AFFINE POISSON STRUCTURES IN ANALYTICAL MECHANICS

Paweł Urbański

Department of Mathematical Methods in Physics, University of Warsaw Hoża 74, 00-682 Warszawa Poland

Abstract

If the space-time is a product of the space and the time the Poisson structure on the phase bundle is used to describe dynamics of mechanical systems. Further it is shown that if the space-time is a fibration over the time, then the Poisson structure has to be replaced by an affine Poisson structure.

1. TIME-DEPENDENT SYSTEMS

1.1. Time Independent Systems

In order to define a time-independent system the space-time has to be the product of space and time represented by the real line \mathbb{R} . For a time-independent system with the configuration manifold Q the infinitesimal dynamics is a submanifold D of TT^*Q . In particular cases D is the image of a vector field. The cotangent bundle TT^*Q with the canonical 2-form ω_Q is a symplectic manifold. The tangent bundle TT^*Q of the cotangent bundle with the tangent 2-form $d_{\mathsf{T}}\omega_Q$ is a symplectic manifold as well^{4, 5}. We say that the system is *Lagrangian* if the dynamics D is a Lagrange submanifold of $(\mathsf{TT}^*Q, d_{\mathsf{T}}\omega_Q)$.

Let us denote by τ_Q the canonical projection $\tau_Q: \mathsf{T}Q \to Q$ and by π_Q the canonical projection $\pi_Q: \mathsf{T}^*Q \to Q$. There are three, fundamental for the analytical mechanics, isomorphisms of vector bundles:

$$\kappa_Q : (\tau_{\mathsf{T}Q} : \mathsf{T}\mathsf{T}Q \to \mathsf{T}Q) \longrightarrow (\mathsf{T}\tau_Q : \mathsf{T}\mathsf{T}Q \to \mathsf{T}Q) \tag{1.1}$$

$$\alpha_Q : (\mathsf{T}\pi_Q : \mathsf{T}\mathsf{T}^*Q \to \mathsf{T}Q) \longrightarrow (\pi_{\mathsf{T}Q} : \mathsf{T}^*\mathsf{T}Q \to \mathsf{T}Q) \tag{1.2}$$

$$\beta_Q : (\mathsf{T}\pi_Q : \mathsf{T}\mathsf{T}^*Q \to \mathsf{T}Q) \longrightarrow (\pi_{\mathsf{T}^*Q}\mathsf{T}^*\mathsf{T}^*Q \to \mathsf{T}^*Q)$$
(1.3)

The mapping α_Q is also a symplectomorfism of $(\mathsf{TT}^*Q, \mathsf{T}\pi_Q)$ and $(\mathsf{T}^*\mathsf{T}Q, \pi_{\mathsf{T}Q})$. The mapping β_Q is a symplectomorphism of $(\mathsf{TT}^*Q, \mathsf{T}\pi_Q)$ and $(\mathsf{T}^*\mathsf{T}^*Q, \pi_{\mathsf{T}^*Q})$.

Let the dynamics D of a system be a Lagrangian submanifold of $(\mathsf{TT}^*Q, \mathsf{T}\pi_Q)$. It follows that $\alpha_Q(D)$ and $\beta_Q(D)$ are Lagrangian submanifolds of $(\mathsf{T}^*\mathsf{T}Q, \pi_{\mathsf{T}Q})$ and $(\mathsf{T}^*\mathsf{T}^*Q, \pi_{\mathsf{T}^*Q})$ respectively. By a theorem of Hörmander $\alpha_Q(D)$ and $\beta_Q(D)$ can be generated (at least locally) by a function (or a Morse family of functions) on a submanifold of $\mathsf{T}Q$ or T^*Q respectively. The generating function on $\mathsf{T}Q$ (if it exists) is called the Lagrangian of the system. The generating function on T^*Q is called the Hamiltonian generating function of the system. In the following we shall consider systems with the dynamics generated by a Lagrangian defined on $\mathsf{T}Q$.

