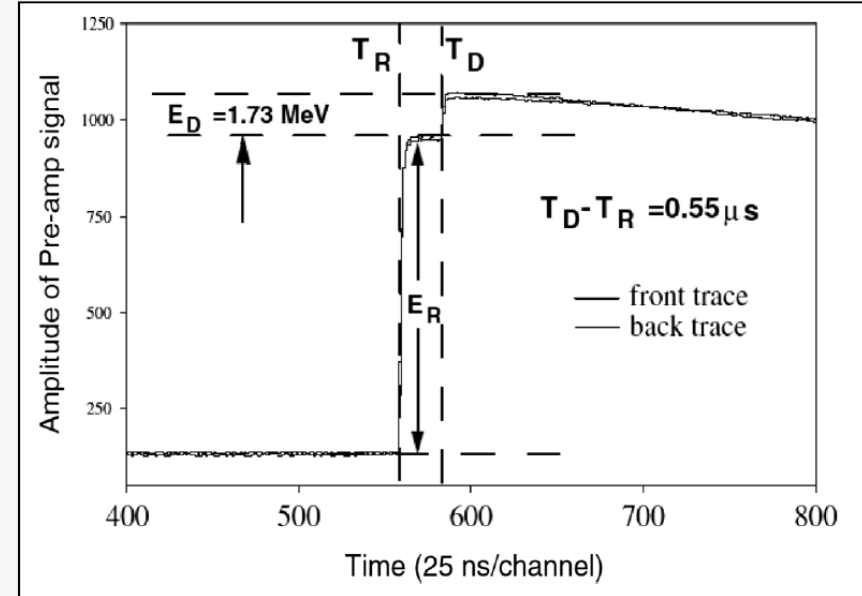


DSSSD – Double Sided Silicon Strip Detector

- Correlations between ion-implantation and its decay in space and in time

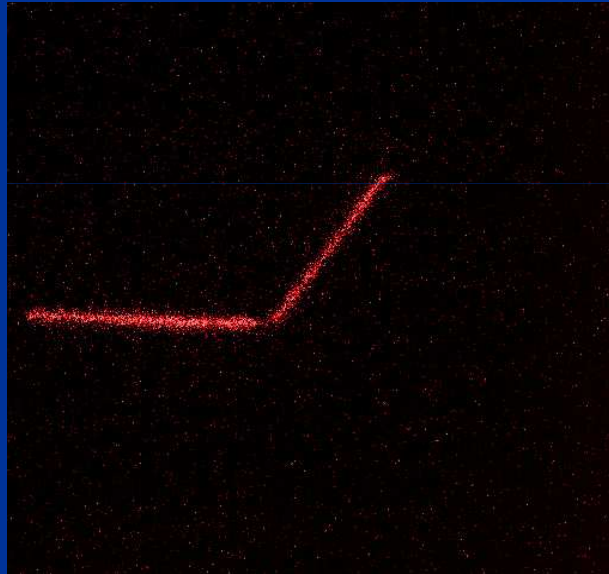


Grzywacz, NIM B 234(2007)

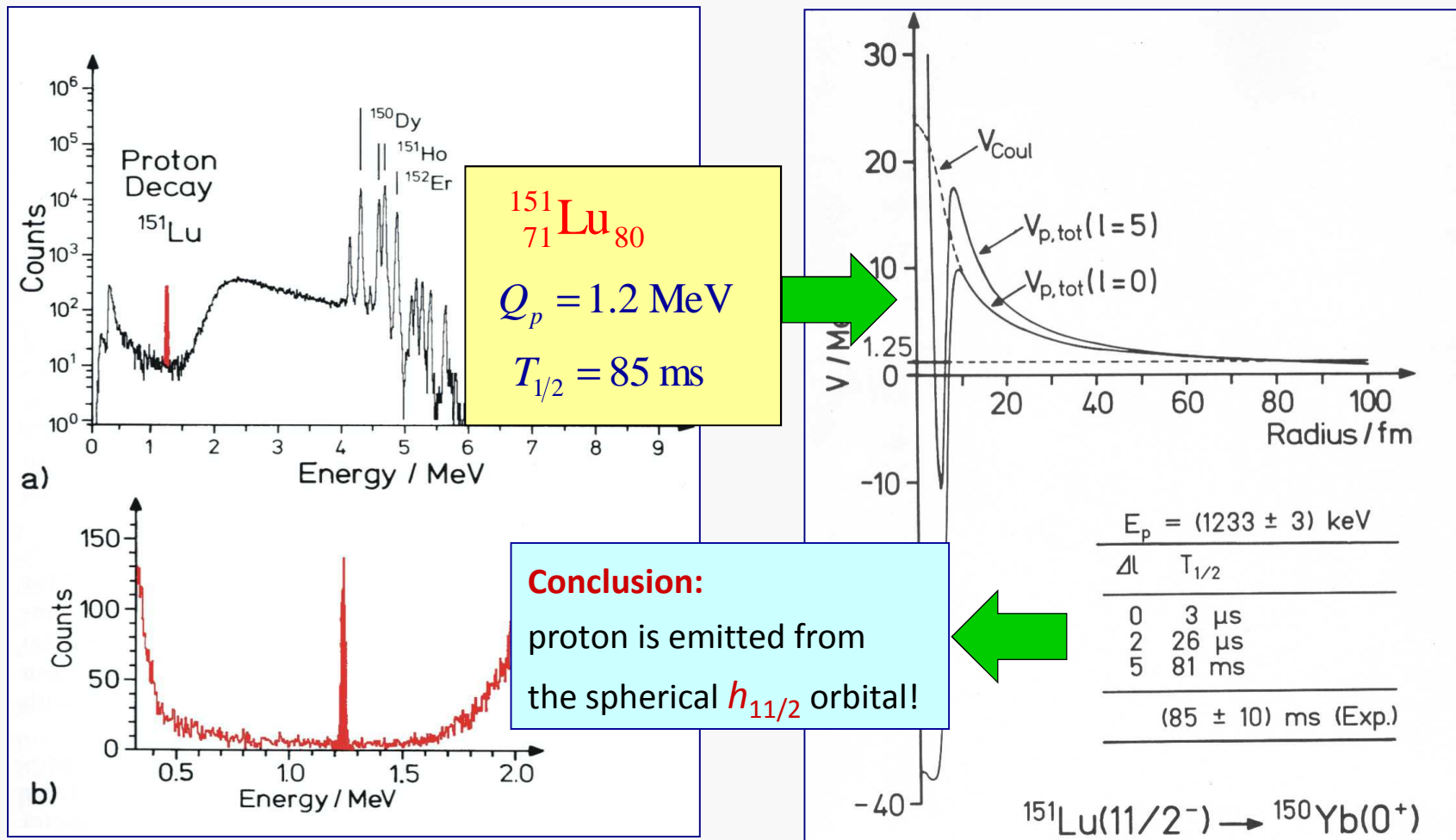


Karny *et al.*, Phys. Rev. Lett. 90 (2003) 012502

# Proton radioactivity



# $p$ radioactivity – the first case



Hofmann et al., Z. Phys. A 305 (1982) 111

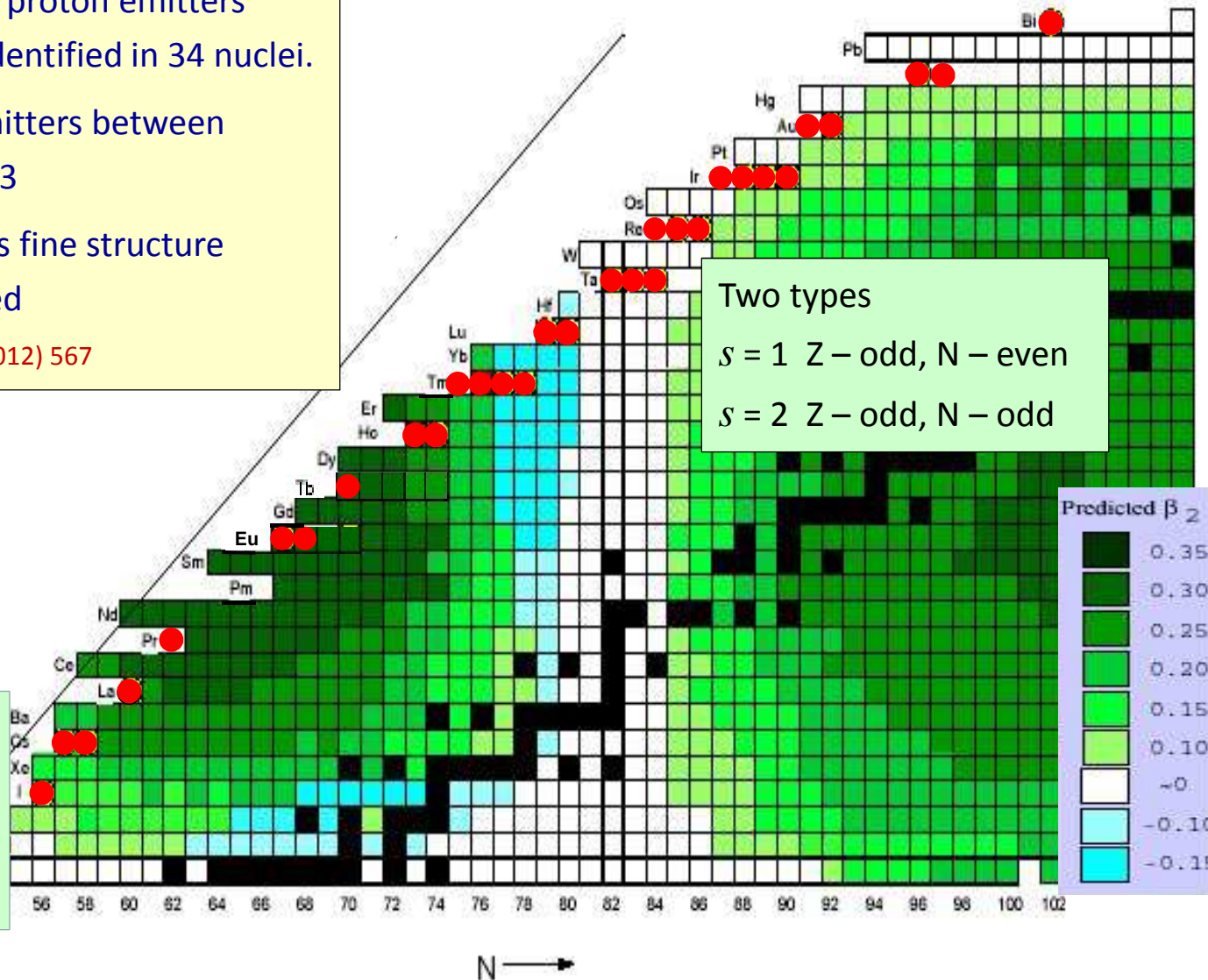
# $p$ radioactivity – the status

- Presently 46 proton emitters (g.s. or m) identified in 34 nuclei.
- All g.s.  $p$ -emitters between  $Z = 50$  and  $83$
- In 7 emitters fine structure was observed

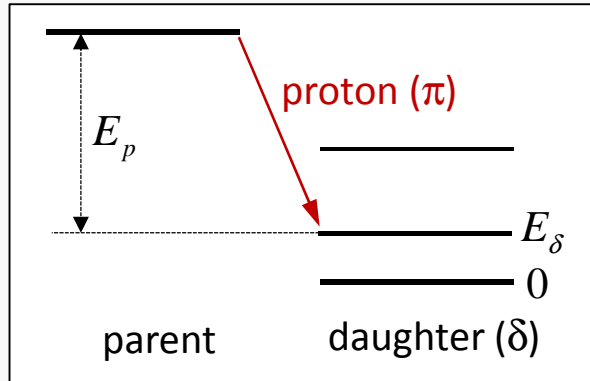
M.P. et al, RMP (2012) 567

Z ↑

$s > 2$   
 $^{53}\text{mCo}$   
 $^{54}\text{mNi}$   
 $^{94}\text{mAg}$



# More general model



- Cluster approximation (proton moves in the (deformed) potential of the daughter)
- Parent wave function:

$$\Psi_{JM}(\vec{r}) = \frac{1}{r} \sum_{\alpha=\{\pi,\delta\}} u_\alpha(r) [\mathcal{Y}_\pi \otimes \Phi_\delta]_{JM}$$

radial part of relative motion      angular part

→ Schrödinger equation, integration over angles:

$$\left[ \frac{d^2}{dr^2} - \frac{l_p(l_p+1)}{r^2} + \frac{2\mu}{\hbar^2} E_p \right] u_\alpha(r) = \frac{2\mu}{\hbar^2} \sum_{\alpha'} (\hat{V}_{\alpha,\alpha'}) u_{\alpha'}(r) \quad \text{(coupled channels)}$$

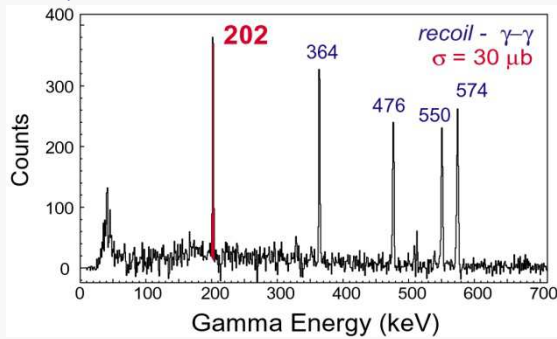
- Solutions must fulfill conditions:

$$u_\alpha(r) \xrightarrow{r \rightarrow 0} 0, \quad u_\alpha(r) \xrightarrow{r \rightarrow \infty} \sim G_{l_p}(\eta, kr) + i F_{l_p}(\eta, kr) \quad \text{(outgoing Coulomb wave)}$$

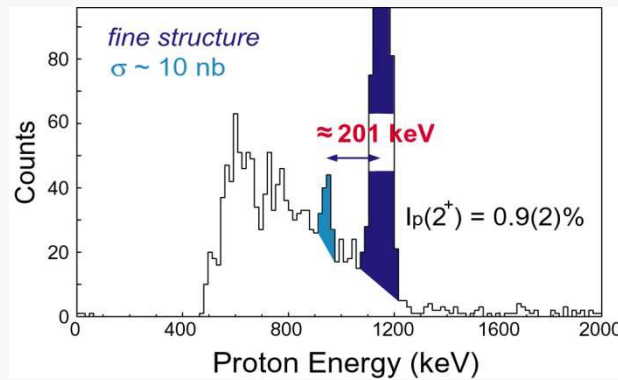
→ From radial functions the width:  $u_\alpha(r) \Rightarrow \Gamma_\alpha$ ,  $\Gamma = \sum_\alpha \Gamma_\alpha$ ,  $\Gamma_\delta = \sum_{\{\pi\}} \Gamma_{\{\pi,\delta\}}$

# Proton emission from deformed $^{141}\text{Ho}$

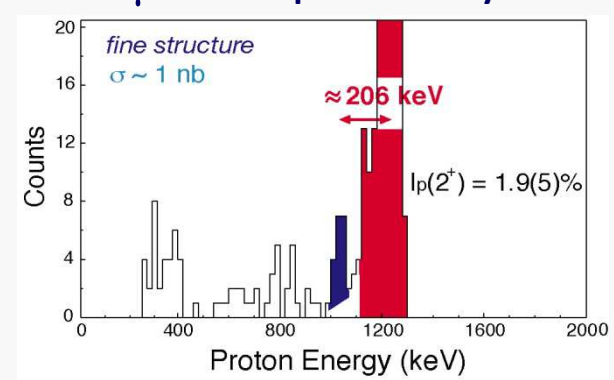
7  $\mu\text{s}$   $^{140\text{m}}\text{Dy}$  isomeric decay



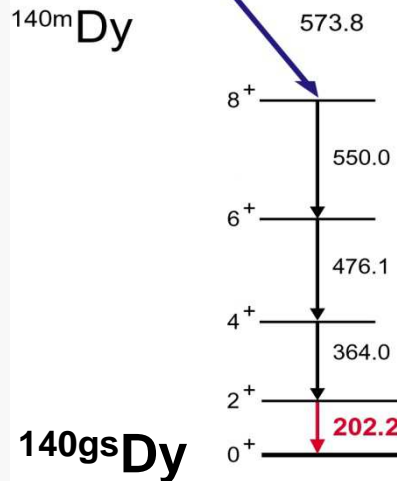
4 ms  $^{141\text{gs}}\text{Ho}$  proton decay



7.2  $\mu\text{s}$   $^{141\text{m}}\text{Ho}$  proton decay



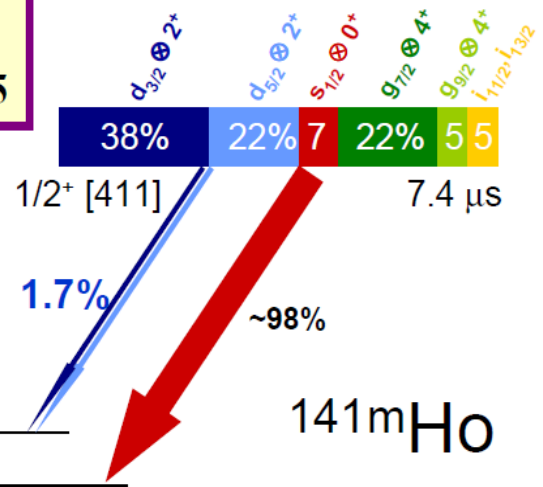
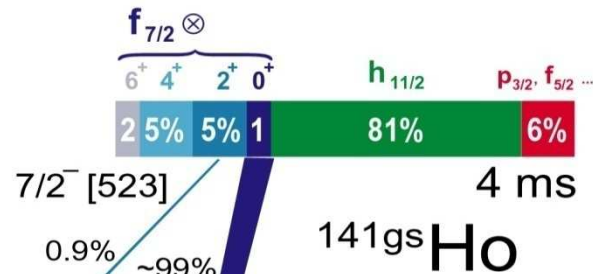
8<sup>-</sup> K-isomer 7  $\mu\text{s}$   
 $\nu 7/2^+ [404] \otimes \nu 9/2^- [514]$



