

Integrable and chaotic generalizations of the rolling of a convex body without slipping on a plane

Mariana Costa Villegas
Joint work with Luis García Naranjo

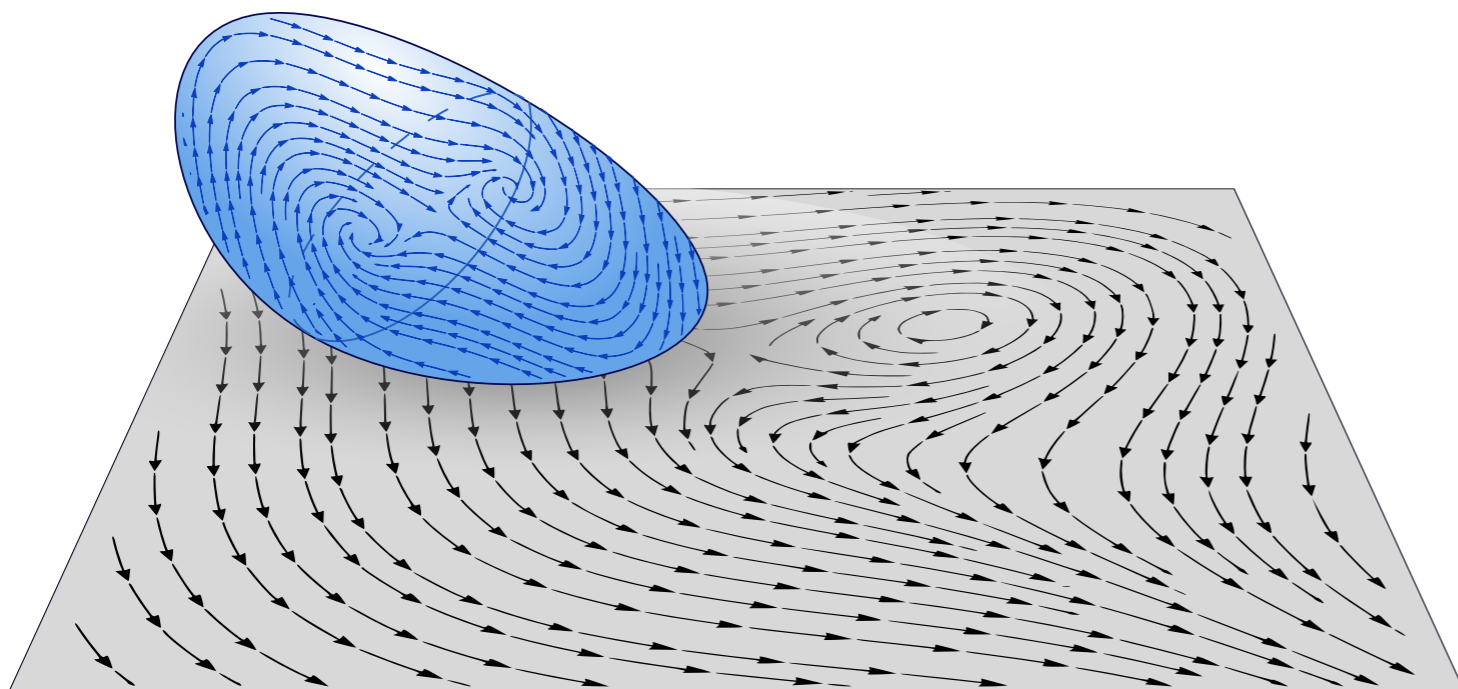


UNIVERSITÀ
DEGLI STUDI
DI PADOVA

 DIPARTIMENTO
MATEMATICA
DIPARTIMENTO DI MATEMATICA "TULLIO LEVI-CIVITA"

**XVIII International Young
Researchers Workshop in Geometry,
Dynamics and Field Theory**

Warsaw, Poland, February 21-23, 2024



Classic convex body rolling without slipping on a plane

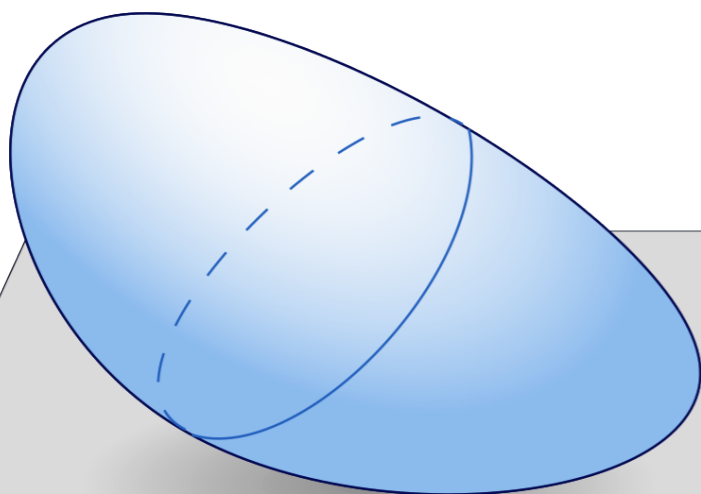
Nonholonomic system

Configuration manifold: $Q = SO(3) \times \mathbb{R}^2$
 $(B, \underline{x}) \in Q$

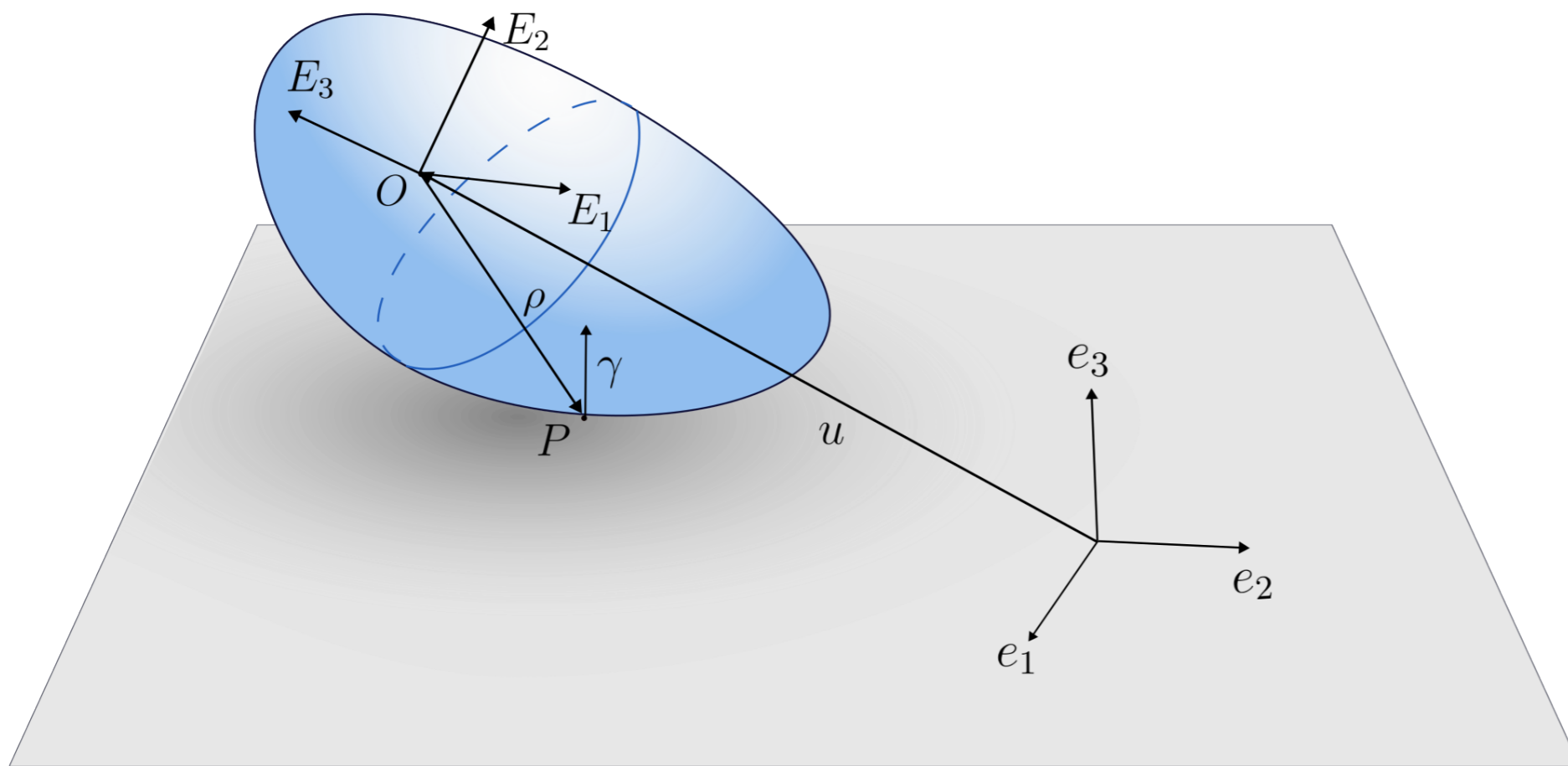
Nonholonomic linear constraint

Phase space: dim 8

$SE(2)$ - symmetry on plane



Classic convex body rolling without slipping on a plane



$$\gamma = B^{-1}e_3$$

$$\rho = \rho(\gamma)$$

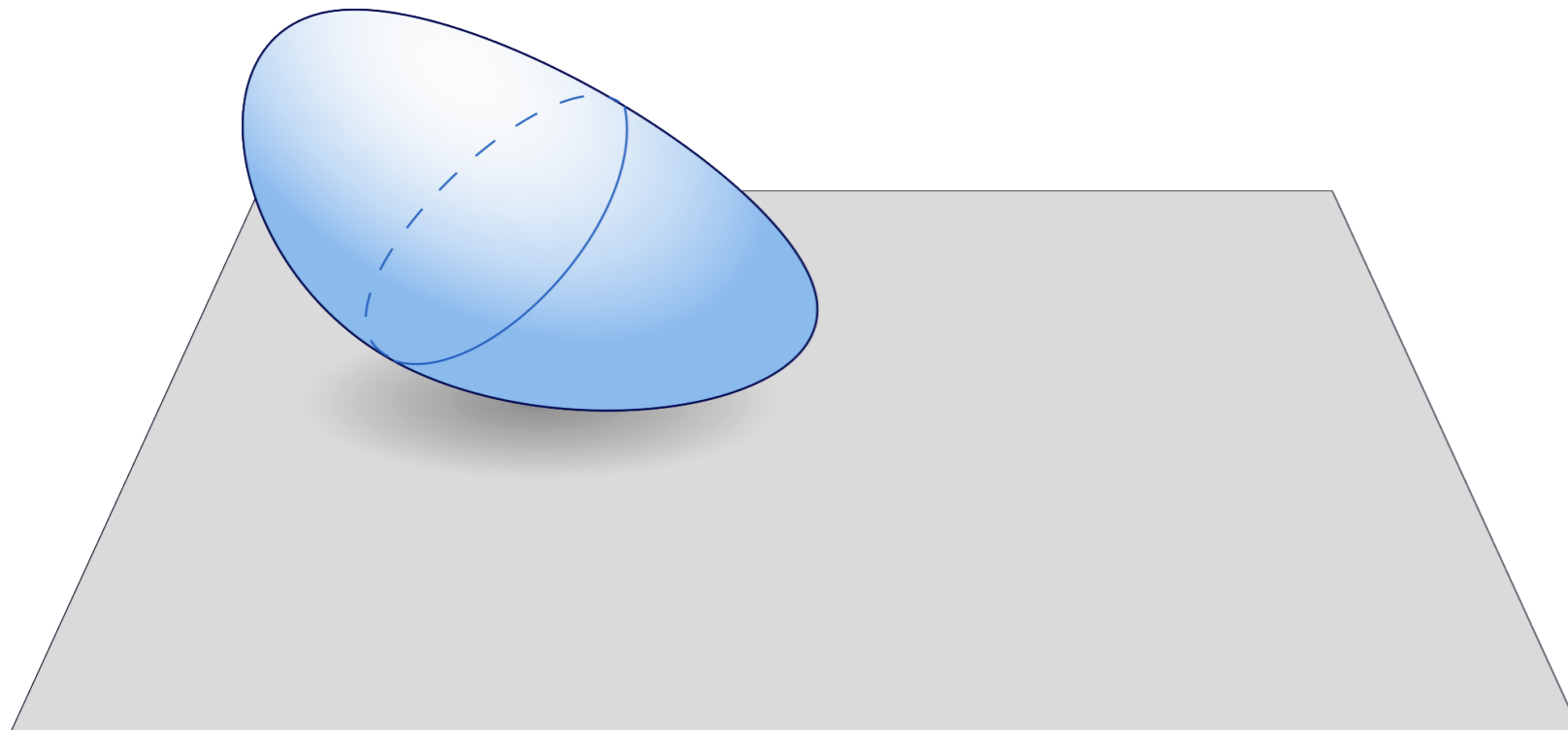
$$\underline{x} = u + B\rho$$

Equations of motion:

$$\dot{M} = M \times \Omega + m\dot{\rho} \times (\Omega \times \rho) + mg\rho \times \gamma$$

$$\dot{\gamma} = \Omega \times \gamma$$

$$M = \mathbb{I}\Omega + m\rho \times (\Omega \times \rho), \quad \rho = \rho(\gamma)$$



Equations of motion:

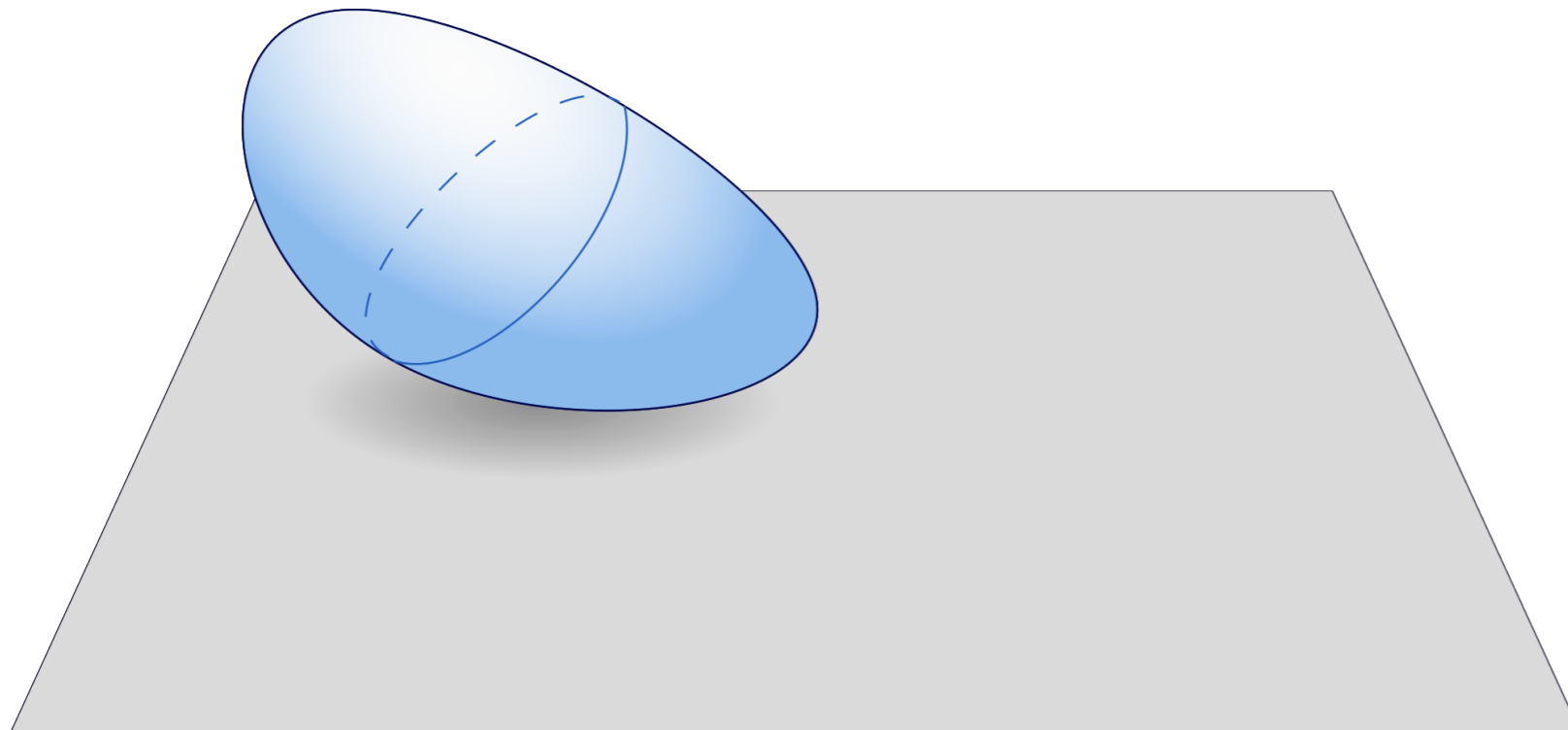
$$\dot{M} = M \times \Omega + m\dot{\rho} \times (\Omega \times \rho) + mg\rho \times \gamma$$

$$\dot{\gamma} = \Omega \times \gamma$$

$$M = \mathbb{I}\Omega + m\rho \times (\Omega \times \rho), \quad \rho = \rho(\gamma)$$

First integrals:

$$H = \frac{1}{2}\langle M, \Omega \rangle - mg\langle \rho, \gamma \rangle, \quad \|\gamma\|^2 = 1$$



Equations of motion:

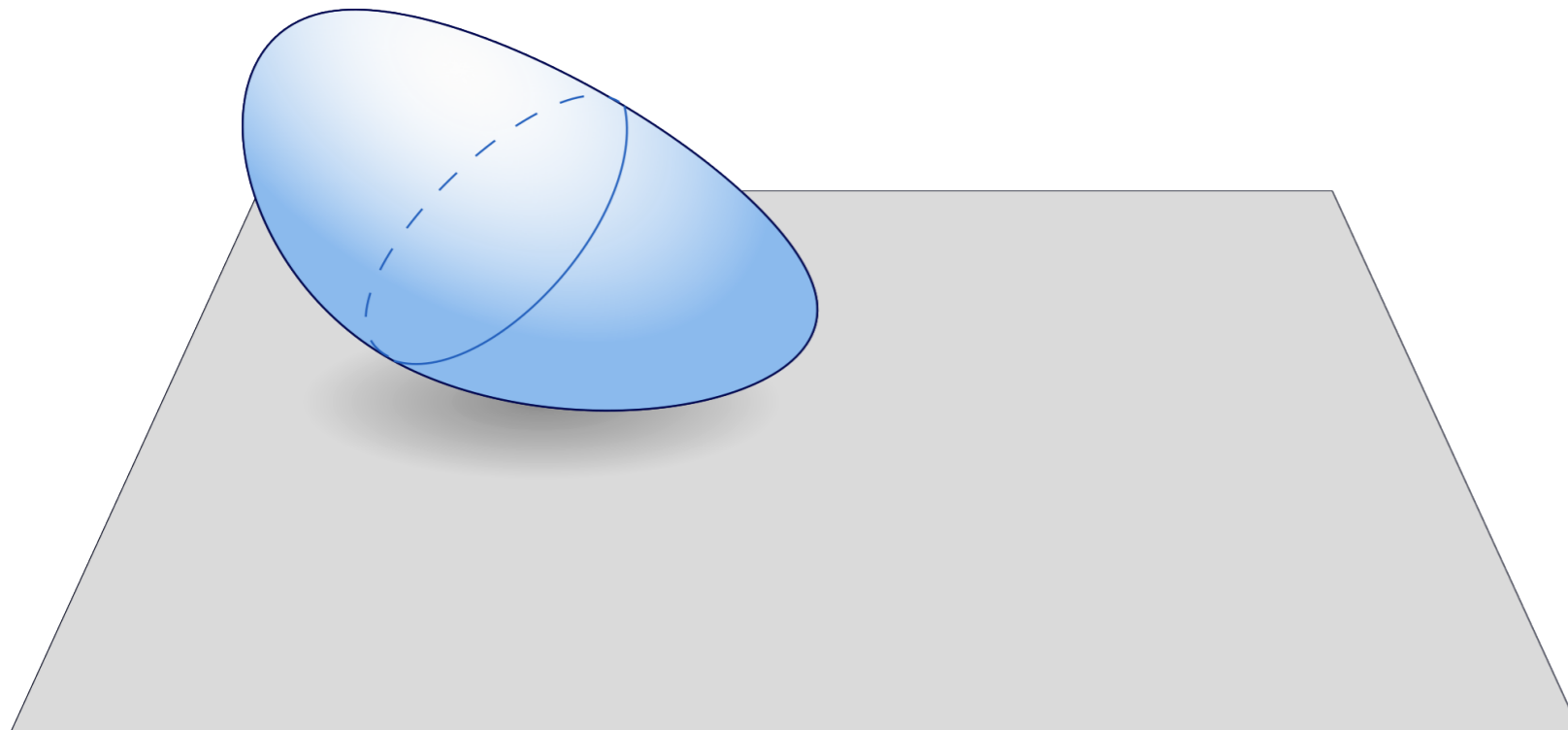
$$\dot{M} = M \times \Omega + m\dot{\rho} \times (\Omega \times \rho) + mg\rho \times \gamma$$
$$\dot{\gamma} = \Omega \times \gamma$$

$$M = I\Omega + m\rho \times (\Omega \times \rho), \quad \rho = \rho(\gamma)$$

First integrals:

$$H = \frac{1}{2}\langle M, \Omega \rangle - mg\langle \rho, \gamma \rangle, \quad \|\gamma\|^2 = 1$$

- No invariant measure
- Generally chaotic



Equations of motion:

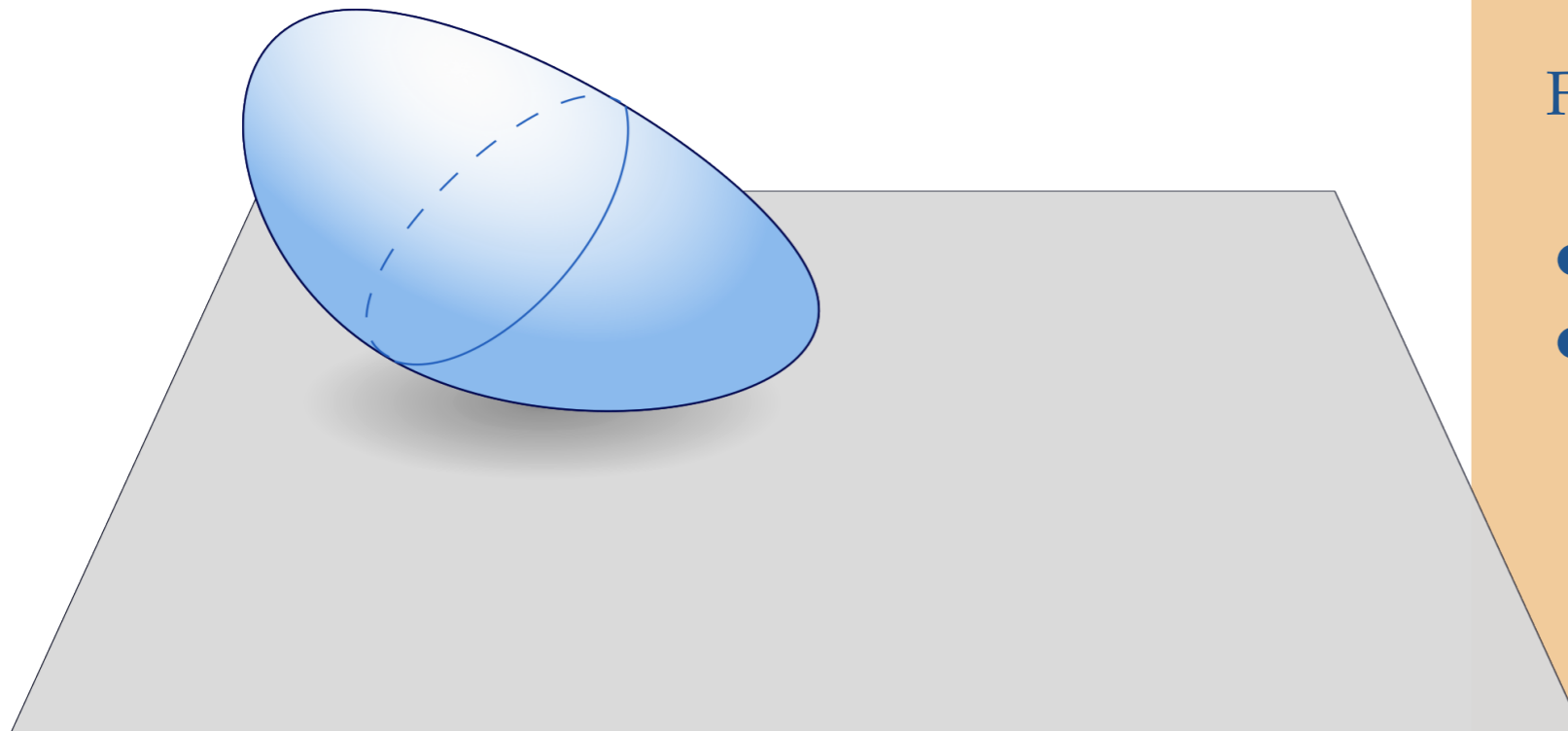
$$\dot{M} = M \times \Omega + m\dot{\rho} \times (\Omega \times \rho) + mg\rho \times \gamma$$
$$\dot{\gamma} = \Omega \times \gamma$$

$$M = I\Omega + m\rho \times (\Omega \times \rho), \quad \rho = \rho(\gamma)$$

First integrals:

$$H = \frac{1}{2}\langle M, \Omega \rangle - mg\langle \rho, \gamma \rangle, \quad \|\gamma\|^2 = 1$$

- No invariant measure
- Generally chaotic

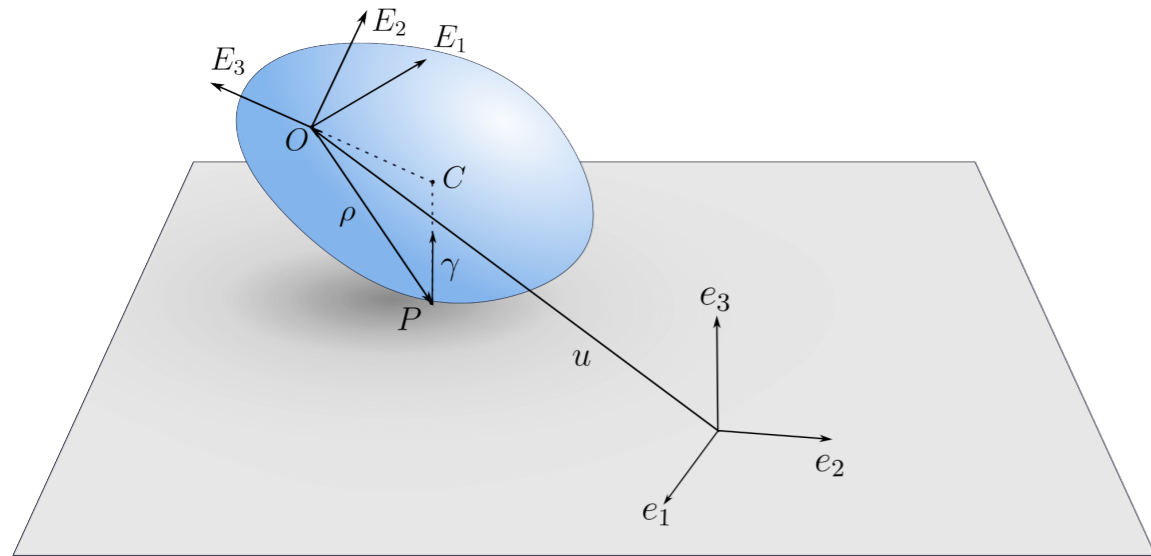


For it to be integrable we need:

- 2 additional first integrals
- invariant measure.