1.2. Time-Dependent Systems. Inhomogeneous formulation.

Let us assume that, as before, the space-time is the product of the space and the time. Let Q be the manifold of space configurations of a system. $M = Q \times \mathbb{R}$ is the manifold of the space-time configurations of the system. Let $\zeta: M \to Q$ be the canonical projection. An infinitesimal configuration is a vector $v \in \mathsf{T}M$ such that $\mathsf{T}\zeta(v) = (\zeta(\tau_M v), \partial_t)$. $\mathsf{T}_1 M$ will denote the bundle of infinitesimal configurations. The phase bundle is the product $P = \mathsf{T}^* Q \times \mathbb{R}$. Let $\eta: P \to \mathbb{R}$ be the canonical projection. For each $t \in \mathbb{R}$ the fiber $P_t = \eta^{-1}(t)$ is a symplectic manifold. An infinitesimal state is a vector $w \in \mathsf{T}P$ such that $\mathsf{T}\eta(w) = (\eta(\tau_P w), \partial_t)$. We denote by T_1P the bundle of infinitesimal states. The dynamics of the system is a submanifold D of T_1P . Let D_t denote a subset of D defined by $D_t = \{D \ni w; \eta(\tau_P w) = t \in \mathbb{R}\}$. Since $P = \mathsf{T}^* Q \times \mathbb{R}$, we have also $\mathsf{T}P = \mathsf{T}\mathsf{T}^*Q \times \mathsf{T}\mathbb{R}$ and $D_t = \overline{D}_t \times (t, \partial_t)$. We say that the system is Lagrangian if \overline{D}_t is a Lagrangian submanifold of $(\mathsf{TT}^*Q, d_\mathsf{T}\omega_0)$ for each $t \in \mathbb{R}$. It follows (see previous section) that \overline{D}_t can be generated by the Lagrangian L_t on $\mathsf{T}Q$ or by the Hamiltonian generating function H_t on T^*Q . Families of functions L_t, H_t define functions L, H on $\mathsf{T}Q \times \mathbb{R}$ and $\mathsf{T}^*Q \times \mathbb{R}$ respectively. A procedure of generating the component \overline{D} of the dynamics can be formulated in terms of the canonical Poisson structure on P. This formulation is equivalent to the described one.

1.3. Time-Dependent Systems. Homogeneous formulation.

In the homogeneous formulation the configuration manifold is the manifold M. Infinitesimal configurations are vectors tangent to M. TM is the manifold of infinitesimal configurations. If the system is Lagrangian, i. e., the dynamics is generated by the Lagrangian L on $TQ \times \mathbb{R}$, we define a function $\hat{L}: TM \to \mathbb{R}$ by the formula

$$L(v) = sL_t(\overline{v})$$

where \overline{v} is the component of v in $\mathsf{T}Q$ and $\mathsf{T}\zeta(v) = (t, s\partial_t)$.

The function \hat{L} generates a Lagrangian submanifold of $\mathsf{T}^*\mathsf{T}M$ and, consequently, of TT^*M (Section 1.2). We denote by \widehat{D} the generated by \widehat{L} submanifold of TT^*M . It can also be generated also by a Hamiltonian.

Proposition 1 The manifold \widehat{D} is generated by a function equal to zero and defined on a submanifold $C \subset \mathsf{T}^*M$

$$C = \{\mathsf{T}^*M \ni (p, t, \epsilon); -\epsilon = H_t(p)\}$$

We can also get D_t from \widehat{D} directly. It is easy to verify that TT^*Q is canonicaly identified with the reduction of TT^*M with respect to the coisotropic submanifold $\mathsf{T}K_t$ where $K_t \subset \mathsf{T}^*M$ defined by

$$K_t = \{\mathsf{T}^* M \ni (p, t', \epsilon); t' = t\}.$$

Let us denote this reduction by $\mathsf{T}\varrho_t$.

Proposition 2

$$D_t = \mathsf{T}\varrho_t(D)$$

2. AFFINE SPACES AND AFFINE BUNDLES

2.1. Principal Affine Spaces

An affine space is a triple (A, V, α) , where A is a set, V is a real vector space of finite dimension and α is a mapping $\alpha: A \times A \to V$ such that

- 1. $\alpha(a_3, a_2) + \alpha(a_2, a_1) + \alpha(a_1, a_3) = 0;$
- 2. the mapping $\alpha(\cdot, a): A \to V$ is bijective for each $a \in A$.

We will write for brevity $a_2 - a_1$ instead of $\alpha(a_2, a_1)$. We will denote by a + v the unique point $a' \in A$ such that a' - a = v.

