

# Experimental security analysis of a four-photon private state

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**Konrad Banaszek**

**Rafał Demkowicz-Dobrzański**

**Michał Karpiński**

*Wydział Fizyki  
Uniwersytet Warszawski*

**Krzysztof Dobek**

*Wydział Fizyki, Uniwersytet Adama  
Mickiewicza w Poznaniu*



**Paweł Horodecki**

*Wydział Fizyki Technicznej i Matematyki  
Stosowanej, Politechnika Gdańsk*

**Karol Horodecki**

*Instytut Informatyki, Uniwersytet Gdańsk*



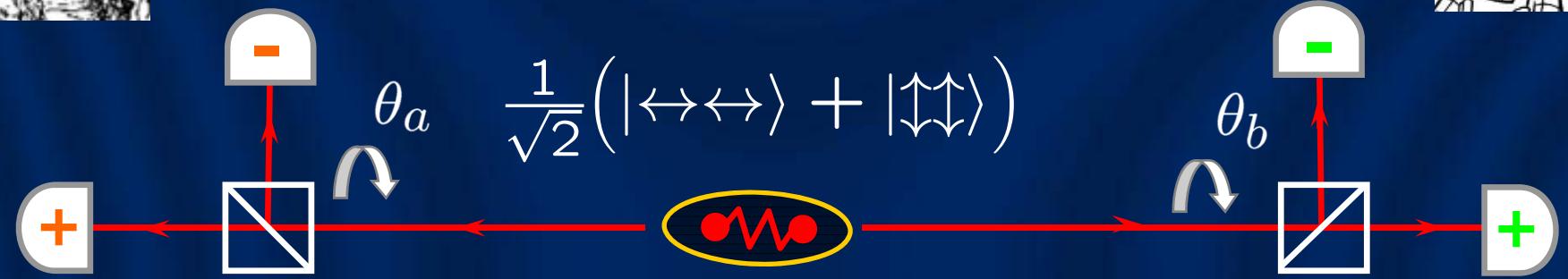
INNOVATIVE ECONOMY  
NATIONAL COHESION STRATEGY

  
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EUROPEAN REGIONAL  
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# Bell's inequalities



$$\mathbf{A}: \theta_a = 45^\circ$$

$$\mathbf{A}': \theta_a = 0^\circ$$

$$\mathbf{B}: \theta_b = 22.5^\circ$$

$$\mathbf{B}': \theta_b = 67.5^\circ$$

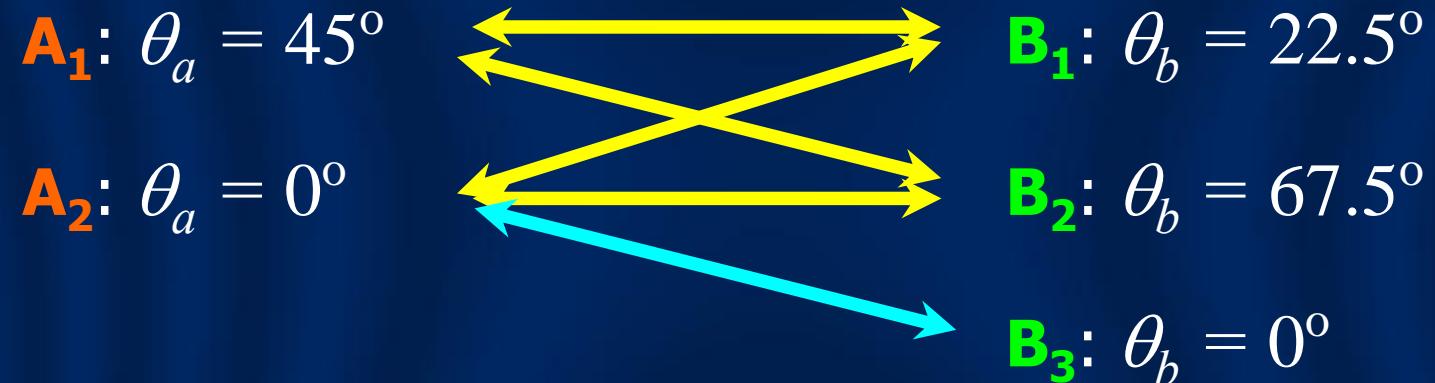
Clauser-Horne-Shimony-Holt inequality is violated!

$$\langle AB \rangle + \langle A'B \rangle + \langle AB' \rangle - \langle A'B' \rangle = 2\sqrt{2}$$

# Quantum cryptography

A. K. Ekert, Phys. Rev. Lett. **67**, 661 (1991)

For each photon pair Alice and Bob select randomly measurement bases...



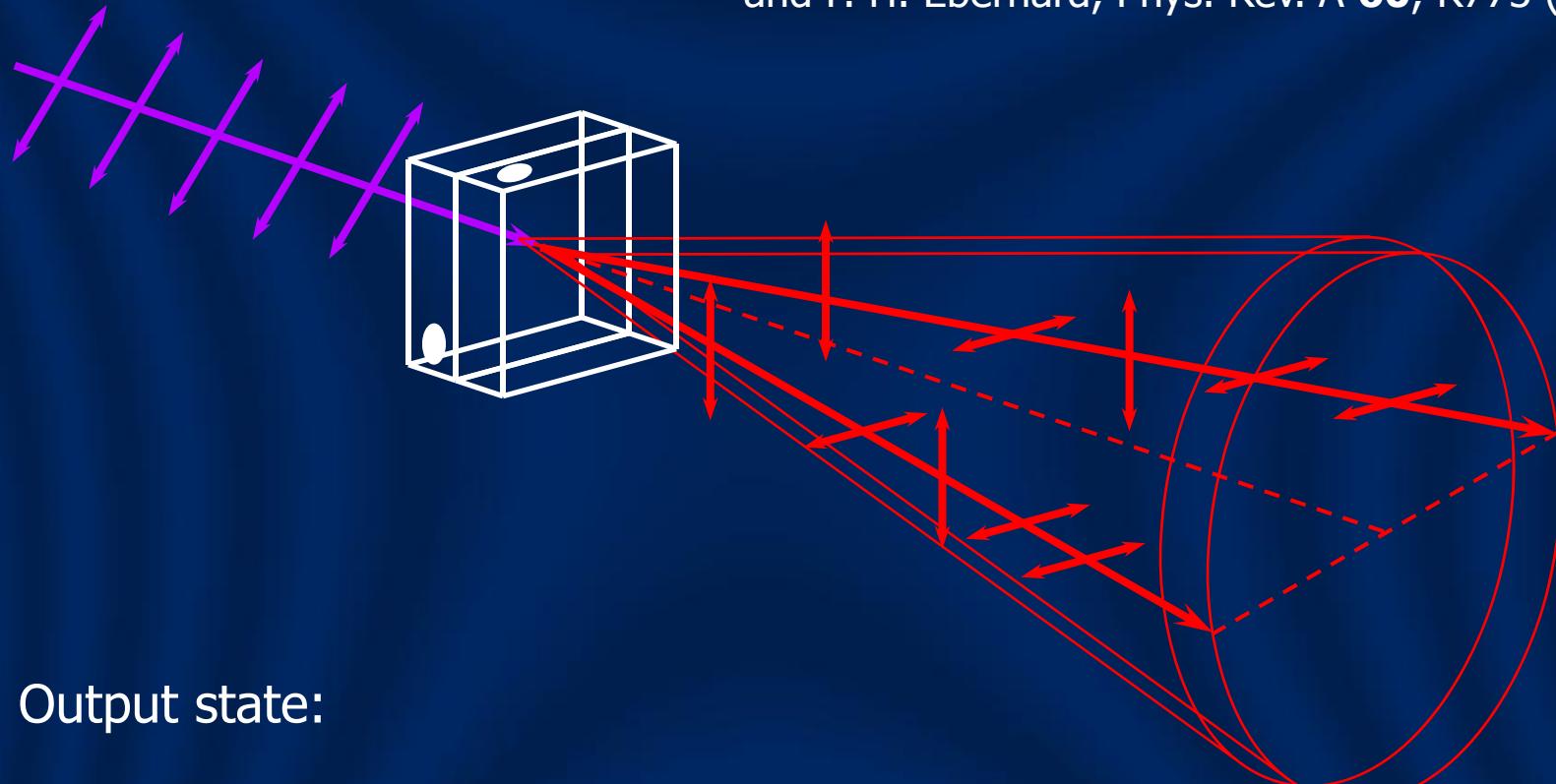
...and compare measurements over a public channel afterwards.