Coupled-channels approach  
Kruppa et al., PRL 84 (2000) 4549

$$\beta_2 = 0.35$$

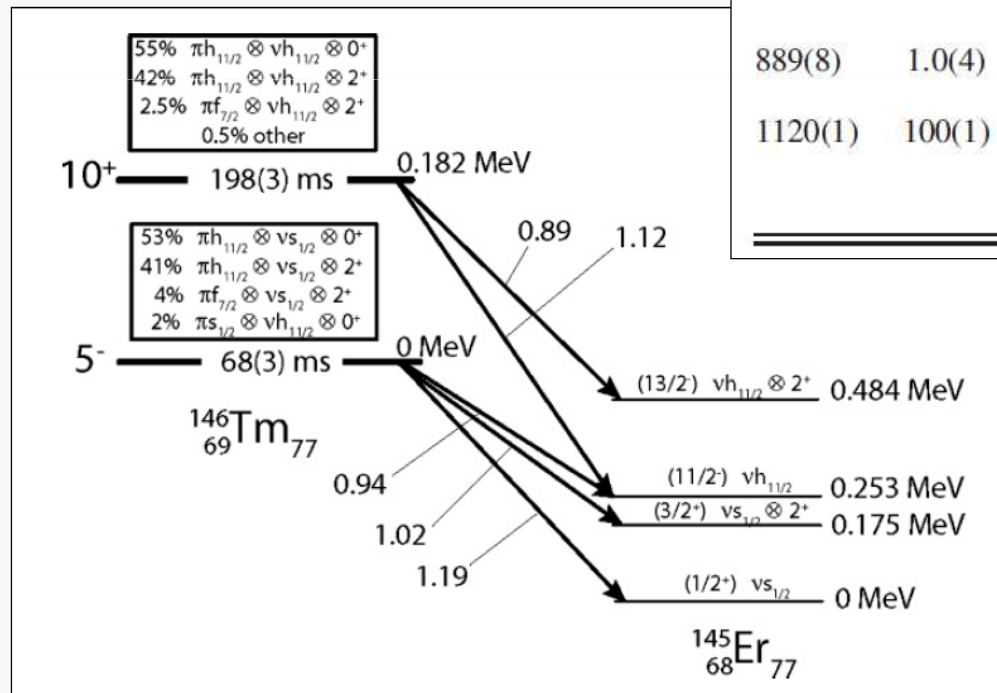
$$\beta_4 = -0.05$$



M. Karny et al., Phys. Lett. B664 (2008) 52

# Odd-odd example: $^{146}\text{Tm}$

- The richest emitter known:  
5 proton lines observed!
- ➔ In Z – odd, N – odd (s=2)  
nucleus the proton radioactivity  
provides data on neutron levels!

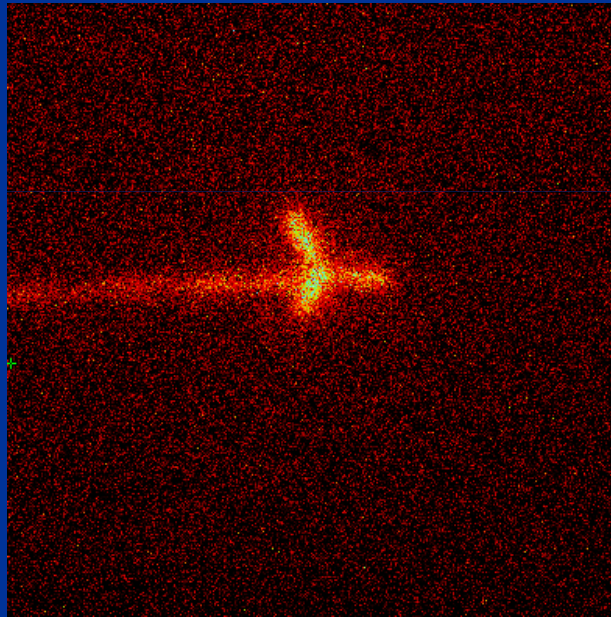


$E_p$	$I_p^{\text{exp}}$ (%)	Wave function composition	$I_p^{\text{cal}}$	$\Delta l$	$E_f$
Ground state					
$I^\pi = 5^-, T_{1/2} = 68(5) \text{ ms}$					
938(4)	13.8(9)	2% $\pi s_{1/2} \otimes \nu h_{11/2} \otimes 0^+$	(15) <sup>a</sup>	0	253
1016(4)	18.3(11)	4% $\pi f_{7/2} \otimes \nu s_{1/2} \otimes 2^+$	15	3	175
		41% $\pi h_{11/2} \otimes \nu s_{1/2} \otimes 2^+$	0.003	5	175
1191(1)	68.1(19)	53% $\pi h_{11/2} \otimes \nu s_{1/2} \otimes 0^+$	70	5	0
Isomeric state					
$I^\pi = 10^+, T_{1/2} = 198(3) \text{ ms}$					
889(8)	1.0(4)	2.5% $\pi f_{7/2} \otimes \nu h_{11/2} \otimes 2^+$	1.2	3	484
		42% $\pi h_{11/2} \otimes \nu h_{11/2} \otimes 2^+$	0.04	5	484
1120(1)	100(1)	55% $\pi h_{11/2} \otimes \nu h_{11/2} \otimes 0^+$	98.6	5	253
		0.1% $\pi h_{9/2} \otimes \nu h_{11/2} \otimes 0^+$	0.2	5	253
		0.4% $\pi(l > 5) \otimes \nu h_{11/2}$			

- Good agreement with the model assuming coupling of a particle to core vibrations (the c-c scheme)

Tantawy *et al.*, Phys. Rev. C 73 (2006) 024316  
Hagino, Phys. Rev. C 64 (2001) 041304(R)

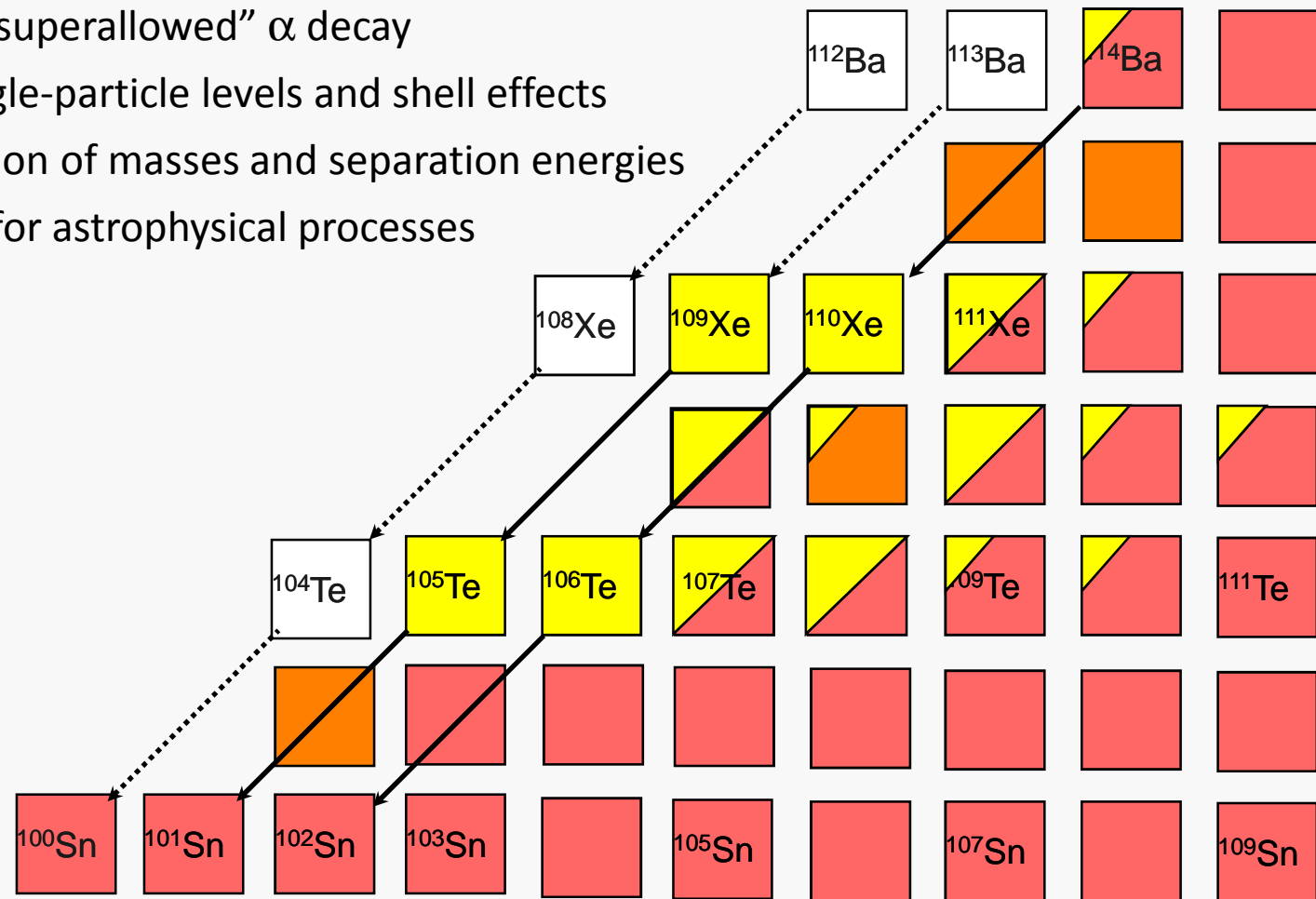
# Alpha radioactivity





# Island of $\alpha$ emitters above $^{100}\text{Sn}$

- Search for „superallowed”  $\alpha$  decay
- Probing single-particle levels and shell effects
- Determination of masses and separation energies
- Conditions for astrophysical processes

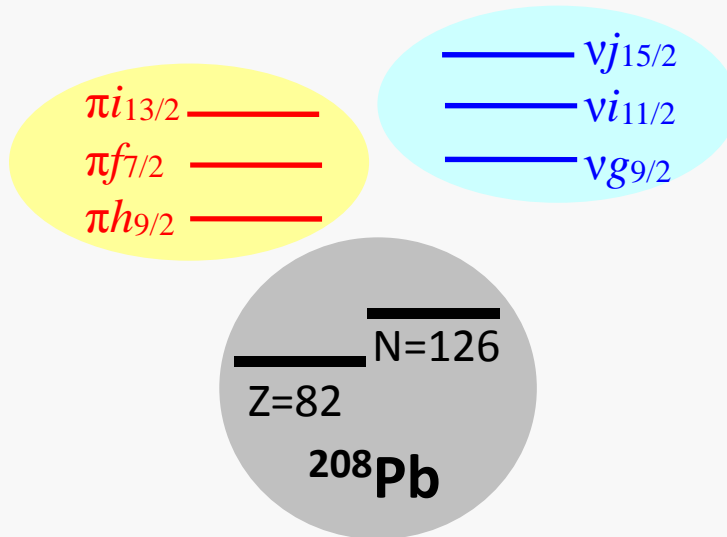


# Superaligned $\alpha$ decay?

- Present  $\alpha$ -decay reference:  $^{212}\text{Po}$



$\alpha$  made of protons and neutrons  
from different orbitals of opposite parity

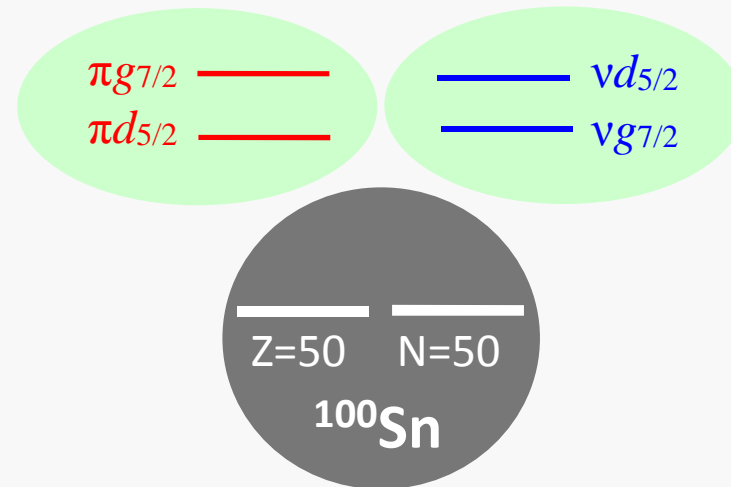


- Expected standard:  $^{104}\text{Te}$

Macfarlane and Siivola, PRL 14 (1965) 114



$\alpha$  formed by protons and neutrons  
in the same orbitals



Predictions for  $^{104}\text{Te} \rightarrow ^{100}\text{Sn}$  decay

$$E_{\alpha} > 5 \text{ MeV}, T_{1/2} < 50 \text{ ns !}$$

# $\alpha$ decay of $^{105}\text{Te}$

► Decay of  $^{105}\text{Te}$  studied:

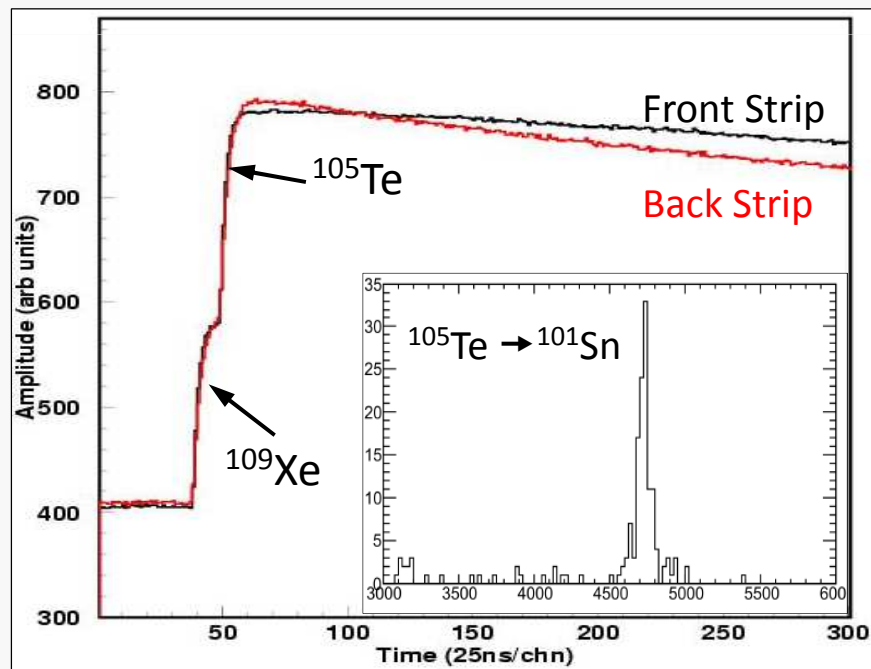
- directly at FMA (Argonne) using  $^{50}\text{Cr}(^{58}\text{Ni}, 3n)^{105}\text{Te}$  and fast recovery electronics
- via decay of  $^{109}\text{Xe}$  at HRIBF (ORNL) by  $^{54}\text{Fe}(^{58}\text{Ni}, 3n)^{109}\text{Xe}$  and DSP

→  $^{105}\text{Te}$  decay:  $E_\alpha = 4.7 \text{ MeV}$ ,  $T_{1/2} = 0.6 \mu\text{s}$

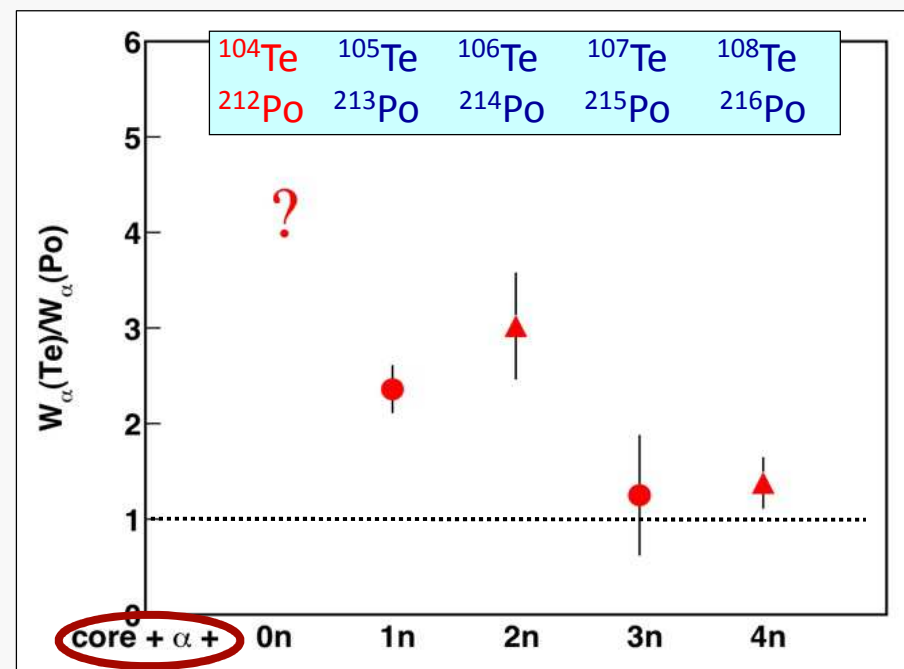
Seweryniak et al., PRC 73(2006) 061301(R)