Integrable cases

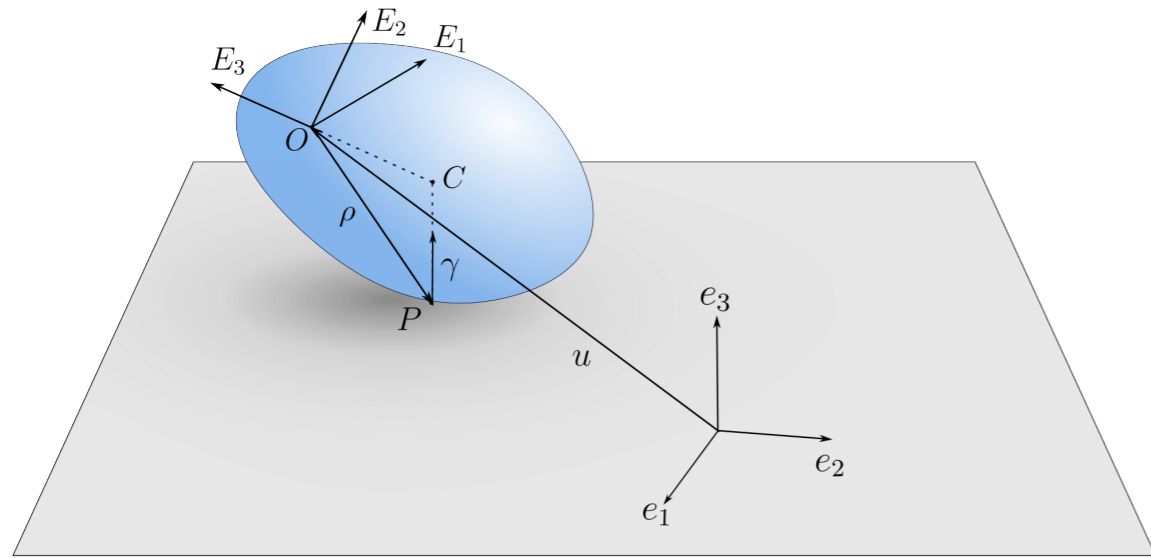
Integrable cases



Body of revolution (Chaplygin, 1897)

$$I_1 = I_2 \neq I_3, O \neq C$$

Integrable cases



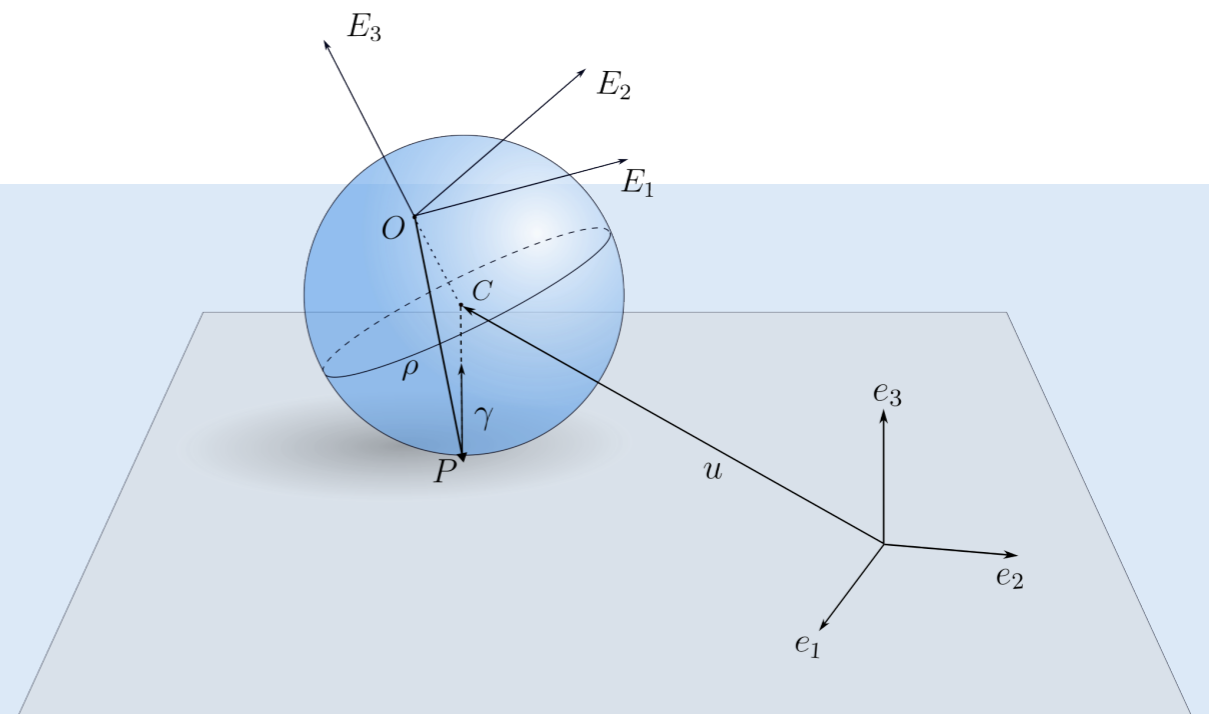
Body of revolution (Chaplygin, 1897)

$$I_1 = I_2 \neq I_3, O \neq C$$

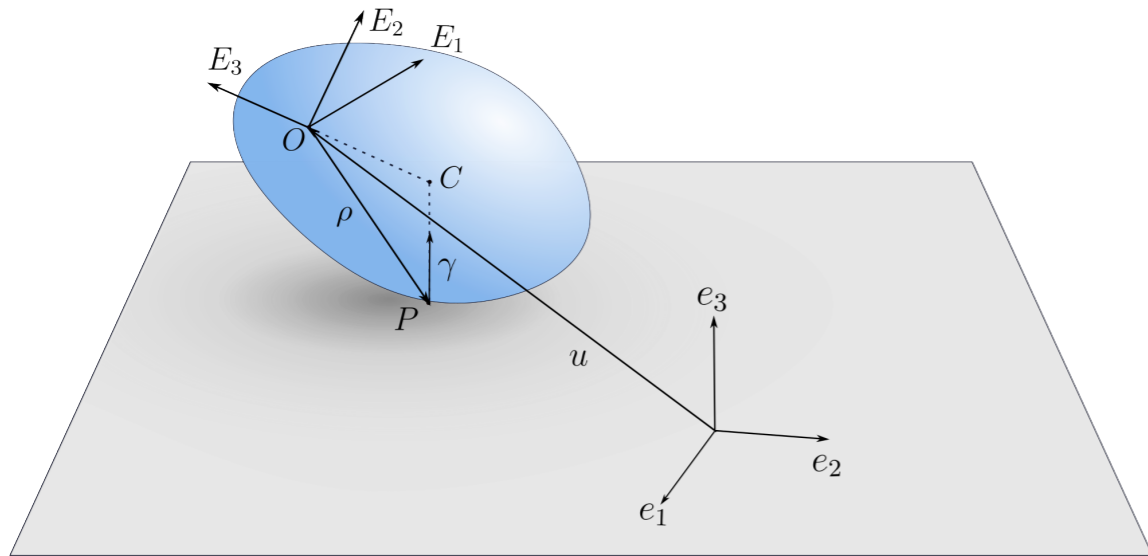
Routh's sphere (Routh, 1884)

$$I_1 = I_2 \neq I_3, O \neq C$$

$$\rho(\gamma) = -r\gamma - lE_3$$



Integrable cases



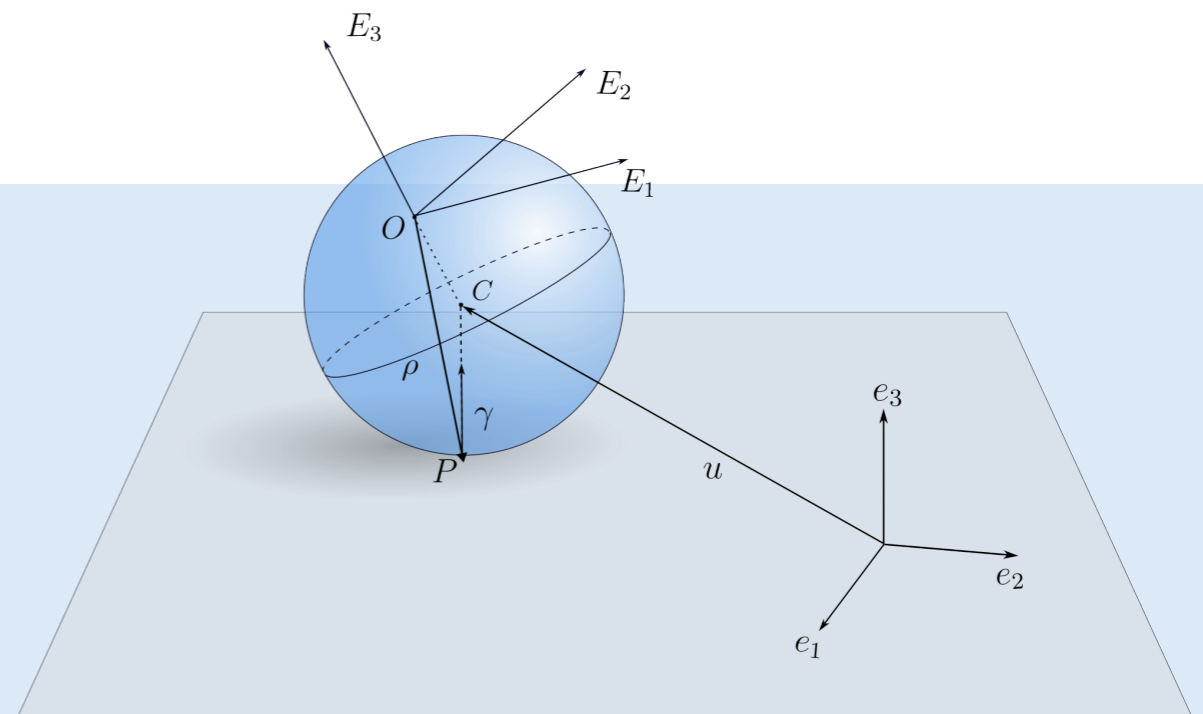
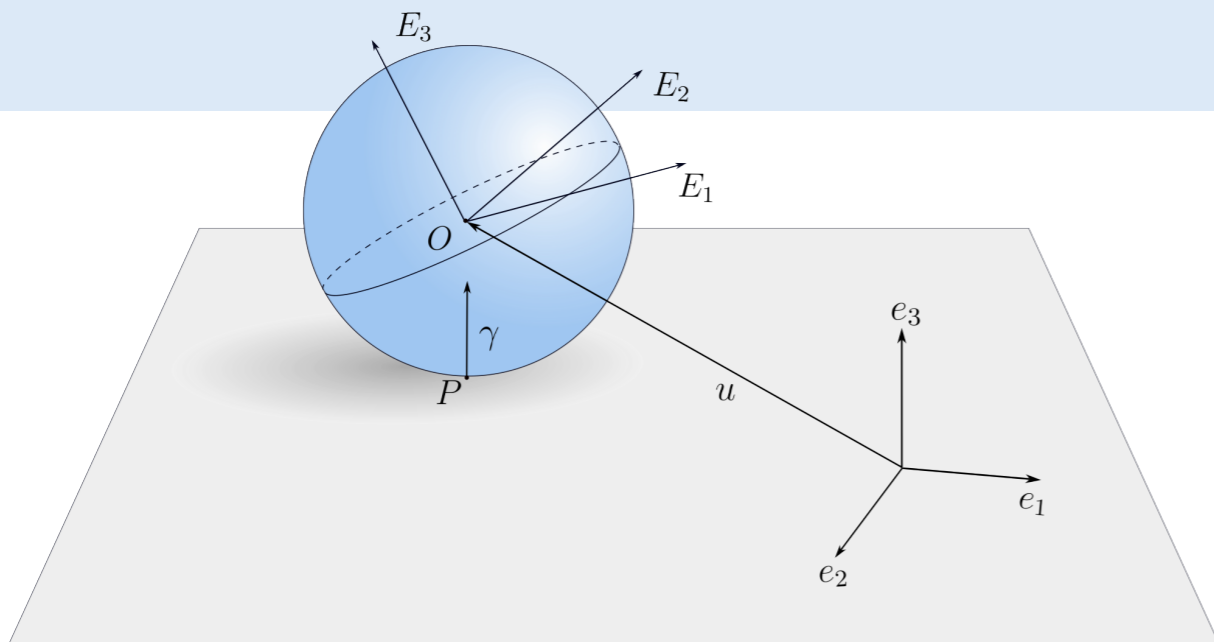
Body of revolution (Chaplygin, 1897)

$$I_1 = I_2 \neq I_3, O \neq C$$

Routh's sphere (Routh, 1884)

$$I_1 = I_2 \neq I_3, O \neq C$$

$$\rho(\gamma) = -r\gamma - lE_3$$



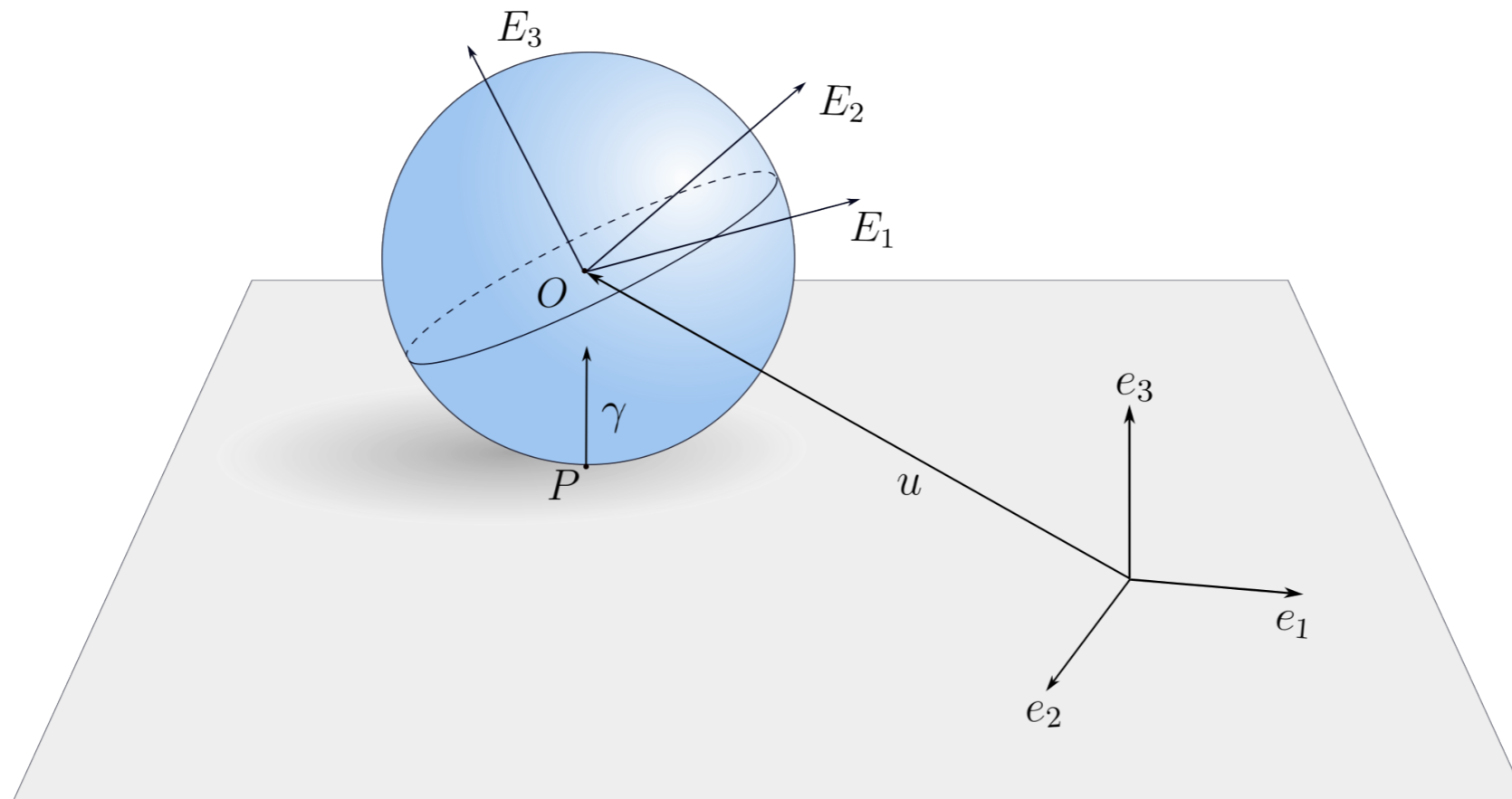
Chaplygin sphere (Chaplygin, 1903)

$$I_1 \neq I_2 \neq I_3, O = C$$

Integrable cases

Homogeneous sphere

$$I_1 = I_2 = I_3, \quad O = C$$

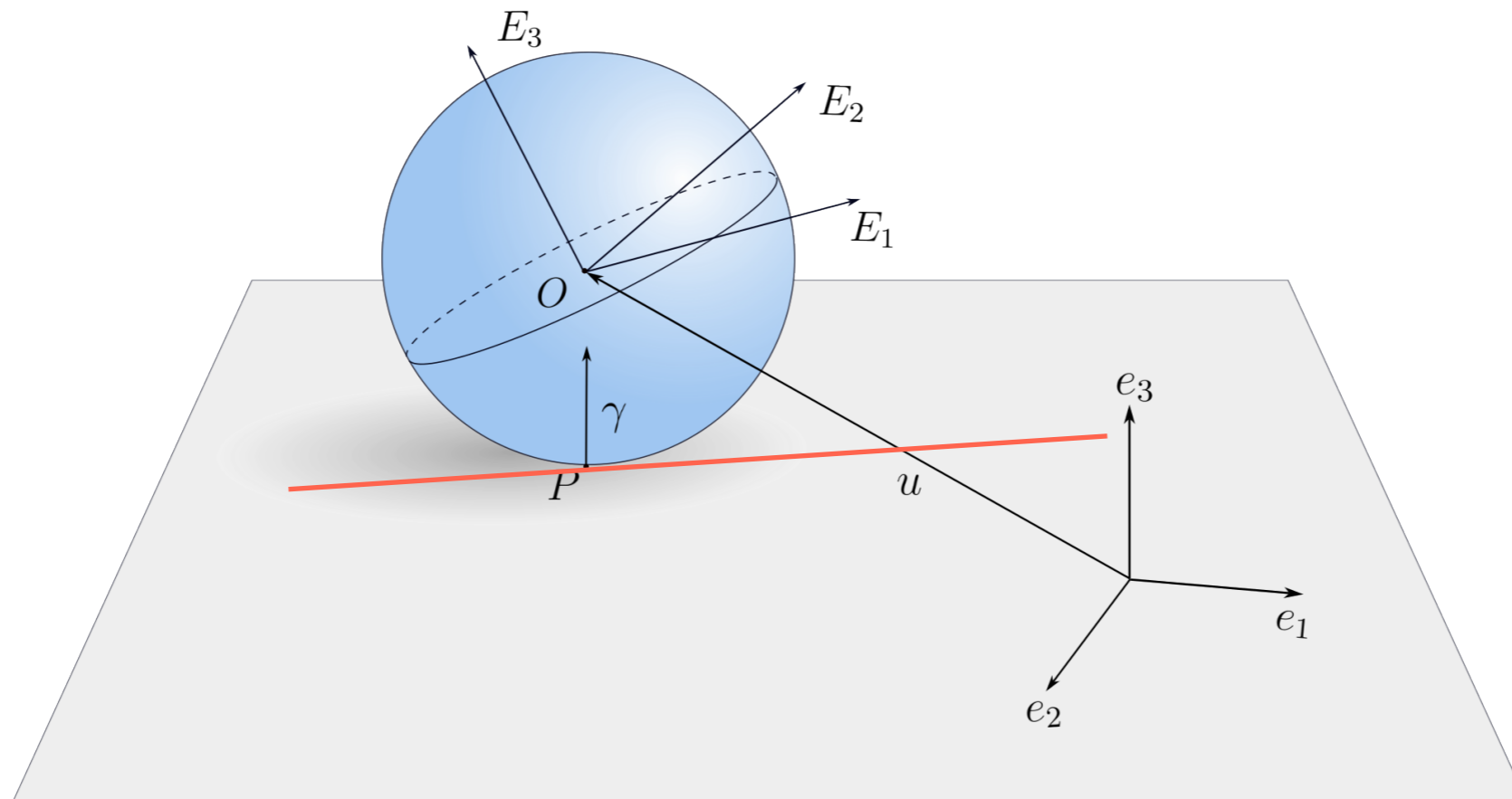


Integrable cases

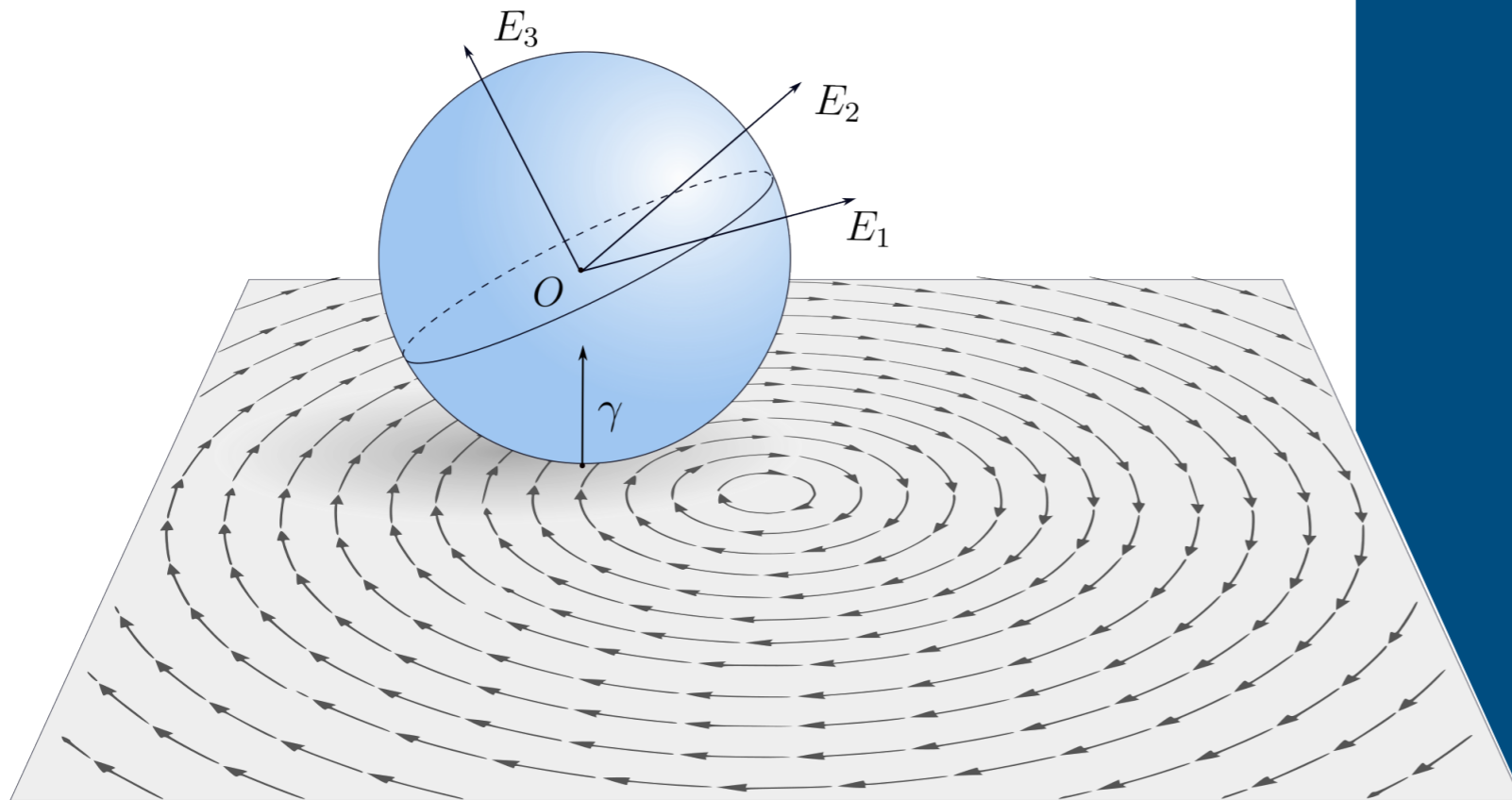
Homogeneous sphere

$$I_1 = I_2 = I_3, \quad O = C$$

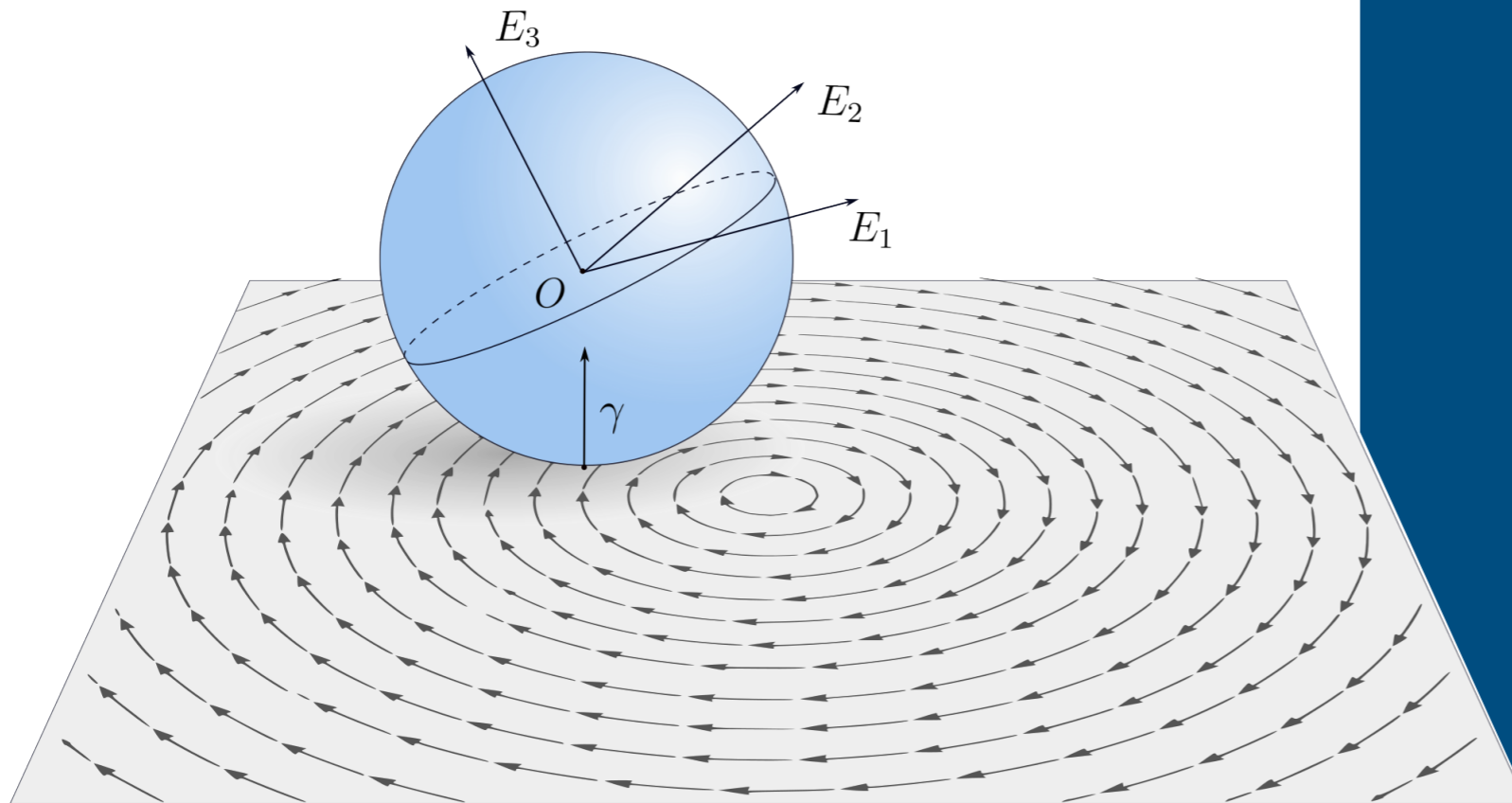
Moves in straight line



What happens if the plane
is moving?