We consider quadruples (A, V, α, v_0) , where (A, V, α) is an affine space and v_0 is a distinguished nonzero vector in the model space V of the affine space (A, V, α) . Such objects will be called *principal affine spaces*. A principal affine mapping from (A, V, α, v_0) to (B, W, β, w_0) is an affine mapping φ from (A, V, α) to (B, W, β) such that $\overline{\varphi}(v_0) = w_0$ ($\overline{\varphi}$ is the linear part of φ . A principal affine space (A, V, α, v_0) can be considered as a principal bundle with the structural group \mathbb{R} and the action

$$(r,a) \mapsto a + sv_0.$$

The category of principal affine spaces has a distiguished object $I = (\mathbb{R}, \mathbb{R}, -, 1)$. The *affine dual* to a principal affine space (A, V, α, v_0) is a principal affine space $(A^{\#}, V^{\#}, \alpha^{\#}, f_0)$ where $A^{\#}$ is the space of all principal affine mappings from (A, V, α, v_0) to $I, V^{\#}$ is the vector space of affine functions on the quotient vector space $V/_{\{v_0\}}, \alpha^{\#}(\varphi - \varphi') = \varphi - \varphi'$ and f_0 is the equal 1 constant function.

2.2. Affine Bundles

Let $\xi: E \to N$ be a vector fibration. An *affine fibration modelled on* ξ is a differential fibration $\eta: A \to N$ and a differentiable mapping $\rho: A \times_N A \to E$ such that

- 1. $\xi \circ \rho = \eta \times_N \eta$,
- 2. $\rho(a_3, a_2) + \rho(a_2, a_1) = \rho(a_3, a_1)$ for each triple $(a_3, a_2, a_1) \in A \times_N A \times_N A$,
- 3. for each local section $\sigma: U \to A$ of η , the mapping $\rho_{\sigma}: \eta^{-1}(U) \to \xi^{-1}(U)$ defined by

$$\rho_{\sigma}(a) = \rho(a, \sigma(\eta(a)))$$

is a diffeomorphism.

A *principal affine fibration* is an affine fibration with e nowhere vanishing section of the model vector fibration. It follows that a fiber of a principal affine fibration is a principal affine space. The affine dual to a principal affine fibration we define in the obvious way.

An affine fibration modelled on the trivial vector fibration $pr_N: N \times \mathbb{R} \to N$ is usually interpreted as a principal fibration with structure group **R**. We denote by **I** the trivial principal fibration $(pr_1: \{1\} \times \mathbb{R} \to \{1\})$.

Let $\mathbf{Z} = (\zeta : Z \to N, \rho : Z \times_N Z \to N \times \mathbb{R})$ be an affine fibration modelled on the trivial fibration $pr_N : N \times \mathbb{R} \to N$. We define an equivalence relation in the set of all pairs (m, σ) , where m is a point in N and σ is a section of ζ . Two pairs (m, σ) and (m', σ') are equivalent if m' = m and $d(\sigma' - \sigma)(m) = 0$. We have identified the section $\sigma' - \sigma$ of pr_N with a function on N for the purpose of evaluating the differential $d(\sigma' - \sigma)(m)$. We denote by $\mathsf{P}Z$ the set of equivalence classes. The class of (m, σ) will be denoted by $d\sigma(m)$ and will be called the *differential* of σ at m. We define a mapping $\mathsf{P}\zeta : \mathsf{P}Z \to N$ by $\mathsf{P}\zeta(d\sigma(m)) = m$. We define a mapping

$$\mathsf{P}\rho:\mathsf{P}Z\times_N\mathsf{P}Z\to\mathsf{T}^*N$$

by

$$\mathsf{P}\rho(\mathrm{d}\sigma_2(m),\mathrm{d}\sigma_1(m)) = \mathrm{d}(\sigma_2 - \sigma_1)(m).$$

The pair $\mathbf{PZ} = (\mathsf{P}\zeta, \mathsf{P}\rho)$ is an affine fibration modelled on the fibration $\pi_N: \mathsf{T}^*N \to N$. This fibration is called the *phase fibration* of $\mathbf{Z} = (\zeta, \rho)$. Let φ be a section of $\mathsf{P}\zeta$ and let σ be a section of ζ . We define the *differential* $d\varphi$ of φ by $d\varphi = d(\varphi - d\sigma)$. Since for two sections σ, σ' of ζ we have $d(d\sigma - d\sigma') = dd(\sigma - \sigma') = 0$ it follows that the definition of the differential does not depend on the choice of σ .