S.N. Liddick et al., PRL 97 (2006) 082501

I.G. Darby et al. PRL 105 (2010) 162502

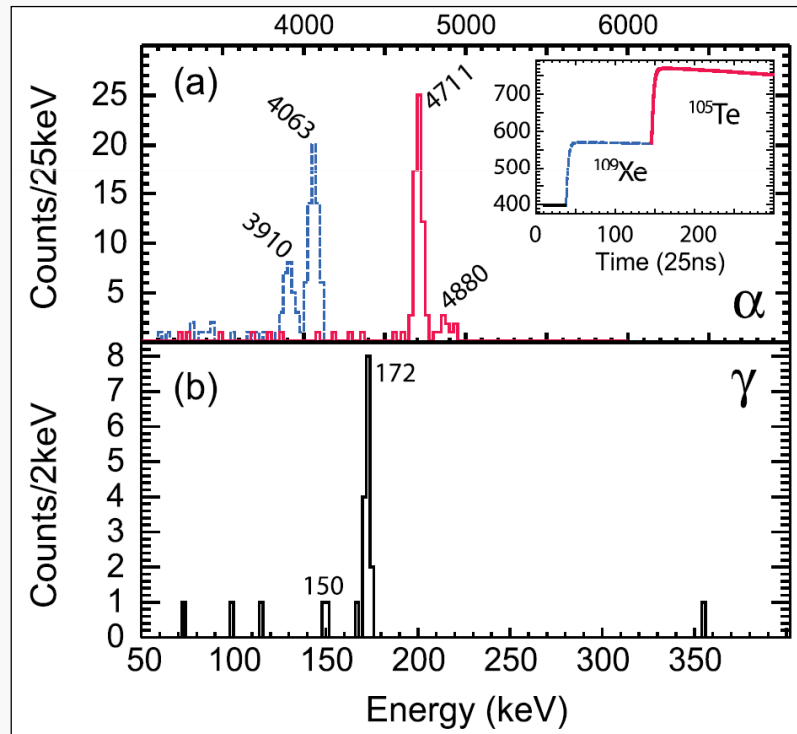


## Renormalized $\alpha$ decay width ( $l = 0$ transitions)

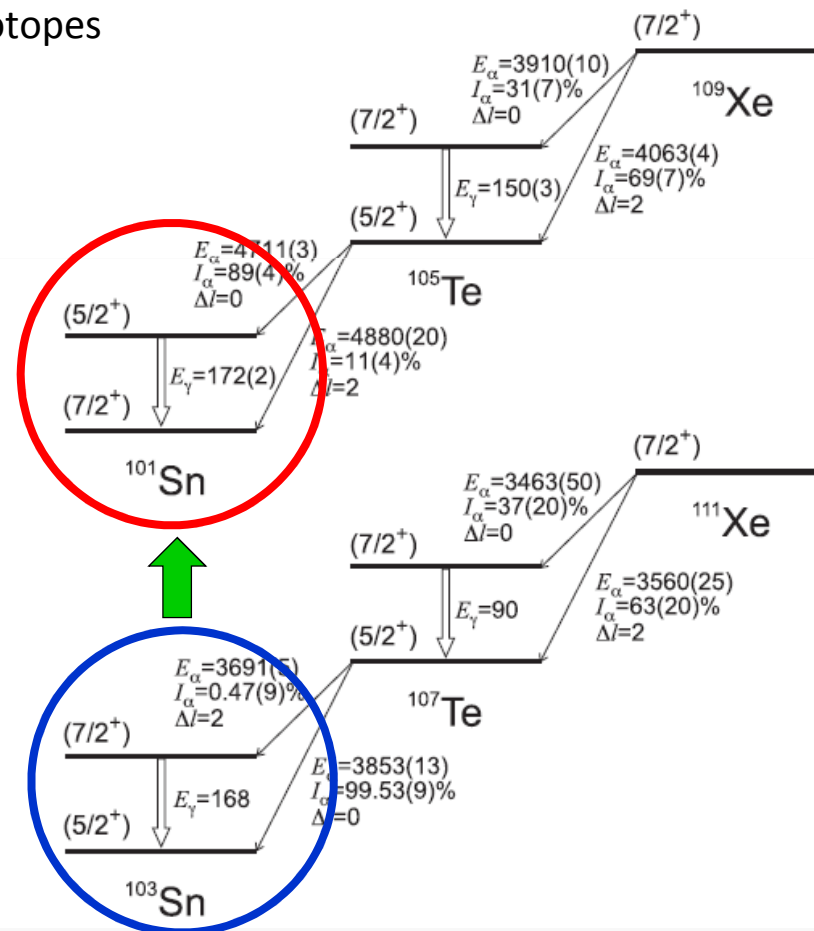


# Single particle states in $^{101}\text{Sn}$

- Details of  $\alpha$  decay of  $^{109}\text{Xe}$  (fine structure) yield surprising result on  $^{101}\text{Sn}$ !  
 $5/2^+$  and  $7/2^+$  levels are reversed between  $^{103}\text{Sn}$  and  $^{101}\text{Sn}$
- ➔ Orbital dependent pairing, stronger for  $(g_{7/2})^2$  than for  $(d_{5/2})^2$ , is responsible for  $5/2^+$  g.s of  $^{103}\text{Sn}$  and heavier odd tin isotopes



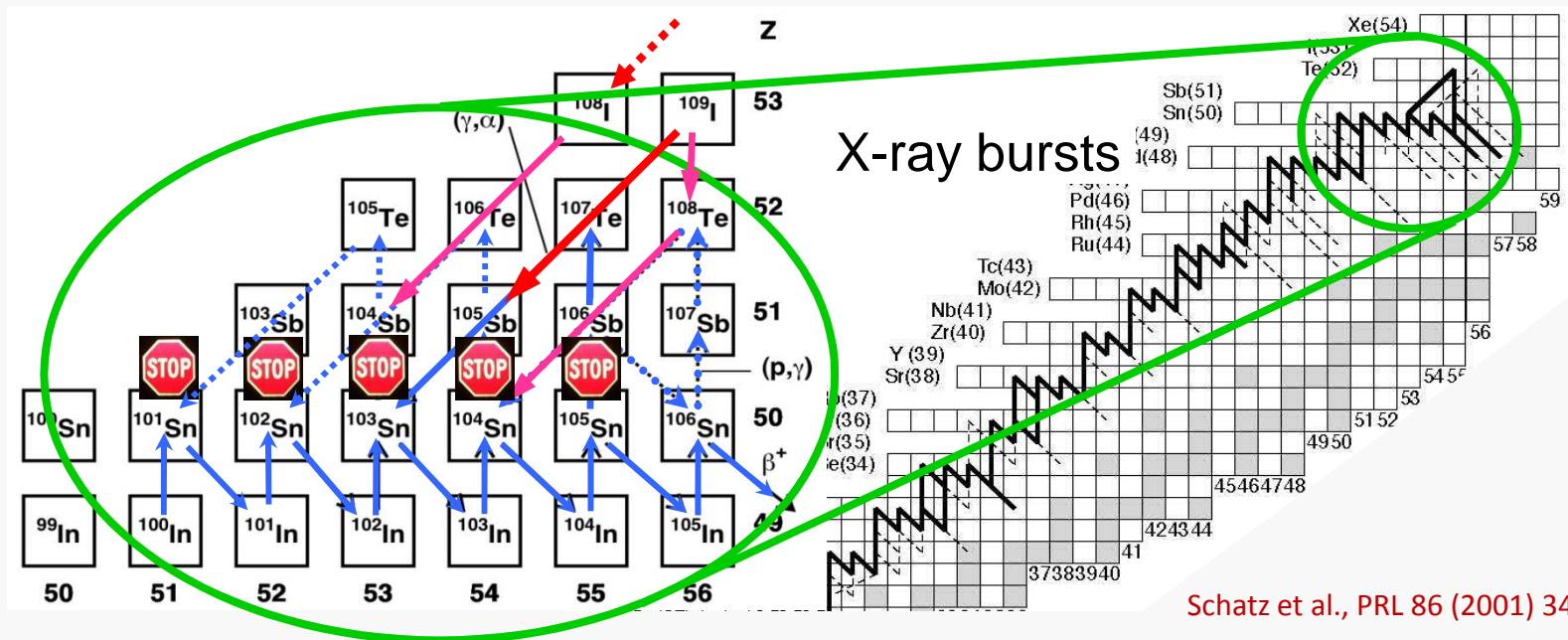
I.G. Darby et al. PRL 105 (2010) 162502



# Termination of $rp$ -process

- In some conditions  $rp$ -process could terminate by Sn-Sb-Te cycle. The details of this cycle depend critically on  $Q_p$  values of Sb-isotopes.
  - $\alpha$ -decay of  $^{109}\text{I}$  yielded the  $Q_p = 356 \text{ keV}$  for  $^{105}\text{Sb}$  solving an old controversy.
  - Similarly, the search for  $\alpha$ -decay of  $^{112}\text{Cs}$  produced a limit  $Q_p > 150 \text{ keV}$  for  $^{104}\text{Sb}$ .
  - The  $Q_p$  values of  $^{106}\text{Sb}$  and heavier isotopes were measured in JYFLTRAP.
- ➔ The  $rp$ -process most likely dies out before reaching the Sn-Sb-Te cycle.

Mazzocchi et al., PRL 98 (2007) 212501    Cartegni et al., PRC 85 (2012) 014312    Elomaa et al., PRL 102 (2009) 252501

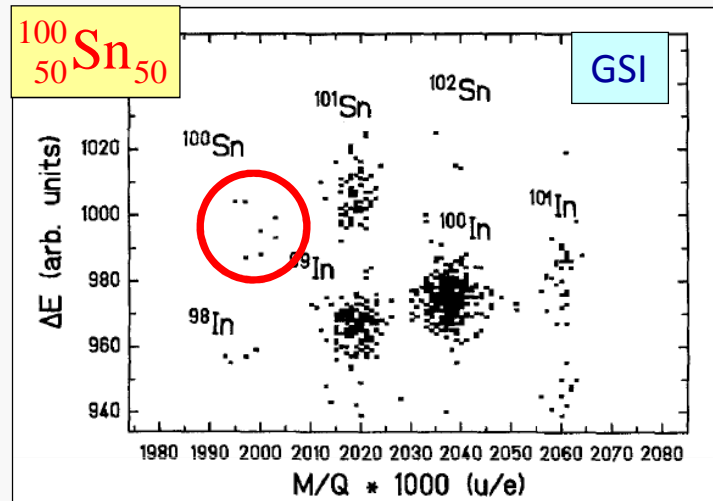


Schatz et al., PRL 86 (2001) 3471

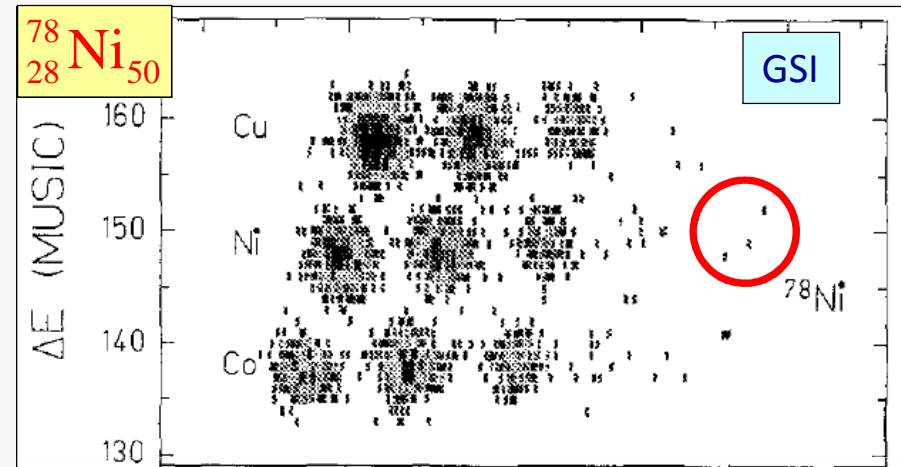
# In-flight technique above Fermi energy



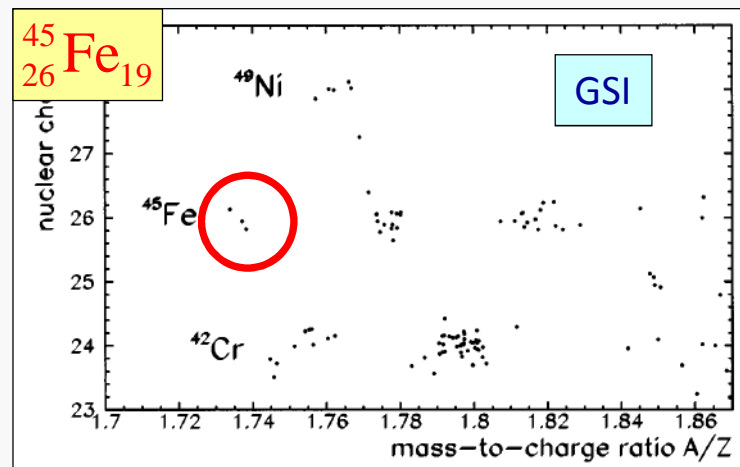
# Fragmentation milestones



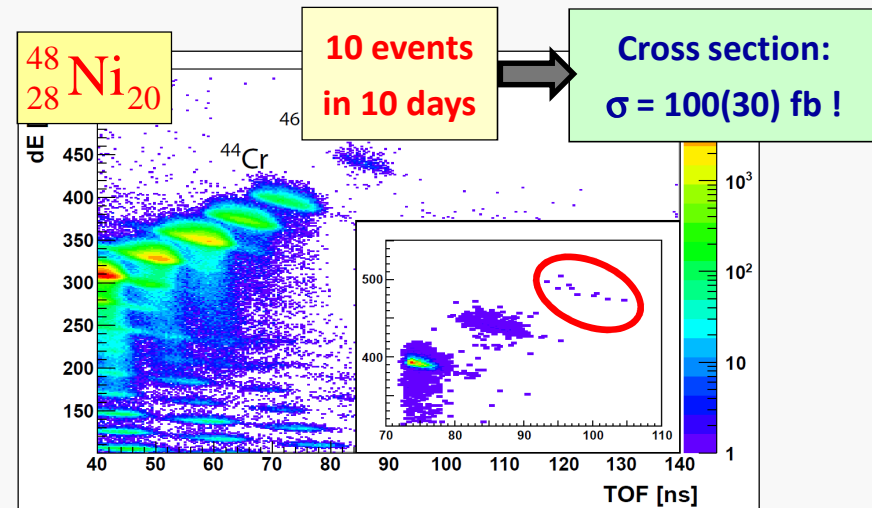
Schneider et al., Z. Phys. A 348 (1994) 241



Engelmann et al., Z. Phys. A352 (1995) 351

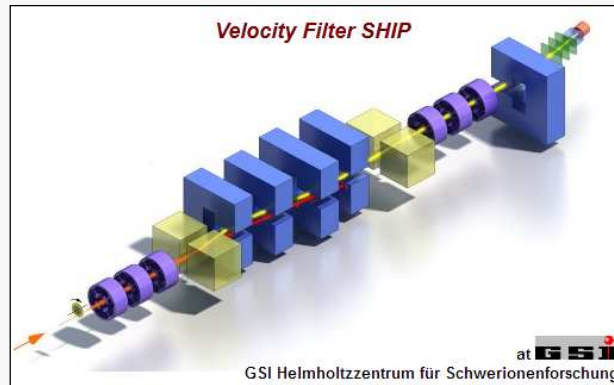


Blank et al., PRL 77 (1996) 2893



Pomorski et al., PRC 83 (2011) 061303(R)

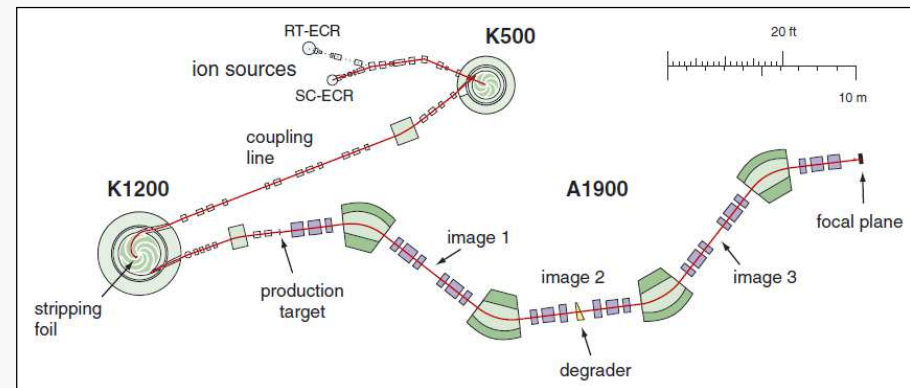
# Sensitivity



- Hunt for element 120 at SHIP



Cross section:	$\approx 90 \text{ fb}$
Beam intensity:	750 pA
Running time (1 event):	$\approx 100 \text{ days}$
Total dose:	$\approx 4 \cdot 10^{19} \text{ part.}$
Target thickness:	$0.5 \text{ mg/cm}^2$