Perfect correlations → one-time pad

Security test

# Entangled photon pairs

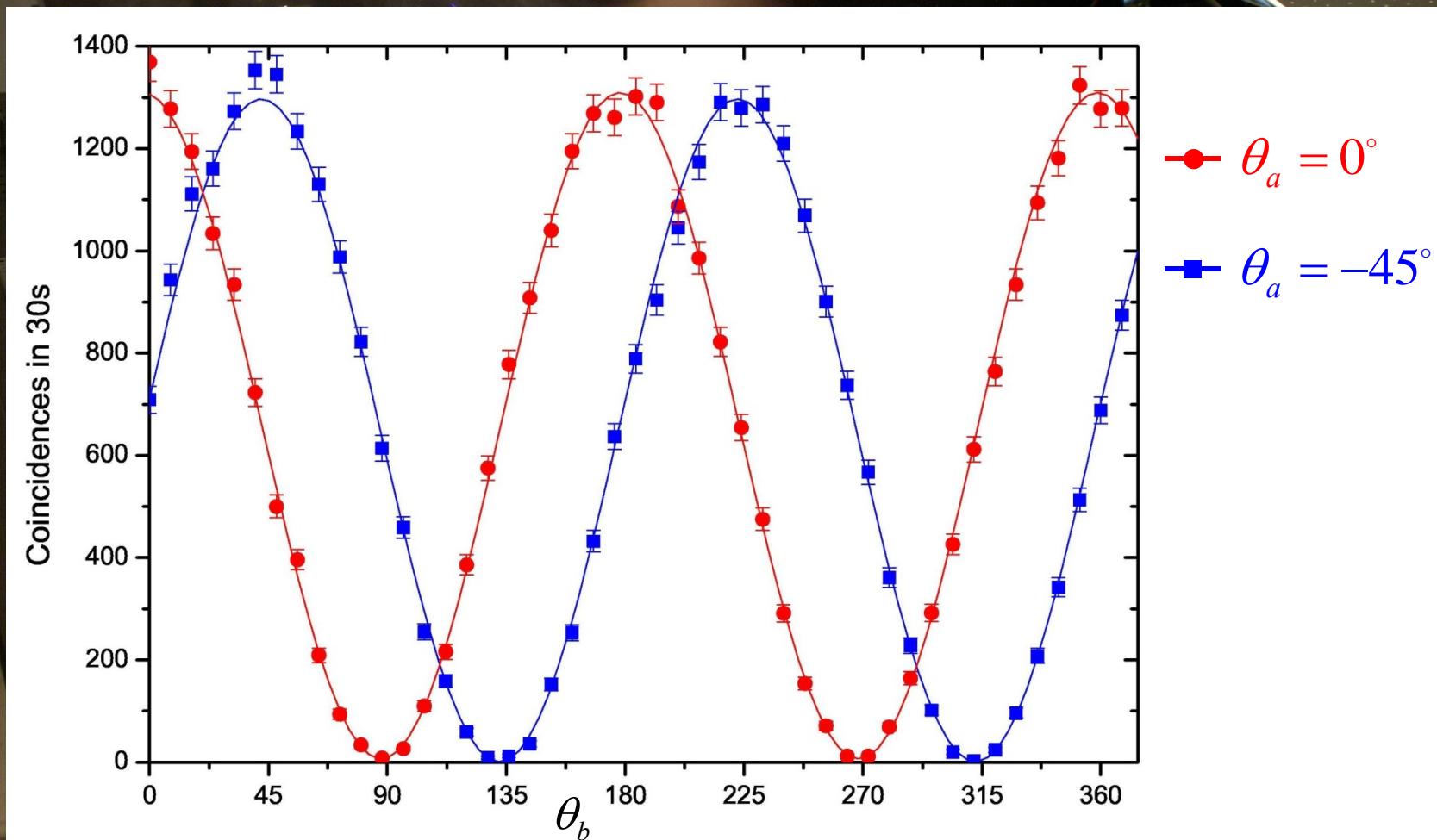
P. G. Kwiat, E. Waks, A. G. White, I. Appelbaum,  
and P. H. Eberhard, Phys. Rev. A **60**, R773 (1999)



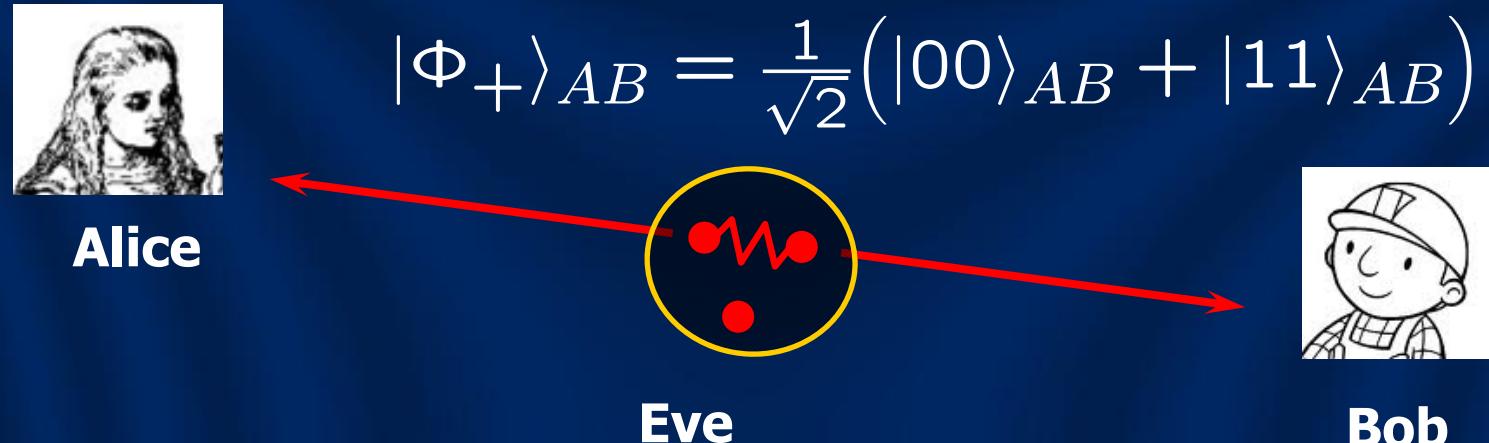
Output state:

$$|\Phi_+\rangle \propto |\leftrightarrow\leftrightarrow\rangle + |\uparrow\downarrow\rangle$$

# Correlation measurements



# Entanglement monogamy



- Even when the pair has been prepared by Eve...
- ...if Alice and Bob verify that the systems arrived in a maximally entangled pure state...
- ...measurement results will be known *only* to Alice and Bob.

# Statistical mixture

Define:  $|\Phi_{\pm}\rangle_{AB} = \frac{1}{\sqrt{2}}(|00\rangle_{AB} \pm |11\rangle_{AB})$

Equally weighted mixture:

$$\begin{aligned} & \frac{1}{2}|\Phi_+\rangle_{AB}\langle\Phi_+| + \frac{1}{2}|\Phi_-\rangle_{AB}\langle\Phi_-| \\ &= \frac{1}{2}(|00\rangle_{AB}\langle 00| + |11\rangle_{AB}\langle 11|) \\ &= \text{Tr}_E(|\Psi\rangle_{ABE}\langle\Psi|) \end{aligned}$$



$$|\Psi\rangle_{ABE} = \frac{1}{\sqrt{2}}(|000\rangle_{ABE} + |111\rangle_{ABE})$$

# Density matrix

Maximally entangled state

$$|\Phi_+\rangle_{AB}\langle\Phi_+|$$

$$\begin{array}{c} |00\rangle \\ |01\rangle \\ |10\rangle \\ |11\rangle \end{array} \quad \left( \begin{array}{cccc} \frac{1}{2} & \cdot & \cdot & \frac{1}{2} \\ \cdot & \cdot & \cdot & \cdot \\ \cdot & \cdot & \cdot & \cdot \\ \frac{1}{2} & \cdot & \cdot & \frac{1}{2} \end{array} \right)$$

Statistical mixture

$$\frac{1}{2}\left(|\Phi_+\rangle_{AB}\langle\Phi_+| + |\Phi_-\rangle_{AB}\langle\Phi_-|\right)$$

$$\left( \begin{array}{cccc} \frac{1}{2} & \cdot & \cdot & \cdot \\ \cdot & \cdot & \cdot & \cdot \\ \cdot & \cdot & \cdot & \cdot \\ \cdot & \cdot & \cdot & \frac{1}{2} \end{array} \right)$$

- Correlations between measurement outcomes in the key basis
- Security tested by the violation of Bell's inequalities  
(If trusting quantum theory, could be also tested by measurements in the  $(|0\rangle \pm |1\rangle)/\sqrt{2}$  basis.)

