Hubbard model with spin dependent disorder

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Collaboration and Discussion

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Introduction: Optical lattices



L. Sanchez-Palencia, M. Lewenstein, Nature Phys. 6, 87 (2010)

$$H = \sum_{ij\sigma} J_{ij\sigma} a_{i\sigma}^{\dagger} a_{j\sigma} + \sum_{i\sigma} V_{i\sigma} a_{i\sigma}^{\dagger} a_{i\sigma} + \frac{1}{2} \sum_{i\sigma\sigma'} U_{i\sigma\sigma'} n_{i\sigma} n_{i\sigma'}$$

- high tunability and control over system parameters

- (almost) perfect realization of model lattice Hamiltonians
- access to systems not seen in solid-state matter

Introduction: Eg. Hexagonal spin-dependent lattice



P. Soltan-Panahi et al., Nature Phys. 7, 434 (2011); D McKay, B DeMarco, NJP 12, 055013 (2010)

$$H = \sum_{ij\sigma} J_{ij\sigma} a_{i\sigma}^{\dagger} a_{j\sigma} + \sum_{i\sigma} V_{i\sigma} a_{i\sigma}^{\dagger} a_{i\sigma} + \frac{1}{2} \sum_{i\sigma\sigma'} U_{i\sigma\sigma'} n_{i\sigma} n_{i\sigma'}$$

- "For neighbouring sites along the vertices of the hexagonal lattice, ..., the local polarization alternates between σ^+ and σ^- ."

- "As atoms in a light field experience a polarization dependent a.c. Stark shift, the potential at these σ^+ and σ^- sites is different for different atomic Zeeman substates labeled by m_F "

Introduction: Eg. Optical lattices with disorder



L. Sanchez-Palencia, M. Lewenstein, Nature Phys. 6, 87 (2010)

$$H = \sum_{ij\sigma} J_{ij\sigma} a^{\dagger}_{i\sigma} a_{j\sigma} + \sum_{i\sigma} V_{i\sigma} a^{\dagger}_{i\sigma} a_{i\sigma} + \frac{1}{2} \sum_{i\sigma\sigma'} U_{i\sigma\sigma'} n_{i\sigma} n_{i\sigma'}$$

- Optical speckle potential or superposition of incommensurate periodic fields creates a random (pseudo random) disorder on top of the optical lattice

Introduction: Eg. Anderson localization



G. Roati et al., Nature 453, 895 (2008)

$$\Psi(x)|^2 \sim e^{-2x/L_{loc}}$$

- Observation of Anderson localization of matter waves made of $^{87}\mathrm{Rb},$ V_R is disorder strength

Introduction: Our motivation

Spin Dependent Disorder

- soon available experiments where disorder acts selectively only on one spin component

- not know theoretically too much ...

Let's do theoretical work ...

Hubbard model with local spin dependent disorder

$$H = \sum_{ij\sigma} t_{ij\sigma} a_{i\sigma}^{\dagger} a_{j\sigma} + \sum_{i\sigma} V_{i\sigma} n_{i\sigma} + U \sum_{i} n_{i\uparrow} n_{i\downarrow},$$

with $V_{i\uparrow} = \epsilon_i$ - random variable, $V_{i\downarrow} = 0$.

$$\mathcal{P}(\epsilon_1, ..., \epsilon_n, ...) = \prod_i P(\epsilon_i)$$

and

$$P(\epsilon) = \begin{cases} \frac{1}{\Delta} & |\epsilon| < \frac{\Delta}{2} \\ 0 & |\epsilon| > \frac{\Delta}{2}. \end{cases}$$

We consider now fermions $({}^{40}K)$ and leave bosons $({}^{87}Rb)$ for later on.

Model A: $\langle n_{\uparrow} \rangle + \langle n_{\downarrow} \rangle = n$ conserved - Today Model B: $\langle n_{\uparrow} \rangle = n_{\uparrow}$, $\langle n_{\downarrow} \rangle = n_{\downarrow}$ independently conserved - Next time

Hubbard model with spin dependent disorder

Simple trick

$$n_i = n_{i\uparrow} + n_{i\downarrow},$$
$$m_i = n_{i\uparrow} - n_{i\downarrow},$$

leads to Hamiltonian

$$H = \sum_{ij\sigma} t_{ij\sigma} a_{i\sigma}^{\dagger} a_{j\sigma} + \sum_{i} \mu_{i} n_{i} + \sum_{i} h_{i} m_{i} + U \sum_{i} n_{i\uparrow} n_{i\downarrow},$$

where we identify local random chemical potential and random magnetic field that are correlated

$$\langle \mu_j h_j \rangle_{\rm dis} = \frac{\Delta^2}{48} \delta_{ij}.$$

- possible simultaneous response in both density and spin channels

Hubbard model with spin dependent disorder -Dynamical Mean-Field Theory

M. Ulmke, V. Janis, D. Vollhardt, Phys. Rev. B 51, 10411 (1995)

$$Z\{\mathcal{G}^{-1},\epsilon_i\} = \int Da^* Da \ e^{\mathcal{A}}$$

$$\mathcal{A} = \sum_{n\sigma} a_{n\sigma}^* (\mathcal{G}^{-1} - \epsilon_i \delta_{\sigma\uparrow}) a_{n\sigma} - U \int_0^\beta d\tau n_\uparrow(\tau) n_\downarrow(\tau)$$

with $\mathcal{G}_{n\sigma}^{-1} = G_{n\sigma}^{-1} + \Sigma_{n\sigma}$ and

$$G_{n\sigma} = \int_{0}^{\beta} d\tau e^{i\omega_n \tau} \langle \langle a_{\sigma}^*(\tau) a_{\sigma}(0) \rangle_{\rm qm,T} \rangle_{\rm dis}$$

and generalization for AB lattices with given DOS No Anderson localization physics