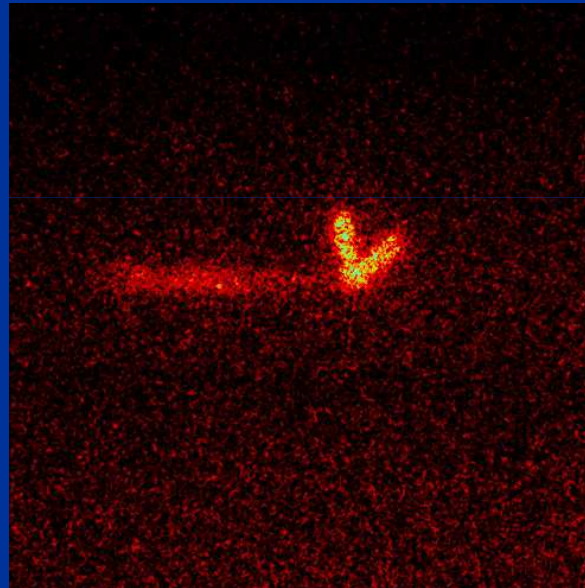
- Production of  $^{48}\text{Ni}$  at NSCL



$\approx 100 \text{ fb}$	38 times weaker
20 pA	100 times shorter
$\approx 1 \text{ day}$	<b>4000 times smaller</b>
$\approx 1 \cdot 10^{16} \text{ part.}$	
$580 \text{ mg/cm}^2$	

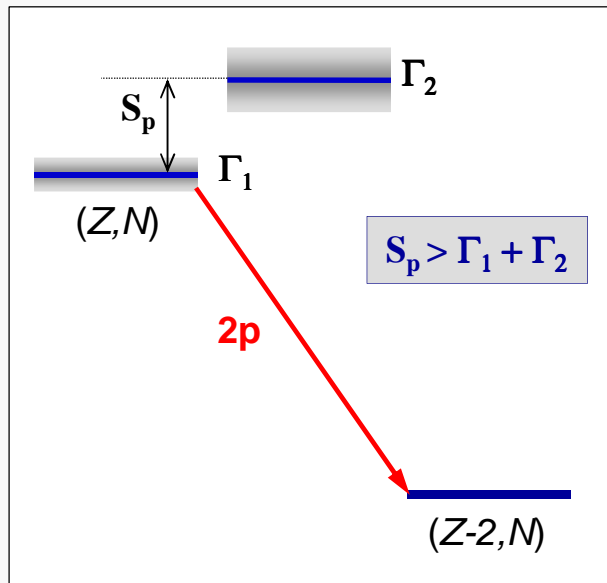


# Two-proton radioactivity



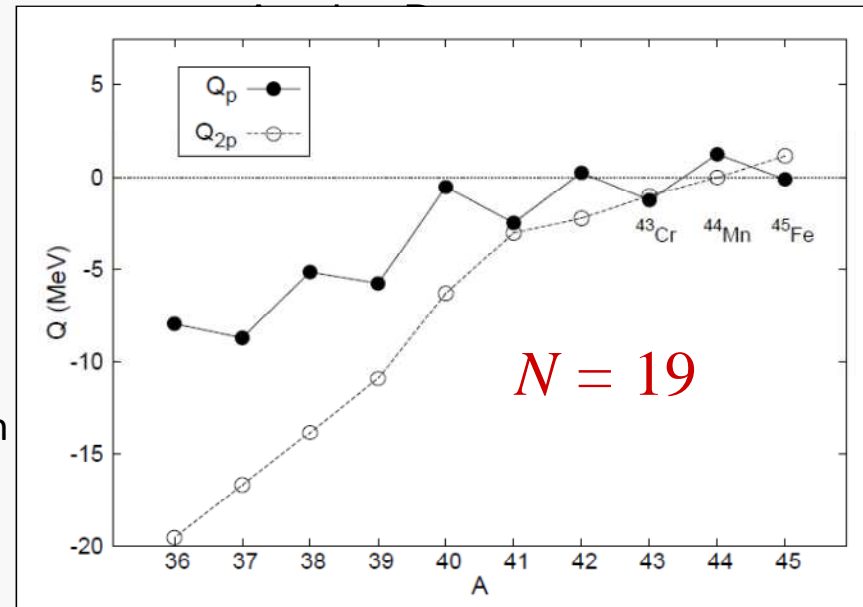
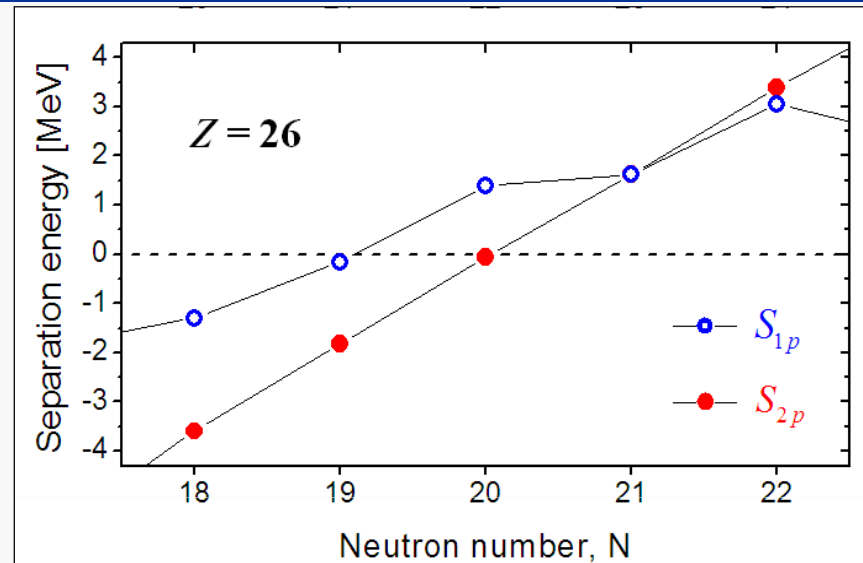
# Two protons can be unbound!

- It is possible that pair of protons is unbound while each of individual proton is bound!



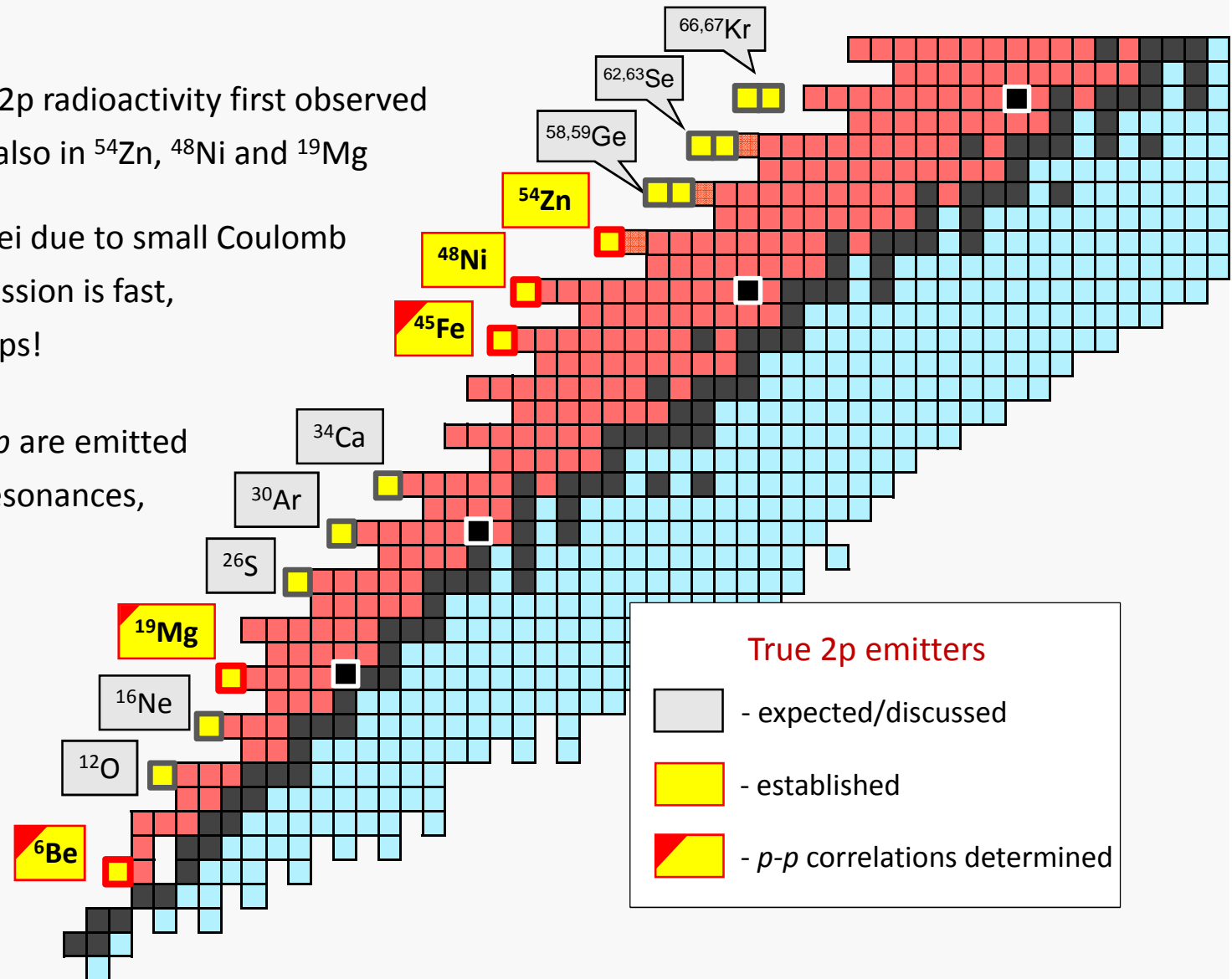
- ➔ True 2p decay is an essentially three-body phenomenon
- ➔ It offers more information: in addition to energy and half-life, there is a distribution of protons' momenta

Goldansky, Nucl. Phys. 19 (1960) 482

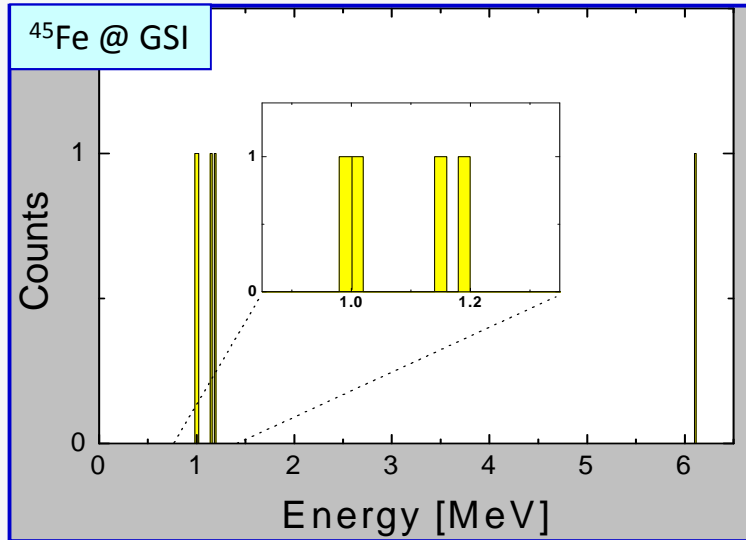


# True 2p emitters

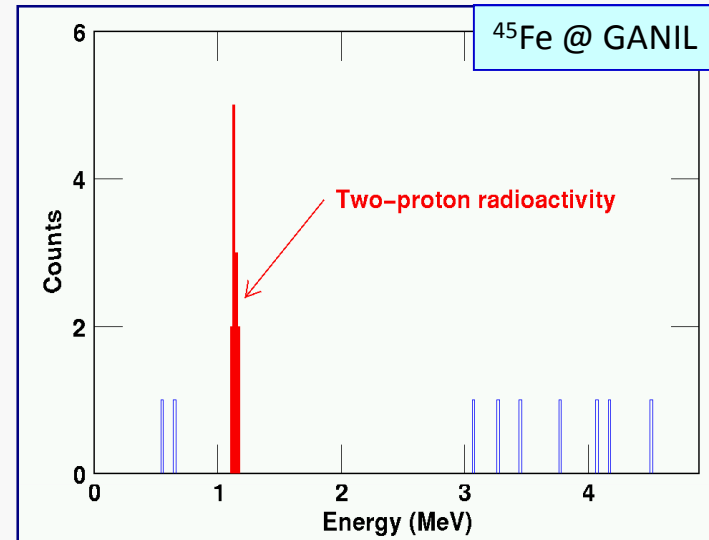
- Ground-state 2p radioactivity first observed in  $^{45}\text{Fe}$ . Later also in  $^{54}\text{Zn}$ ,  $^{48}\text{Ni}$  and  $^{19}\text{Mg}$
- In lighter nuclei due to small Coulomb barrier 2p emission is fast,  $T_{1/2}(^{19}\text{Mg}) = 4 \text{ ps!}$
- Below  $^{19}\text{Mg}$  2p are emitted from broad resonances, like  $^6\text{Be}$