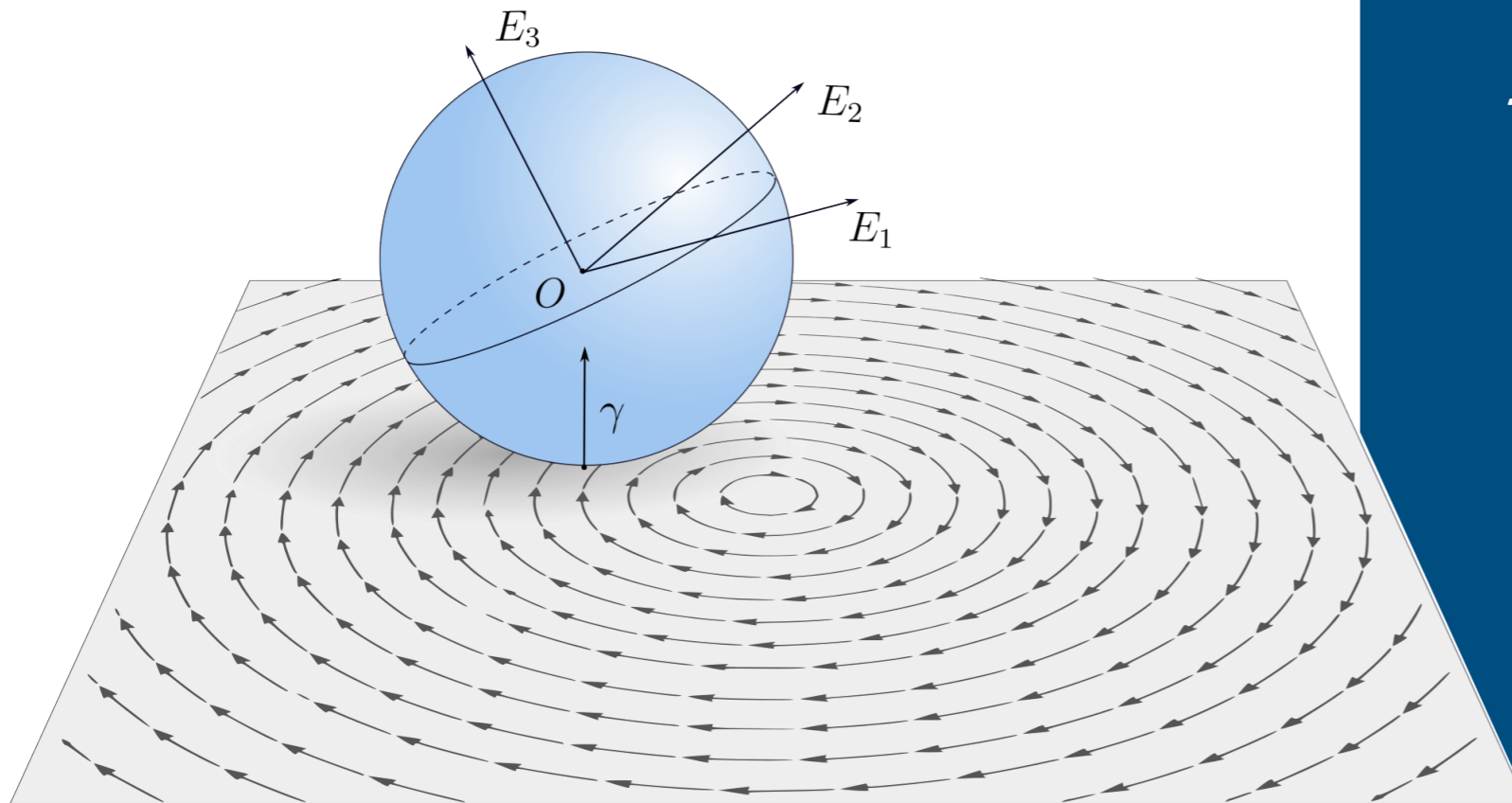


What happens if the plane
is moving?



- Lose symmetry on the plane
- Lose energy first integral

What happens if the plane
is moving?



- Lose symmetry on the plane
- Lose energy first integral
- Moving energy
(Fassò, Sansonetto, 2015)

Homogeneous sphere on rotating plane

Equations of motion:

$$\begin{aligned}\dot{M} &= M \times \Omega \\ \dot{U} &= -r(\gamma \times \Omega) + U \times \Omega - \eta \gamma \times U \\ \dot{\gamma} &= \gamma \times \Omega\end{aligned}$$

$$M = I\Omega + mr^2\gamma \times (\Omega \times \gamma) + \eta\gamma \times (\gamma \times U)$$

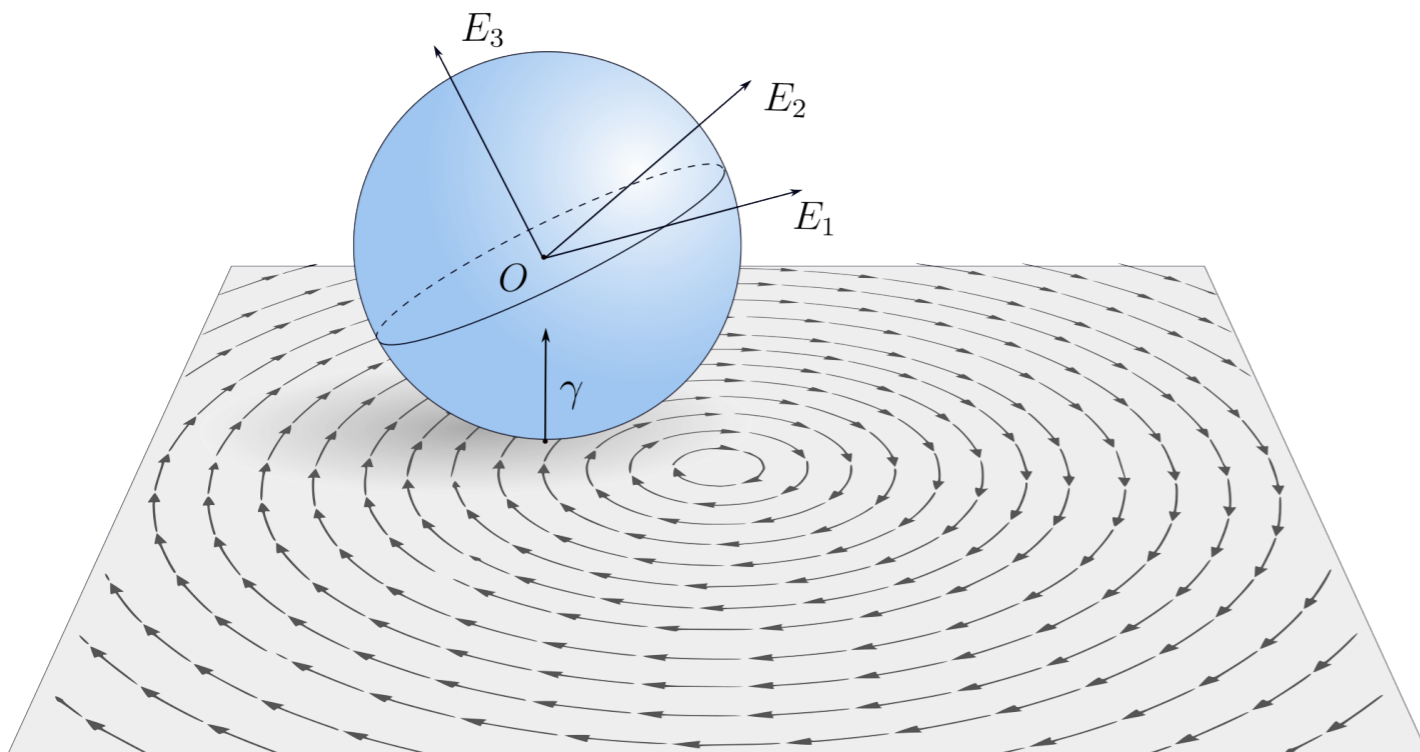
First integrals

$$\|M\|^2, \quad \langle M, \gamma \rangle, \quad \|\gamma\|^2, \quad \langle U, \gamma \rangle = 0,$$

$$E_{mov}, \quad \|\Omega + \eta U\|^2,$$

Invariant measure

$$dM dx dy$$



Homogeneous sphere on rotating plane

Integrable

Equations of motion:

$$\dot{M} = M \times \Omega$$

$$\dot{U} = -r(\gamma \times \Omega) + U \times \Omega - \eta \gamma \times U$$

$$\dot{\gamma} = \gamma \times \Omega$$

$$M = I\Omega + mr^2\gamma \times (\Omega \times \gamma) + \eta\gamma \times (\gamma \times U)$$

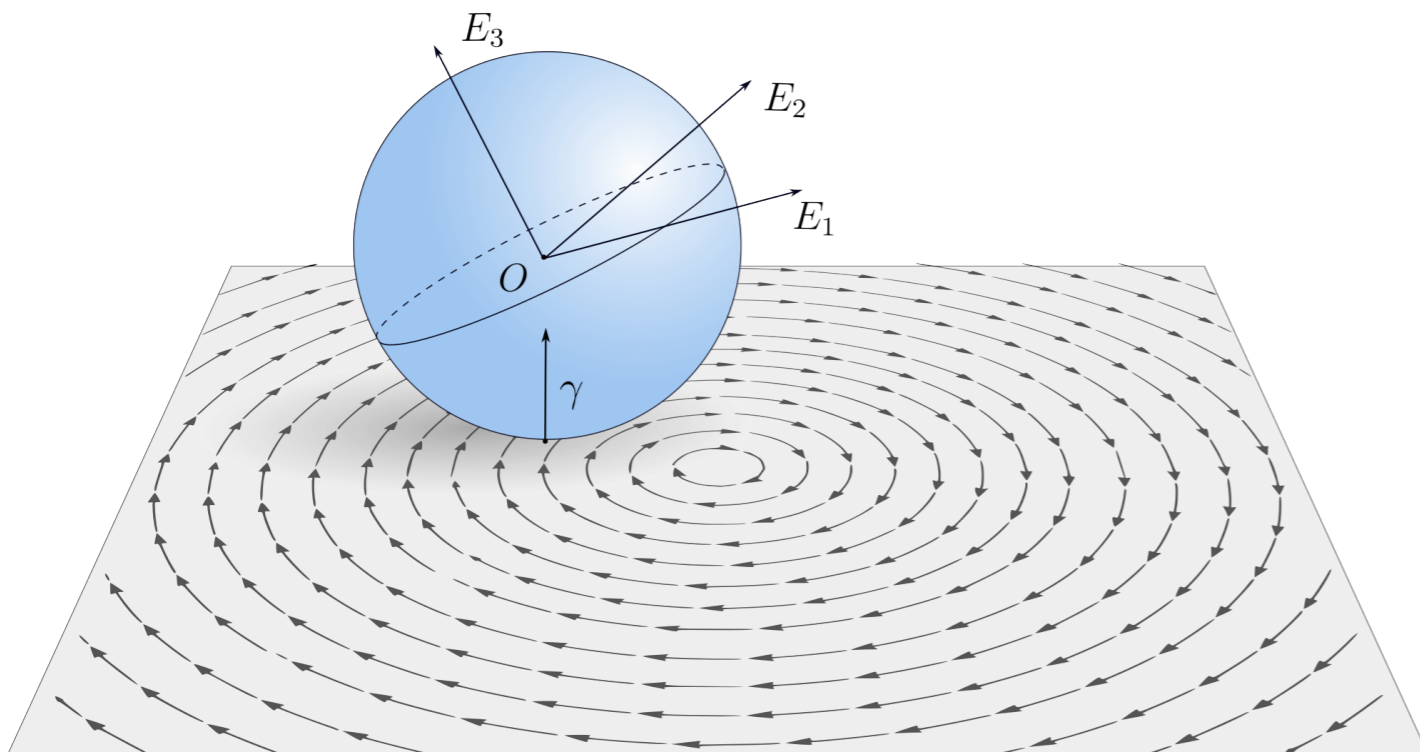
First integrals

$$\|M\|^2, \quad \langle M, \gamma \rangle, \quad \|\gamma\|^2, \quad \langle U, \gamma \rangle = 0,$$

$$E_{mov}, \quad \|\Omega + \eta U\|^2,$$

Invariant measure

$$dM dx dy$$



Homogeneous sphere on rotating plane

Integrable

Moves in circle

Equations of motion:

$$\dot{M} = M \times \Omega$$

$$\dot{U} = -r(\gamma \times \Omega) + U \times \Omega - \eta \gamma \times U$$

$$\dot{\gamma} = \gamma \times \Omega$$

$$M = I\Omega + mr^2\gamma \times (\Omega \times \gamma) + \eta\gamma \times (\gamma \times U)$$

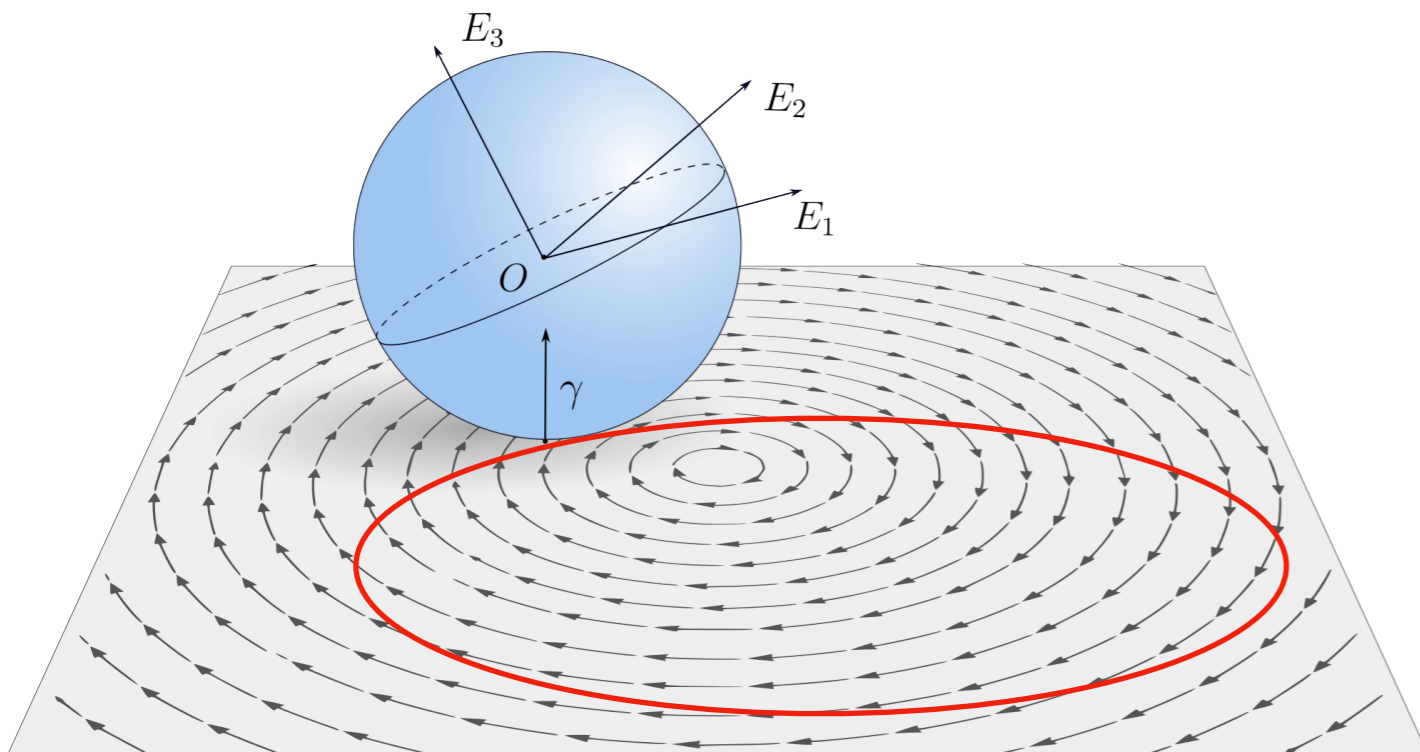
First integrals

$$\|M\|^2, \quad \langle M, \gamma \rangle, \quad \|\gamma\|^2, \quad \langle U, \gamma \rangle = 0,$$

$$E_{mov}, \quad \|\Omega + \eta U\|^2,$$

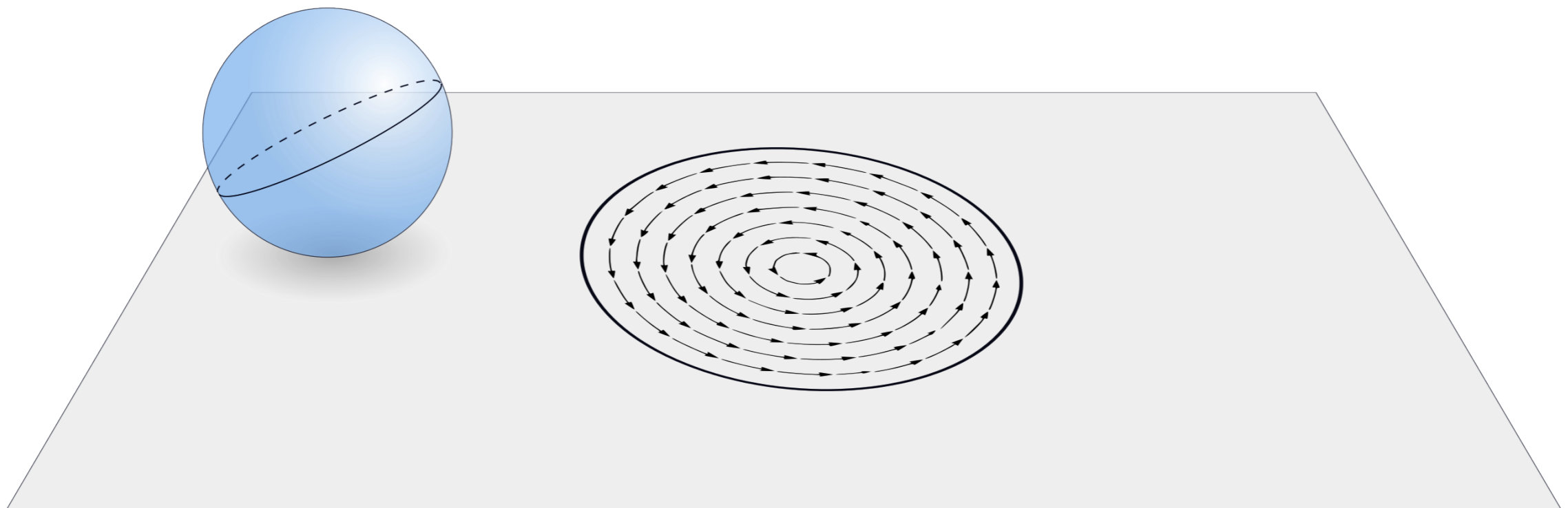
Invariant measure

$$dM dx dy$$



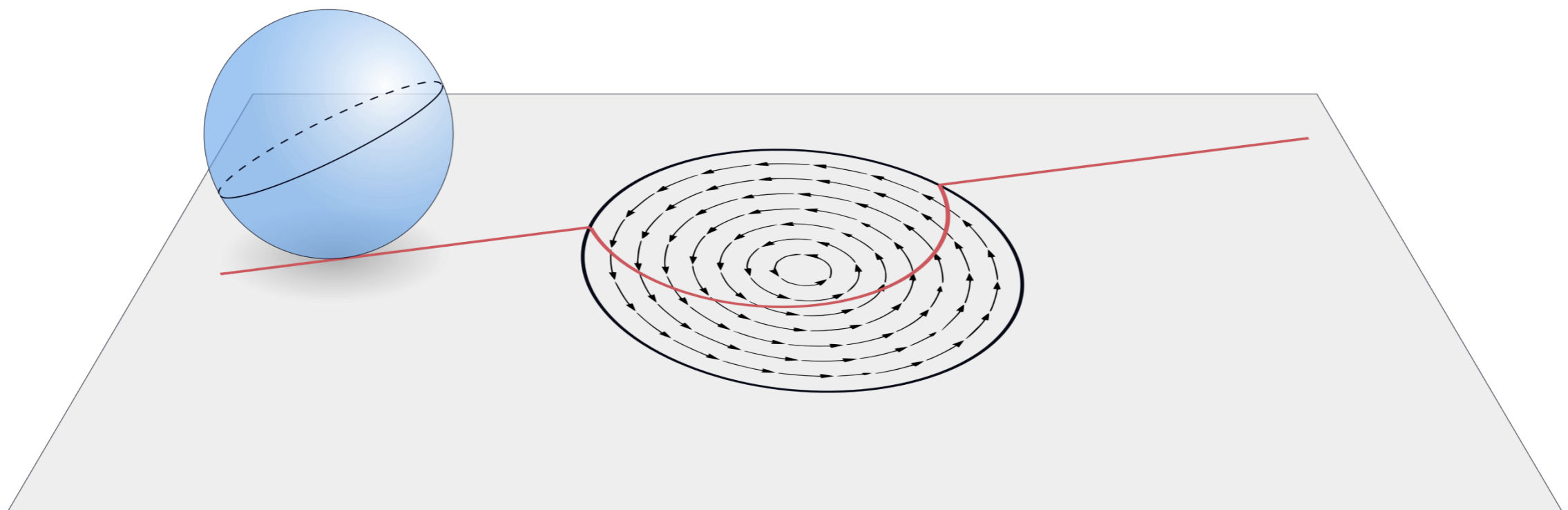
Homogeneous sphere on rotating plane

Anais-billiard phenomenon
(Lévy- Leblond, 1986)



Homogeneous sphere on rotating plane

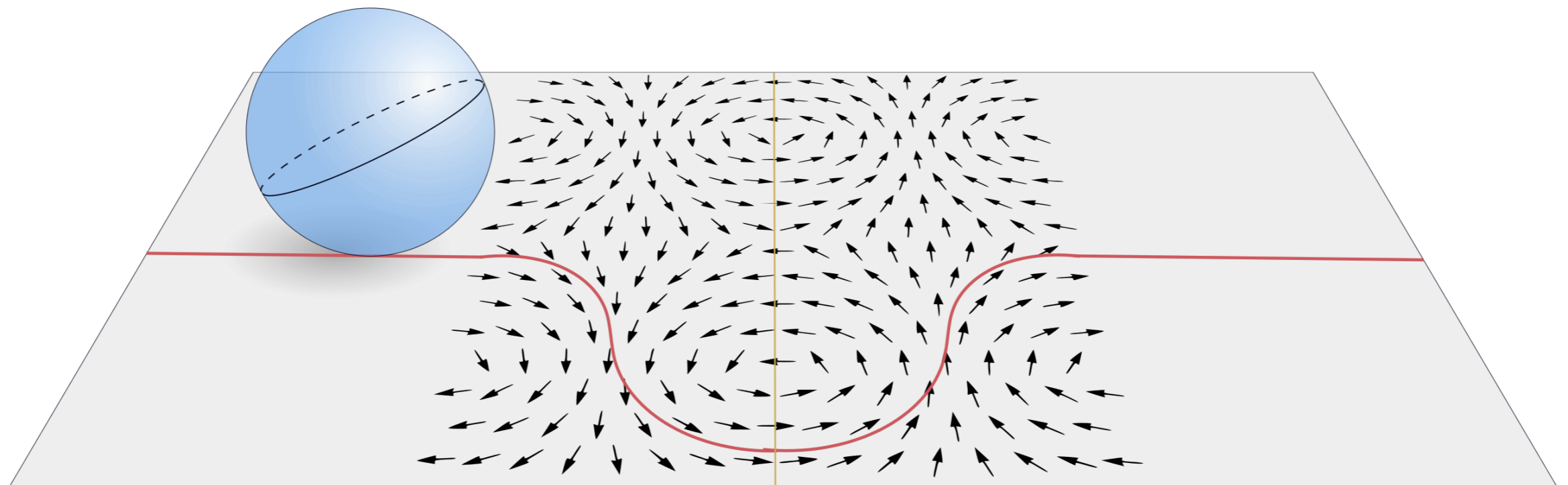
Anais-billiard phenomenon
(Lévy- Leblond, 1986)



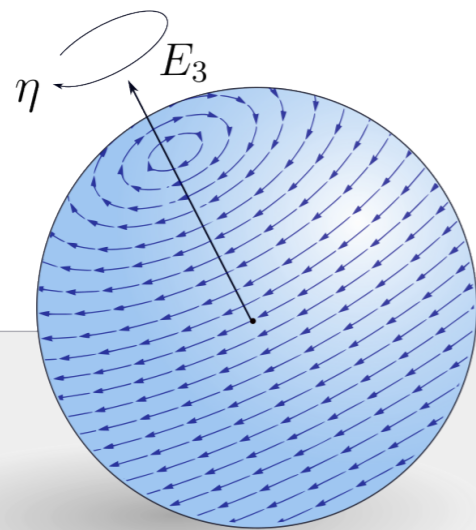
Homogeneous sphere in plane with symmetric vector field

Anais-billiard phenomenon
(Lévy- Leblond, 1986)

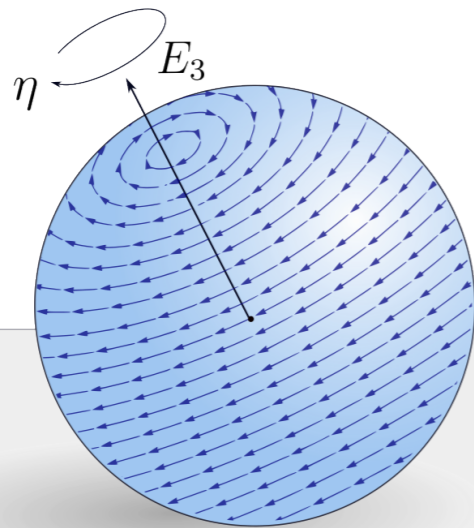
$$\begin{pmatrix} V_1(x, y) \\ V_2(x, y) \end{pmatrix} = \begin{pmatrix} V_1(-x, y) \\ -V_2(-x, y) \end{pmatrix}$$



What happens if the shell of the sphere is moving?