For each **Z** the manifold PZ is a symplectic manifold⁶.

3. AFFINE POISSON STRUCTURES

3.1. Homogeneous Formulation of the Dynamics

In the first section we have assumed that the space-time is the product of the space and the time (represented by the real line \mathbb{R}). This assumption implies that we have chosen a reference frame. In this section we formulate the dynamics of a nonrelativistic system in a frame-independent way. We represent the time by the real line. The space-time is a fibration over the time. It follows that the manifold of space-time configurations of a system is a fibration

$$\zeta: M \to \mathbb{R}.$$

Let us denote by Q_t the fiber over $t \in \mathbb{R}$ of the fibration. Infinitesimal configurations are vectors tangent to M. $\mathsf{T}M$ is the manifold of infinitesimal configurations. The phase bundle is the cotangent bundle T^*M . Let $\hat{\eta}: \mathsf{T}^*M \to \mathbb{R}$ be the canonical projection $\hat{\eta} = \zeta \circ \pi_M$. The dynamics of a system is a submanifold \widehat{D} of $\mathsf{T}\mathsf{T}^*M$. We say that the system is *lagrangian* if D is the Lagrangian submanifold of $(\mathsf{T}\mathsf{T}^*M, d_\mathsf{T}\omega_M)$. Let L be Lagrangian generating function of D. For a nonrelativistic system L is a homogeneous function on $\mathsf{T}M$. It follows that the Hamiltonian generating function is the zero function on a submanifold C of T^*M .

3.2. Inhomogeneous Formulation of the Dynamics

In the formulation of the dynamics presented in Section 1.2 the existence of Lagrangian and Hamiltonian generating functions was possible because the space-time was assumed to be the product of the space and the time. Let $\zeta: M \to \mathbb{R}$ be the configuration manifold of a system, fibered over the time. By M_t we denote a fiber of the fibration ζ , $M_t = (\zeta)^{-1}(t)$. An infinitesimal configuration of the system is a vector $v \in \mathsf{T}M$ such that $\mathsf{T}\zeta(v) = (\zeta(\tau_M v), \partial_t)$. T_1M will denote the bundle of infinitesimal configurations. For each $t \in \mathbb{R}$ a submanifold T_1M_t of T_1M is defined by

$$\mathsf{T}_1 M_t = \{\mathsf{T}_1 M \ni w; \ \eta(\tau_M w) = t\}.$$

The phase bundle is a fibration $\eta: P \to \mathbb{R}$ with $P_t = (\eta)^{-1}(t) = \mathsf{T}^* M_t$. For each $t \in \mathbb{R}$ the fiber P_t can be considered as the reduction of $\mathsf{T}^* M$ with respect to a coisotropic submanifold $K_t = \{\mathsf{T}^* \ni p; \zeta(\pi_M p) = t\}$. An infinitesimal state is a vector $w \in \mathsf{T} P$ such that $\mathsf{T} \eta(w) = (\eta(\tau_P w), \partial_t)$. We denote by $\mathsf{T}_1 P$ the bundle of infinitesimal states . For each $t \in \mathbb{R}$ a submanifold $\mathsf{T}_1 P_t$ of $\mathsf{T}_1 P$ is defined by

$$\mathsf{T}_1 P_t = \{\mathsf{T}_1 P \ni w; \ \eta(\tau_P w) = t\}.$$

Proposition 3 A submanifold T_1P_t is the reduction of $(TT^*M, d_T\omega_M)$ with respect to a coisotropic submanifold T_1K_t defined by

$$\mathsf{T}_1 K_t = \{\mathsf{T}\mathsf{T}^* M \ni w; \ \mathsf{T}\pi_M(w) \in \mathsf{T}_1 M \text{ and } \zeta(\tau_{\mathsf{T}^* M} \circ \pi_M(w)) = t\}$$

It follows from this proposition that T_1P_t is a symplectic manifold. The dynamics of the system is a submanifold D of T_1P . Let D_t denote a subset of D defined by

$$D_t = D \cap \mathsf{T}_1 P_t.$$

The system is Lagrangian if for each $t \in \mathbb{R}$ the dynamics D_t is a Lagrangian submanifold of $\mathsf{T}_1 P_t$. The existence of a Lagrangian generating function follows from the theorem