# Noisy entanglement

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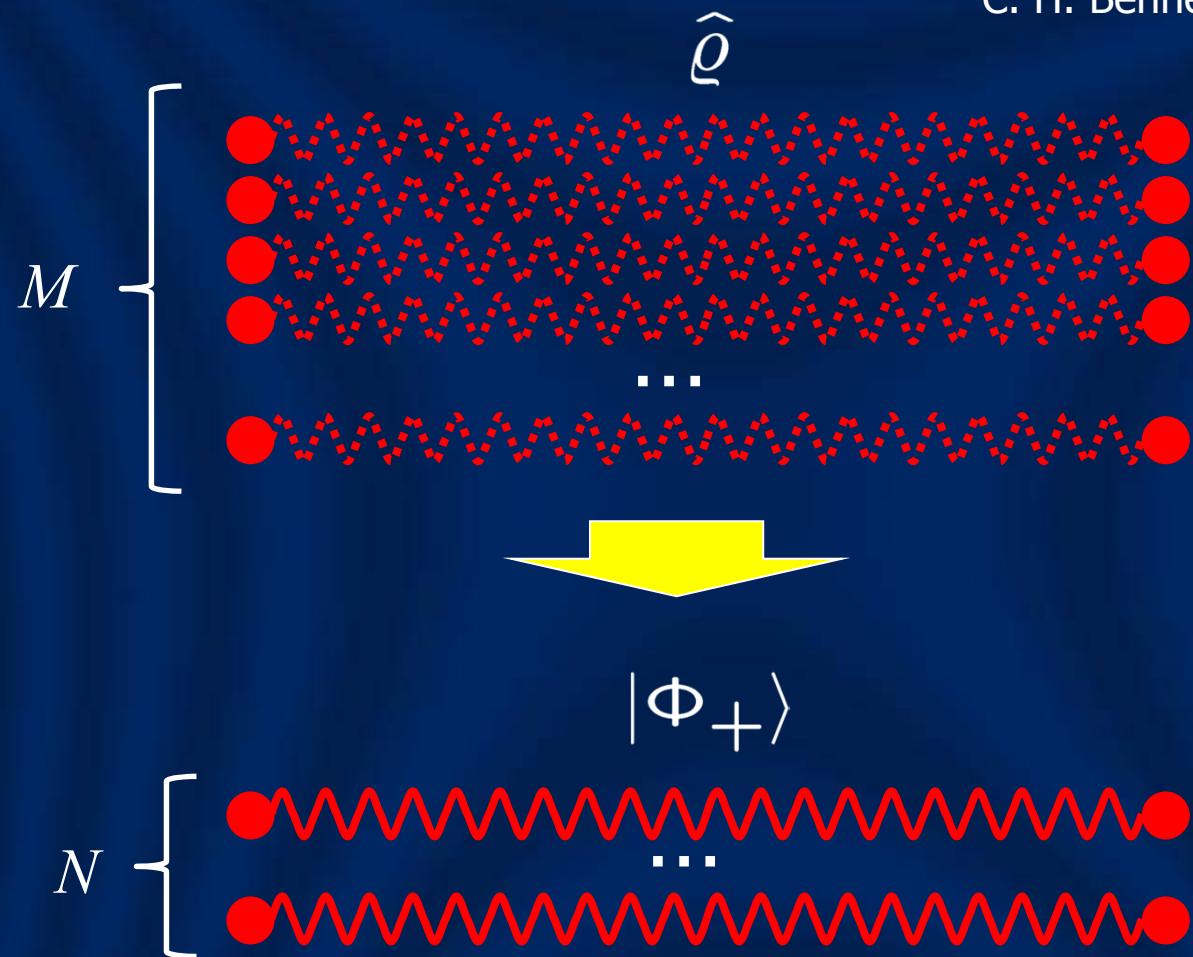
$$\frac{1}{4}|\Phi_+\rangle_{AB}\langle\Phi_+| + \frac{3}{4}|\Phi_-\rangle_{AB}\langle\Phi_-|$$

$$\begin{pmatrix} \frac{1}{2} & \cdot & \cdot & \frac{1}{4} \\ \cdot & \cdot & \cdot & \cdot \\ \cdot & \cdot & \cdot & \cdot \\ \frac{1}{4} & \cdot & \cdot & \frac{1}{2} \end{pmatrix}$$

*How much secure key can be extracted from a noisy state?*

# Distillation

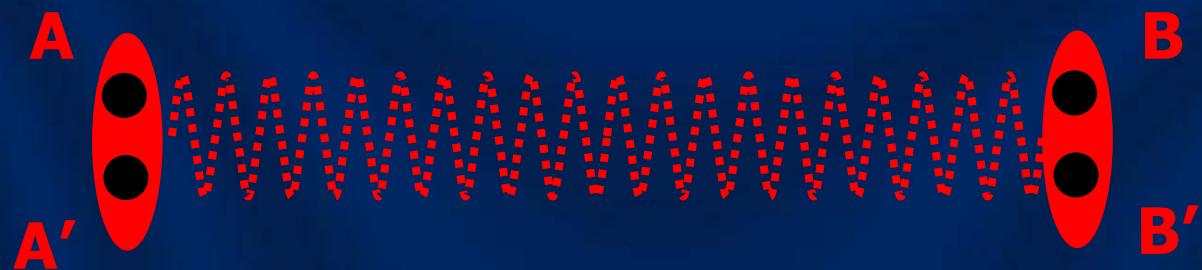
C. H. Bennett *et al.*, Phys. Rev. Lett.  
**76**, 722 (1996)



Distillable entanglement:  $E_D(\hat{\rho}) = \lim_{M \rightarrow \infty} \frac{N}{M}$

# Example I

$$|\Phi_{\pm}\rangle_{AB} = \frac{1}{\sqrt{2}}(|00\rangle_{AB} \pm |11\rangle_{AB})$$



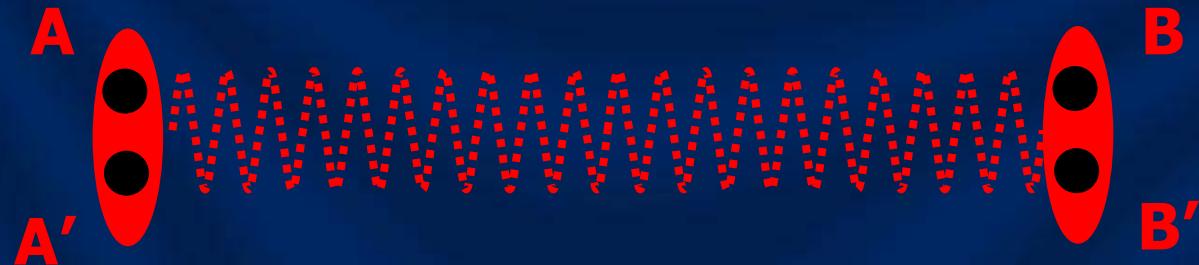
$$\hat{\varrho}_{AA'BB'} = \frac{3}{4}|\Phi_+\rangle_{AB}\langle\Phi_+| \otimes \hat{\varrho}_{A'B'}^{(+)} + \frac{1}{4}|\Phi_-\rangle_{AB}\langle\Phi_-| \otimes \hat{\varrho}_{A'B'}^{(-)}$$

Shield states:

$$\hat{\varrho}_{A'B'}^{(+)} = |00\rangle_{A'B'}\langle 00|, \quad \hat{\varrho}_{A'B'}^{(-)} = |11\rangle_{A'B'}\langle 11|$$

enable Alice and Bob to distinguish locally  $|\Phi_{\pm}\rangle_{AB}$  and generate the key using the standard strategy. Hence  $E_D = 1$

## Example II



What if

$$\hat{\varrho}_{A'B'}^{(+)} = \frac{1}{3}(\hat{\mathbb{I}} - |\Psi_-\rangle_{A'B'}\langle\Psi_-|)$$