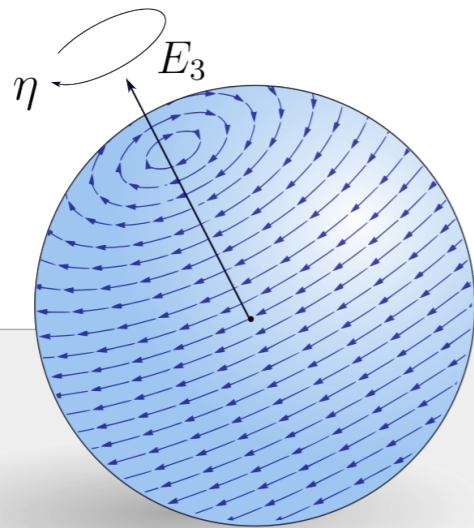


What happens if the shell of the sphere is moving?



- Lose symmetry on the sphere
- Lose energy first integral

What happens if the shell of the sphere is moving?



- Lose symmetry on the sphere
- Lose energy first integral
- Moving energy

Homogeneous sphere with rotating shell

Equations of motion:

$$\dot{M} = M \times \Omega$$

$$\dot{\gamma} = \gamma \times \Omega$$

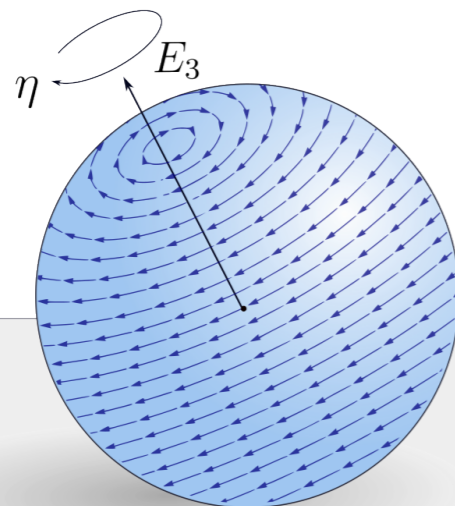
$$M = I\Omega + mr^2\gamma \times (\Omega \times \gamma) - mr^2\eta_1\gamma \times (\gamma \times E_3)$$

First integrals:

$$\|M\|^2, \quad \langle M, \gamma \rangle, \quad \|\gamma\|^2, \quad E_{mov}$$

Invariant measure:

$$dMdxddy$$



Homogeneous sphere with rotating shell

Integrable

Equations of motion:

$$\dot{M} = M \times \Omega$$

$$\dot{\gamma} = \gamma \times \Omega$$

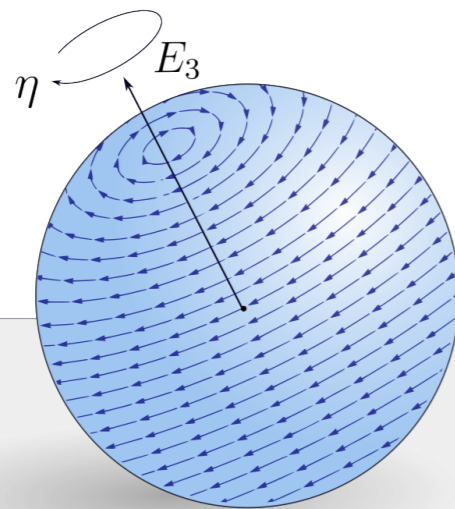
$$M = I\Omega + mr^2\gamma \times (\Omega \times \gamma) - mr^2\eta_1\gamma \times (\gamma \times E_3)$$

First integrals:

$$\|M\|^2, \quad \langle M, \gamma \rangle, \quad \|\gamma\|^2, \quad E_{mov}$$

Invariant measure:

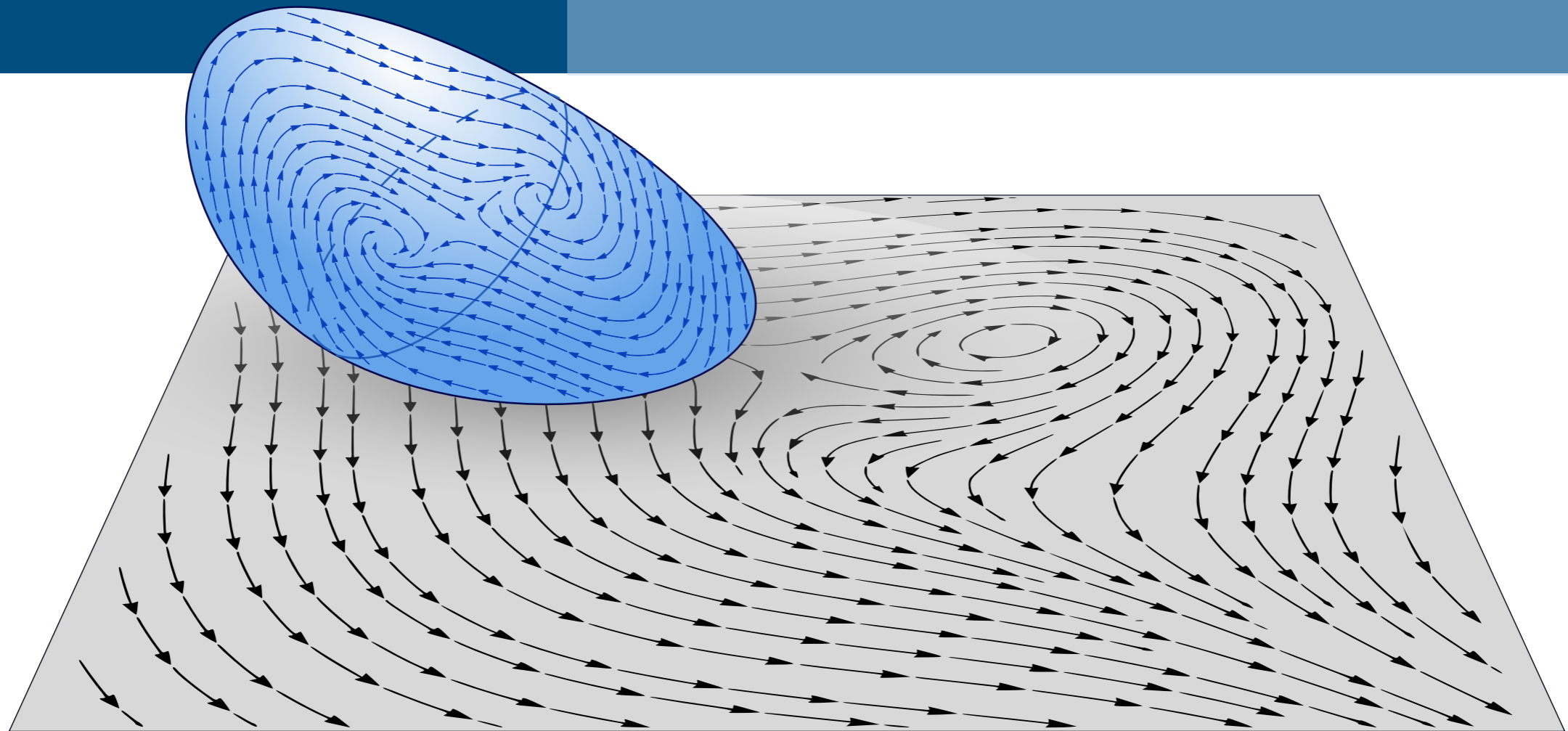
$$dMdxddy$$



Generalization

Convex body rolling on a plane

Add general vector fields to plane and body shell

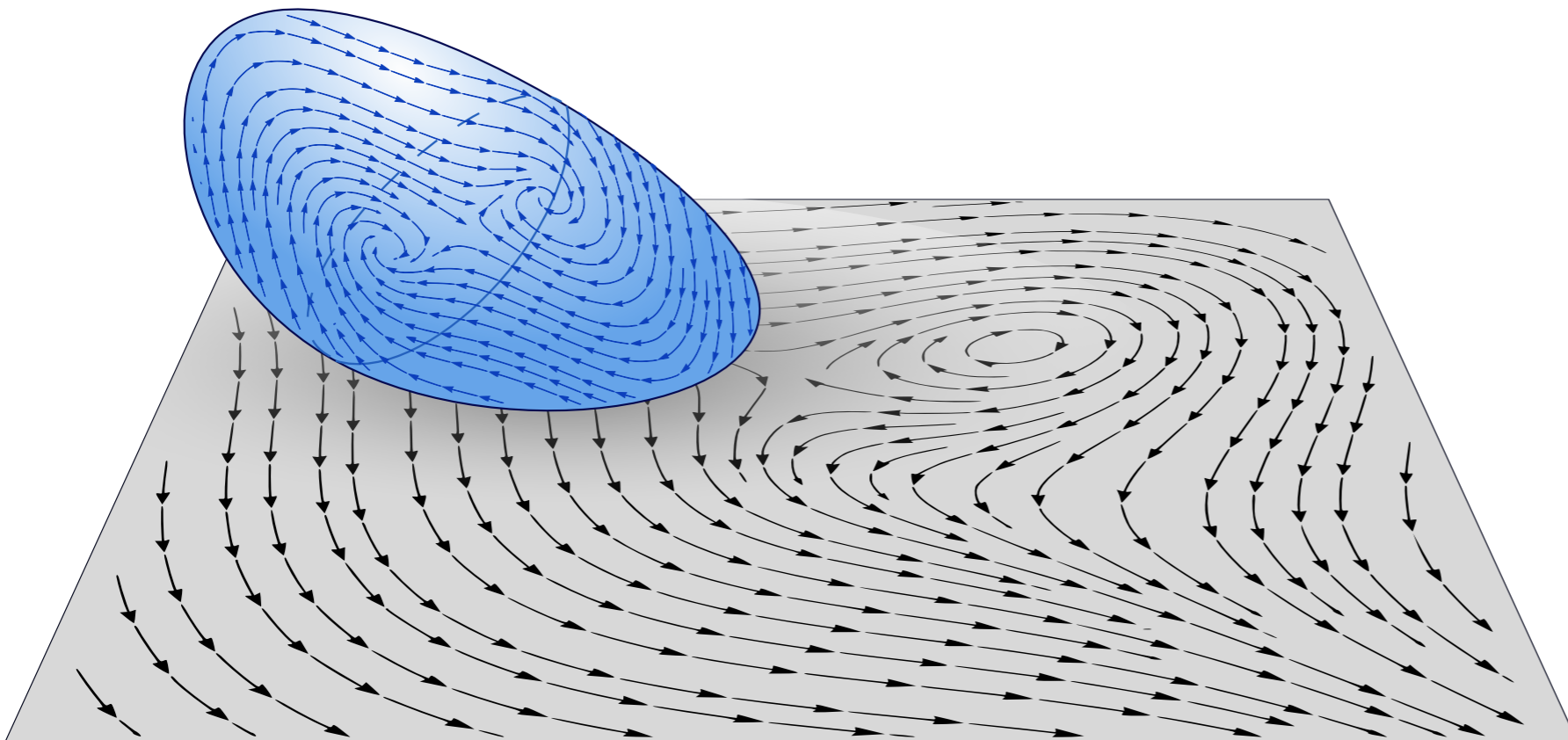


Generalization

Goal:

Identify general principles leading to

- First integrals
- Moving energy
- Invariant measure



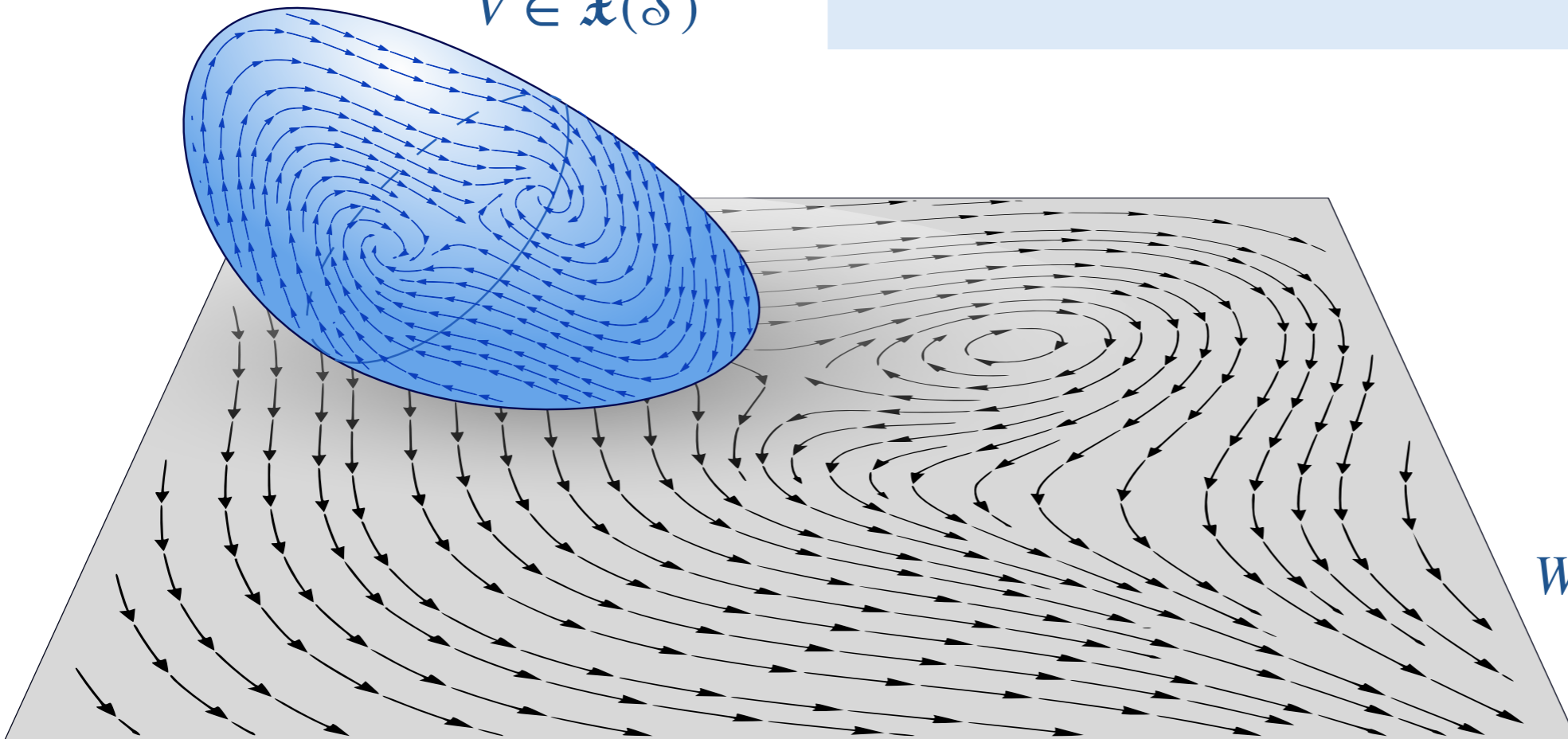
Generalization

Configuration manifold: $Q = SO(3) \times \mathbb{R}^2$
 $(B, \underline{x}) \in Q$

Affine nonholonomic constraint:
 $\dot{i} = B(\rho \times \Omega) + BV + W$

$V \in \mathfrak{X}(\mathcal{S})$

$W \in \mathfrak{X}(\mathbb{R}^2)$



Generalization

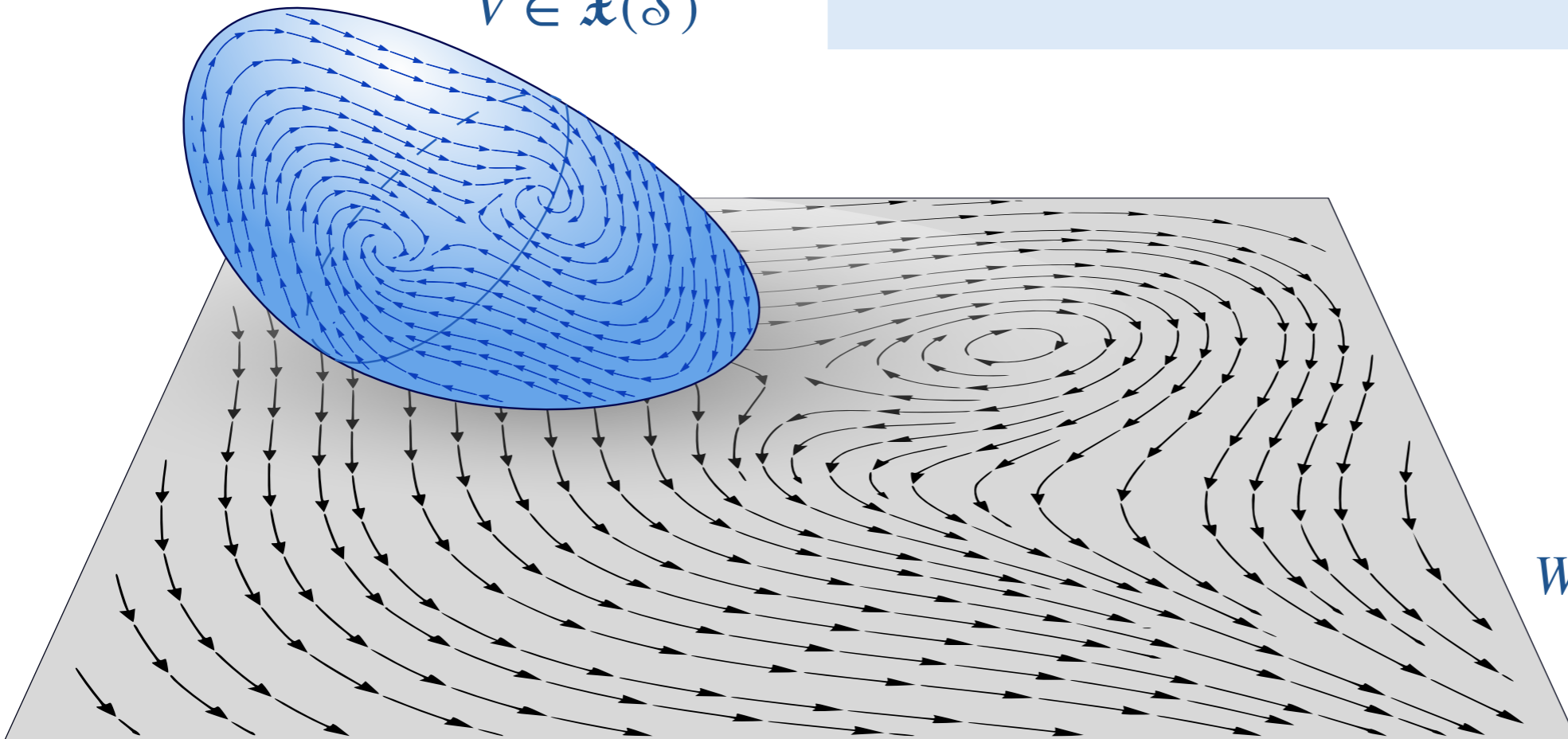
No symmetries on
plane or body

Configuration manifold: $Q = SO(3) \times \mathbb{R}^2$
 $(B, \underline{x}) \in Q$

Affine nonholonomic constraint:
 $\dot{i} = B(\rho \times \Omega) + BV + W$

$V \in \mathfrak{X}(\mathcal{S})$

$W \in \mathfrak{X}(\mathbb{R}^2)$



General system

Theorem (C., García Naranjo, 2023):

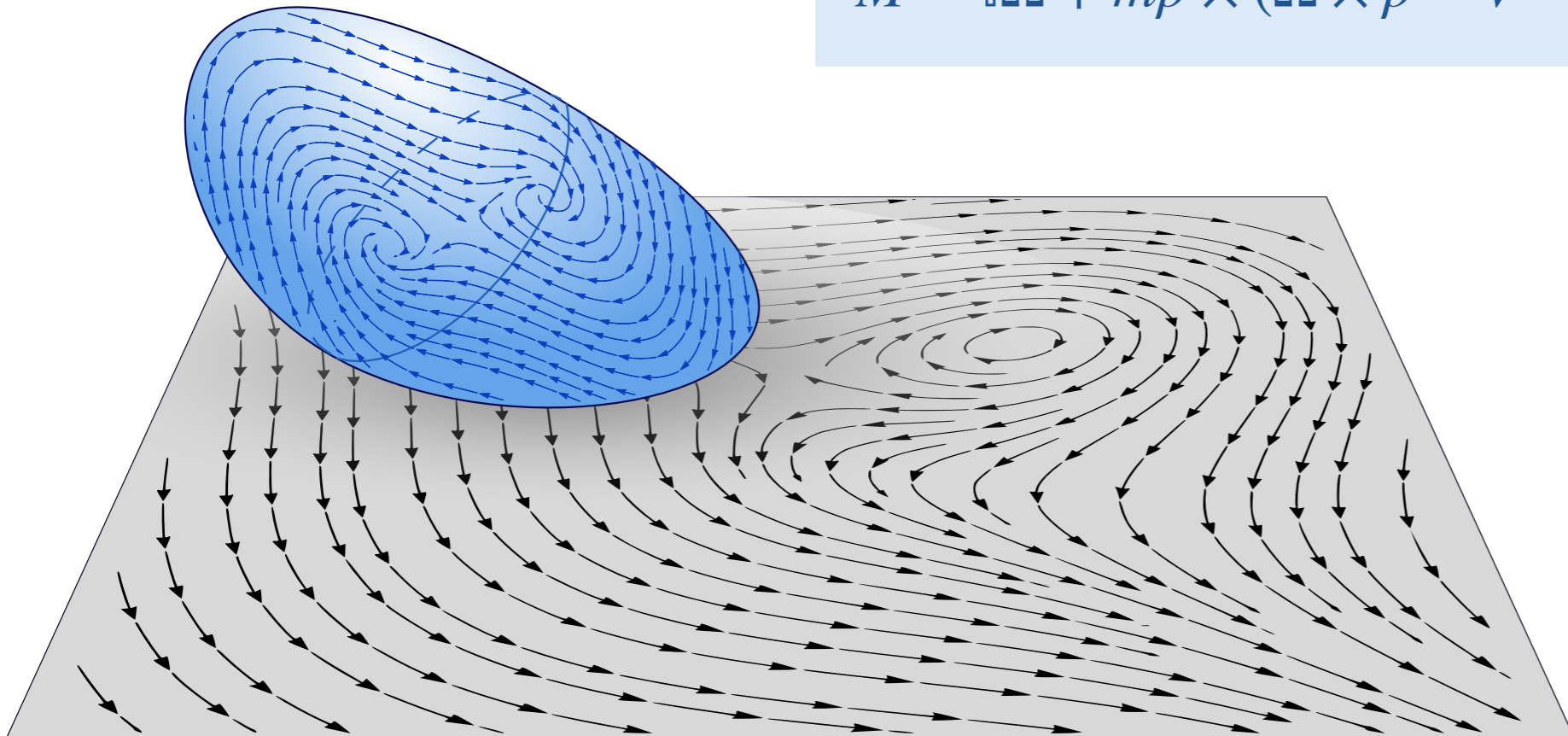
Equations of motion :

$$\dot{M} = M \times \Omega + m\dot{\rho} \times (\Omega \times \rho) + mg\rho \times \gamma \\ + m(V + B^{-1}W) \times (\dot{\rho} + \Omega \times \rho)$$