Preliminary results obtained within HF-QMC for

 $\begin{array}{l} n=0.5,\ 0.7,\ {\rm and}\ 1.0\\ U=0.1,\ 1.0,\ 2.0,\ {\rm and}\ 3.0\\ \Delta=0.0,\ 1.0,\ 3.0,\ {\rm and}\ 5.0 \end{array}$

model semi-elliptic density of states with bandwidth W = 2.

Computed: FM susceptibility, AF susceptibility, compressibility, DOS at Fermi level, double occupancy (local moment), magnetization as functions of temperatures.

$$d = \langle n_{i\uparrow} n_{i\downarrow} \rangle = \langle n_{i\uparrow} + n_{i\downarrow} \rangle - \frac{2}{3} \langle \mathbf{S}_i^2 \rangle$$

Almost non-interacting system away from half-filling

n=0.5, U=0.1, W=2.0

0.06 0.055 0.05 $\chi_{FM}^{}(T)$ (L) ۲ × 0.5 ¥ 0.045 $\Delta = 0.0$ $\Delta = 1.0$ 0.04 $\Delta = 3.0$ -2 $\Delta = 5.0$ 0.035 -1 0.03 n -3 0.2 °0 0.2 0.4 0.6 0.8 0 0.2 0.4 0.6 0.8 0 0.2 0.4 0.6 0.8 ٥ 0.4 0.6 0.8 10 3.5 0.25 8 3 1.5 0.2 $1/\chi_{FM}(T)$ $1/\chi_{\rm AF}(T)$ 6 2.5 DOS E 0.15 2 1.5 0.1 0.5 0.05 0.5 0 0 Δ 0 0.2 0.4 0.6 0 0.2 0.8 0.8 0 0.2 0.4 0.6 0.8 0.2 0.4 0.6 0.8 0.4 0.6 Ό 0 т т т т

- divergence in FM susceptibility and compressibility
- negative sign of FM susceptibility and compressibility
- no singularity in other thermodynamic quantities
- finite magnetization, increasing local moments

n=0.5, U=0.1, W=2.0

Hubbard model with normal disorder - comparison

Almost non-interacting system away from half-filling

n=0.5, U=0.1

n=0.5, U=0.1



- no singularity in FM susceptibility and (?) compressibility

- no singularity in other thermodynamic quantities
- zero magnetization, decreasing local moments

Almost non-interacting system at half-filling

n=1.0, U=0.1, W=2.0

n=1.0, U=0.1, W=2.0



- singularity in AF susceptibility killed by disorder
- no singularity FM susceptibility and compressibility
- no singularity in other thermodynamic quantities
- zero magnetization, increasing local moments

Intermediate interacting system away from half-filling

n=0.5, U=1.0, W=2.0

n=0.5, U=1.0, W=2.0



- divergence in FM susceptibility and compressibility
- negative sign of FM susceptibility and compressibility
- no singularity in other thermodynamic quantities
- finite magnetization, increasing local moments

Hubbard model with normal disorder - comparison

Intermediate interacting system away from half-filling

n=0.5, U=1.0

n=0.5, U=1.0, W=2.0



- no singularity in FM susceptibility and (??) compressibility
- no singularity in other thermodynamic quantities
- zero magnetization, decreasing local moments

Intermediate interacting system at half-filling

n=1.0, U=1.0, W=2.0





- divergence in AF susceptibility
- negative sign of AF susceptibility
- no singularity in other thermodynamic quantities
- zero magnetization, increasing local moments

Strongly interacting system away from half-filling

n=0.5, U=2.0, W=2.0

n=0.5, U=2.0, W=2.0



- divergence in FM susceptibility and compressibility
- negative sign of FM susceptibility and compressibility
- no singularity in other thermodynamic quantities
- finite magnetization, increasing local moments

Hubbard model with normal disorder - comparison

Strongly interacting system away from half-filling

n=0.5, U=4.0, W=2.0

n=0.5, U=4.0, W=2.0



- singularity in FM susceptibility and AF susceptibility
- susceptibility crossing
- no singularity in other thermodynamic quantities
- zero magnetization, decreasing/increasing local moments

Strongly interacting system at half-filling

n=1.0, U=2.0, W=2.0



- divergence in AF susceptibility
- negative sign of AF susceptibility
- no singularity in other thermodynamic quantities
- zero magnetization, increasing local moments

Conclusions and outlook

- 1. Wanted better understanding
- 2. Why susceptibilities/compressibility change signs
- 3. First order transitions, instabilities toward which phases
- 4. Any suggestions