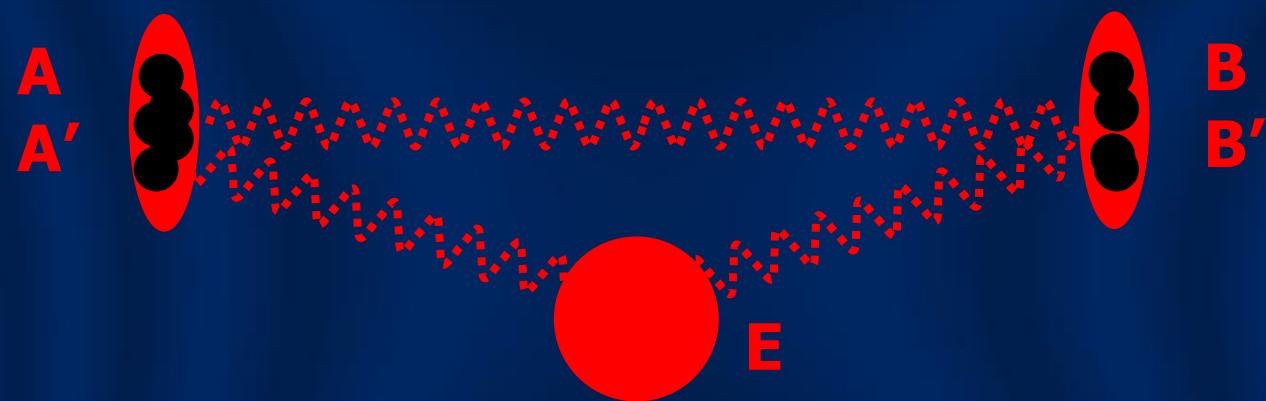
$$\hat{\varrho}_{A'B'}^{(-)} = |\Psi_-\rangle_{A'B'}\langle\Psi_-| \quad |\Psi_-\rangle = \frac{1}{\sqrt{2}}(|01\rangle - |10\rangle)$$

- States  $\hat{\varrho}_{A'B'}^{(+)}$  and  $\hat{\varrho}_{A'B'}^{(-)}$  cannot be discriminated unambiguously using local operations and classical communication.
- Distillable entanglement bounded by log-negativity:

$$E_D \leq \log_2 3 - 1 \approx 0.585$$

# Eavesdropping

K. Horodecki, M. Horodecki, P. Horodecki, and J. Oppenheim,  
Phys. Rev. Lett. **94**, 160502 (2005)



The worst case scenario: all the noise is controlled by Eve

$$\hat{\rho}_{AA'BB'} = \text{Tr}_E(|\Psi\rangle_{AA'BB'E}\langle\Psi|)$$

# Alice → Eve channel

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Alice measures an outcome  $a$  with a probability

$$p_a = \text{Tr}_{A'BB'E} \left( {}_A \langle a | \Psi \rangle \langle \Psi | a \rangle_A \right)$$



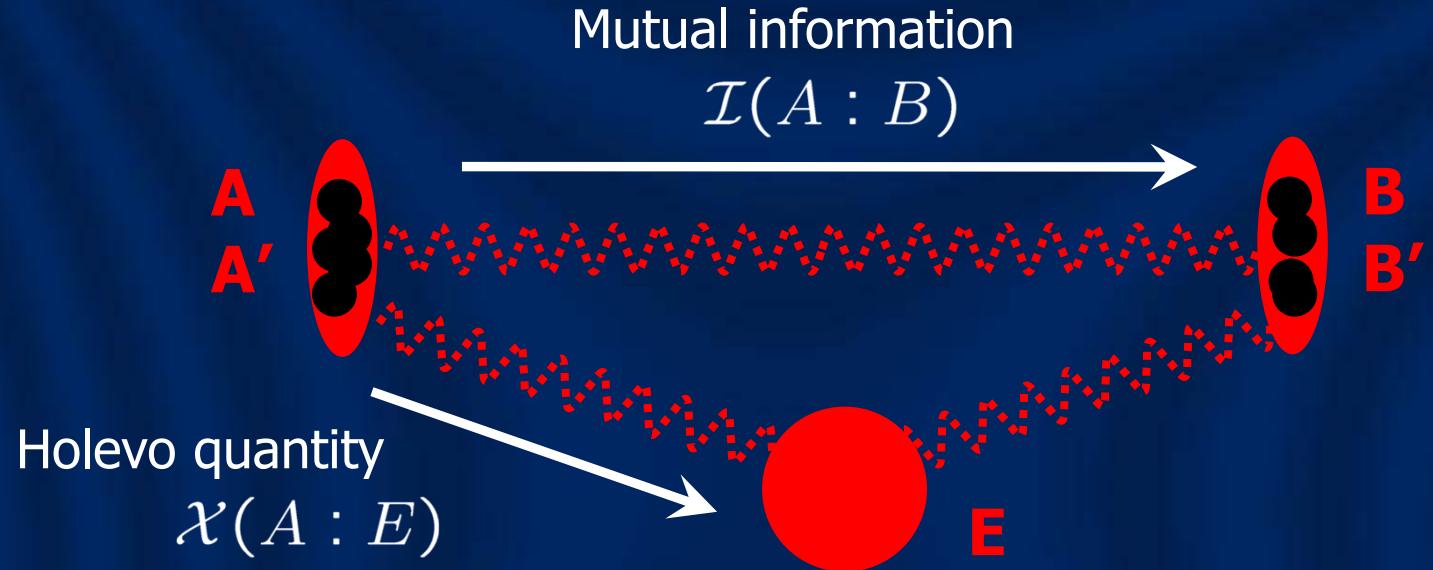
Eve infers  $a$  from the conditional state of her subsystem  $E$ :

$$\hat{\varrho}_E^{(a)} = \frac{1}{p_a} \text{Tr}_{AA'BB'} [|\Psi\rangle\langle\Psi| (|a\rangle_A\langle a| \otimes \hat{I}_{A'BB'})]$$

Holevo quantity:

$$\mathcal{X}(A : E) = S \left( \sum_a p_a \hat{\varrho}_E^{(a)} \right) - \sum_a p_a S(\hat{\varrho}_E^{(a)})$$

# Key rate



$$\text{Key rate} \quad K_D \geq \mathcal{I}(A : B) - \mathcal{X}(A : E)$$

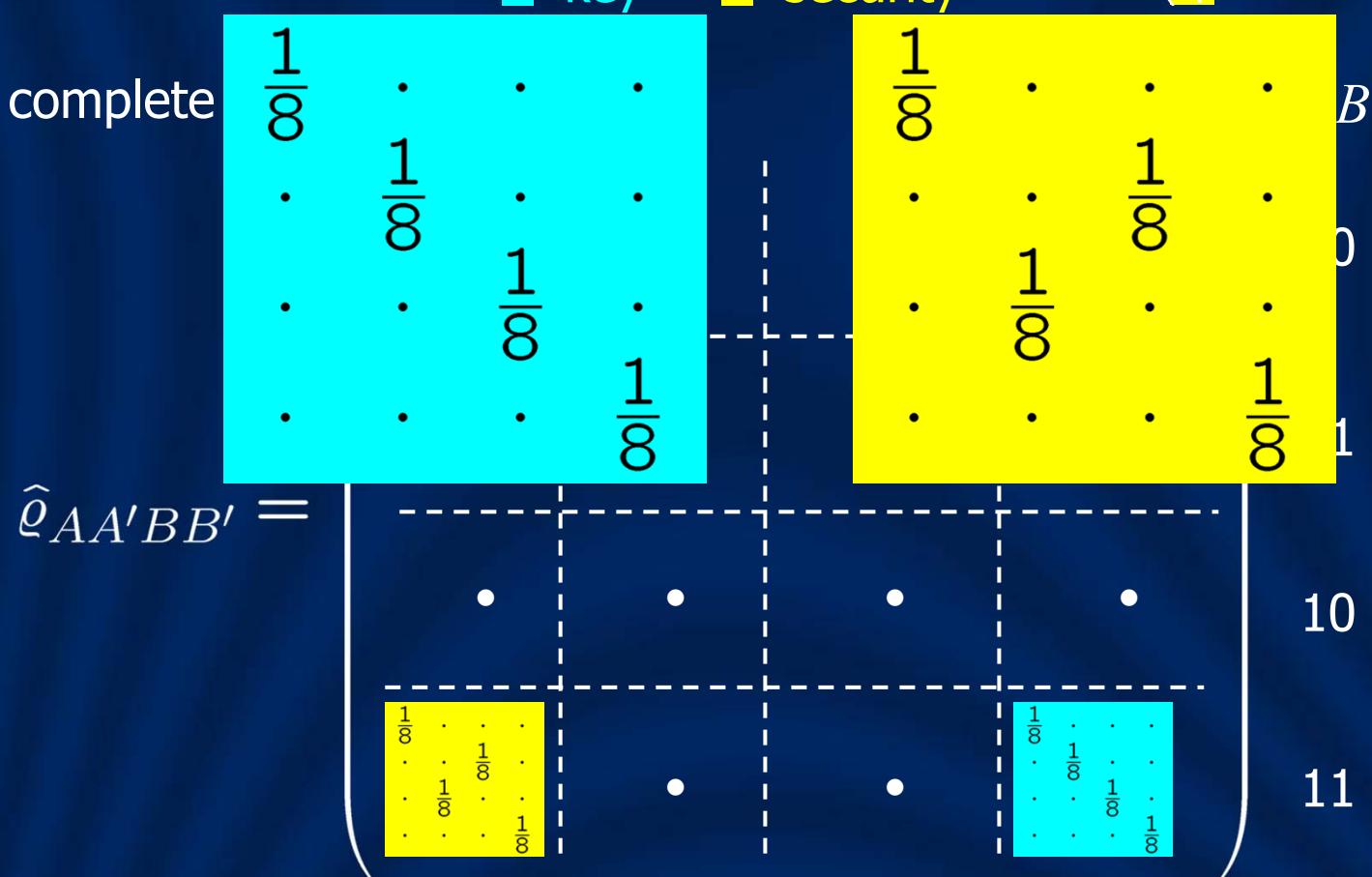
For Example II, Eve's subsystem contains no information about outcomes of Alice's measurement on her qubit, hence  $K_D = 1$ .