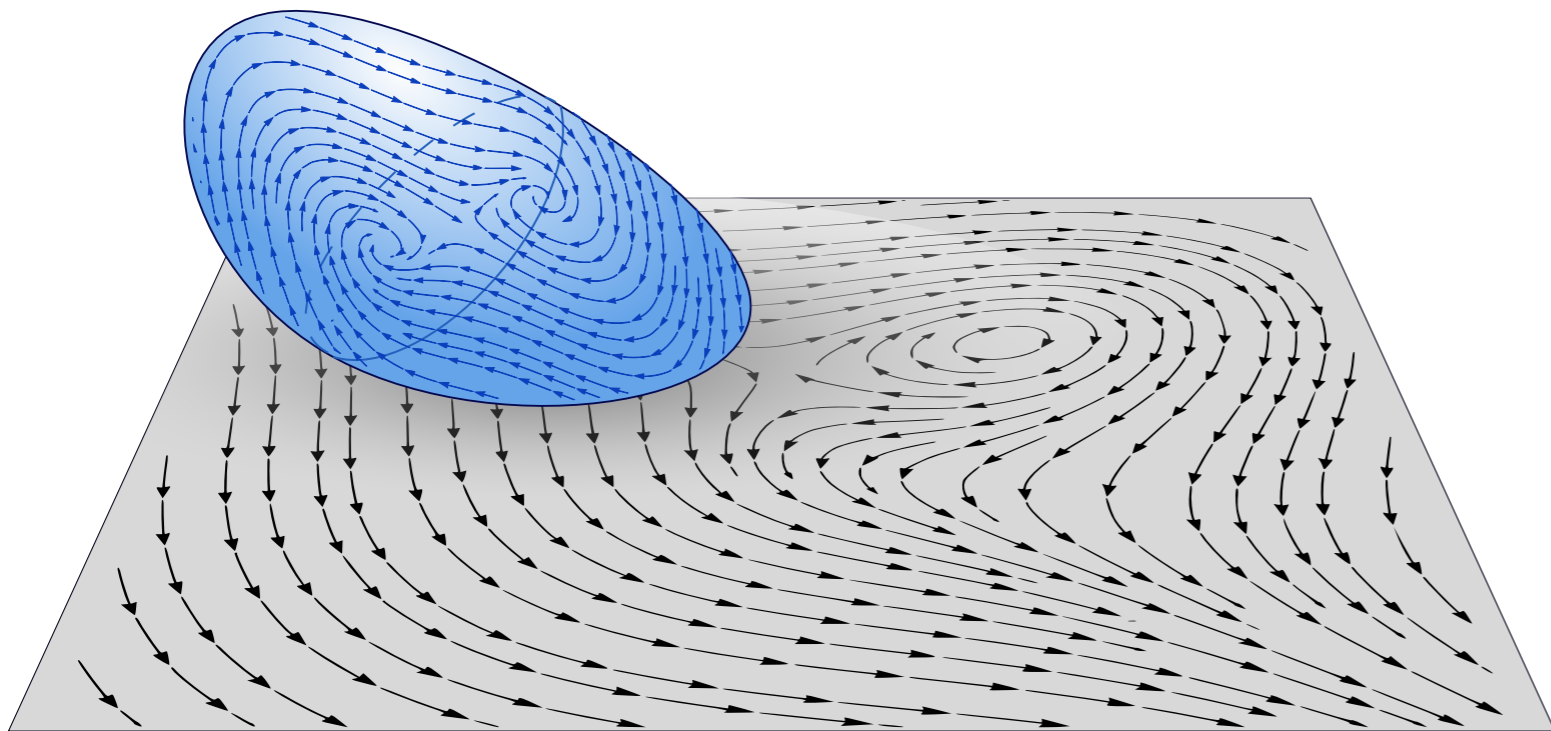
$$\dot{u} = B(\rho \times \Omega) + BV + W$$

$$\dot{B} = B\hat{\Omega}$$

$$M = \mathbb{I}\Omega + m\rho \times (\Omega \times \rho - V - B^{-1}W)$$

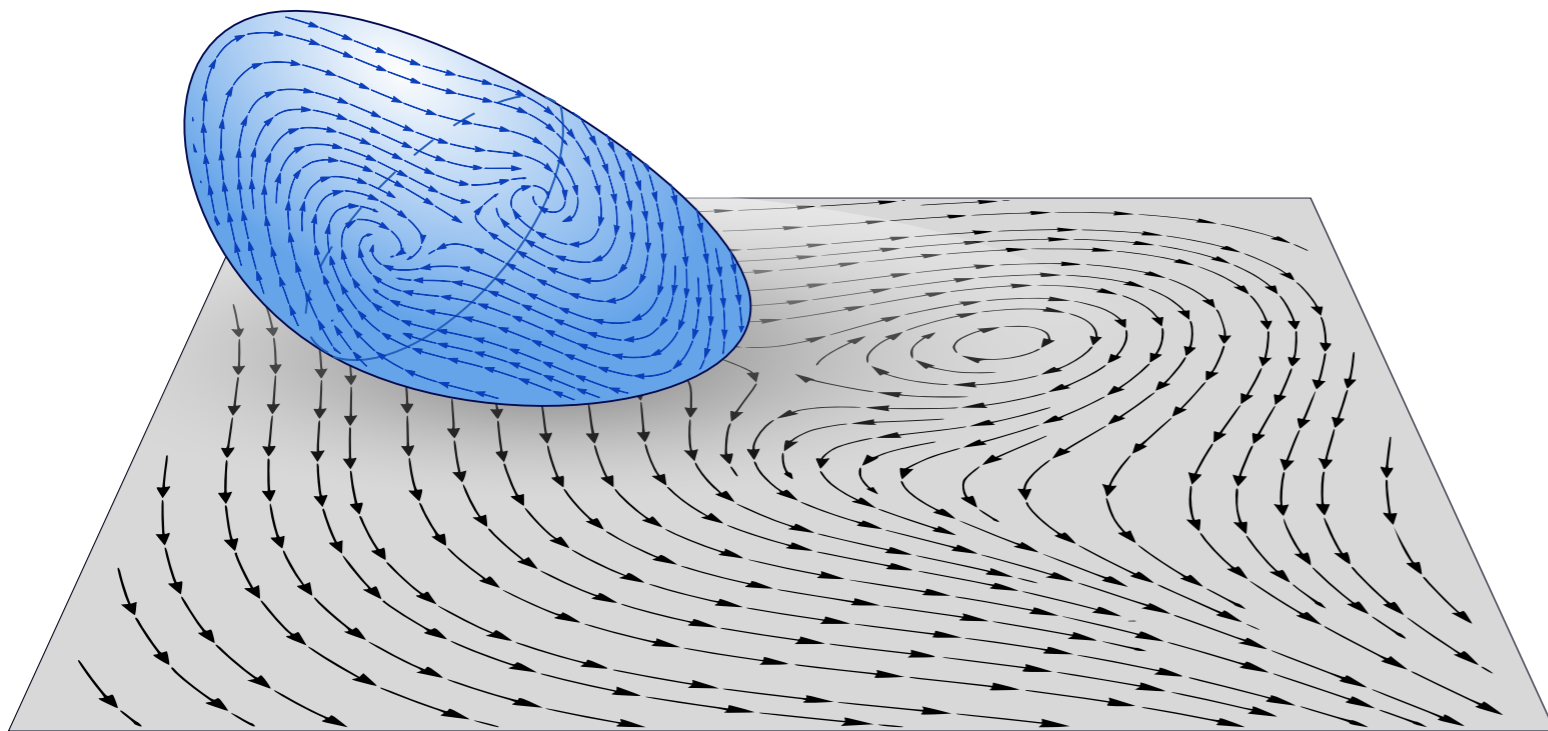


General system



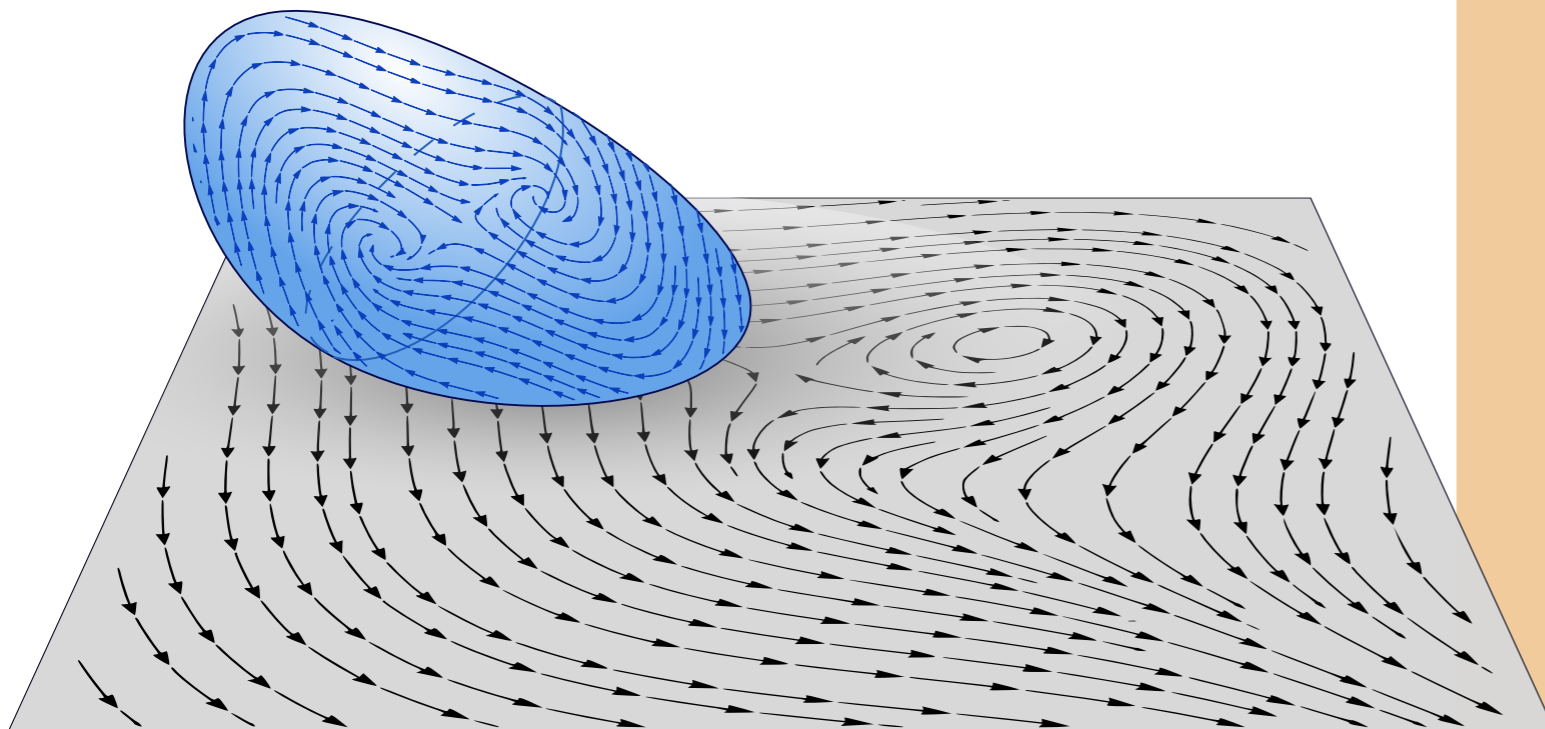
General system

- No invariant measure
- Generally chaotic



General system

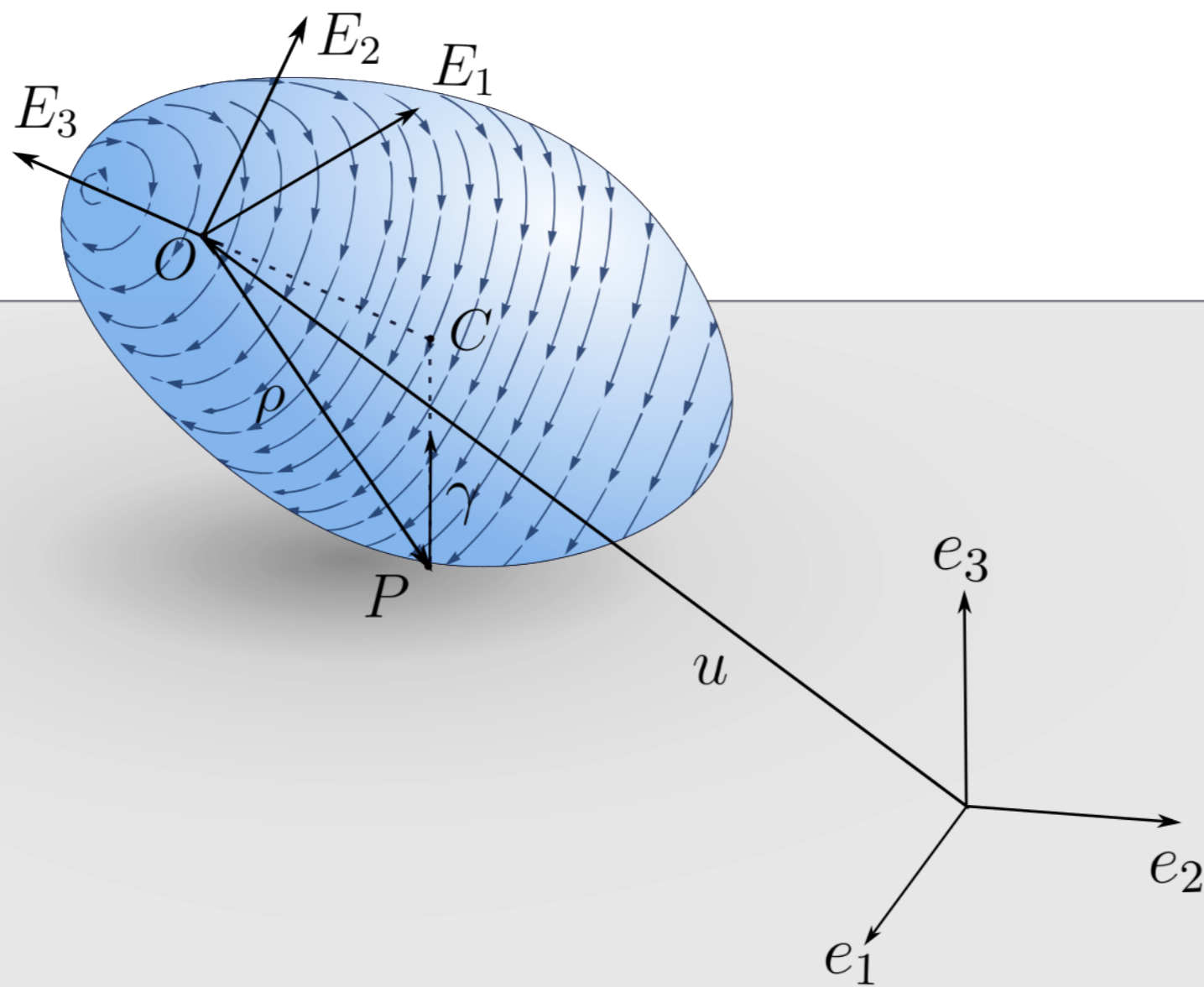
- No invariant measure
- Generally chaotic



For it to be integrable we need extra symmetries that lead to a sufficient number of:

- First integrals
- Invariant measure

Body of revolution with rotating shell



Body of revolution with rotating shell

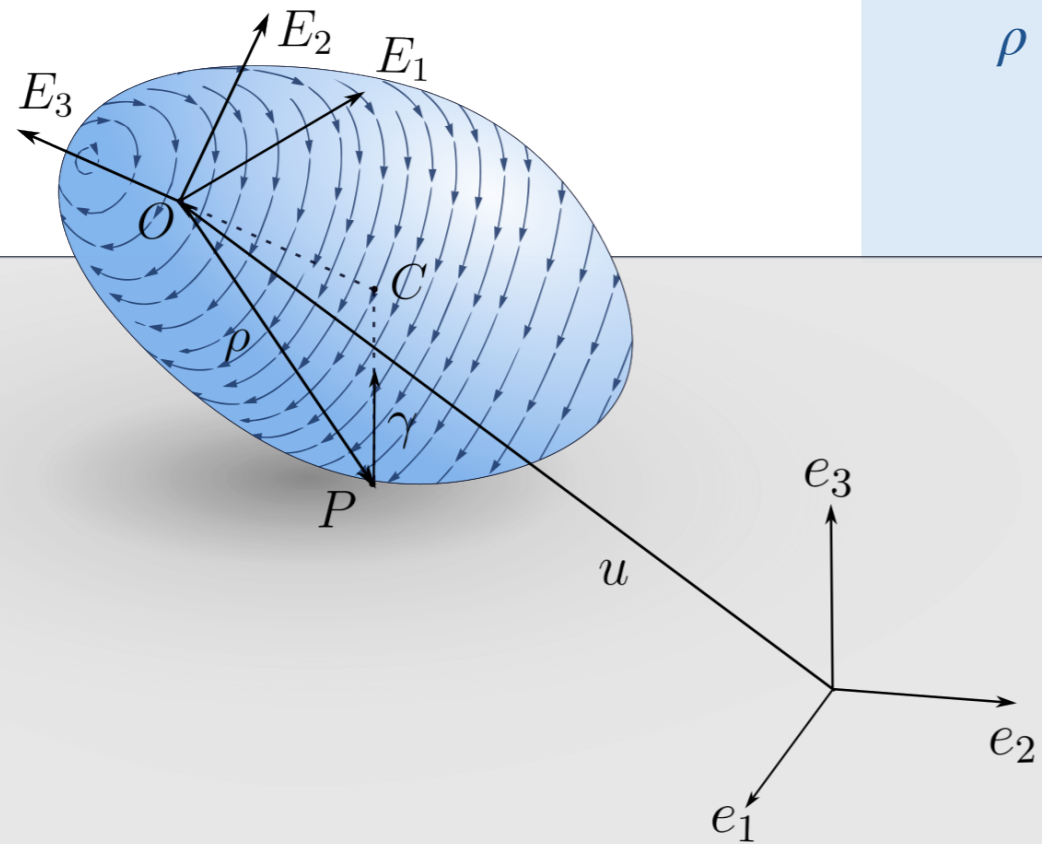
Equations of motion:

$$\dot{M} = M \times \Omega + m\dot{\rho} \times (\Omega \times \rho) + mg\rho \times \gamma \\ + m\eta(\rho \times E_3) \times (\dot{\rho} + \Omega \times \rho)$$

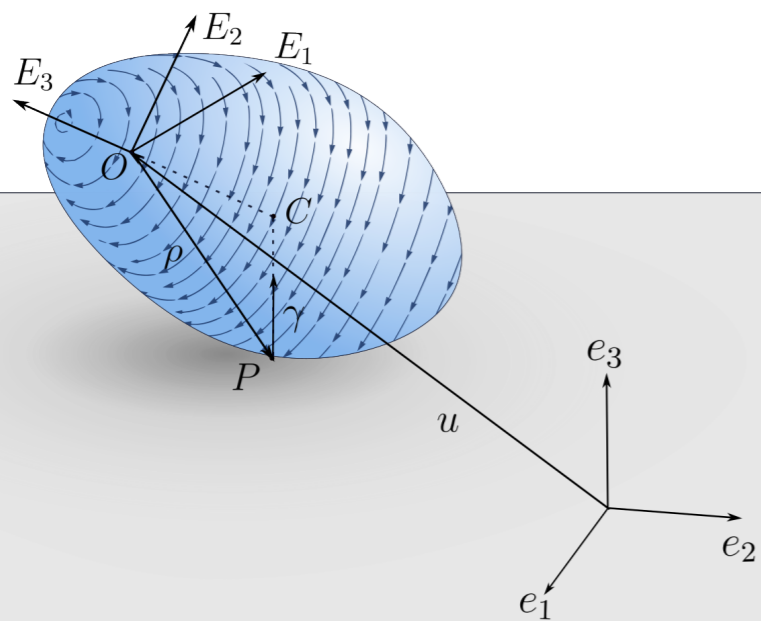
$$\dot{\gamma} = \gamma \times \Omega$$

$$M = \mathbb{I}\Omega + mr^2\gamma \times (\Omega \times \gamma) + mr\eta\gamma \times (\rho \times E_3)$$

$$\rho = (f_1(\gamma_3)\gamma_1, f_1(\gamma_3)\gamma_2, f_2(\gamma_3))$$



Body of revolution with rotating shell



First integrals:

$$\begin{pmatrix} J_1 \\ J_2 \end{pmatrix} = U^{-1} \left(\begin{pmatrix} K_1 \\ K_2 \end{pmatrix} - \eta u \right),$$

$$K_1 = \frac{\langle M, \rho \rangle}{f_1}, \quad K_2 = \frac{\Omega_3}{\sqrt{I_1 I_3 + m \langle \rho, \mathbb{I} \rho \rangle}},$$

U solution matrix of $U' = G(\gamma_3)U$, $U(0) = Id$

u solution of $u' = Gu + b$.

$$G = G(\gamma_3), \quad b = b(\gamma_3)$$

Body of revolution with rotating shell

Theorem (C., García Naranjo, 2023):

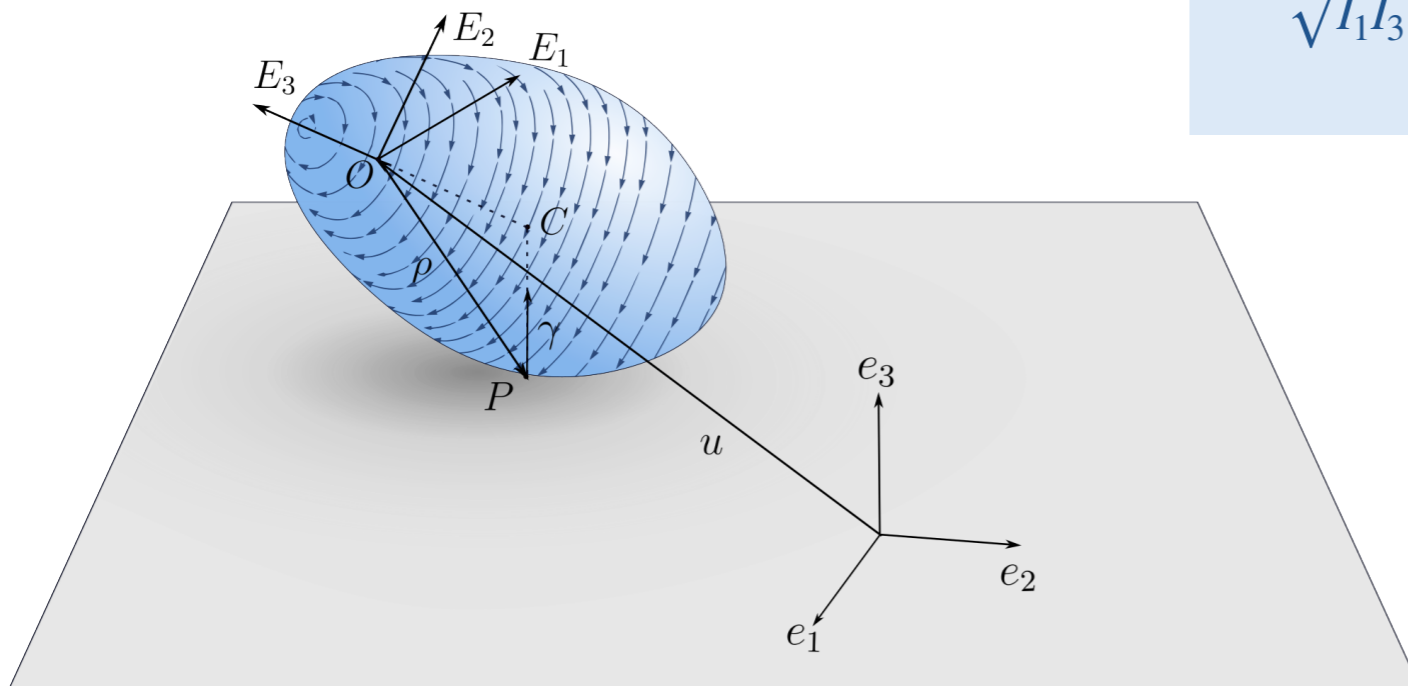
First integrals:

$$\|\gamma\|^2, J_1, J_2$$

$$E_{mov} = \frac{1}{2} \langle \mathbb{I}(\Omega + \eta E_3), \Omega + \eta E_3 \rangle + \frac{m}{2} \|\rho \times (\Omega + \eta E_3)\|^2 - mr \langle \rho, \gamma \rangle$$

Invariant measure:

$$\frac{1}{\sqrt{I_1 I_3 + m \langle \rho, \mathbb{I} \rho \rangle}} dM d\gamma$$



Body of revolution with rotating shell

Integrable

Theorem (C., García Naranjo, 2023):