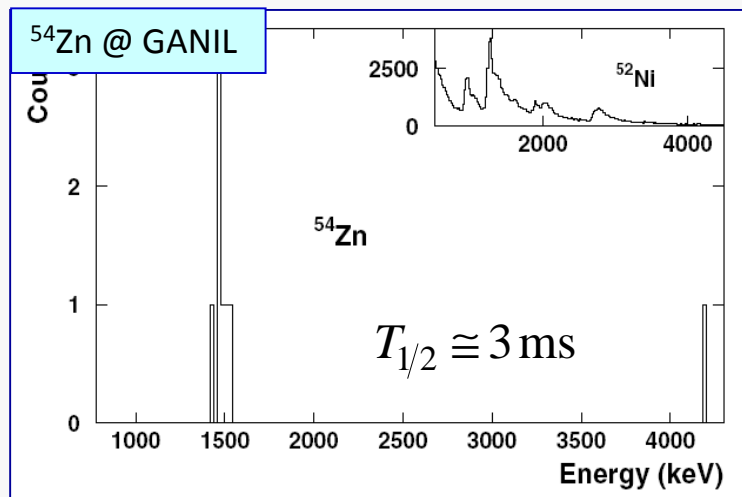
# First, with silicon detectors



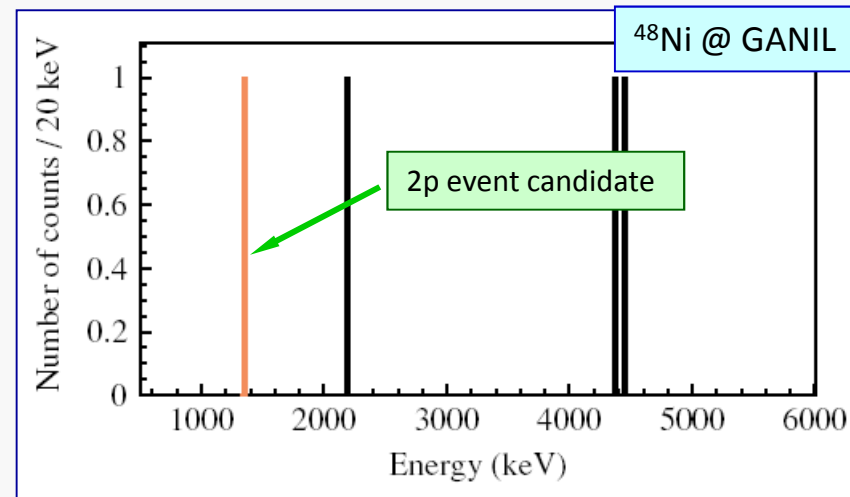
MP et al., EPJ A 14 (2002) 279



Giovinazzo et al., PRL 89 (2002) 102501



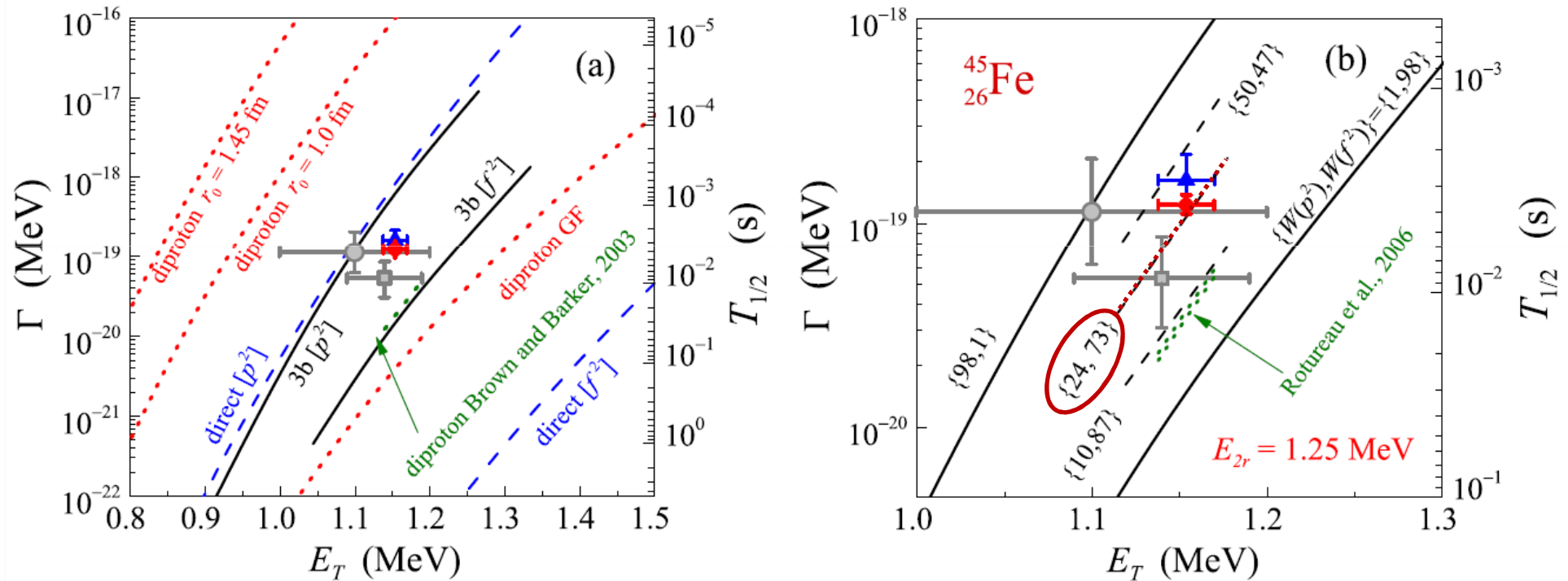
Blank et al., PRL 94 (2005) 232501



Dossat et al., PRC 72 (2005) 054315

# Decay energy and time

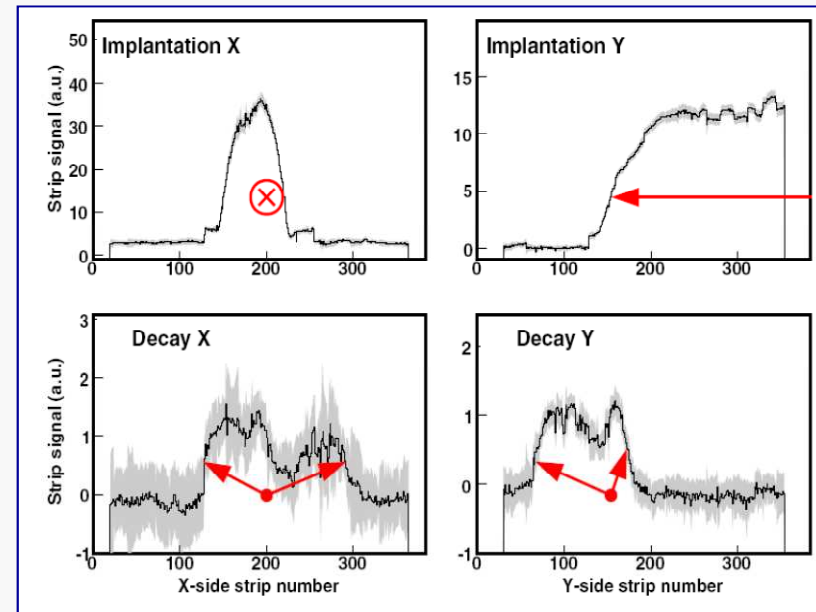
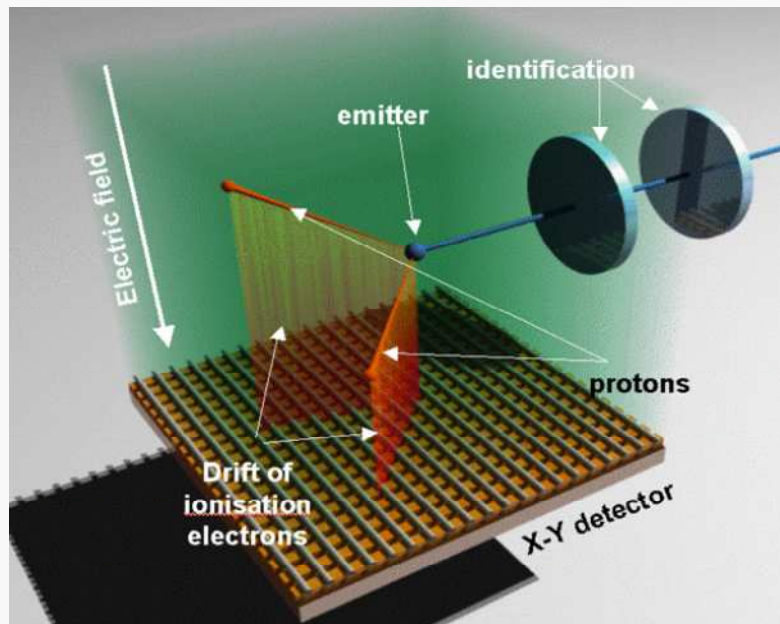
- The decay energy and the lifetime are enough to establish the 2p decay. Most models used for comparison, however, are based on two-body approximations.



- ➔ To explore fully the physics of the process, the correlations between proton's momenta must be determined! New detection technique is needed and a model capturing the three-body kinematics.

# TPC principle

- A „classical” Time Projection Chamber (TPC) constructed at CEN Bordeaux.  
It has fully electronic readout. The position on the x-y plane is detected by two orthogonal sets of 768 strips readout by ASIC-type electronics.



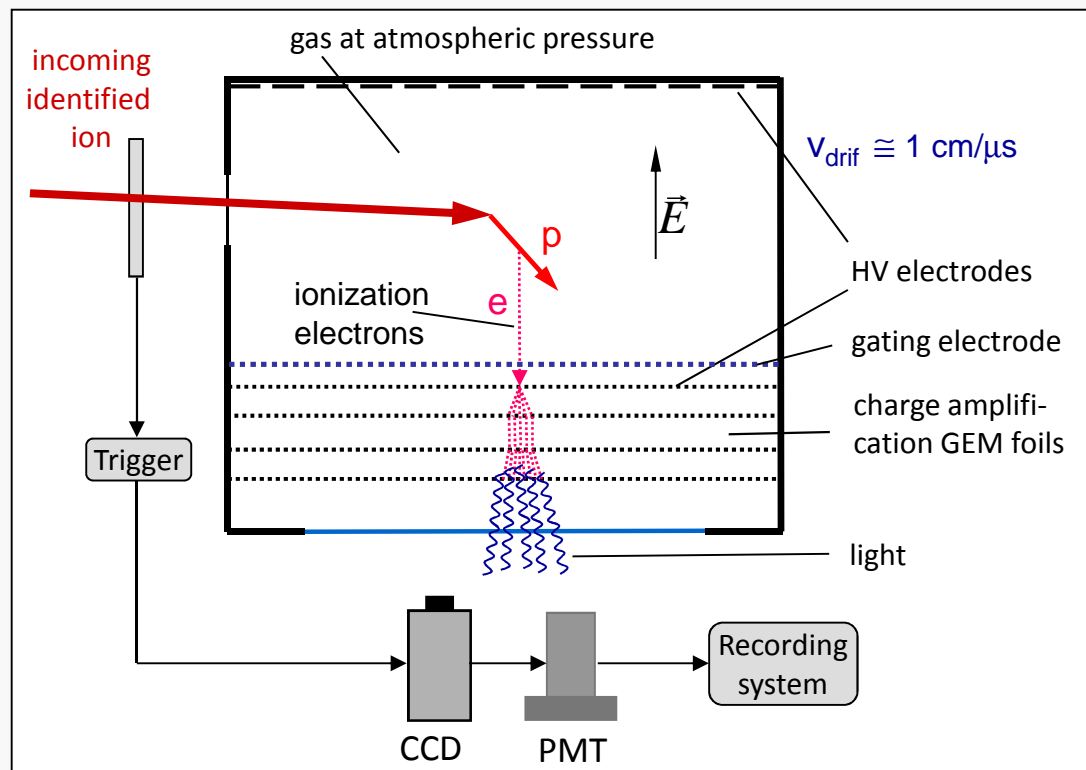
A decay event of  $^{45}\text{Fe}$

- ➔ Expensive and difficult to handle. Problems with information on z coordinate

Giovinazzo et al., PRL 99 (2007) 102501

# New idea – TPC with optical readout

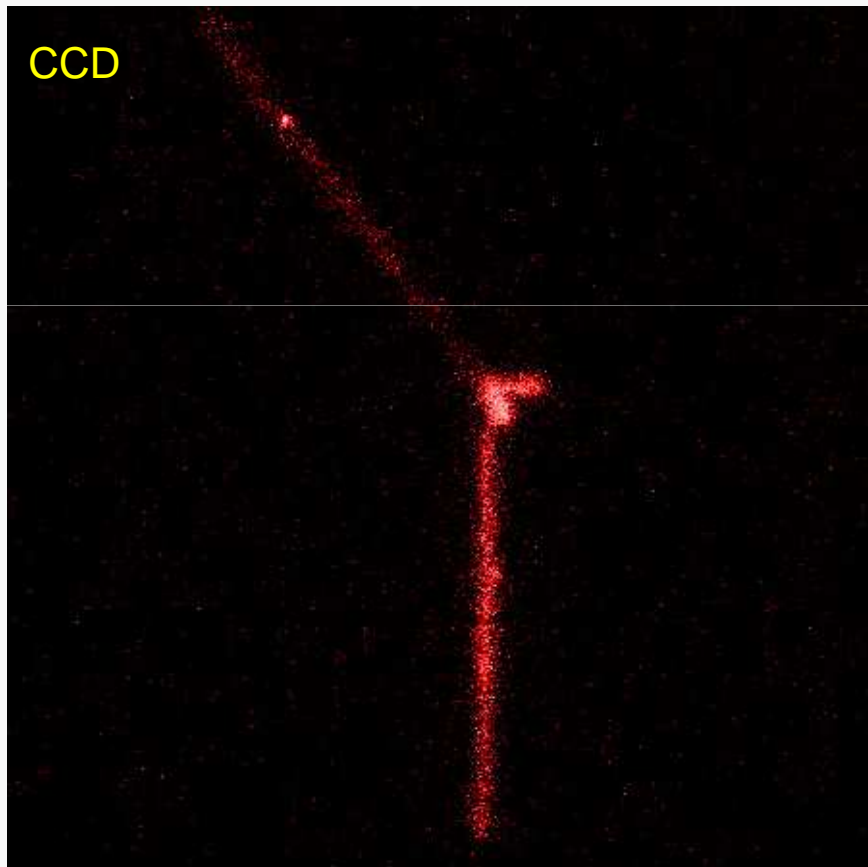
## OTPC – Optical Time Projection Chamber



Miernik et al., NIM A581 (2007) 194

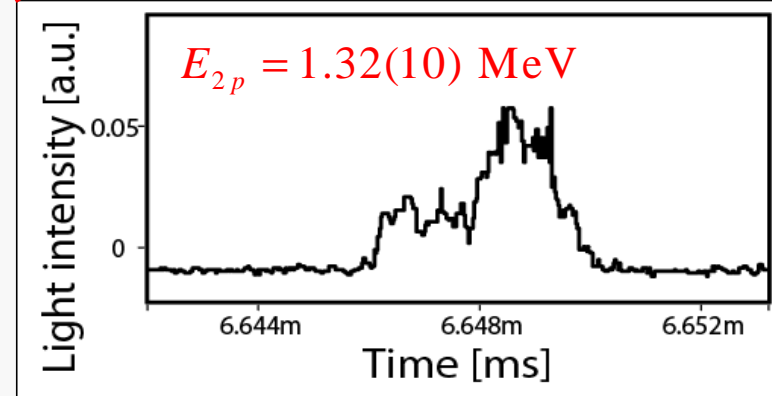
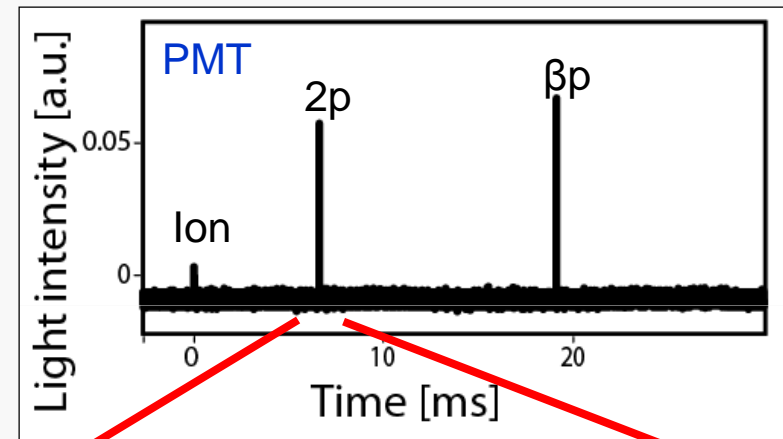
# OTPC operation

- The CCD picture yields 2D projection of tracks of particles



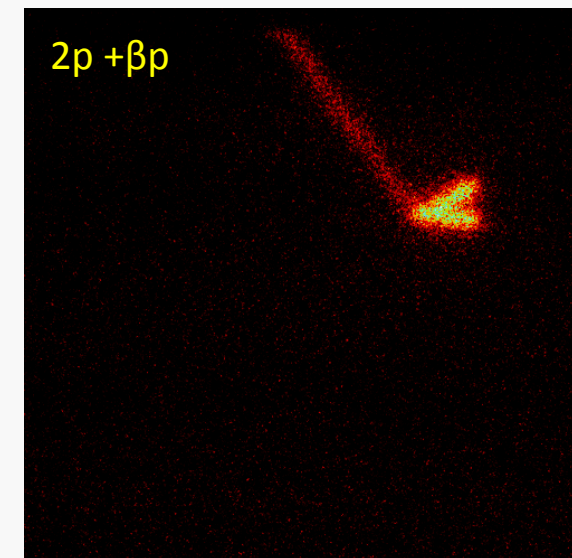
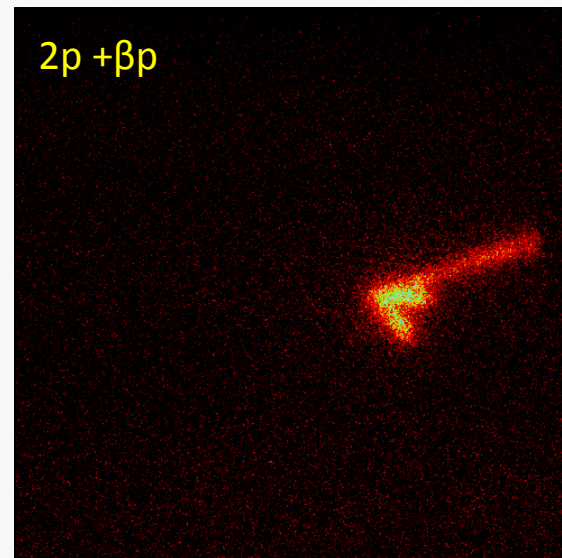
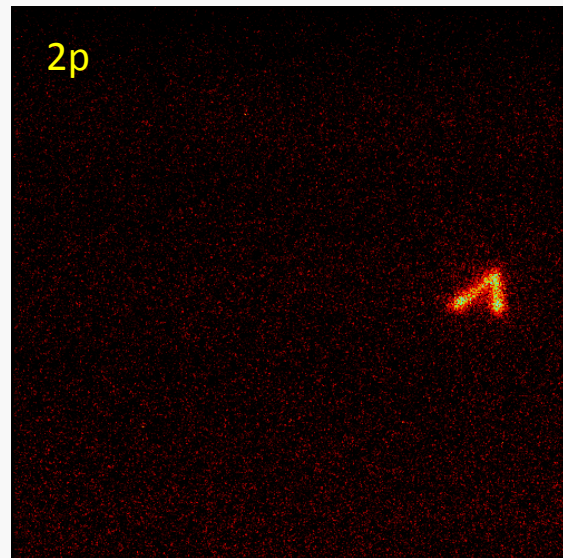
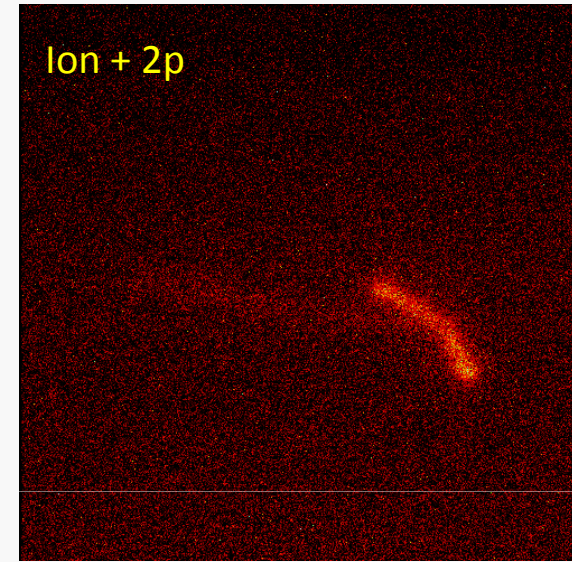
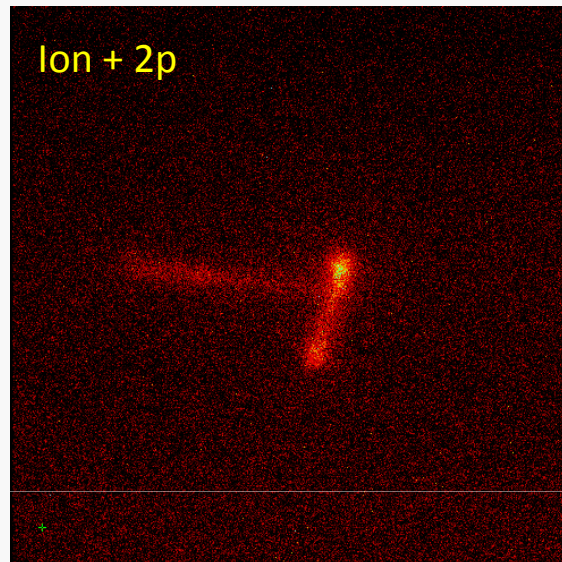
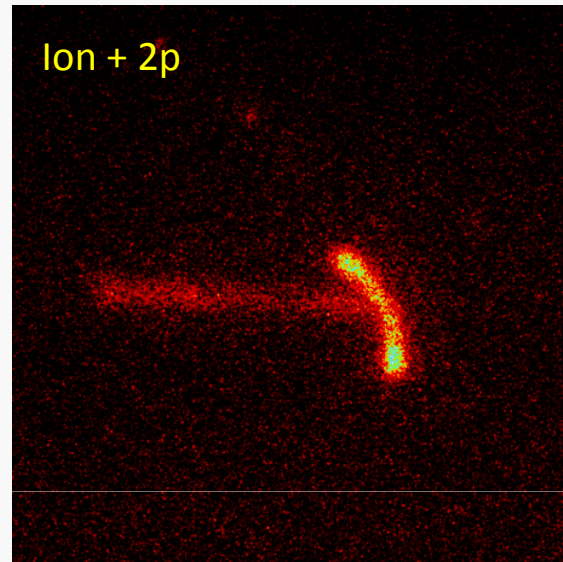
Pomorski et al., PRC 83 (2011) 061303(R)