# Shield

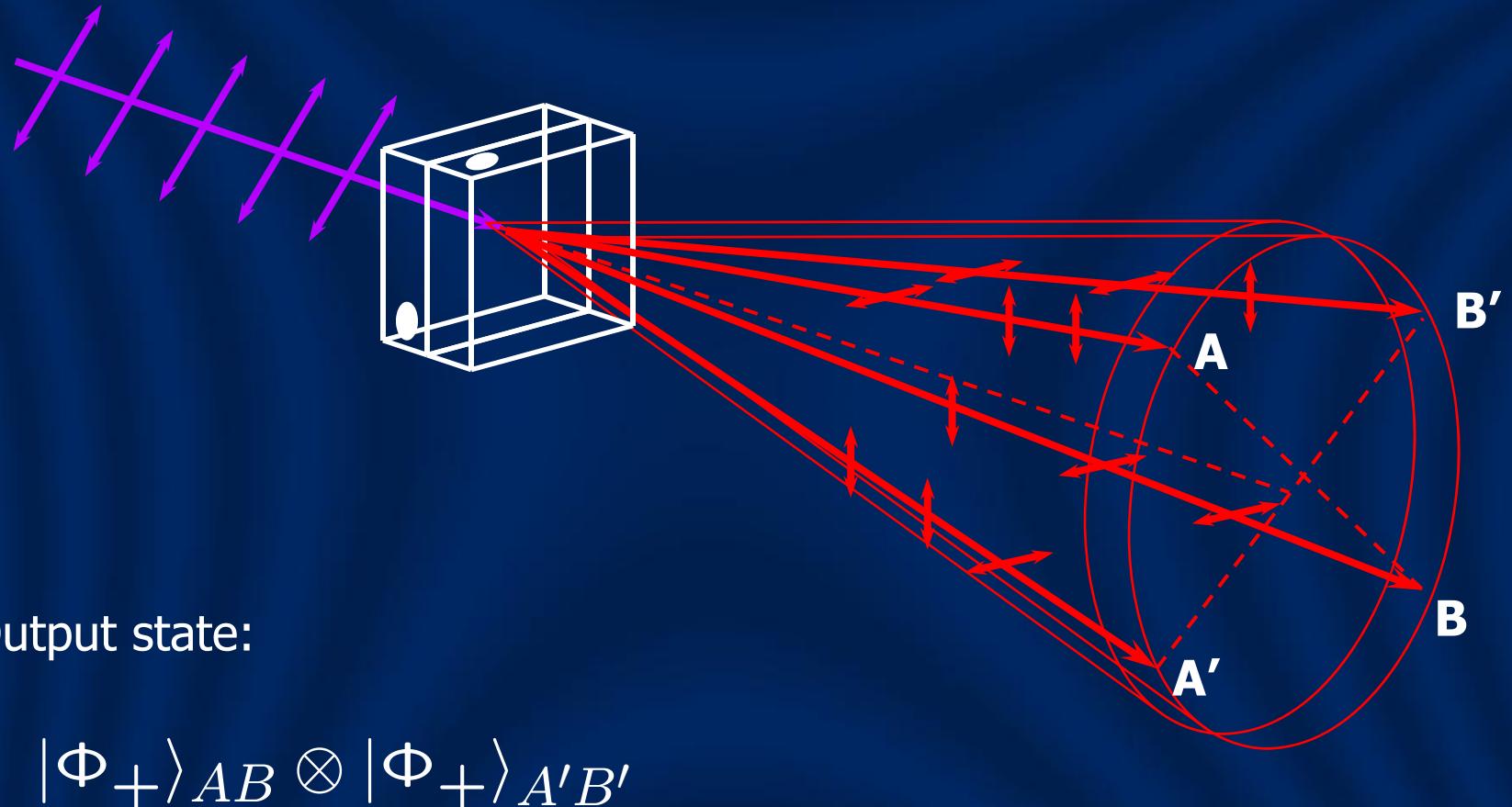
## Without the shield:

$$\text{Tr}_{A'B'}(\hat{\varrho}_{AA'BB'}) = \begin{pmatrix} \frac{1}{2} & & & & \frac{1}{4} \\ \vdots & \ddots & \ddots & & \vdots \\ \vdots & & \ddots & \ddots & \vdots \\ \frac{1}{4} & & & \ddots & \frac{1}{2} \end{pmatrix}$$

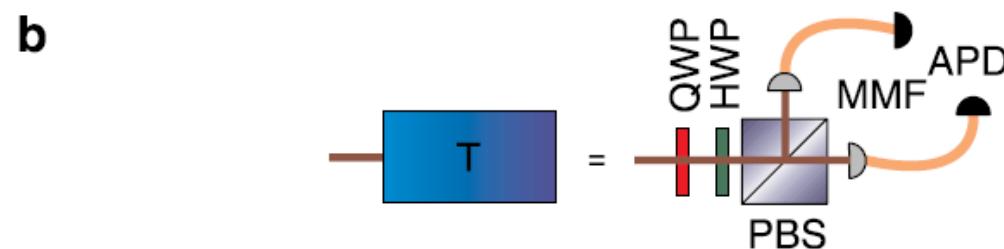
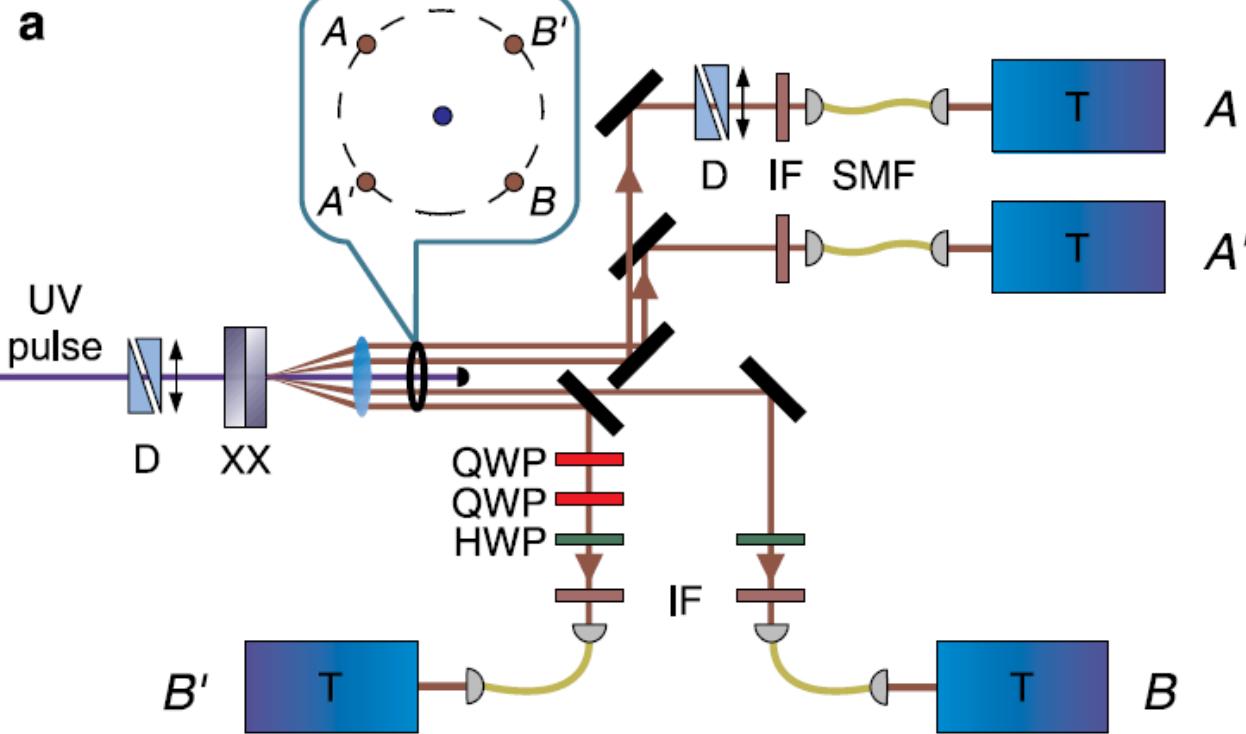
# The complete