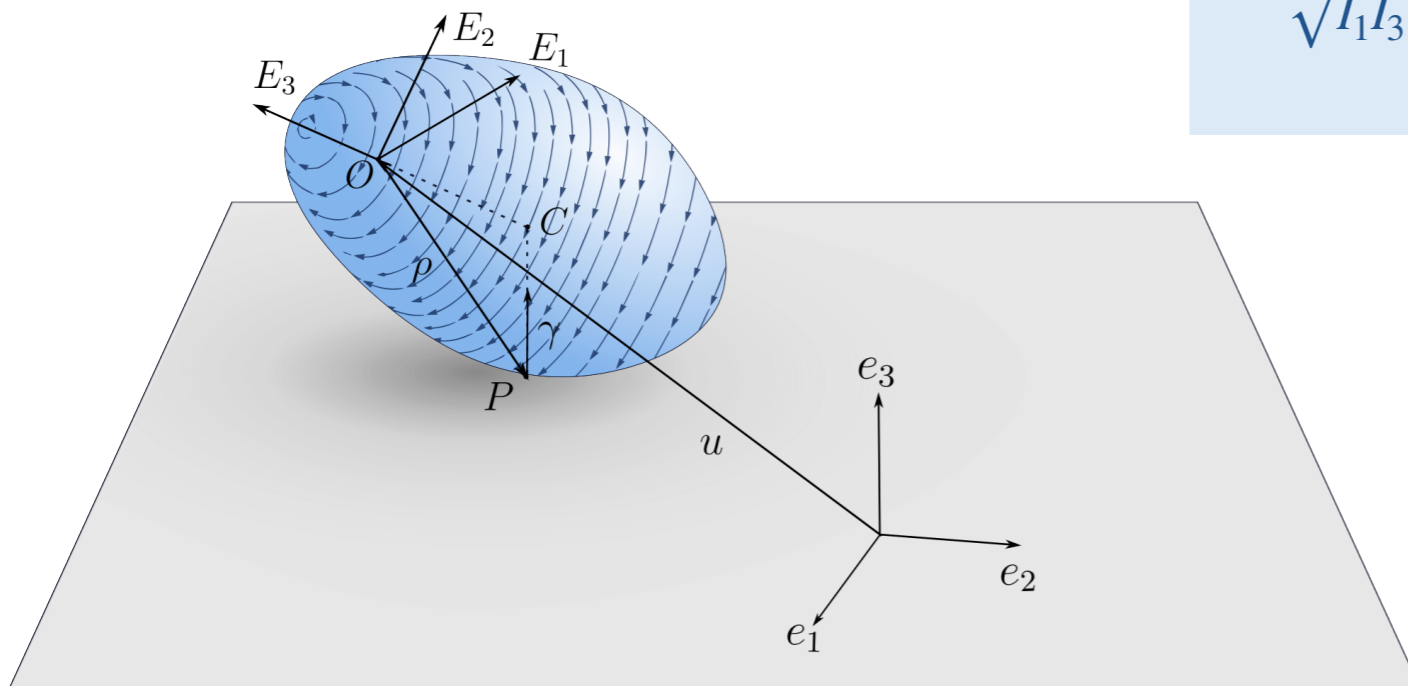
First integrals:

$$\|\gamma\|^2, J_1, J_2$$

$$E_{mov} = \frac{1}{2} \langle \mathbb{I}(\Omega + \eta E_3), \Omega + \eta E_3 \rangle + \frac{m}{2} \|\rho \times (\Omega + \eta E_3)\|^2 - mr \langle \rho, \gamma \rangle$$

Invariant measure:

$$\frac{1}{\sqrt{I_1 I_3 + m \langle \rho, \mathbb{I} \rho \rangle}} dM d\gamma$$



Routh's sphere with rotating shell

$$\rho(\gamma) = -r\gamma - lE_3$$

First integrals

$$||\gamma||^2,$$

$$J_1 = \frac{1}{r}\langle M, \rho \rangle,$$

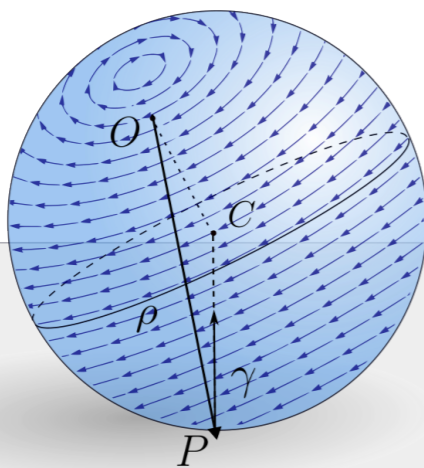
$$J_2 = \Omega_3 + \eta \left(\frac{I_1 I_3 \sqrt{m}}{(I_1 - I_3)^{3/2}} \arctan \left(\frac{\mu'}{r\sqrt{m(I_1 - I_3)}} \right) - \frac{I_1}{I_1 - I_3} \mu \right),$$

$$\mu = \sqrt{I_1 I_3 + m\langle \rho, \mathbb{1}\rho \rangle}$$

$$E_{mov} = \frac{1}{2}\langle \mathbb{1}(\Omega + \eta E_3), \Omega + \eta E_3 \rangle + \frac{m}{2} ||\rho \times (\Omega + \eta E_3)||^2 - mr\langle \rho, \gamma \rangle$$

Invariant measure

$$\frac{1}{\sqrt{I_1 I_3 + m\langle \rho, \mathbb{1}\rho \rangle}} dM d\gamma$$



Routh's sphere with rotating shell

Integrable

$$\rho(\gamma) = -r\gamma - lE_3$$

First integrals

$$\|\gamma\|^2,$$

$$J_1 = \frac{1}{r}\langle M, \rho \rangle,$$

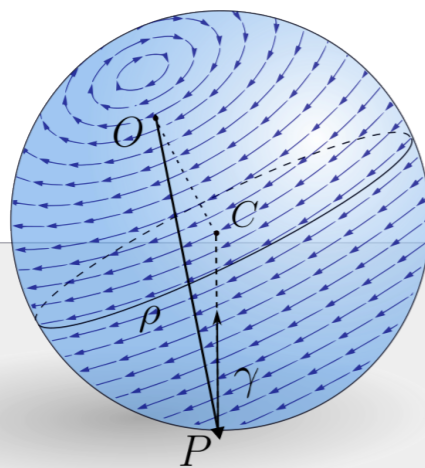
$$J_2 = \Omega_3 + \eta \left(\frac{I_1 I_3 \sqrt{m}}{(I_1 - I_3)^{3/2}} \arctan \left(\frac{\mu'}{r\sqrt{m(I_1 - I_3)}} \right) - \frac{I_1}{I_1 - I_3} \mu \right),$$

$$\mu = \sqrt{I_1 I_3 + m\langle \rho, \mathbb{I}\rho \rangle}$$

$$E_{mov} = \frac{1}{2}\langle \mathbb{I}(\Omega + \eta E_3), \Omega + \eta E_3 \rangle + \frac{m}{2} \|\rho \times (\Omega + \eta E_3)\|^2 - mr\langle \rho, \gamma \rangle$$

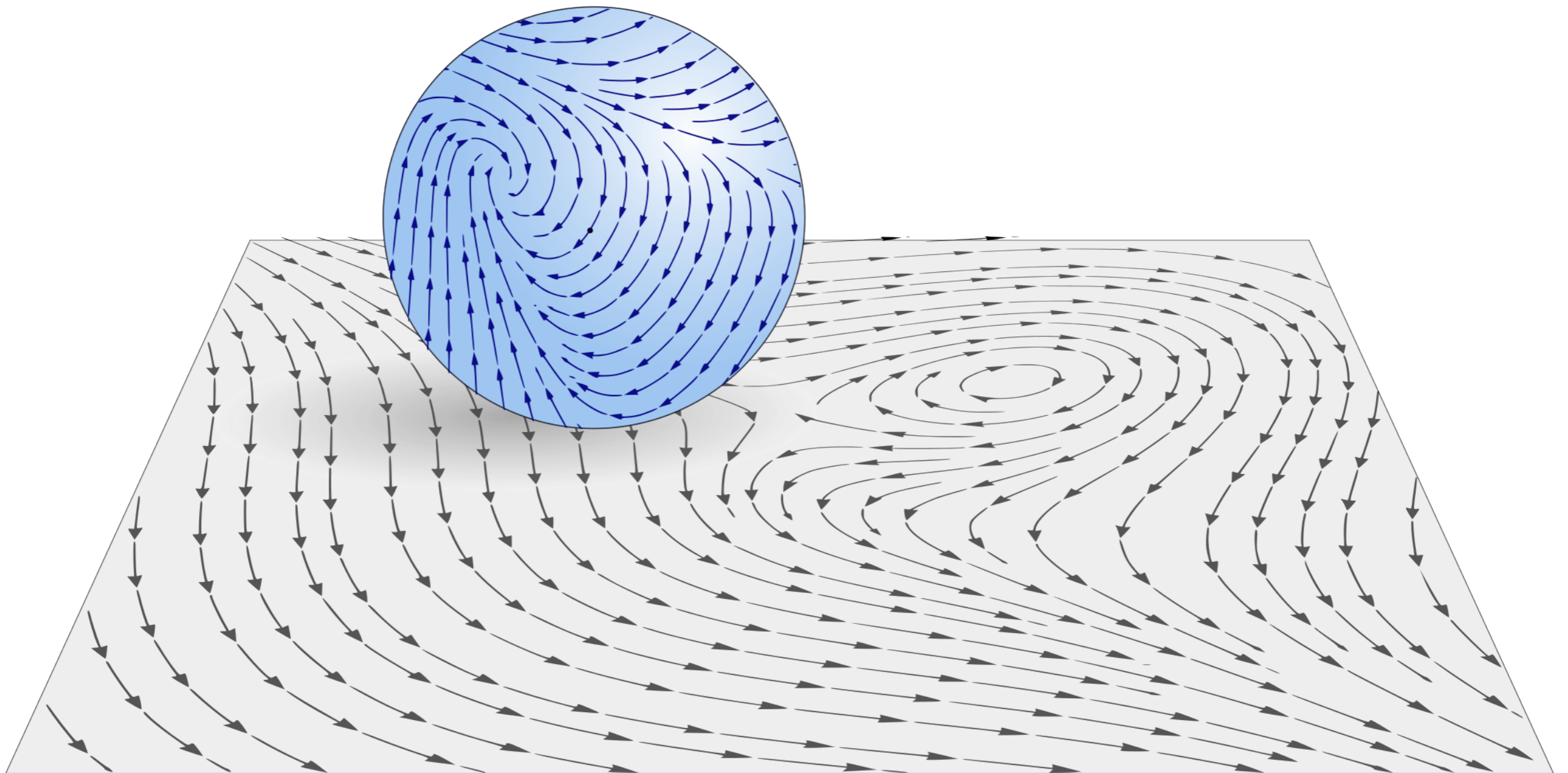
Invariant measure

$$\frac{1}{\sqrt{I_1 I_3 + m\langle \rho, \mathbb{I}\rho \rangle}} dM d\gamma$$



Chaplygin sphere

$$I_1 \neq I_2 \neq I_3$$



Chaplygin sphere

$$I_1 \neq I_2 \neq I_3$$

Equations of motion:

$$\dot{M} = M \times \Omega$$

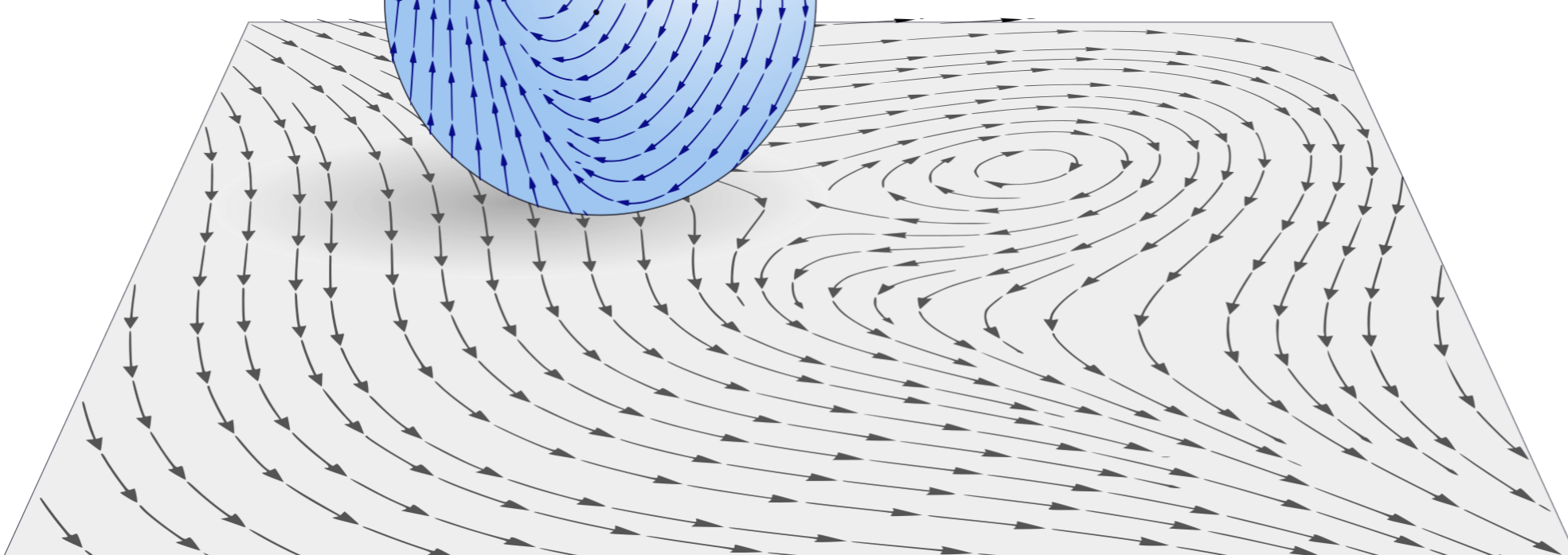
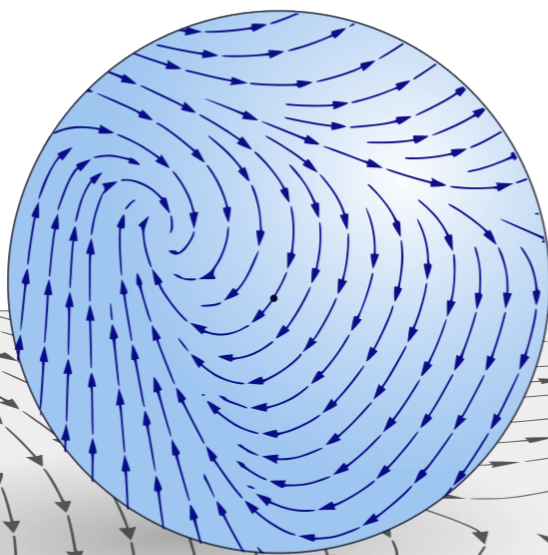
$$\dot{u} = -rB(\gamma \times \Omega) + BV + W$$

$$\dot{B} = B\hat{\Omega}$$

Theorem (C. García Naranjo, 2023):

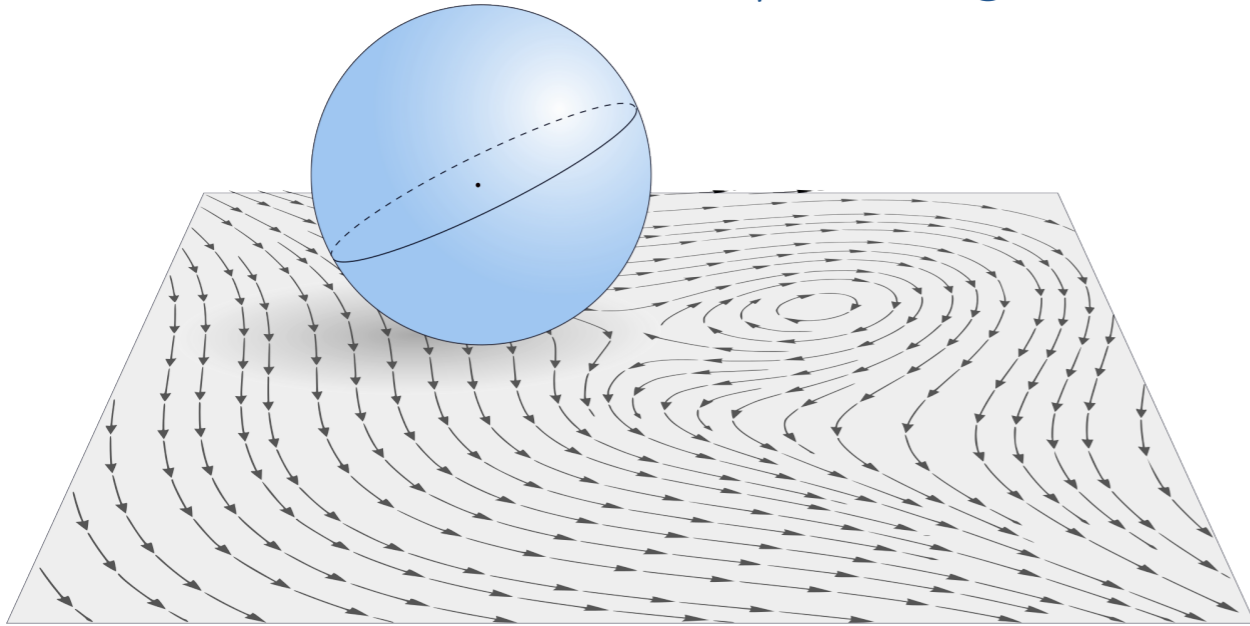
$$M = \mathbb{I}\Omega + mr^2\gamma \times (\Omega \times \gamma) + mr\gamma \times (V + B^{-1}W)$$

is constant in space coordinates

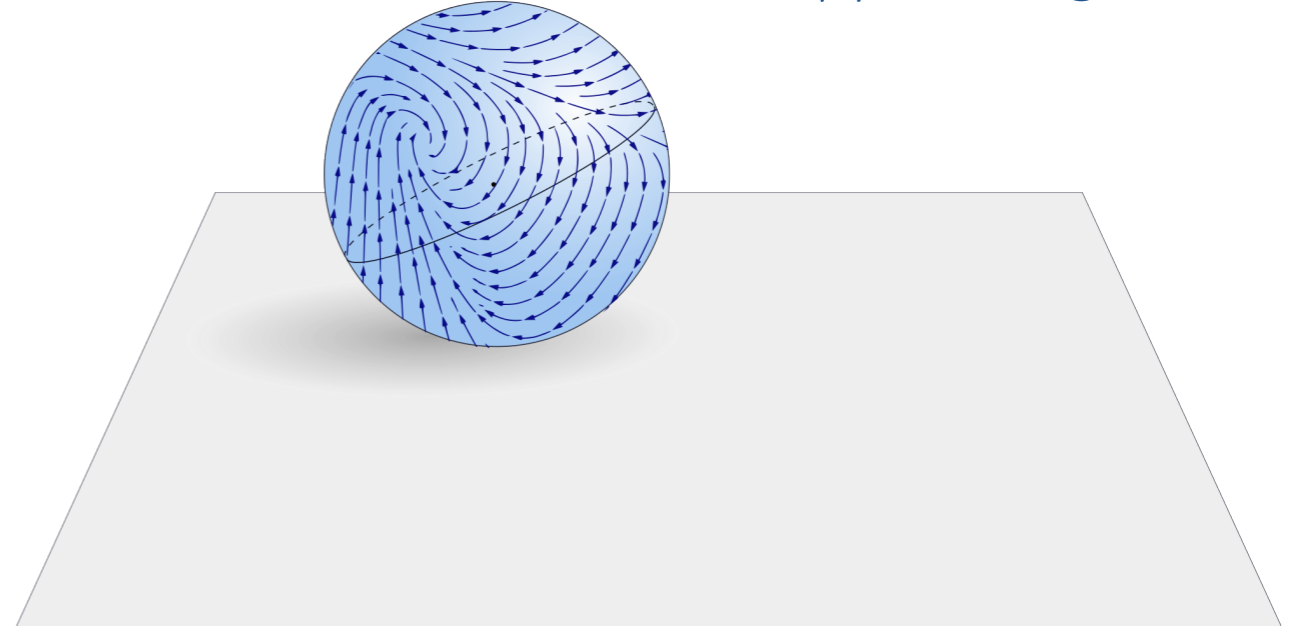


Chaplygin sphere

$$V = 0$$

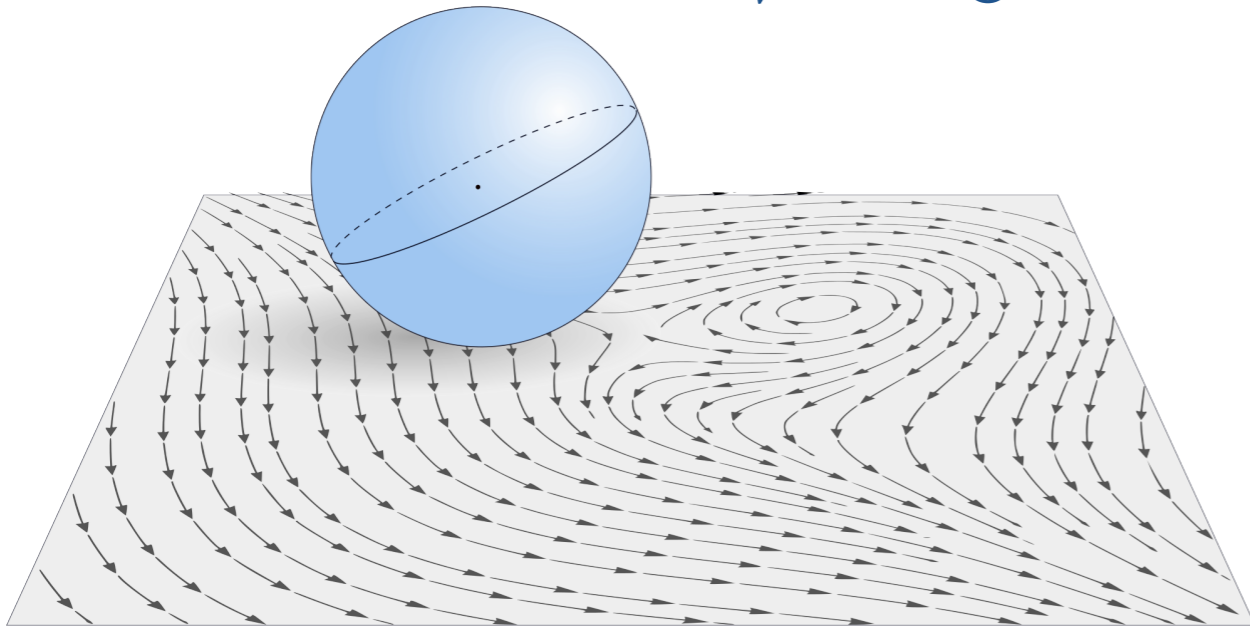


$$W = 0$$

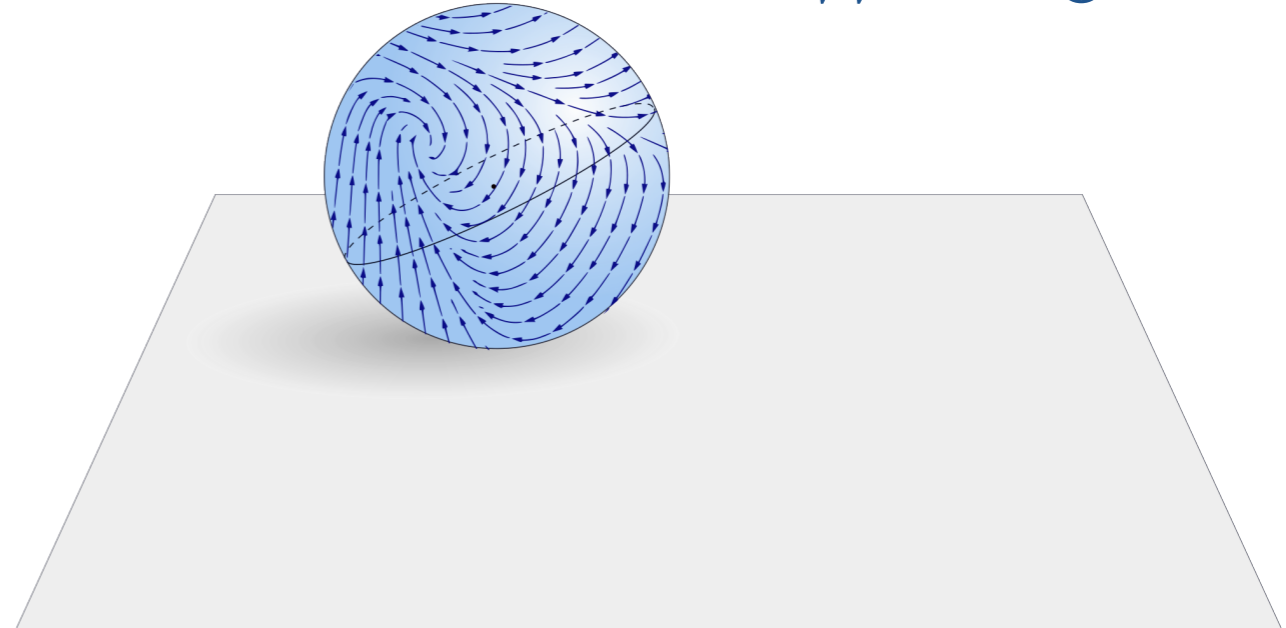


Chaplygin sphere

$$V = 0$$



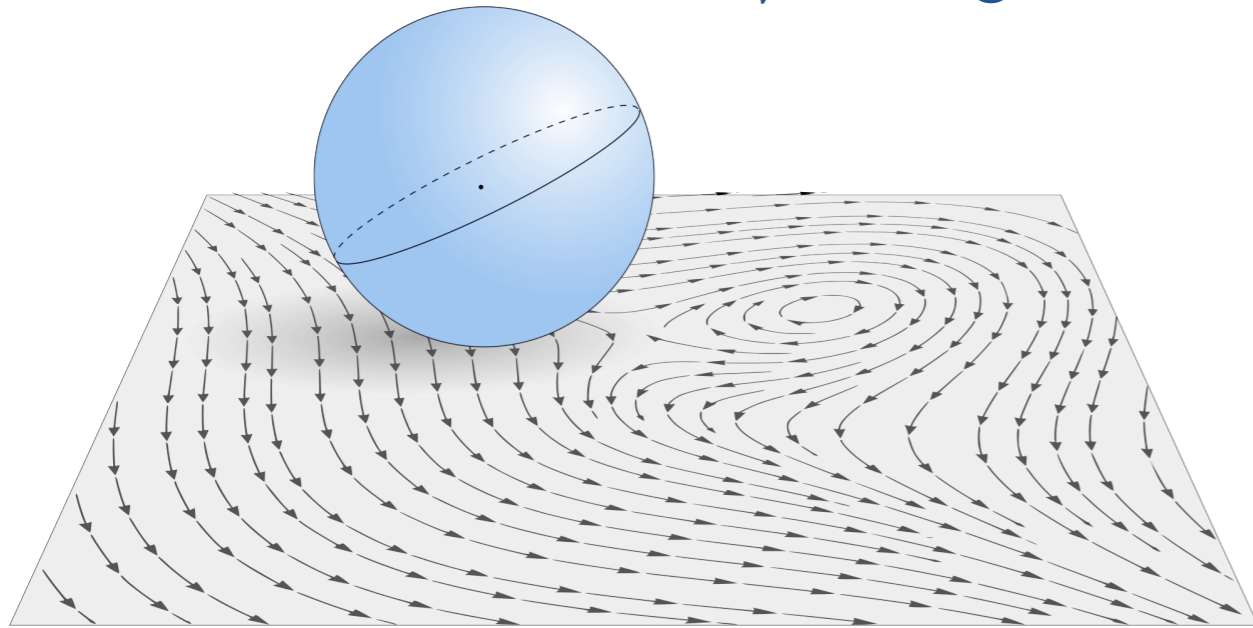
$$W = 0$$



First integrals
Invariant measure

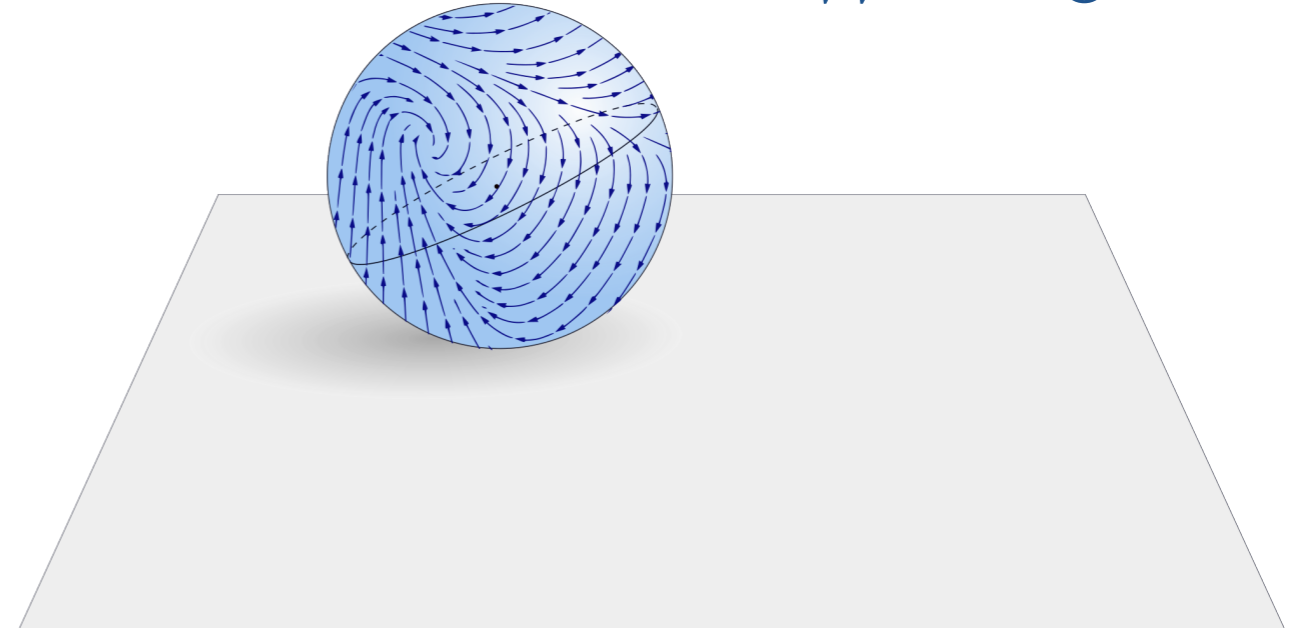
Chaplygin sphere

$$V = 0$$



First integrals
Invariant measure

$$W = 0$$



SE(2)-symmetry
First integrals

Chaplygin sphere

$$V = 0$$

Examples:

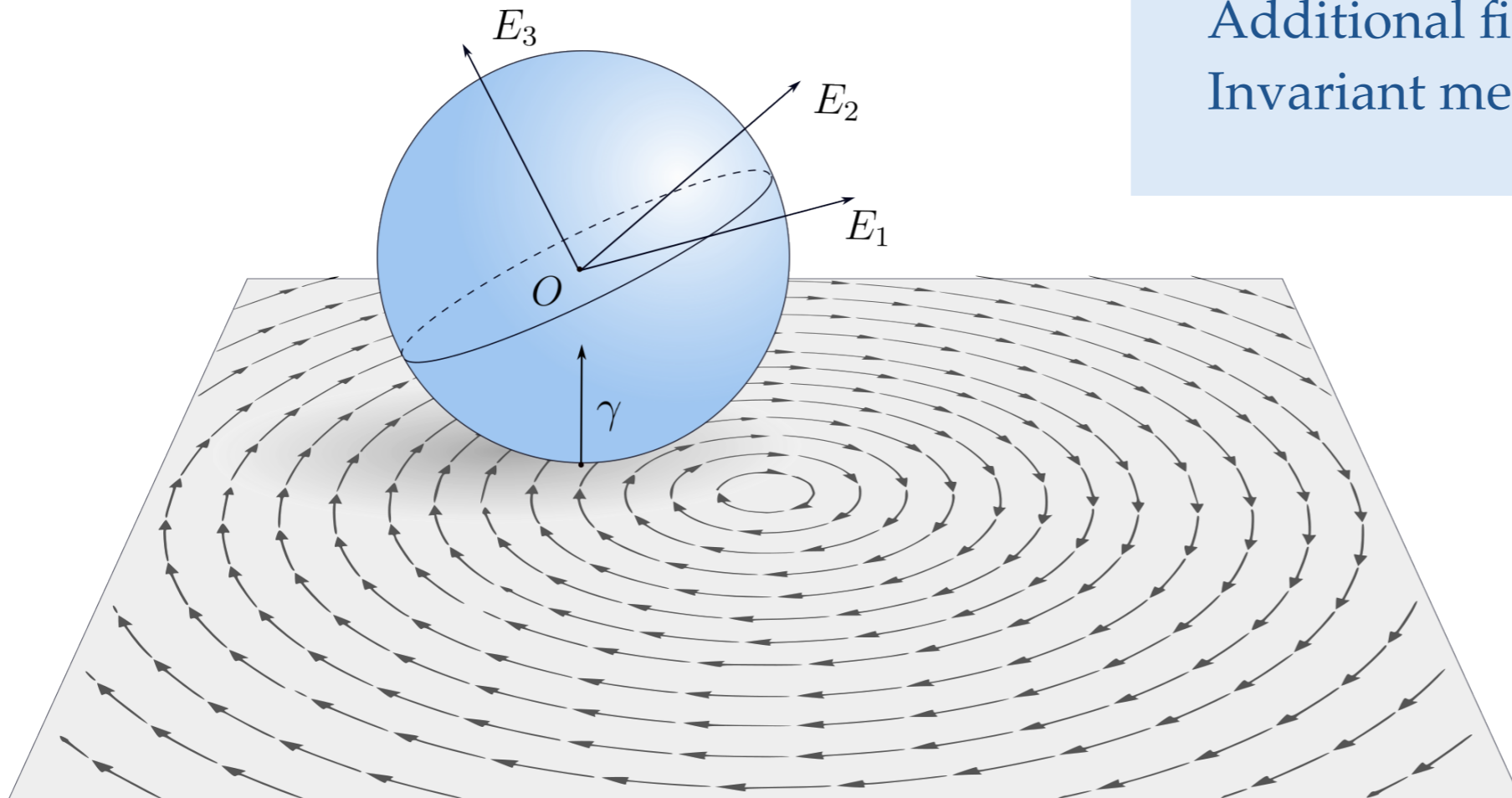
Rotating plane:

Bizyaev, Borisov, Mamaev, 2018

Vibrating plane:

Kilin, Pivovarova, 2021

Additional first integral- moving energy
Invariant measure



Chaplygin sphere

$$V = 0$$

Generally chaotic

Examples:

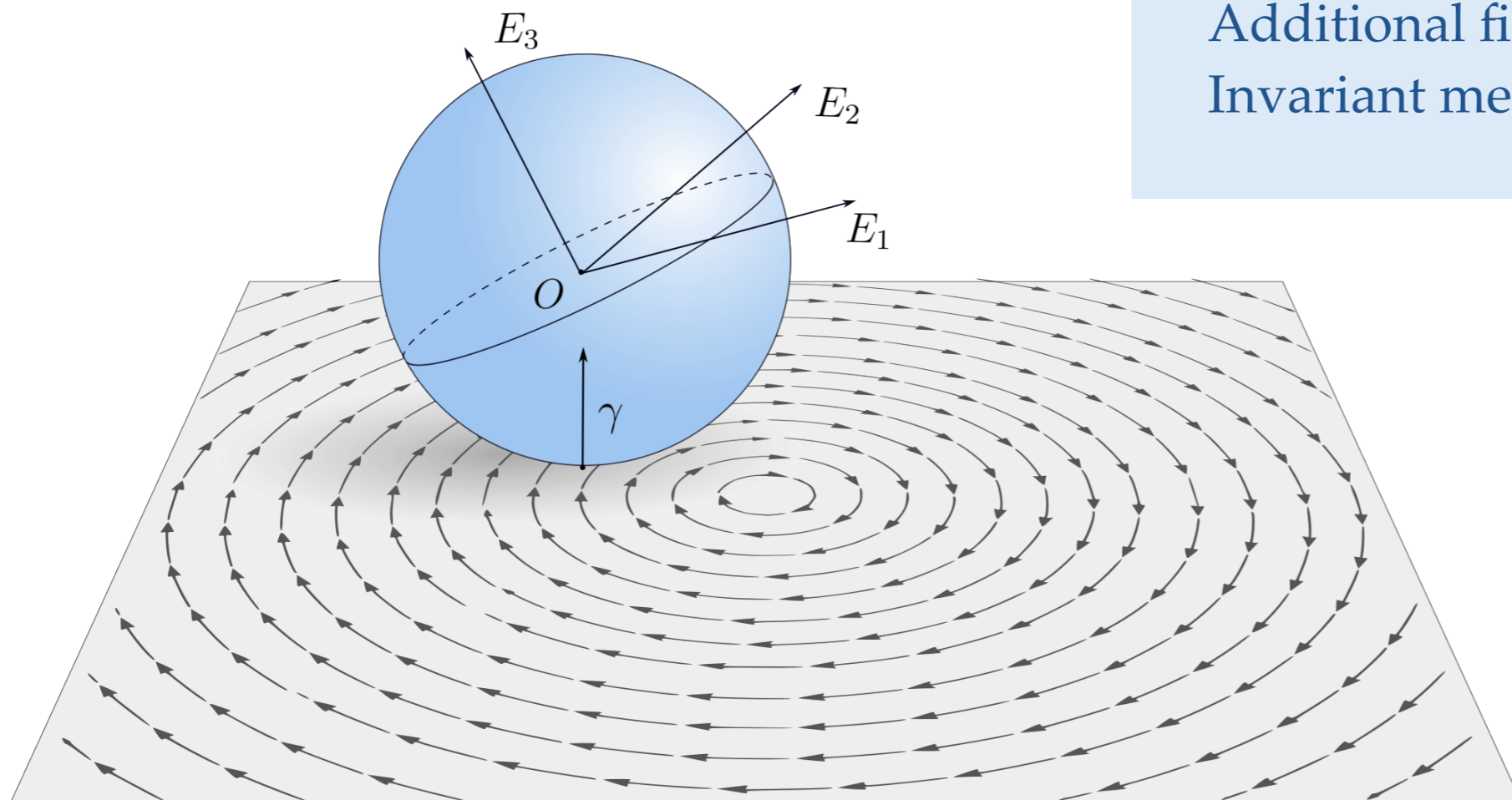
Rotating plane:

Bizyaev, Borisov, Mamaev, 2018

Vibrating plane:

Kilin, Pivovarova, 2021

Additional first integral- moving energy
Invariant measure



Chaplygin sphere

$$V = 0$$

$$W \in \mathfrak{X}(\mathbb{R}^2)$$

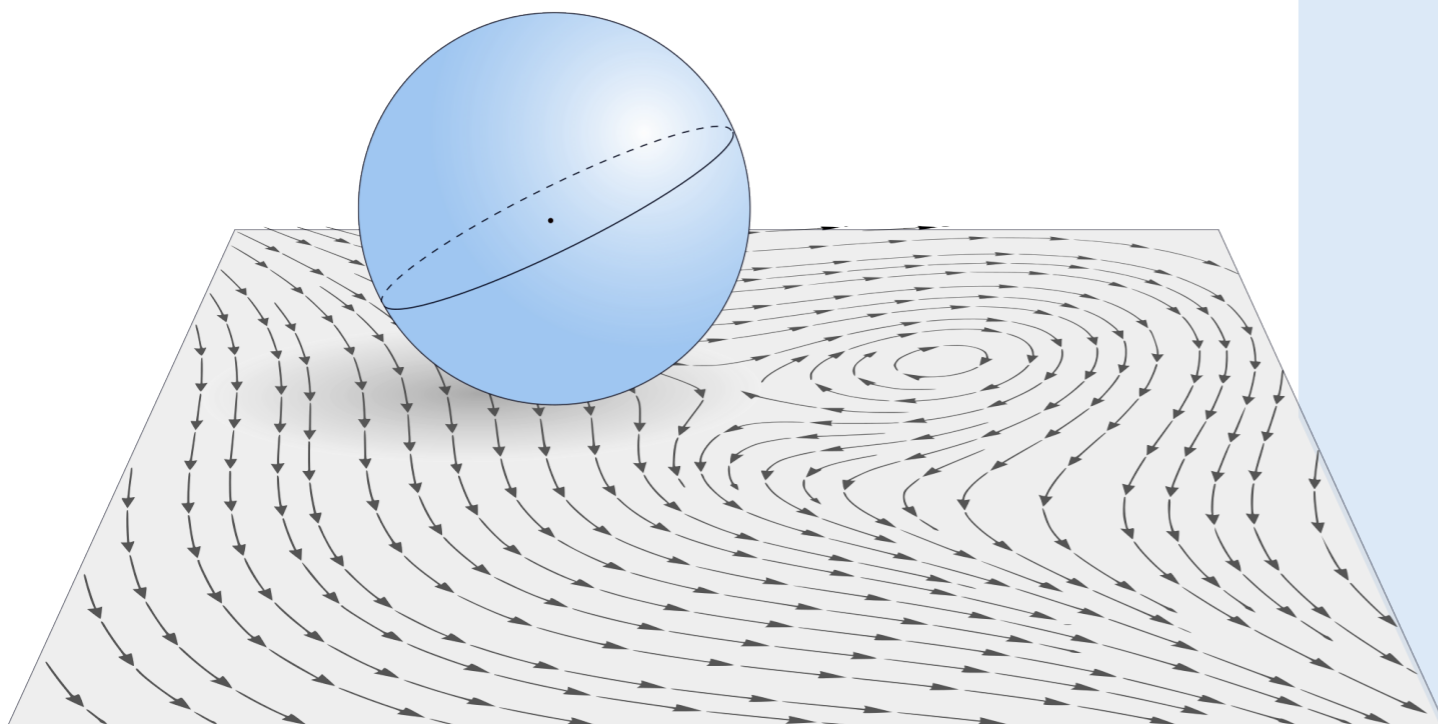
Equations of motion:

$$\dot{M} = M \times \Omega$$

$$\dot{u} = -rB(\gamma \times \Omega) + W$$

$$\dot{B} = B\hat{\Omega}$$

$$M = \mathbb{I}\Omega + mr^2\gamma \times (\Omega \times \gamma) + mr\gamma \times (B^{-1}W)$$



Chaplygin sphere

$$V = 0$$

$$W \in \mathfrak{X}(\mathbb{R}^2)$$

Equations of motion:

$$\dot{M} = M \times \Omega$$

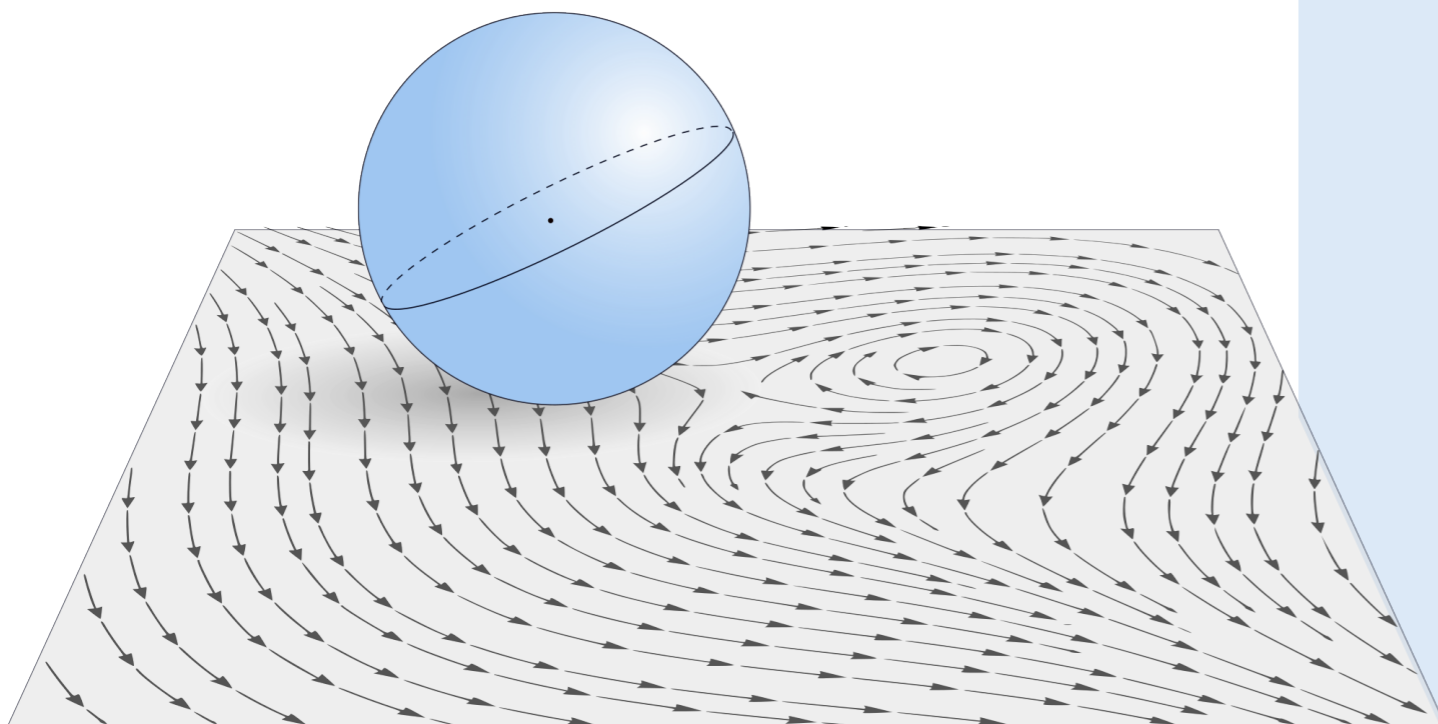
$$\dot{u} = -rB(\gamma \times \Omega) + W$$

$$\dot{B} = B\hat{\Omega}$$

$$M = \mathbb{I}\Omega + mr^2\gamma \times (\Omega \times \gamma) + mr\gamma \times (B^{-1}W)$$

First integrals

$$\langle M, \alpha \rangle, \langle M, \beta \rangle, \langle M, \gamma \rangle$$



Chaplygin sphere

$$V = 0$$

$$W \in \mathfrak{X}(\mathbb{R}^2)$$

Equations of motion:

$$\begin{aligned}\dot{M} &= M \times \Omega \\ \dot{u} &= -rB(\gamma \times \Omega) + W \\ \dot{B} &= B\hat{\Omega}\end{aligned}$$

$$M = \mathbb{I}\Omega + mr^2\gamma \times (\Omega \times \gamma) + mr\gamma \times (B^{-1}W)$$

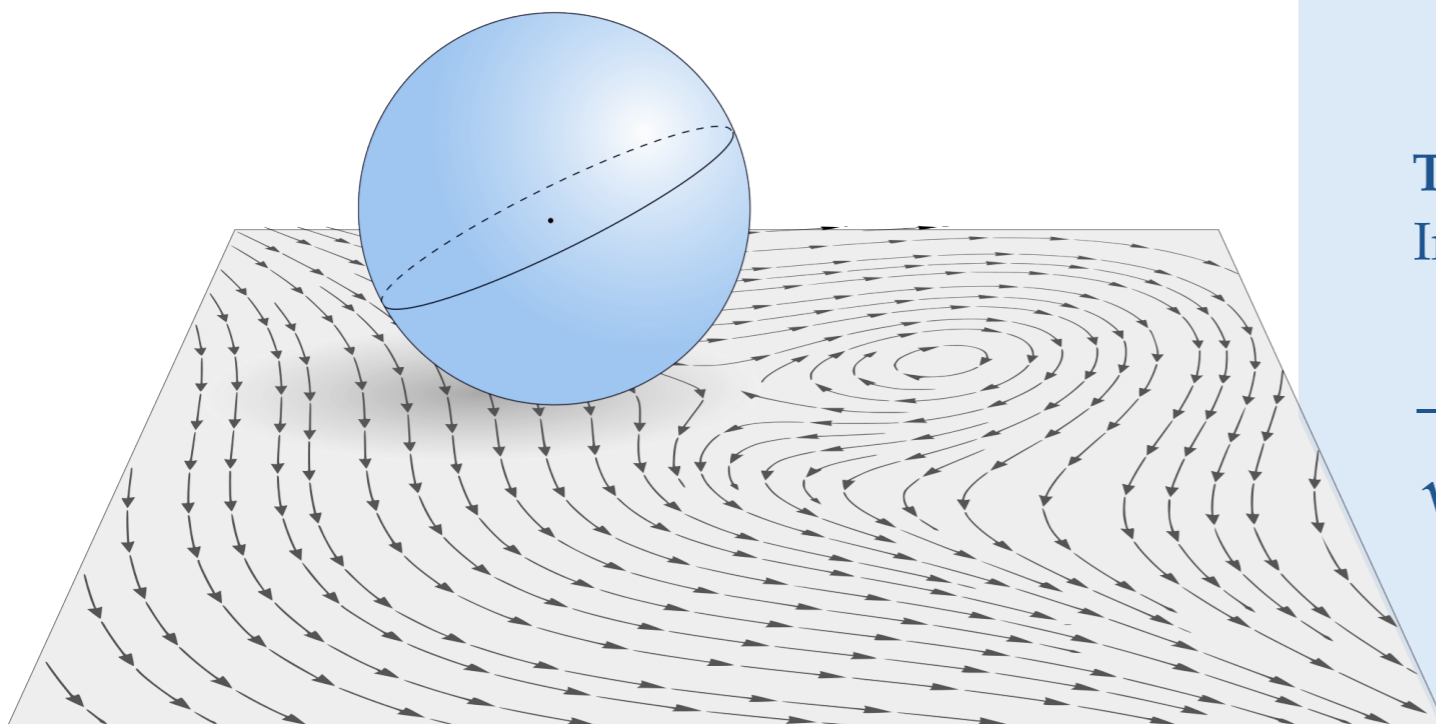
First integrals

$$\langle M, \alpha \rangle, \langle M, \beta \rangle, \langle M, \gamma \rangle$$

Theorem (C. García Naranjo, 2023):

Invariant measure if $\operatorname{div}_{\mathbb{R}^2} W = 0$

$$\frac{1}{\sqrt{1 - mr^2\langle \gamma, (\mathbb{I} + mr^2)^{-1}\gamma \rangle}} dM dx dy d\alpha d\beta dy$$



Chaplygin sphere

$$W = 0$$

Equations of motion:

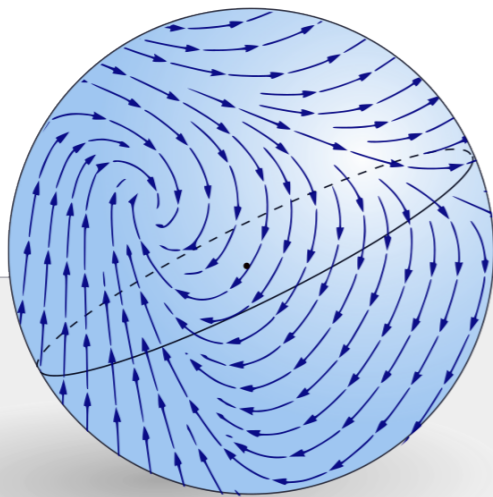
$$\dot{M} = M \times \Omega$$

$$\dot{\gamma} = \gamma \times \Omega$$

$$M = \mathbb{I}\Omega + mr^2\gamma \times (\Omega \times \gamma) + mr\gamma \times V$$

First integrals

$$\|\gamma\|^2, \quad \|M\|^2, \quad \langle M, \gamma \rangle$$



Chaplygin sphere

$$W = 0$$

No invariant measure
No moving energy

Equations of motion:

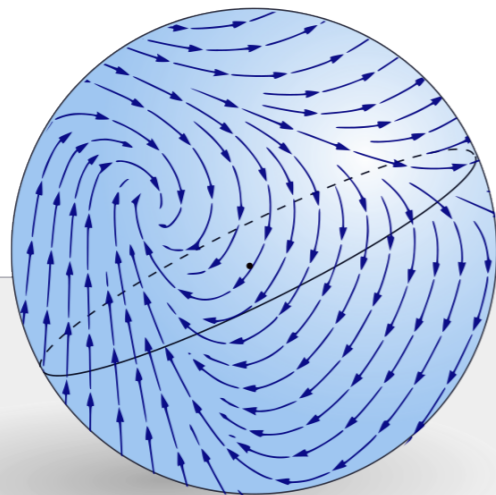
$$\dot{M} = M \times \Omega$$

$$\dot{\gamma} = \gamma \times \Omega$$

$$M = \mathbb{I}\Omega + mr^2\gamma \times (\Omega \times \gamma) + mr\gamma \times V$$

First integrals

$$\|\gamma\|^2, \quad \|M\|^2, \quad \langle M, \gamma \rangle$$



Chaplygin sphere with rotating shell

$$V = -r\eta\gamma \times E_3$$

Equations of motion:

$$\dot{M} = M \times \Omega$$

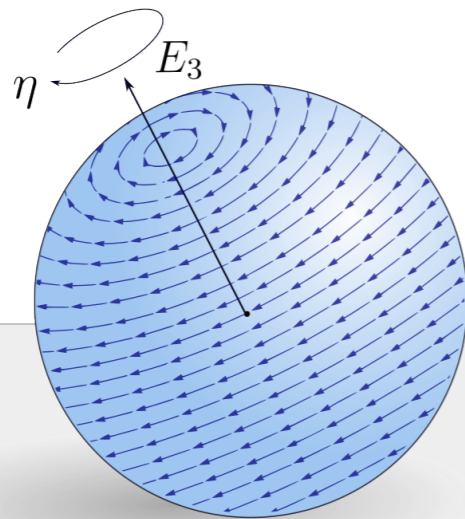
$$\dot{\gamma} = \gamma \times \Omega$$

$$M = \mathbb{I}\Omega + mr^2\gamma \times (\Omega \times \gamma) - m^2\eta\gamma \times (\gamma \times E_3)$$