- The PMT provides information on sequence, timing, and z-coordinate





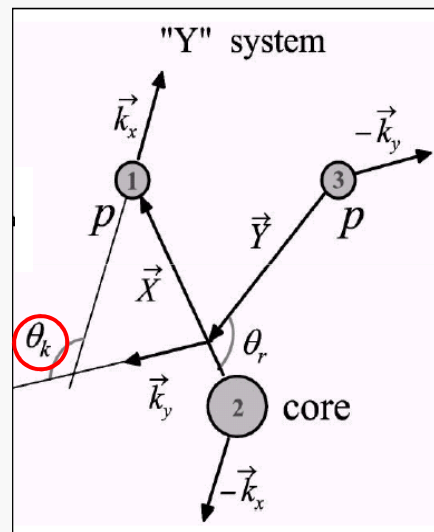
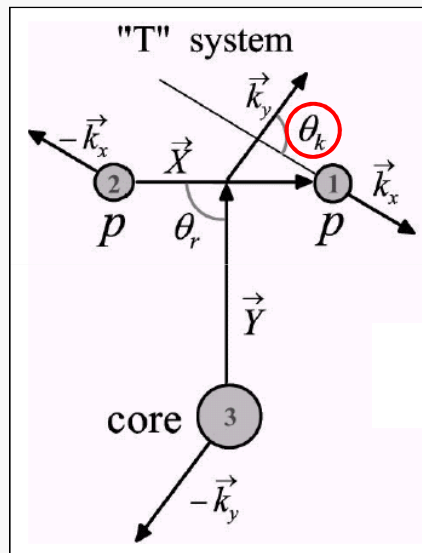
# $2p$ decays of $^{45}\text{Fe}$



# Jacobi coordinates

- ▶ Three-body kinematics is simpler in Jacobi coordinates.

The  $p$ - $p$  correlations are fully described by two variables:  $\varepsilon = E_X/E_T$  and  $\theta_k$



Transition from CM to  
Jacobi „T” system

$\vec{k}_1, \vec{k}_2$  - protons' momenta in CM

$$E_X = \frac{(\vec{k}_1 - \vec{k}_2)^2}{4 m_p}$$

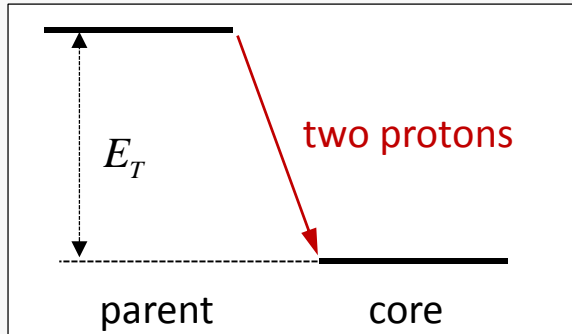
$\theta_k$  is the angle between vectors:

$$(\vec{k}_1 - \vec{k}_2) \text{ and } (\vec{k}_1 + \vec{k}_2)$$

- ▶ In place of radius and solid angle we introduce the *hyperradius* and *hyper solid angle*:

$$r, \Omega \rightarrow \rho, \Omega_5 \quad \Omega_5 = \{\theta_\rho, \Omega_X, \Omega_Y\} \quad \rho = \frac{A_1 A_2 A_3}{A_1 + A_2 + A_3} \left( \frac{\vec{r}_{12}^2}{A_3} + \frac{\vec{r}_{23}^2}{A_1} + \frac{\vec{r}_{31}^2}{A_2} \right)$$

# Three-body model



- Cluster approximation (two protons and the core)

- Parent wave function:

$$\Psi_{JM}(\rho, \Omega_5) = \frac{1}{\rho^{5/2}} \sum_{\alpha=\{K, \dots\}} \chi_{\alpha}(\rho) \mathcal{J}_{\alpha}^{JM}(\Omega_5)$$

radial functions      hyperspherical harmonics

→ Schrödinger equation, integration over angles:

$$\left[ \frac{d^2}{d\rho^2} - \frac{\mathcal{L}_K(\mathcal{L}_K + 1)}{\rho^2} + \frac{2\mu}{\hbar^2} E_T \right] \chi_{\alpha}(\rho) = \frac{2\mu}{\hbar^2} \sum_{\alpha'} (\hat{V}_{\alpha, \alpha'}) \chi_{\alpha'}(\rho) \quad \text{(coupled channels)}$$

- Main problem: the asymptotic form of radial functions is not known!

Solution of **Grigorenko and Zhukov**:

$$\chi_{\alpha}(\rho) \xrightarrow{\rho \rightarrow \infty} \sim \sum_{\alpha'} \hat{A}_{\alpha, \alpha'} \left[ G_{\mathcal{L}_0}(\eta_{\alpha}, \kappa\rho) + i F_{\mathcal{L}_0}(\eta_{\alpha}, \kappa\rho) \right]$$

$\hat{A}_{\alpha, \alpha'}$  is the matrix which diagonalizes the Coulomb part of  $\hat{V}_{\alpha, \alpha'}$

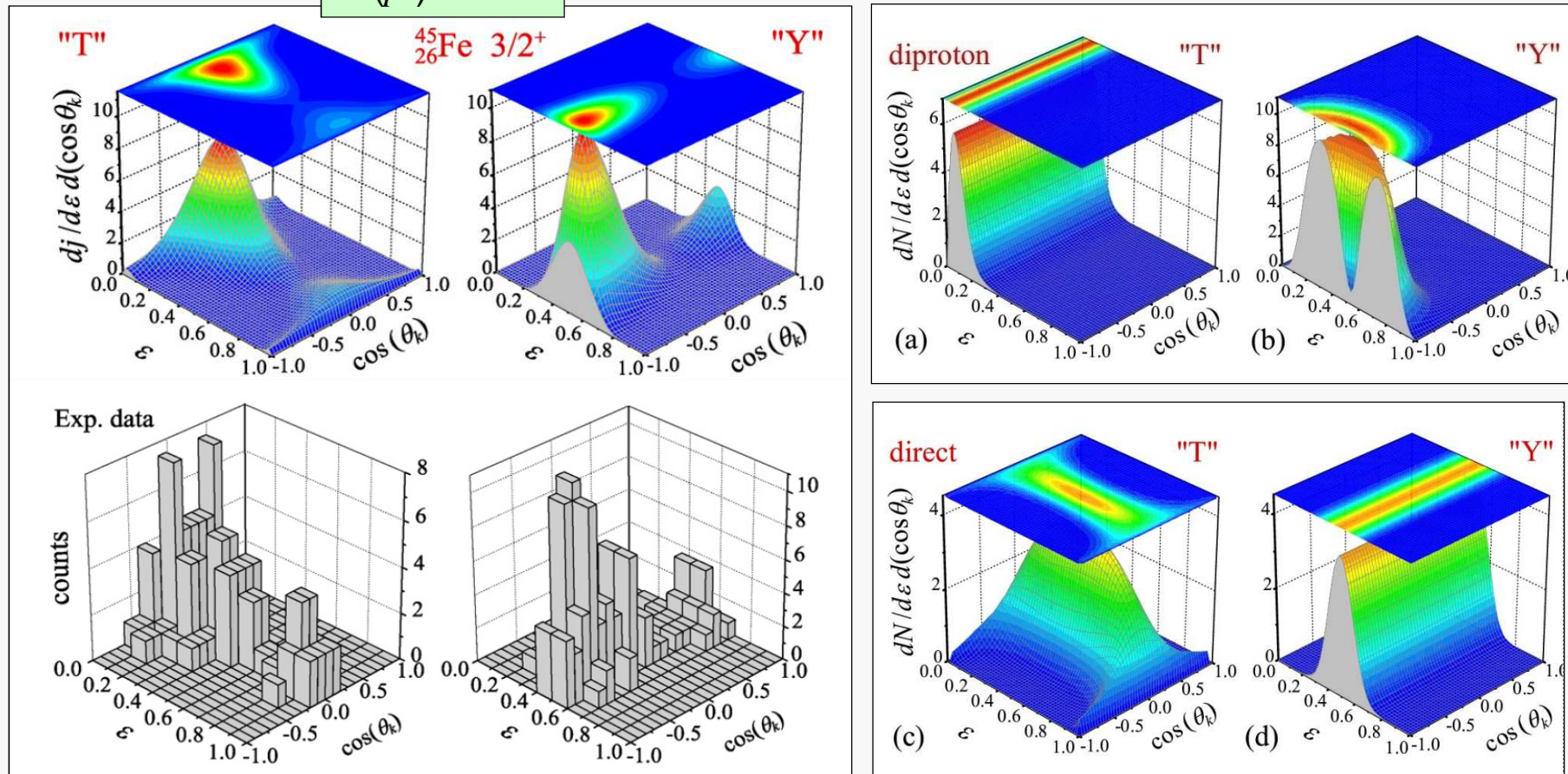
→ From radial functions the width and correlations :

$$\chi_{\alpha}(\rho) \Rightarrow \Gamma, \quad dj/d\varepsilon d(\cos \theta_k)$$

Grigorenko and Zhukov, PRC 68 (03) 054005  
M.P. et al, RMP (2012) 567

# Full picture for $^{45}\text{Fe}$

$W(p^2) = 24\%$



➤ All observables are simultaneously well reproduced by the 3-body model

Grigorenko *et al.*, PLB 677 (2009) 30

# Full picture for ${}^6\text{Be}$

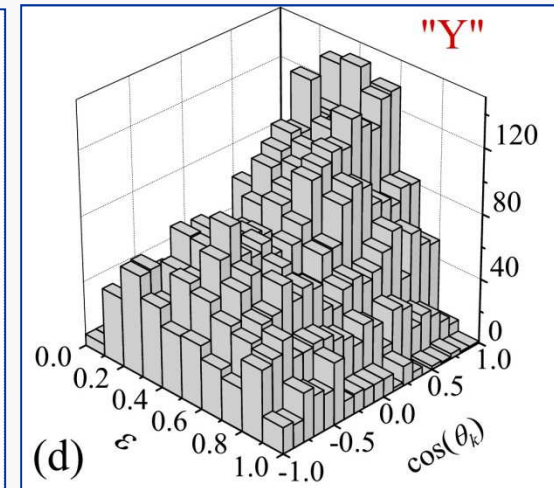
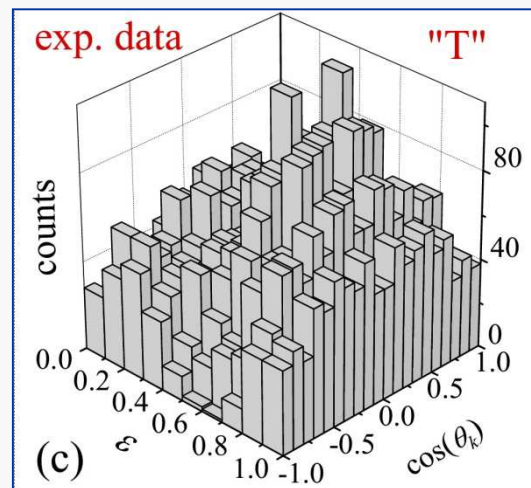
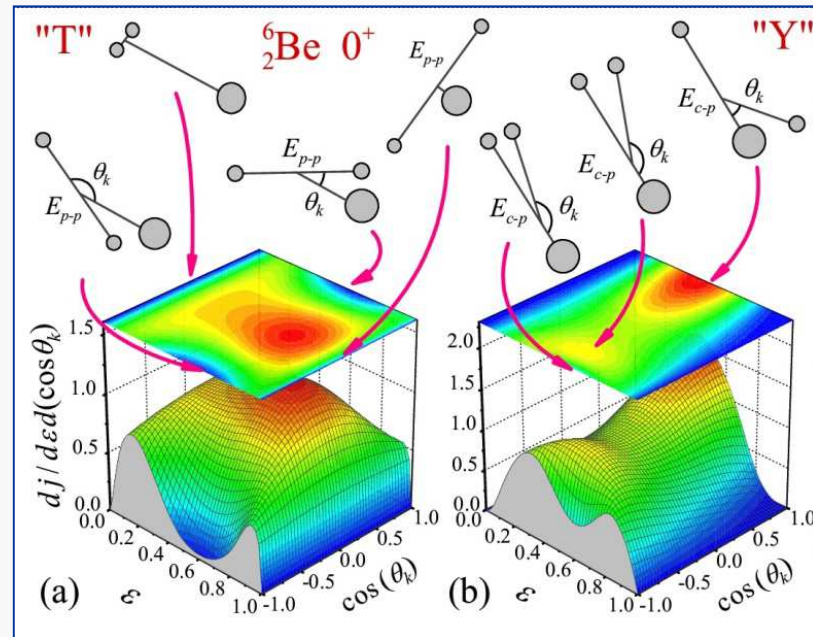
➤ Radioactive beam experiment at Texas A&M University

${}^{10}\text{C}$  inelastic scattering

- ①  $p({}^{10}\text{B}, {}^{10}\text{C})n$  @15 MeV/u
- ② 11 MeV/u  ${}^{10}\text{C} + \text{C/Be} \rightarrow {}^{10}\text{C}^*$
- ③  ${}^{10}\text{C}^* \rightarrow {}^6\text{Be} + \alpha$

Mercurio et al., PRC 78 (08) 031602(R)