# Double photon pairs



# Experimental setup



# Quantum state tomography

Projective qubit measurements:  $\hat{\sigma}_x, \hat{\sigma}_y, \hat{\sigma}_z$

Four-qubit POVM:

$$\hat{M}_i = |\pm_{i_A}\rangle \otimes |\pm_{i_{A'}}\rangle \otimes |\pm_{i_B}\rangle \otimes |\pm_{i_{B'}}\rangle$$

$3^4 = 81$  measurement bases

$3^4 \times 2^4 = 1296$  event types

Probability of an outcome  $i$ :

$$p(i|\hat{\rho}) = Tr(\hat{M}_i \hat{\rho})$$

$n_i$ : number of events  $i$



Density matrix  
estimate  $\hat{\rho}$

$$\sum_i n_i \approx 5 \times 10^5$$

# Maximum likelihood reconstruction

Probability of an outcome  $i$ :

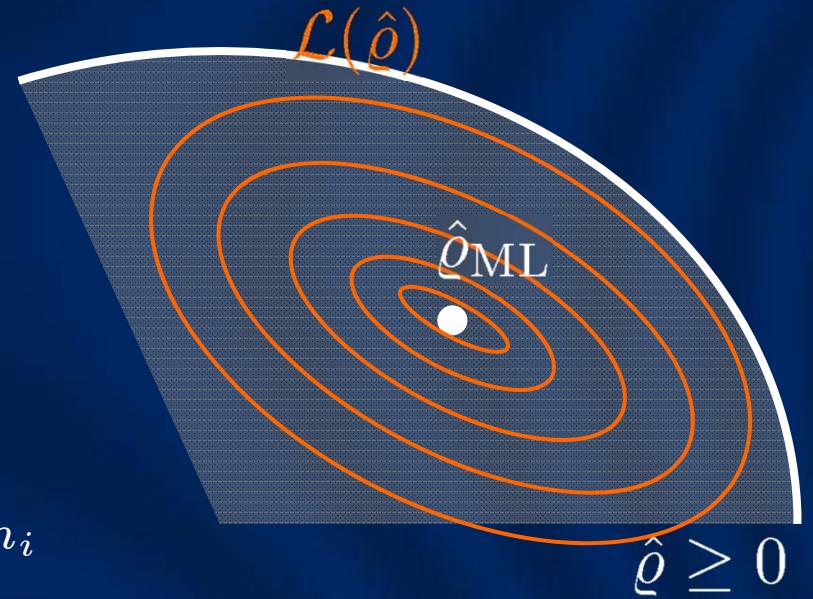
$$p(i|\hat{\rho}) = \text{Tr}(\hat{M}_i \hat{\rho})$$

$n_i$  – number of events  $i$

Likelihood function:

$$\mathcal{L}(\hat{\rho}) = p(\{n_i\}|\hat{\rho}) = \prod_i [p(i|\hat{\rho})]^{n_i}$$

Maximum-likelihood estimate  $\hat{\rho}_{\text{ML}}$   
maximizes  $\mathcal{L}(\hat{\rho})$



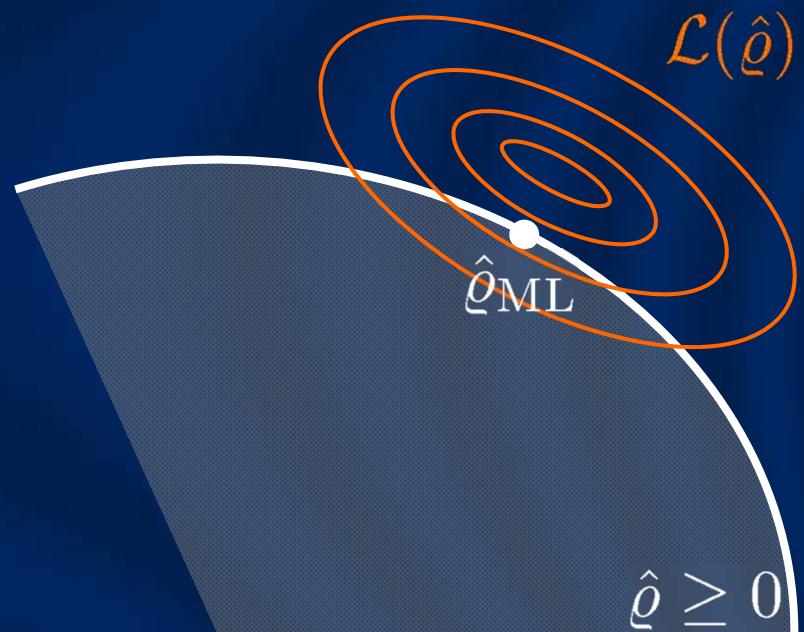
# ML: Parametrisation

K. Banaszek, G. M. D'Ariano, M. G. A. Paris, and M. F. Sacchi,  
Phys. Rev. A **61**, 010304(R) (1999)

Ensuring positivity:  $\hat{\rho} = \hat{T}^\dagger \hat{T}$ ,  $\hat{T} = \nabla$

Task: maximize  $\log \mathcal{L}(\hat{T}^\dagger \hat{T})$

with a constraint  $\text{Tr}(\hat{T}^\dagger \hat{T}) = 1$



**PRO:** - Guaranteed positivity

**CON:** - Impractical in higher dimensions (>6 qubits)  
- Underestimates errors, difficult to include uncertainty of the measuring device (Monte Carlo simulations)  
- Biased towards low-rank matrices for undersampled data

# Bayesian approach

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K. Audenaert and S. Scheel, New J. Phys. **11**, 023028 (2009)

A priori distribution  $p(\hat{\varrho})$

A posteriori:  $p(\hat{\varrho}|\{n_i\}) \propto p(\{n_i\}|\hat{\varrho})p(\hat{\varrho})$

Estimate:  $\hat{\varrho}_{\text{Bayes}} = \int d\varrho \hat{\varrho} p(\hat{\varrho}|\{n_i\})$