Theorem 1 Let T_1TM_t be a submanifold of TTM defined by

$$\mathsf{T}_1\mathsf{T} M_t = \{\mathsf{T}\mathsf{T} M \ni w; \mathsf{T}\tau_M w \in \mathsf{T}_1 M \text{ and } \tau_{\mathsf{T} M} w \in \mathsf{T} M_t\}.$$

There are canonical isomorpisms of vector bundles

$$\widetilde{\kappa}_{M_t} : (\tau_{\mathsf{T}_1 M_t} : \mathsf{T}\mathsf{T}_1 M_t \to \mathsf{T}_1 M_t) \longrightarrow (\mathsf{T}\tau_M : \mathsf{T}_1 \mathsf{T} M_t \to \mathsf{T}_1 M_t)$$
(3.1)

$$\widetilde{\alpha}_{M_t}: (\mathsf{T}\pi_M: \mathsf{T}_1\mathsf{T}^*M_t \to \mathsf{T}_1M) \longrightarrow (\pi_{\mathsf{T}_1M_t}: \mathsf{T}^*\mathsf{T}M \to \mathsf{T}M)$$
(3.2)

Proof. The bundle $\mathsf{T}\tau_M: \mathsf{T}_1\mathsf{T}M_t \to \mathsf{T}_1M_t$ is defind as a subbundle of the bundle $\mathsf{T}\tau_M: \mathsf{T}\mathsf{T}M \to \mathsf{T}M$.

Also the bundle $\tau_{\mathsf{T}_1M_t}: \mathsf{T}\mathsf{T}_1M_t \to \mathsf{T}_1M_t$ can be considered as a subbundle of the bundle $\tau_{\mathsf{T}M}: \mathsf{T}\mathsf{T}M \to \mathsf{T}M$. It is an easy exercise to verify that κ_M restricted to $\mathsf{T}\mathsf{T}_1M_t$ gives the required isomorphism.

The isomorphism α_M is defined as the dual to κ_M . We define the isomorphism $\tilde{\alpha}_{M_t}$ as the dual to $\tilde{\kappa}_{M_t}$ as well. Since α_M is a symplectomorphism we conclude that also $\tilde{\alpha}_{M_t}$ is a symplectomorphism.

It follows that a Lagrangian system can be generated by a Lagrangian generating function defined on T_1M . The Hamiltonian formulation of a dynamics is more complicated and requires affine structures.

3.3. Affine Poisson Structures

This paragraph is based on the relation between Legendre transformation and the affine duality⁷. A Lagrange bundle is the trivial line bundle $\hat{\mathsf{T}}_1 M = \mathsf{T}_1 M \times \mathbb{R}$ over $\mathsf{T}_1 M$. By ξ we denote the canonical projection

$$\xi \colon \widehat{\mathsf{T}}_1 M \to \mathsf{T}_1 M.$$

Lagrangians are sections of the fibration ξ . With respect to the projection

$$\tau_M \circ \xi \colon \widehat{\mathsf{T}}_1 M \to M$$

 $\widehat{\mathsf{T}}_1 M$ is a special affine bundle. The affine dual $\widehat{\mathsf{T}}_1^{\#} M$ to this bundle is a Hamiltonian bundle.

Proposition 4 The special affine bundle $\widehat{\mathsf{T}}_1^{\#} M$ is isomorphic to the cotangent bundle $\pi_M: \mathsf{T}^* M \to M$. The distinguished covector field ϑ is defined by $\langle v, \vartheta \rangle = 0$ for $\mathsf{T} \zeta v = 0$ and $\langle \partial_t, \vartheta \rangle = 1$.

Proof. Let us fix $m \in M$. Elements of $\widehat{\mathsf{T}}_1^{\#}M$ over m are affine functions on $\mathsf{T}_{1,m}M$. An affine function on $\mathsf{T}_{1,m}M$ has the unique extension to a linear function on T_mM , i. e., to an element of T^*M . The distinguished element of $\widehat{\mathsf{T}}_1^{\#}M$ at m is the equal to 1 constant function. The liner extension of this function is a ζ -vertical covector, equal to 1 on vectors which project onto ∂_t .