Grigorenko et al., PLB 677 (2009) 30



# Case of $^{19}\text{Mg}$

- Decay in-flight and tracking for very short-lived  $2p$  decays at GSI

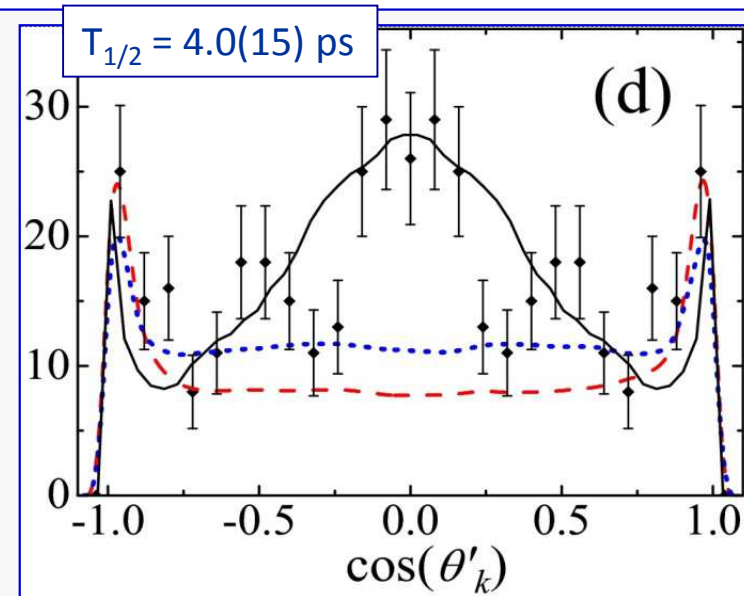
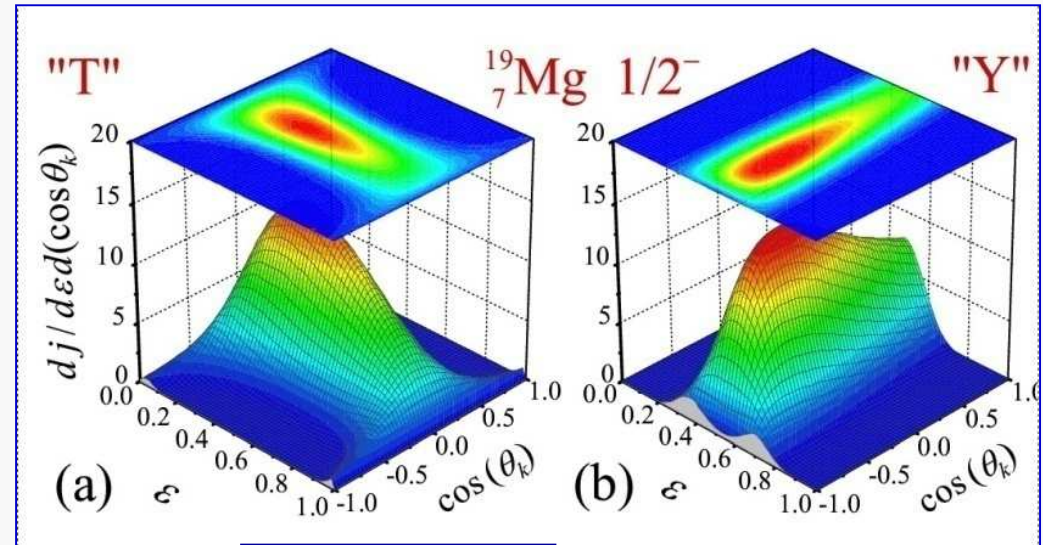
Radioactive beam experiment

- ➊  $^{24}\text{Mg}$  @ 600 MeV/u + Be  $\rightarrow$   $^{20}\text{Mg}$
- ➋  $^{20}\text{Mg}$  + Be  $\rightarrow$   $^{19}\text{Mg}$

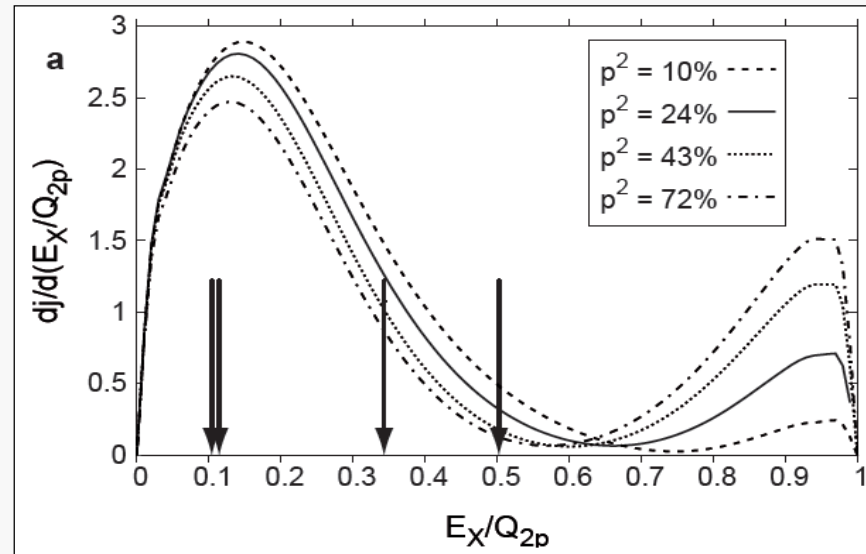
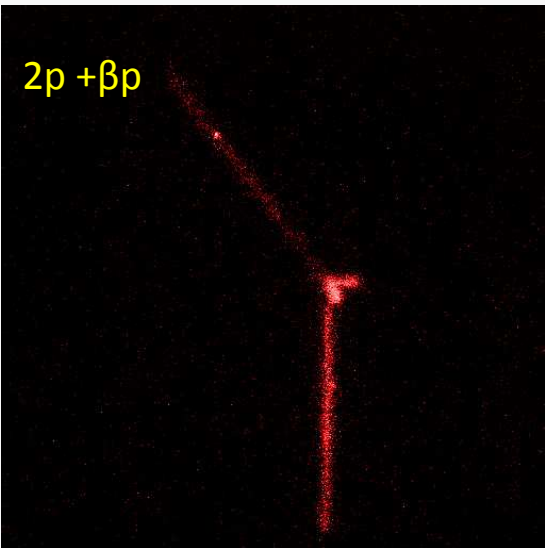
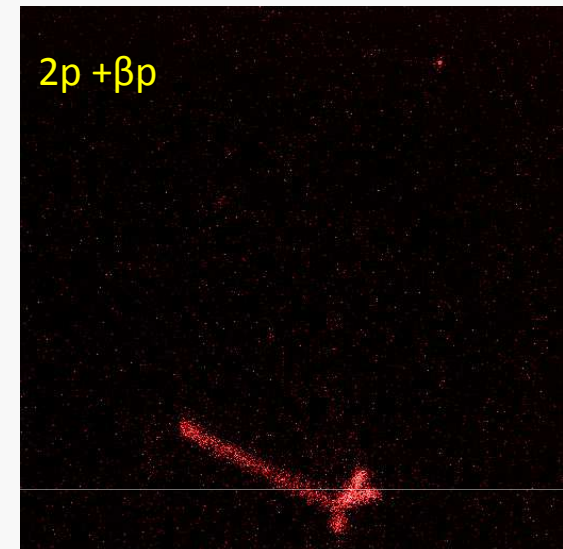
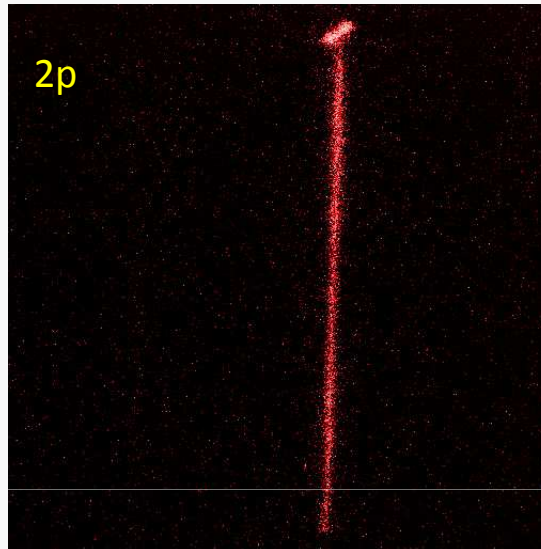
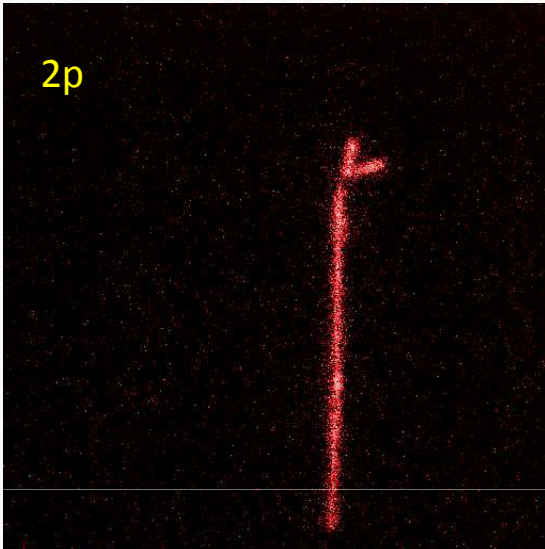
- Only projection of proton's momenta on the plane could be determined.

Mukha et al., PRL. 99 (2007) 182501

Mukha et al., PR C 77 (2008) 061303(R)

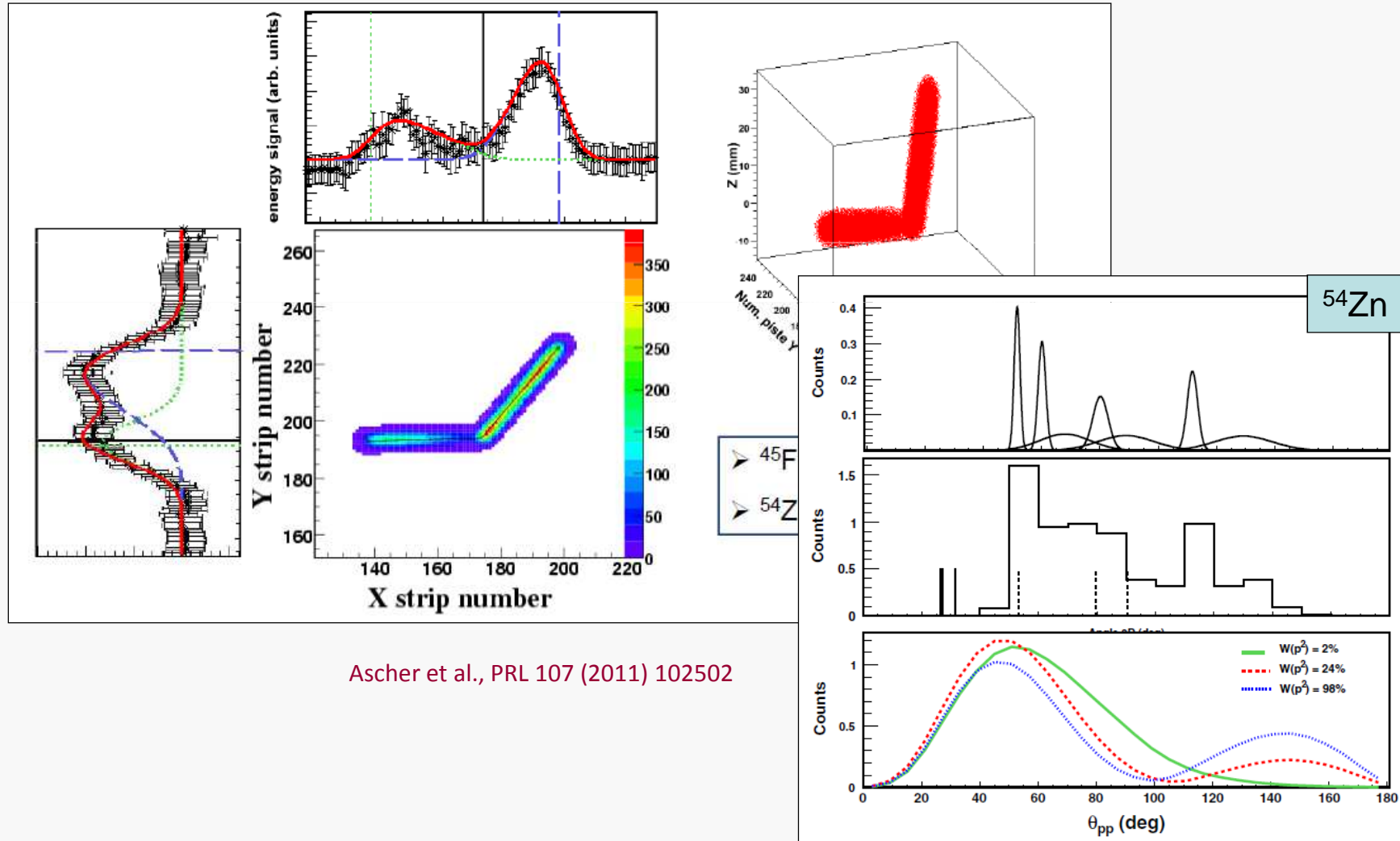


# 2p decays of $^{48}\text{Ni}$



# $p$ - $p$ correlations in $^{54}\text{Zn}$

►  $^{54}\text{Zn}$  studied at GANIL with the Bordeaux TPC. Seven events reconstructed in 3D



Ascher et al., PRL 107 (2011) 102502



# Range of lifetimes

➤ The three-body model seems to work in the range of half-lives covering 18 orders of magnitude!

- Invariant mass method for broad resonances

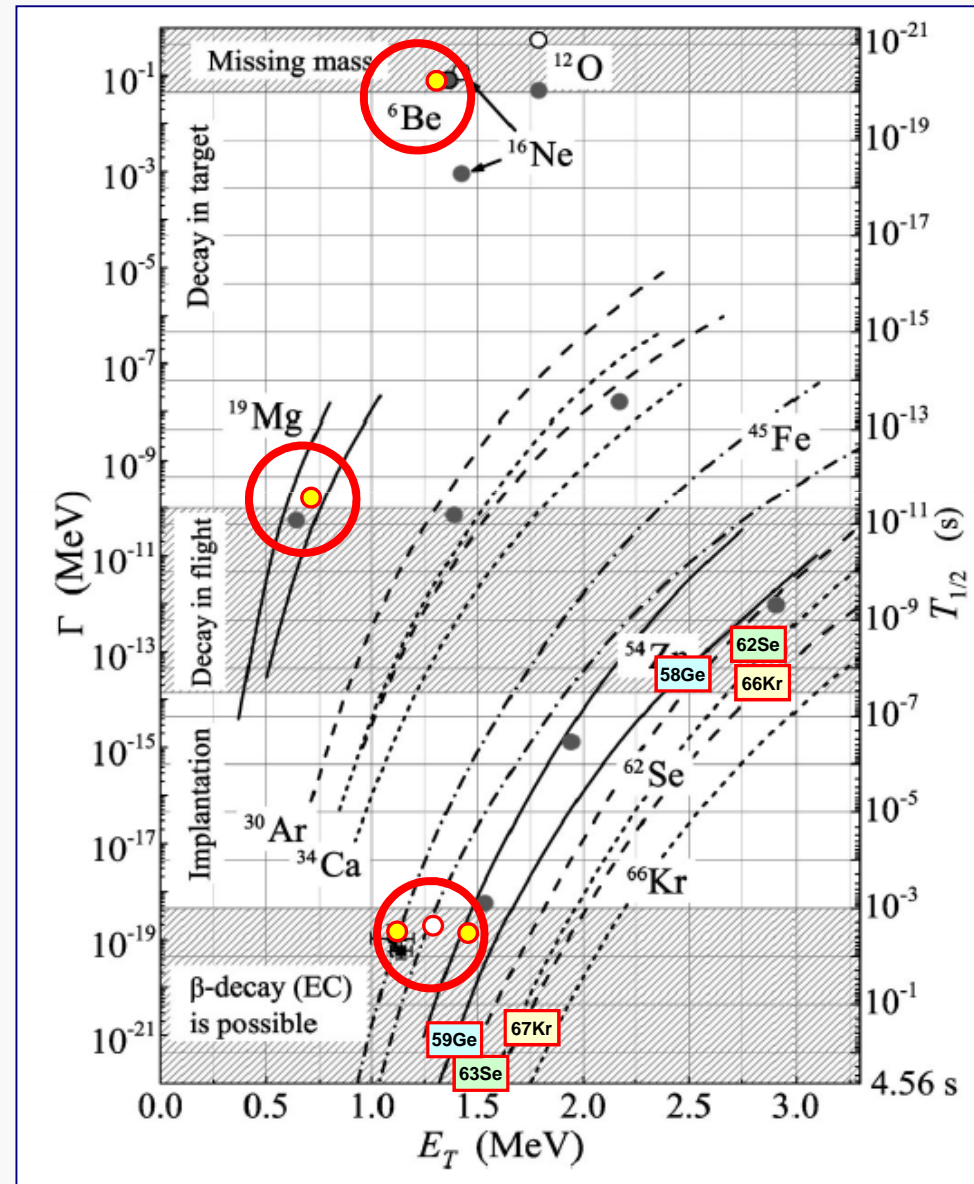
$$T_{1/2} \leq 10^{-19} \text{ s}$$

- In-flight decays

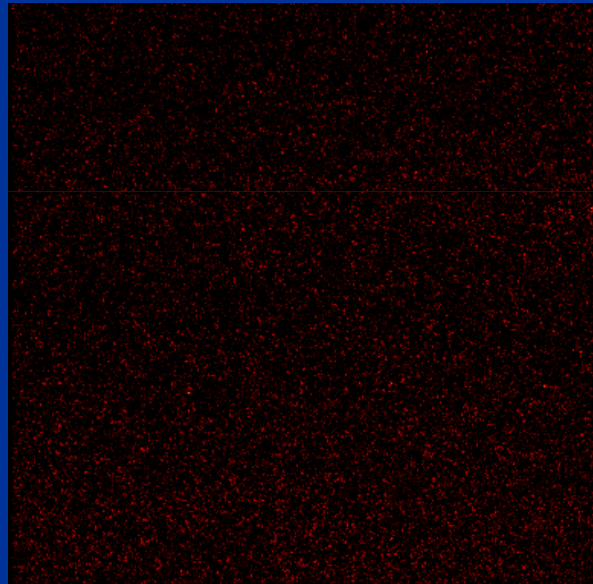
$$T_{1/2} = 1 \text{ ps} - 50 \text{ ns}$$

- Implantation method

$$T_{1/2} > 50 \text{ ns}$$

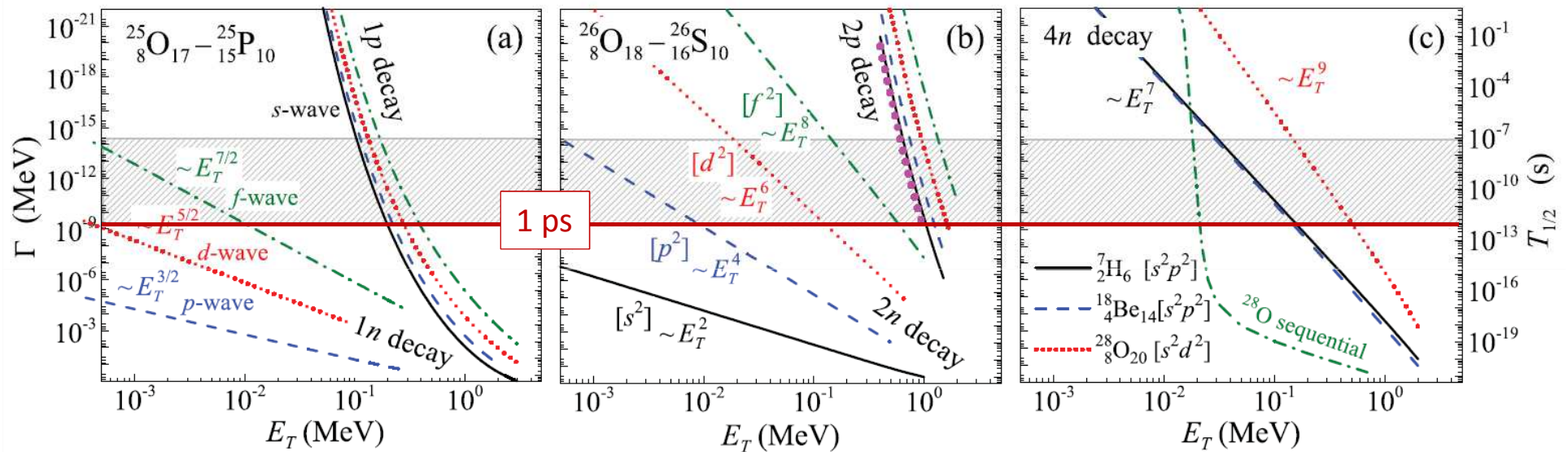


# Neutron radioactivity



# n, 2n, or 4n?

- The xn emission estimated by a simplified version of 3-body model (direct decay model) and compared to proton emission



- Extremely small decay energy needed for a measurable decay time of 1n emission.