- Gaussian approximation
- Truncated to positive definite density operators

**PRO:**

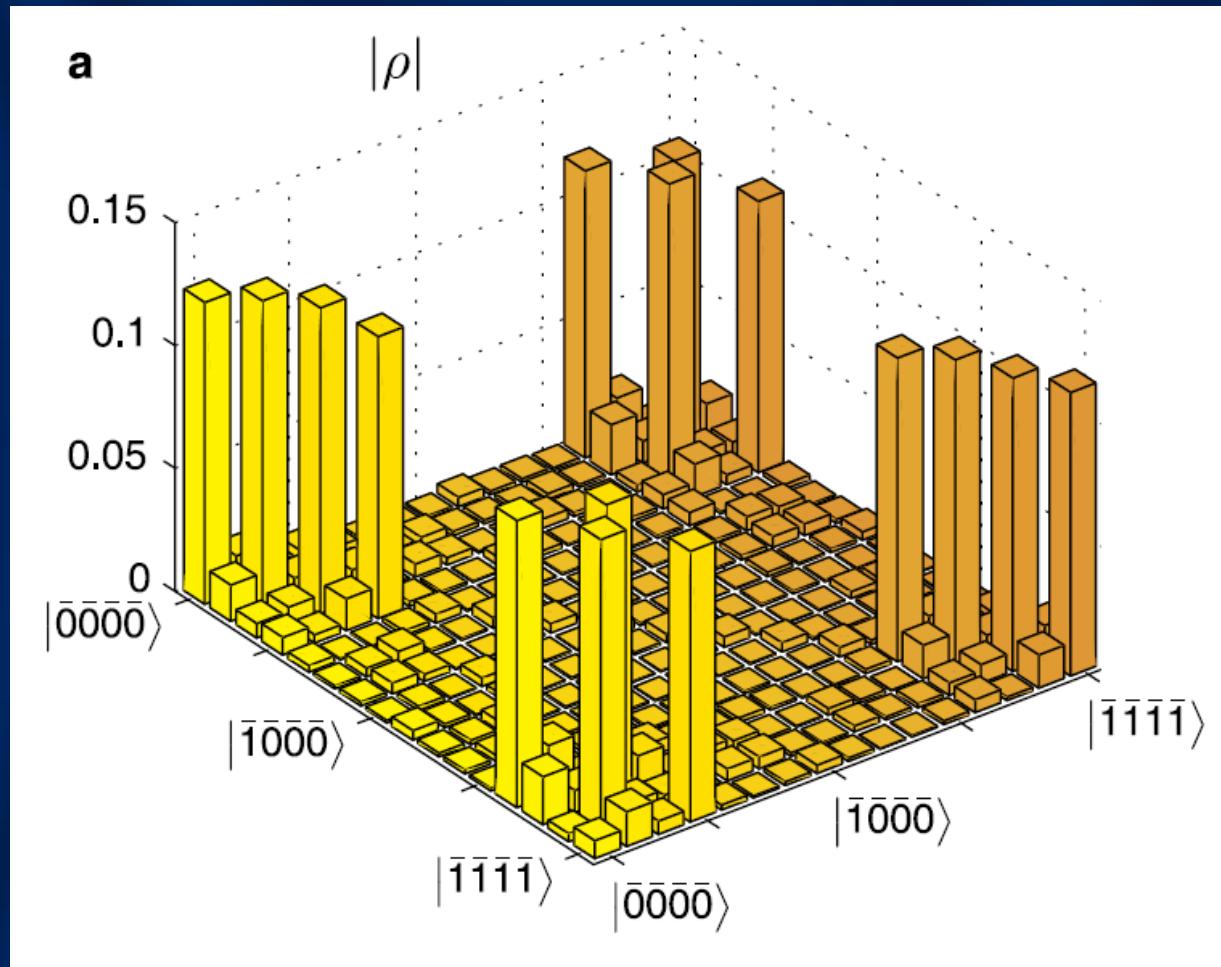
- Clear statistical interpretation
- Provides uncertainty
- No numerical optimisation

**CON:**

- Difficult to normalise probability distribution
- A priori distribution not well defined

# State reconstruction

K. Dobek. M. Karpiński, R. Demkowicz-Dobrzański, K. Banaszek,  
and P. Horodecki, Phys. Rev. Lett. **106**, 030501 (2011)



# Privacy characterisation

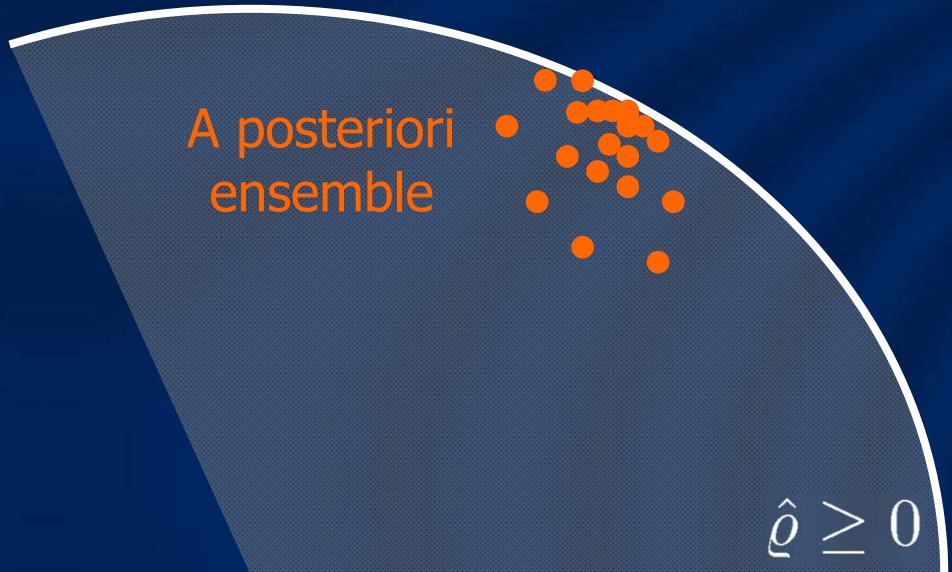
Warning:

$$\langle f(\hat{\rho}) \rangle \neq f(\langle \hat{\rho} \rangle)$$

More conservative

Distillable entanglement:  $E_D \leq 0.581(4)$

Key (cqq scenario):  $K \geq 0.690(7)$



# Distillation protocol

$$\begin{aligned}\hat{\varrho}_{AA'BB'} = & \frac{1}{4}|\Phi_+\rangle_{AB}\langle\Phi_+| \otimes \left(\hat{\mathbb{I}}_{A'B'} - |\Psi_-\rangle_{A'B'}\langle\Psi_-|\right) \\ & + \frac{1}{4}|\Phi_-\rangle_{AB}\langle\Phi_-| \otimes |\Psi_-\rangle_{A'B'}\langle\Psi_-|\end{aligned}$$

Measure qubits  $A'B'$  in the same basis.



Identical outcomes

$$\hat{\varrho}_{AB} = |\Phi_+\rangle\langle\Phi_+|$$

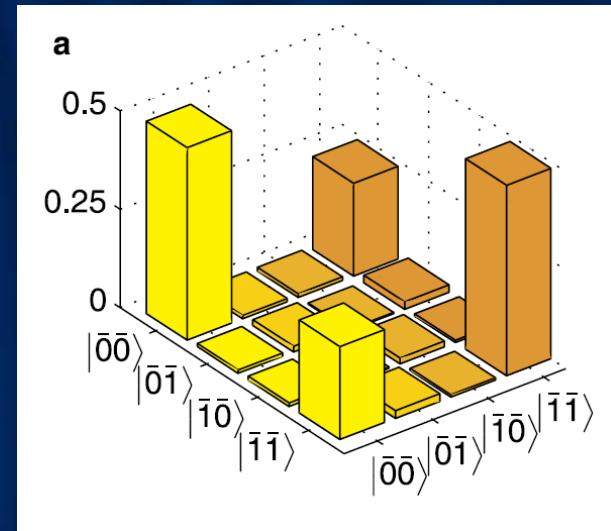
Opposite outcomes

$$\hat{\varrho}_{AB} = \frac{1}{2}(|00\rangle\langle 00| + |11\rangle\langle 11|)$$

# Single-copy distillation

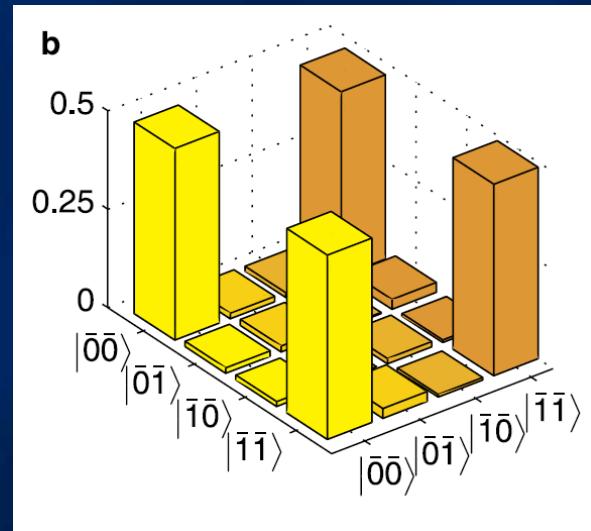
Reduced density matrix  $\hat{\rho}_{AB}$