With this isomorphism the line bundle structure of $\widehat{\mathsf{T}}_1^{\#}M$ is given by the canonical projection

$$\chi: \mathsf{T}^* M \to P.$$

Theorem 2 There is a canonical isomorphism of affine bundles

$$\widetilde{\beta}_{M_t}: (\tau_P: \mathsf{T}_1 P_t \to P_t) \longrightarrow (\mathsf{P}\chi: \mathsf{P}(\widehat{\mathsf{T}}_1^{\#} M_t) \to P_t),$$

which is also a symplectomorphism.

Proof. Let $\gamma: P_t \to \mathsf{T}^* M$ be a section of the fibration χ . We define a function $\tilde{\gamma}$ on $\chi^{-1}(P_t)$ by the formula

$$\widetilde{\gamma}(\gamma(p) + s\vartheta) = -s.$$

We define a relation R from $C^{\infty}(\mathsf{T}^*M)$ to the space of sections of χ over P_t :

 $\gamma \in R(f)$ if $f = \tilde{\gamma}$ on $\chi^{-1}(P_t)$.

The relation R defines a relation

$$\mathrm{d}R: \mathsf{T}^*M \to \mathsf{P}(\widehat{\mathsf{T}}_1^{\#}M_t).$$

It is easy to verify that the composition $dR \circ \beta_M$ projects to an isomorphism $\widetilde{\beta}_{M_t}$ of affine bundles

$$\widetilde{\beta}_{M_t}: (\tau_P: \mathsf{T}_1 P_t \to P_t) \longrightarrow (\mathsf{P}\chi: \mathsf{P}(\widehat{\mathsf{T}}_1^{\#} M_t) \to P_t),$$

and that this isomorphism is a symplectomorphism.

It follows that the dynamics of a Lagrangian system can be generated by a section of the fibration $\chi: \mathsf{T}_1 P \to P$. This section we call the Hamiltonian generating section. It is easy to verify that the image of the Hamiltonian generating section is the submanifold C of $\mathsf{T}^* M$ we mentioned in Section 3.1. The collection $(\tilde{\beta}_{M_t})$ of isomorphisms defines a morphism Λ of affine bundles

$$\Lambda: (\mathsf{P}\chi: \mathsf{P}(\widehat{\mathsf{T}}_1^{\#}M) \to P) \longrightarrow (\tau_P\mathsf{T}_1P: P).$$

Let $\Gamma(\chi)$ be the afine space of sections of the fibration χ . With the morphism Λ we define an *affine Poisson bracket* as a mapping

$$\{,\}: \Gamma(\chi) \times C^{\infty}(P) \to C^{\infty}(P)$$

defined by

$$\{\gamma, f\}(p) = \Lambda(d_p\gamma)(f).$$

The bracket $\{,\}$ has the following properties:

- it is affine with respect to the first and linear with respect to the second argument,
- the linear part is a linear Poisson bracket,
- for each section γ the mapping $C^{\infty}(P) \to C^{\infty}(P)$: $f \mapsto \{\gamma, f\}$ defines a canonical vector field on P.

A discussion on the concept of an affine Poisson structure will be given in a separate publication.

References

- S. Benenti, Symplectic relations in analytical mechanics, in: "Modern Developments in Analytical Mechanics", S. Benenti, M. Francaviglia and A. Lichnerowicz, ed., Atti Accad. Sci. Torino, Suppl. 117 (1983)
- P. Libermann and C.-M. Marle, "Symplectic Geometry and Analytical Mechanics", D. Reidel, Dordrecht (1987)
- C.-M. Marle, Lie group actions on a canonical manifold, *in*: "Symplectic Geometry",
 A. Crumeyrolle and J. Grifone , ed., Pitman-Research Notes in Math. 80 (1983)
- 4. G. Pidello and W. M. Tulczyjew, "Derivations of differential forms on jet bundles", Ann. Matem. Pura Appl., 147 (1987)
- 5. W. M. Tulczyjew, "Geometric Formulations of Physical Theories", Bibliopolis, Naples (1989)
- W. M. Tulczyjew and P. Urbański, An Affine framework for the dynamics of charged particles, *in*: "La 'Mécanique Analytique' de Lagrange et son Héritage -II", *Atti Accad. Sci. Torino*, Suppl. 126 (1992)
- 7. P. Urbański, Affine duality and the Legendre transformation in mechanics, (in preparation).