Very unlikely to find a candidate in the  $s$ - $d$  shell.

- Broader energy window thus higher chances to find a good case.

$^{26}\text{O}$  could be a candidate!

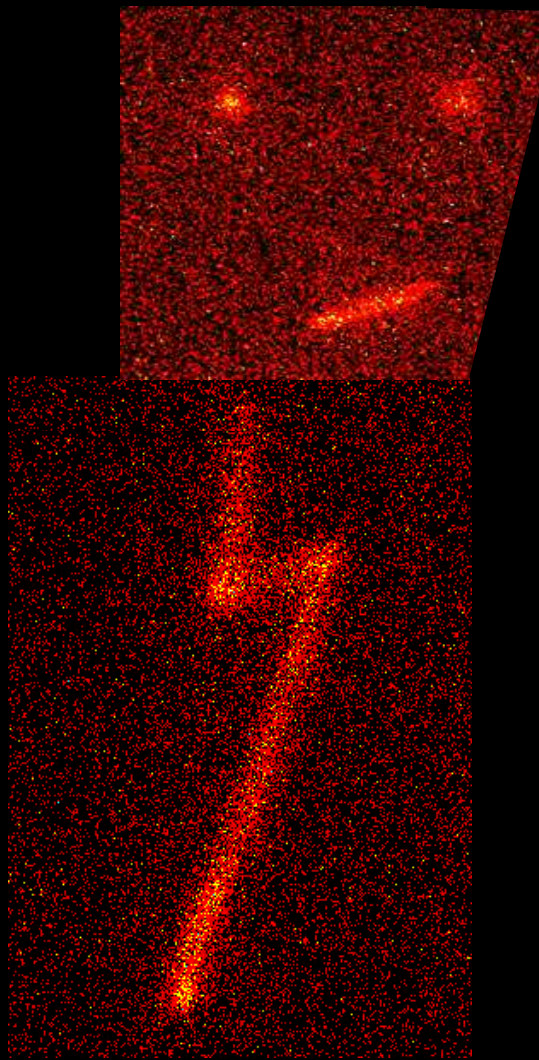
- Special energy configuration required (only  $S_{4n} < 0$ ) but not impossible.

$^7\text{H}$  and  $^{28}\text{O}$  are not excluded!

# Summary

- The particle radioactivity ( $p$ ,  $\alpha$ ) at the proton drip-line is very **efficient tool** in nuclear spectroscopy. Yield masses (separation energies) of very exotic systems, provides stringent tests for models of nuclear structure.
- More than **40 proton-emitting states** are known. 7 emitters exhibit **fine structure**. Observation of g.s. proton radioactivity for  $Z < 50$  remains an experimental challenge.
- The observation of „**superallowed**”  $\alpha$ -decay  $^{104}\text{Te} \rightarrow ^{100}\text{Sn}$  is approaching.
- The direct ground-state **2p emission established** for  $^6\text{Be}$ ,  $^{19}\text{Mg}$ ,  $^{45}\text{Fe}$ ,  $^{48}\text{Ni}$ , and  $^{54}\text{Zn}$ . The hunt for other cases continues.  $^{30}\text{Ar}$  and  $^{59}\text{Ge}$  will be tried soon.
- The observation of **full p-p correlation picture** in  $^6\text{Be}$  and  $^{45}\text{Fe}$  was the major breakthrough in the field. The 3-body model of Grigorenko and Zhukov was confirmed and the **influence of nuclear structure** on the 2p emission was demonstrated. Further theoretical work is needed to elucidate this issue.
- Observation of **two-neutron radioactivity** is probable in nuclei accesible already now.

# Thank you!



# Radioactive decays

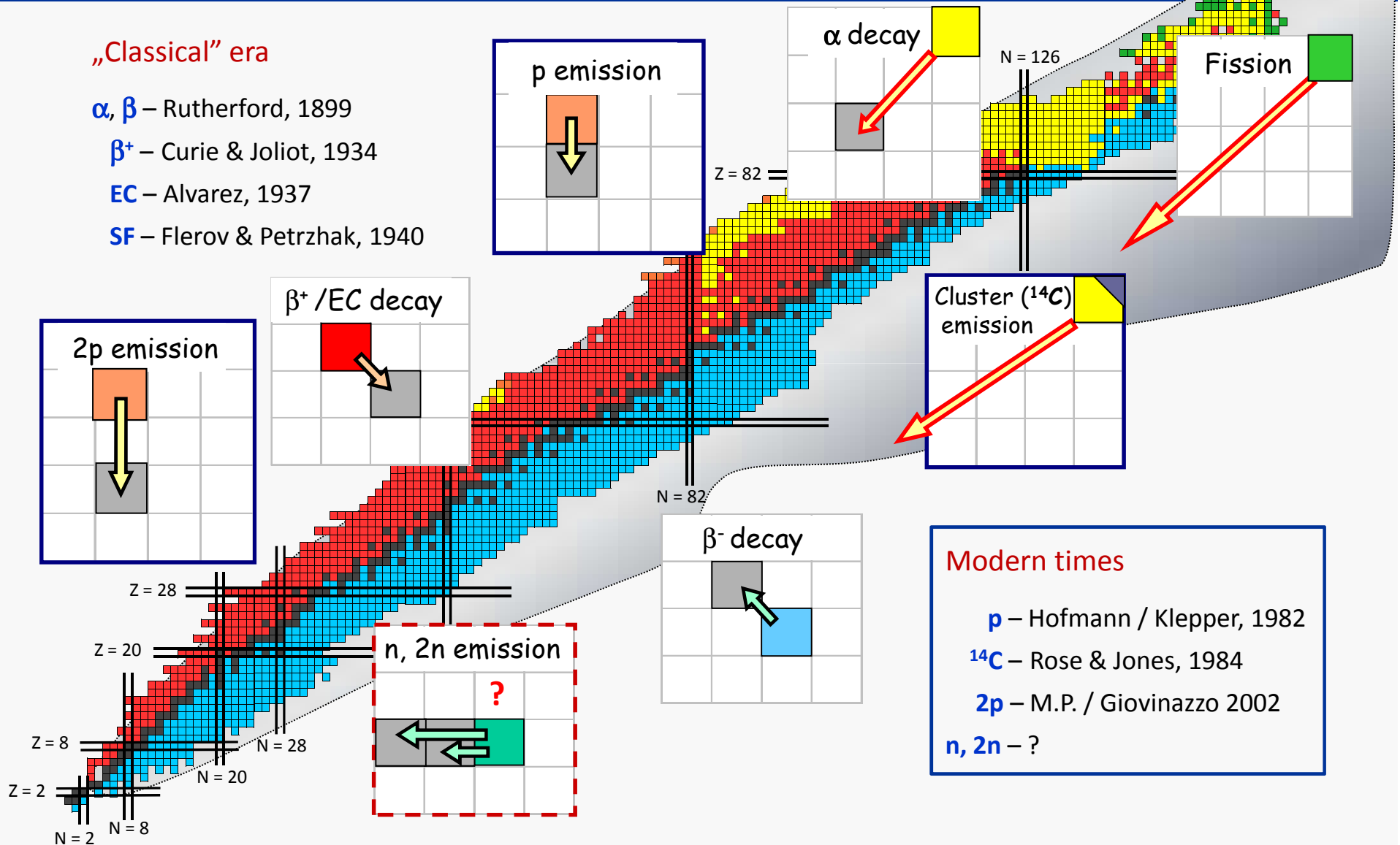
## „Classical” era

$\alpha$ ,  $\beta$  – Rutherford, 1899

$\beta^+$  – Curie & Joliot, 1934

EC – Alvarez, 1937

SF – Flerov & Petrzhak, 1940



## Modern times

**p** – Hofmann / Klepper, 1982

$^{14}\text{C}$  – Rose & Jones, 1984

**2p** – M.P. / Giovinazzo 2002

**n, 2n** – ?

# What is radioactive?

- What is plotted on the chart? Present practice: all systems we **know something about**.
- ➔ Suggestion: plot only long-lived, i.e. **stable** and **radioactive** (those which **exist!**)

## Radioactivity

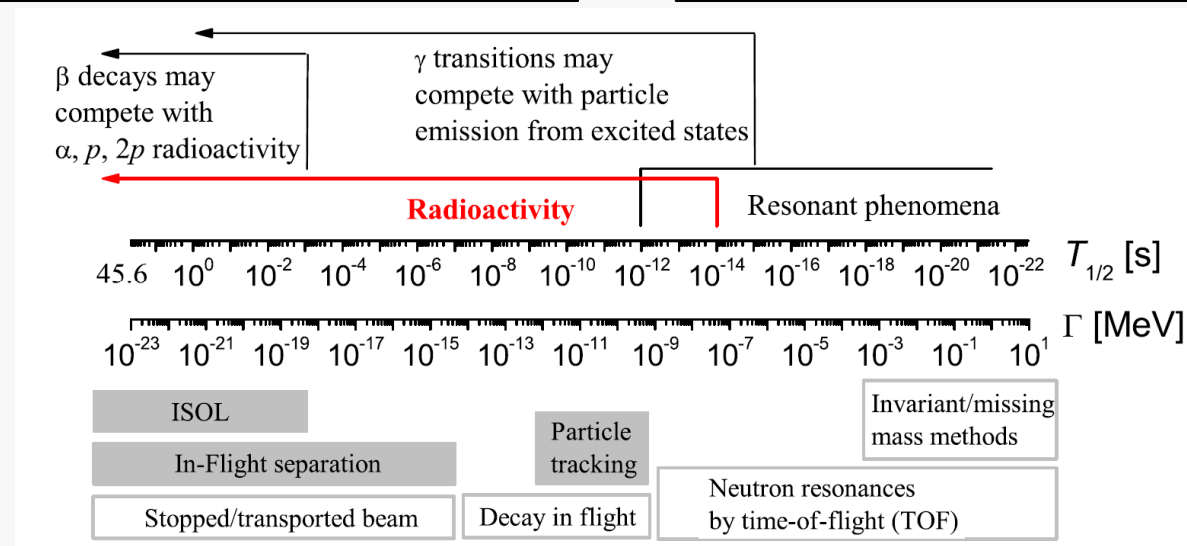
- Slow enough to form neutral atoms
- Characteristic **time** measured directly
- Independent of formation mechanism

$$T_{1/2} \geq 10^{-14} \text{ s}$$

## Reactions/Resonances

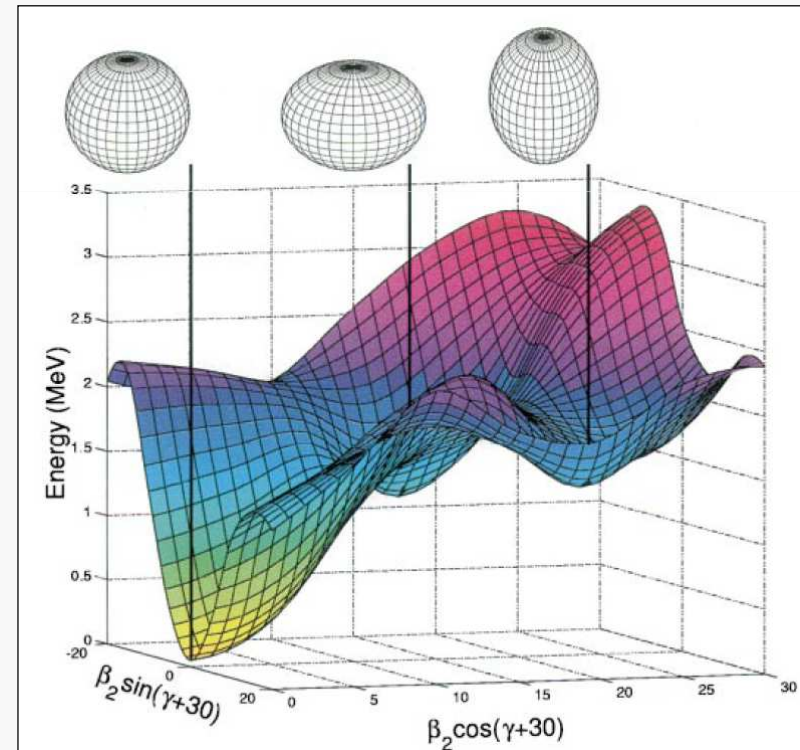
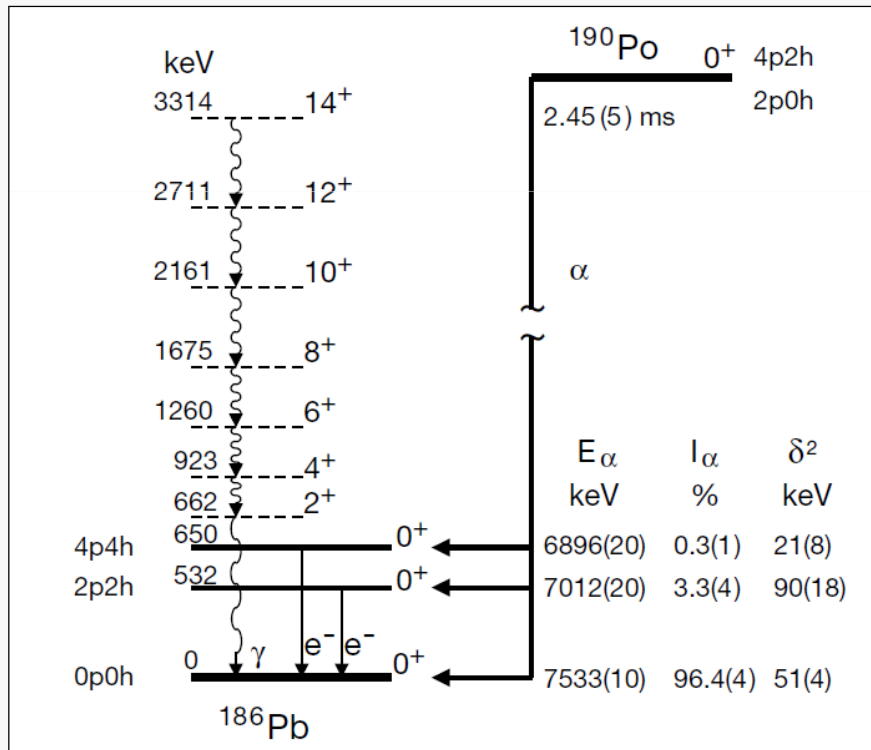
- Fast on atomic scale
- Characteristic **width** measured directly
- Influenced by reaction mechanism

$$\Gamma \geq 1 \text{ meV}$$



# Shape coexistence in $^{186}\text{Pb}$

- Alpha spectroscopy revealed surprising nature of lowest excitations in  $^{186}\text{Pb}$ : the three  $0^+$  levels with different deformation. A unique case of interaction between proton  $p-h$  excitations across the closed shell ( $Z=82$ ) and large number of mid-shell neutrons,  $N=(82+126)/2$ .



Andreyev *et al.*, *Nature* 405 (2000) 430