average

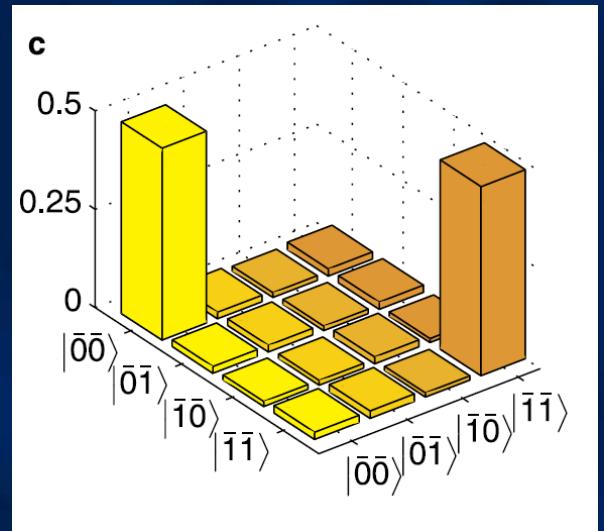


$K = 0$

conditional  
*identical outcomes A'B'*    *opposite outcomes A'B'*



$K = 0.354(5)$



# Cryptographic key

## Optimal strategy

Raw key: 3716 bits



Error correction

2726 bits



Privacy amplification

Secure key: 2164 bits

## Distillation-based approach

Raw key: 1859 bits



$\approx 1300$  bits



Secure key: < 650 bits

# Witnessing privacy

$$\hat{\varrho}_{AA'BB'} = \begin{pmatrix} \hat{A}_{00,00} & \hat{A}_{00,01} & \hat{A}_{00,10} & \hat{A}_{00,11} \\ \hat{A}_{01,00} & \hat{A}_{01,01} & \hat{A}_{01,10} & \hat{A}_{01,11} \\ \hat{A}_{10,00} & \hat{A}_{10,01} & \hat{A}_{10,10} & \hat{A}_{10,11} \\ \hat{A}_{11,00} & \hat{A}_{11,01} & \hat{A}_{11,10} & \hat{A}_{11,11} \end{pmatrix}$$

$$K(\hat{\varrho}_{AA'BB'}) \geq K(\hat{\sigma}_{AB})$$

where

$$\hat{\sigma}_{AB} = \frac{1}{2} \begin{pmatrix} p_+ & \cdot & \cdot & c_+ \\ \cdot & p_- & c_- & \cdot \\ \cdot & c_- & p_- & \cdot \\ c_+ & \cdot & \cdot & p_+ \end{pmatrix} \quad \begin{aligned} p_+ &= \|\hat{A}_{00,00} + \hat{A}_{11,11}\| \\ c_+ &= \|\hat{A}_{00,11} + \hat{A}_{11,00}\| \end{aligned}$$

# Single witness

K. Banaszek, K. Horodecki, and P. Horodecki,  
Phys. Rev. A **85**, 012330 (2012)

Suppose we have measured

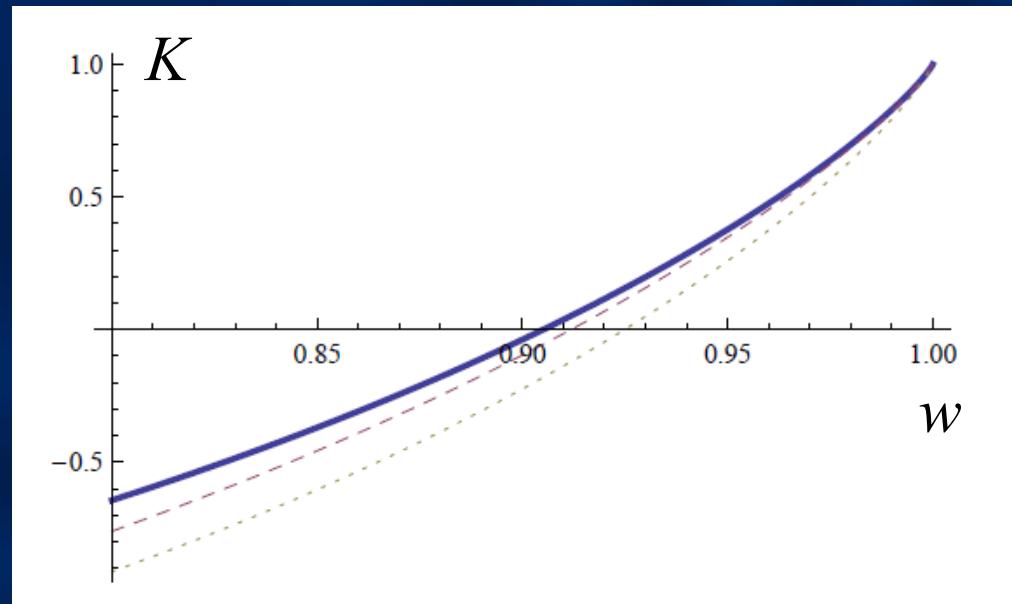
$$w = \left| \langle (\sigma_A^x \otimes \sigma_B^x - \sigma_A^y \otimes \sigma_B^y) \otimes \hat{U}_{A'B'} \rangle \right|$$

where  $\hat{U}^\dagger \hat{U} \leq \hat{I}$

We have:

$$p_+ \geq c_+ \geq w$$

Take the worst-case scenario for  $p_-, c_-$



# Two observables

K. Banaszek, K. Horodecki, and P. Horodecki,  
Phys. Rev. A **85**, 012330 (2012)

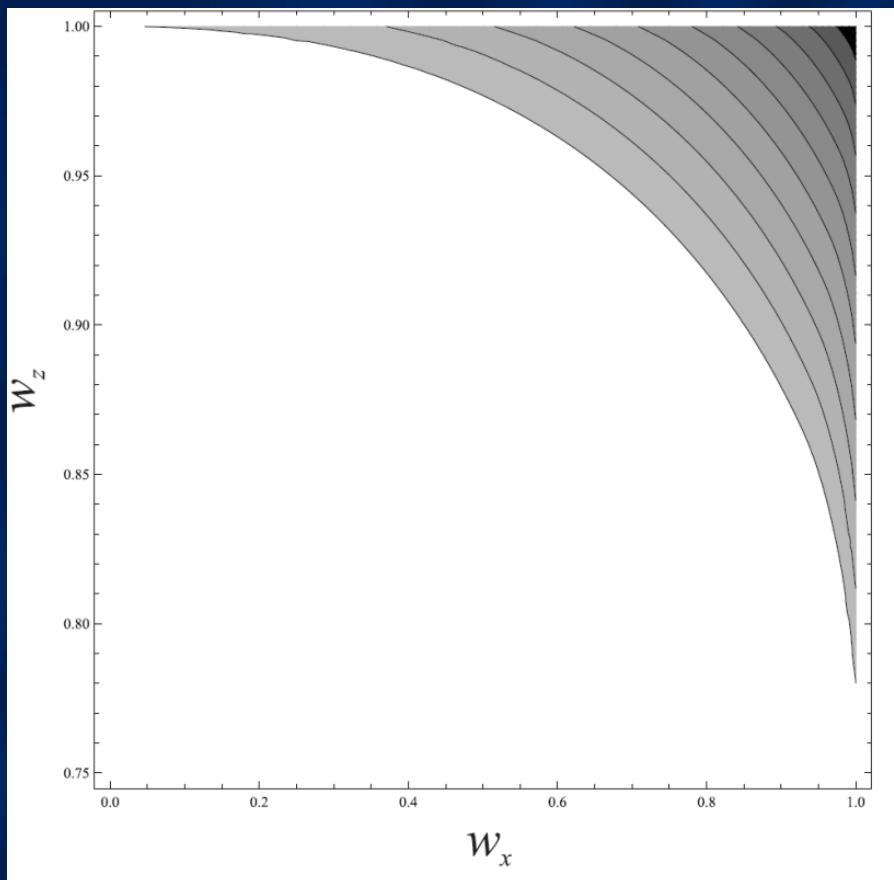
$$w_x = \left| \langle \sigma_A^x \otimes \sigma_B^x \otimes \hat{U}_{A'B'} \rangle \right| \quad w_z = \left| \langle \sigma_A^z \otimes \sigma_B^z \otimes \hat{I}_{A'B'} \rangle \right|$$

We have:

$$c_+ + c_- \geq w_x$$

$$p_{\pm} = \frac{1}{2}(1 \pm w_z)$$

$$c_- \leq p_-$$



# Conclusions

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- Experimental demonstration of the separation between distillable entanglement and cryptographic key contents
- Practical comparison of quantum state reconstruction methods for a noisy multiqubit state
- Full privacy analysis based on the reconstructed state
- Evaluation of highly non-linear information theoretic quantities
- Implementation of a simple entanglement distillation protocol
- Witnessing privacy with few observables
- Multiple degrees of freedom?