First integrals:

$$\|\gamma\|^2, \|M\|^2, \langle M, \gamma \rangle$$

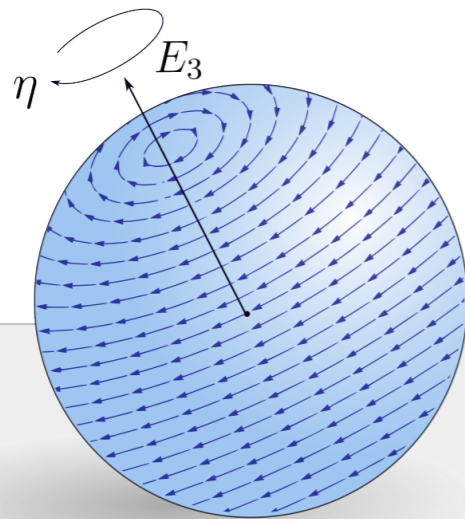
3 dimensional level sets



Chaplygin sphere with rotating shell

$$V = -r\eta\gamma \times E_3$$

Generally chaotic



Equations of motion:

$$\dot{M} = M \times \Omega$$

$$\dot{\gamma} = \gamma \times \Omega$$

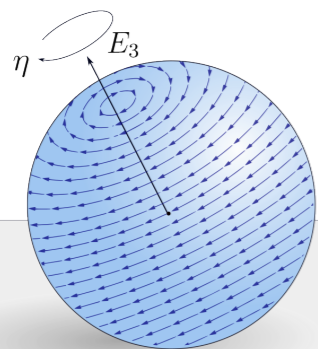
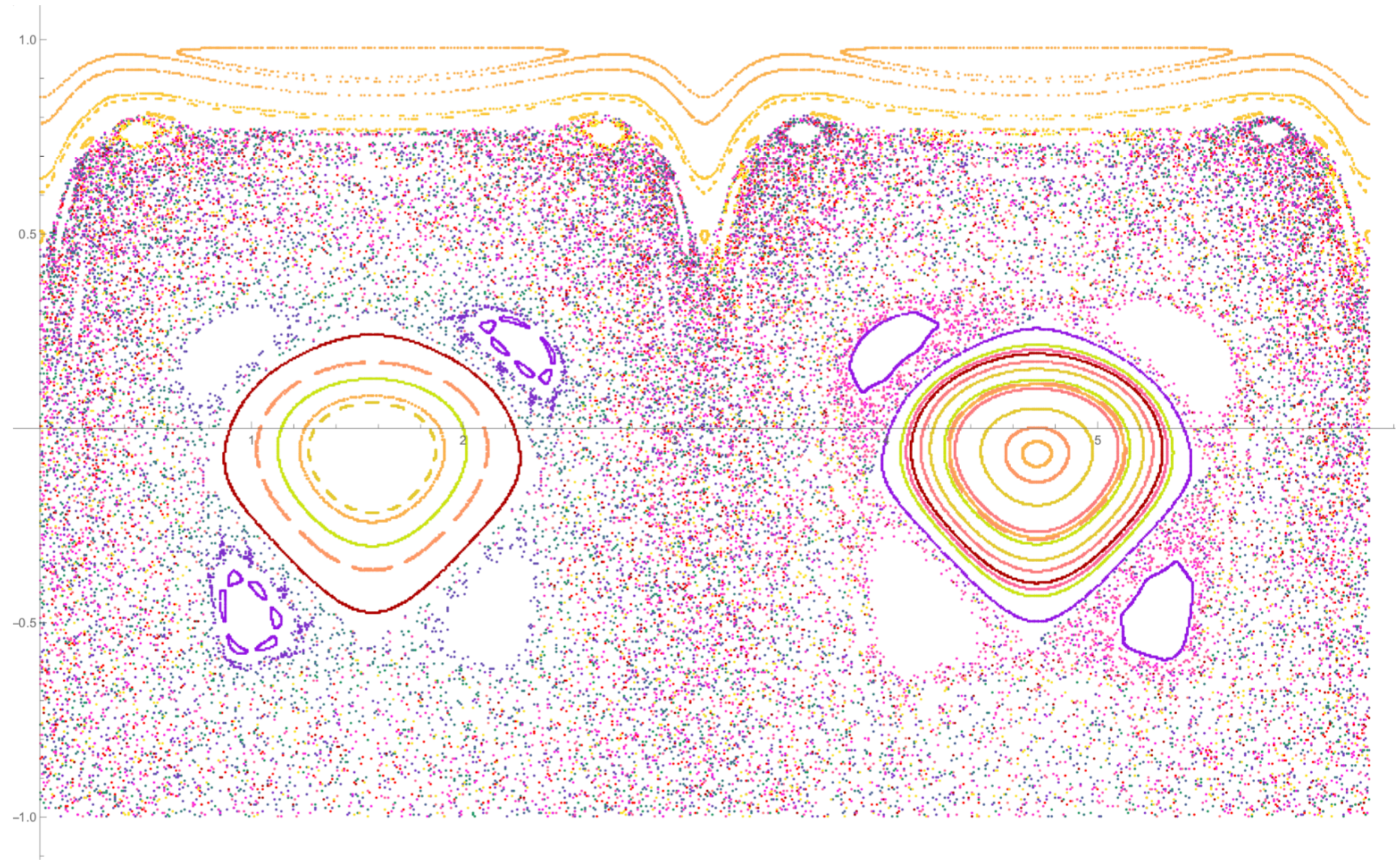
$$M = \mathbb{I}\Omega + mr^2\gamma \times (\Omega \times \gamma) - m^2\eta\gamma \times (\gamma \times E_3)$$

First integrals:

$$\|\gamma\|^2, \|M\|^2, \langle M, \gamma \rangle$$

3 dimensional level sets

Poincaré map



Poincaré map

Equations of motion:

$$\dot{M} = \varepsilon (kM \times \Omega_c) + M \times \Omega_a$$

$$\dot{\gamma} = \varepsilon(k\gamma \times \Omega_c) + \gamma \times \Omega_a$$

Poincaré map

$$\varepsilon = \frac{\|M\|}{mr^2\eta}$$

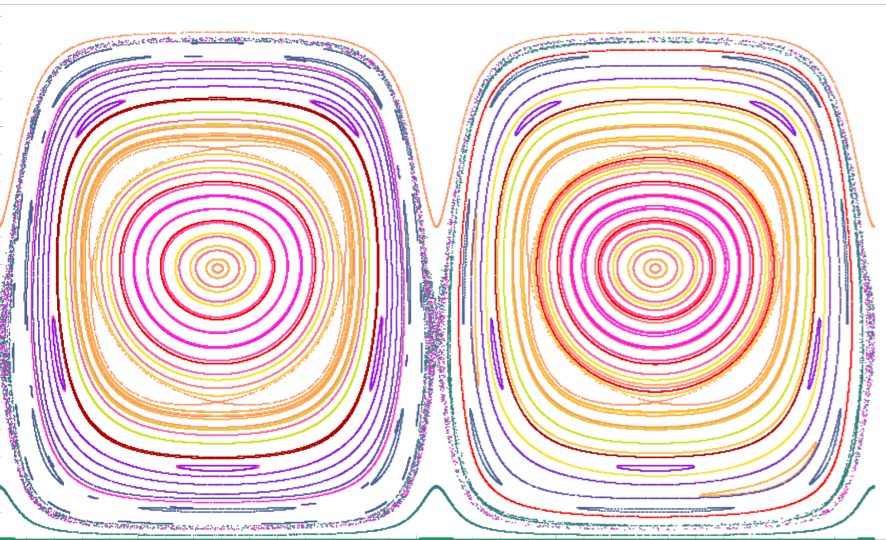
Equations of motion:

$$\dot{M} = \varepsilon (kM \times \Omega_c) + M \times \Omega_a$$

$$\dot{\gamma} = \varepsilon(k\gamma \times \Omega_c) + \gamma \times \Omega_a$$

Poincaré map

$$\varepsilon \gg 1$$



$$\varepsilon = \frac{\|M\|}{mr^2\eta}$$

Equations of motion:

$$\dot{M} = \varepsilon (kM \times \Omega_c) + M \times \Omega_a$$

$$\dot{\gamma} = \varepsilon (k\gamma \times \Omega_c) + \gamma \times \Omega_a$$

Poincaré map

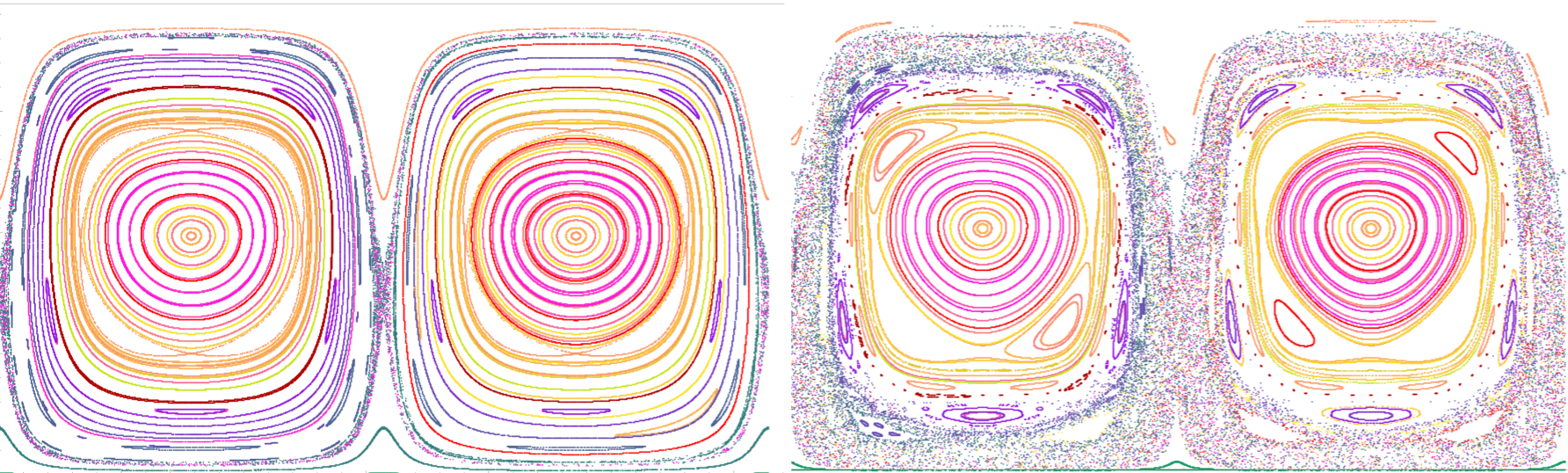
Equations of motion:

$$\dot{M} = \varepsilon (kM \times \Omega_c) + M \times \Omega_a$$

$$\dot{\gamma} = \varepsilon(k\gamma \times \Omega_c) + \gamma \times \Omega_a$$

$$\varepsilon = \frac{\|M\|}{mr^2\eta}$$

$\varepsilon \gg 1$



Poincaré map

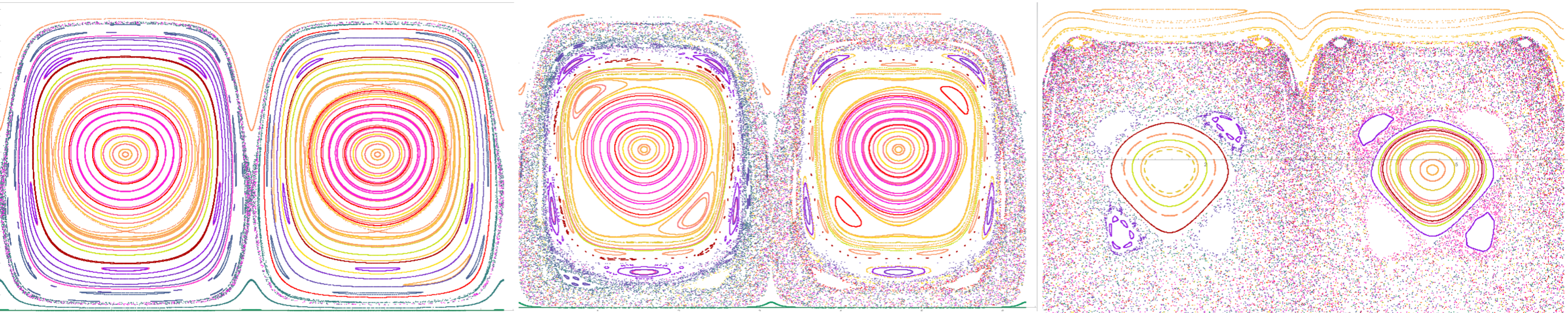
$$\varepsilon = \frac{\|M\|}{mr^2\eta}$$

$$\varepsilon \gg 1$$

Equations of motion:

$$\dot{M} = \varepsilon (kM \times \Omega_c) + M \times \Omega_a$$

$$\dot{\gamma} = \varepsilon(k\gamma \times \Omega_c) + \gamma \times \Omega_a$$



Poincaré map

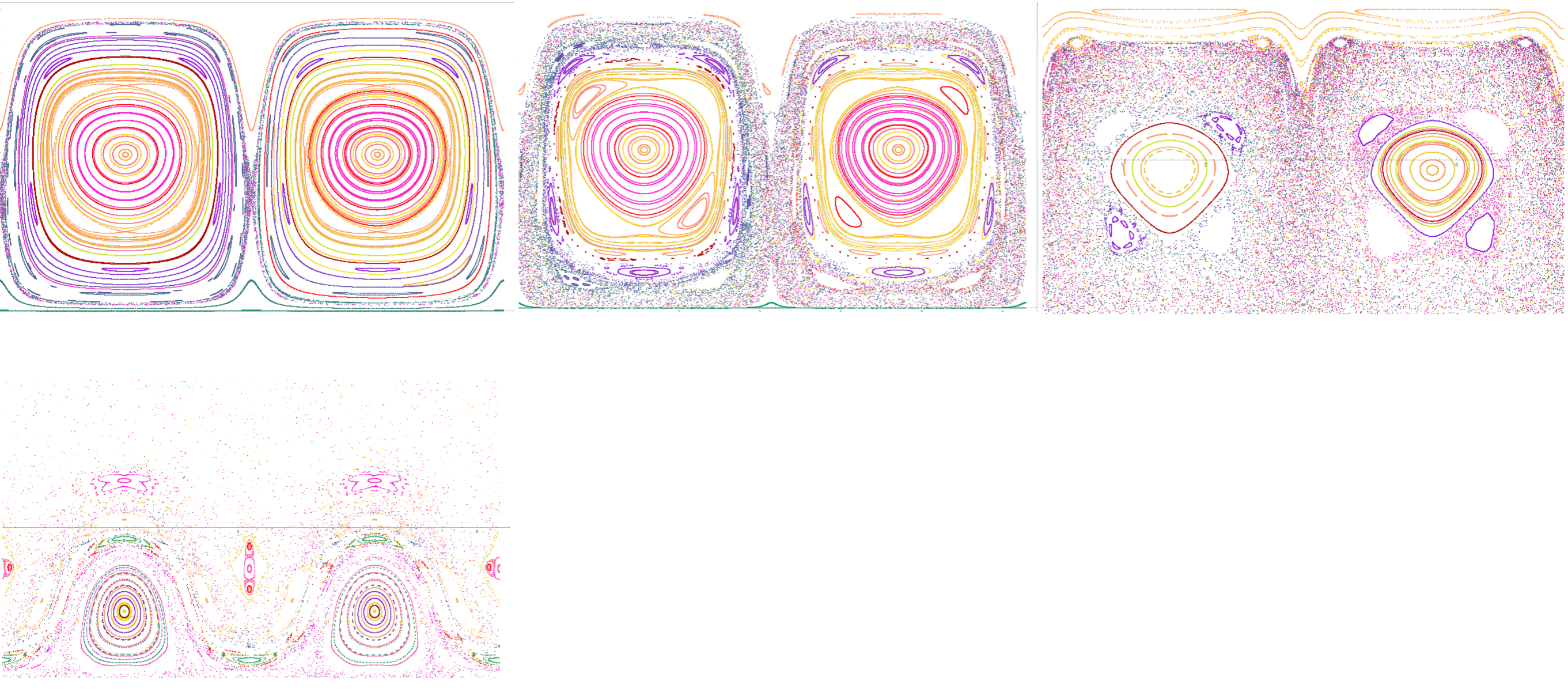
$$\varepsilon = \frac{\|M\|}{mr^2\eta}$$

Equations of motion:

$$\dot{M} = \varepsilon (kM \times \Omega_c) + M \times \Omega_a$$

$$\dot{\gamma} = \varepsilon(k\gamma \times \Omega_c) + \gamma \times \Omega_a$$

$\varepsilon \gg 1$



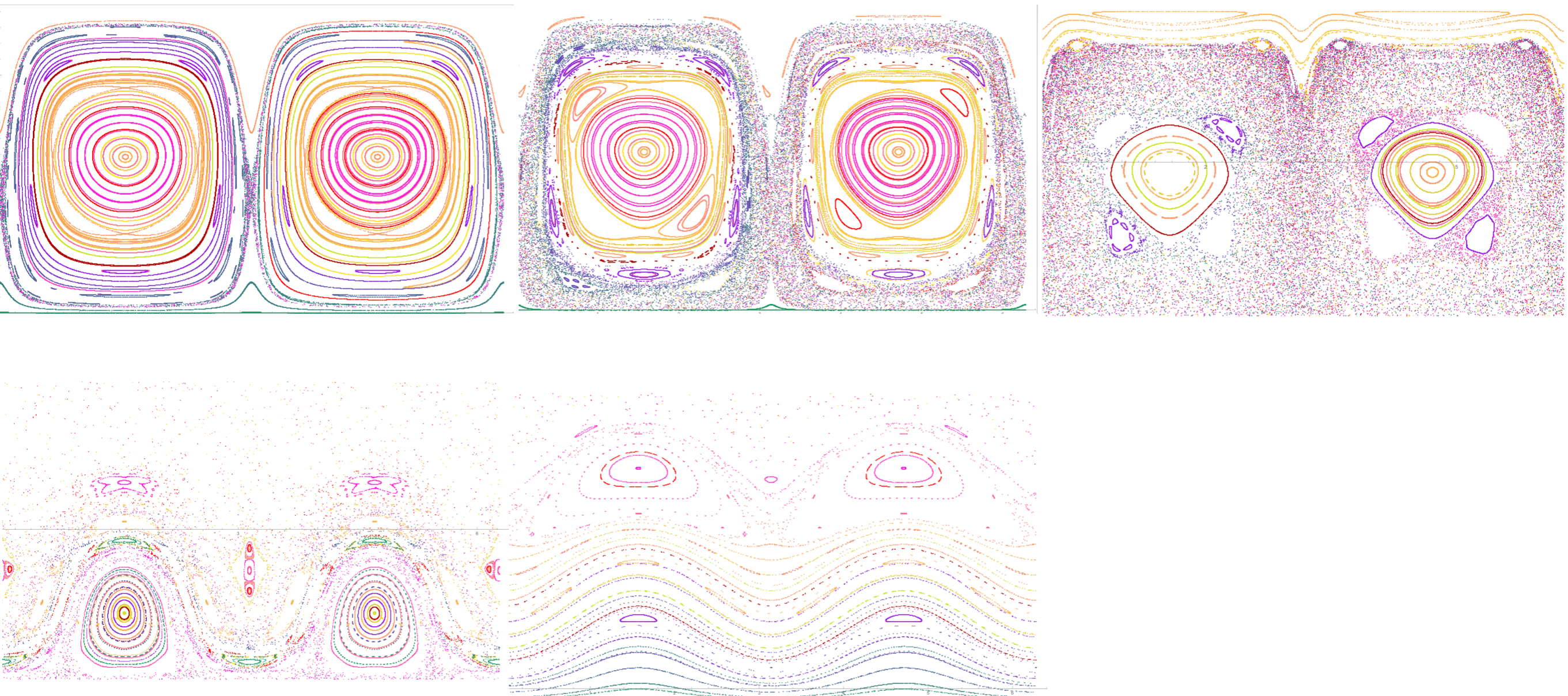
Poincaré map

$$\varepsilon = \frac{\|M\|}{mr^2\eta}$$

Equations of motion:

$$\begin{aligned}\dot{M} &= \varepsilon (kM \times \Omega_c) + M \times \Omega_a \\ \dot{\gamma} &= \varepsilon (k\gamma \times \Omega_c) + \gamma \times \Omega_a\end{aligned}$$

$\varepsilon \gg 1$



Poincaré map

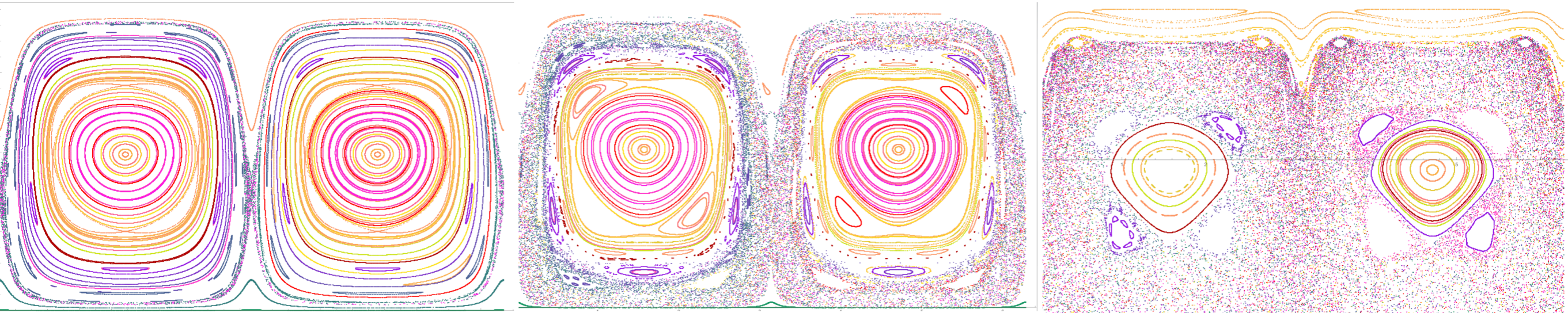
$$\varepsilon = \frac{\|M\|}{mr^2\eta}$$

Equations of motion:

$$\dot{M} = \varepsilon (kM \times \Omega_c) + M \times \Omega_a$$

$$\dot{\gamma} = \varepsilon (k\gamma \times \Omega_c) + \gamma \times \Omega_a$$

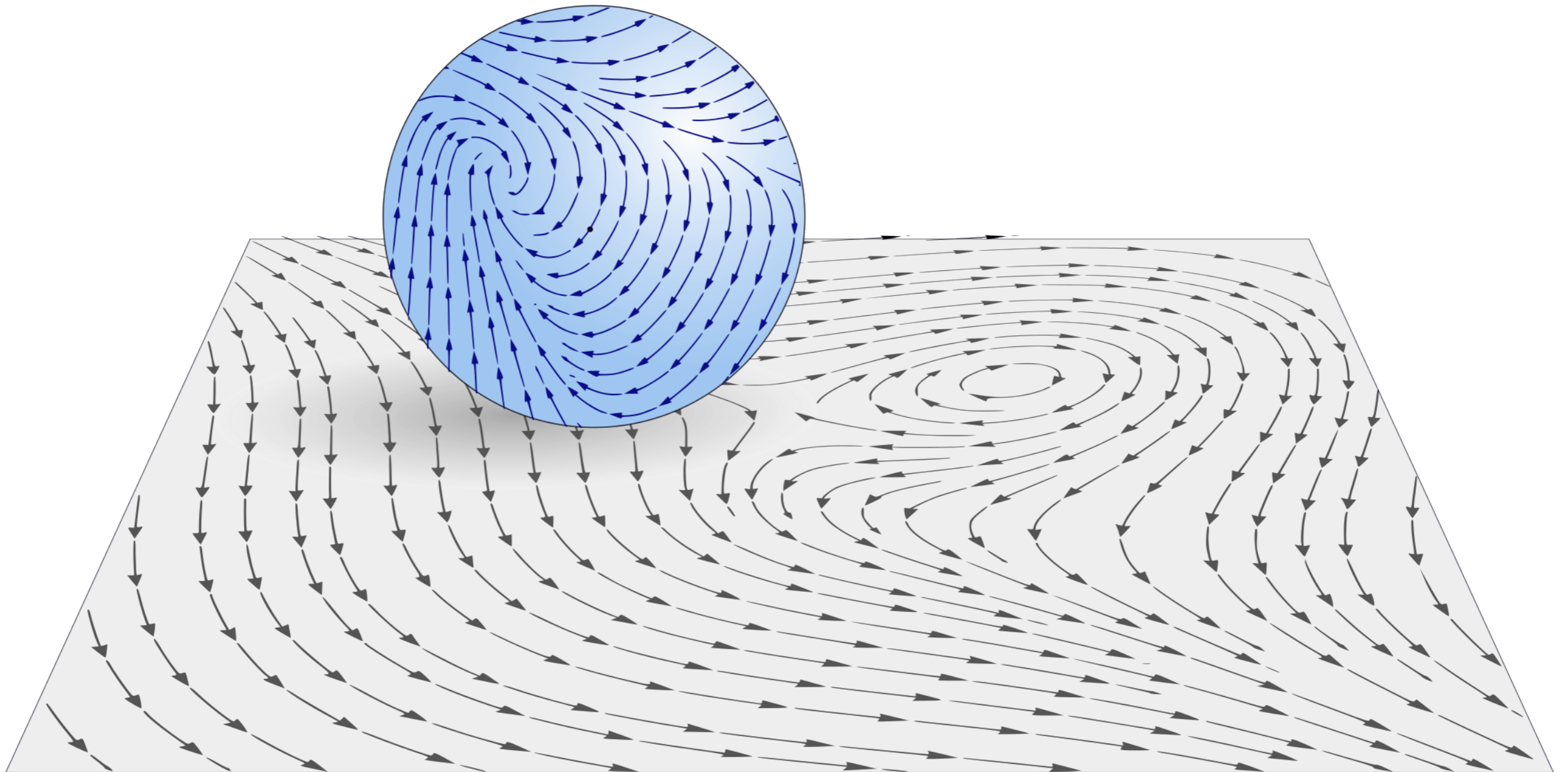
$\varepsilon \gg 1$



$\varepsilon \ll 1$

Homogeneous sphere

$$I_1 = I_2 = I_3$$



Homogeneous sphere

$$I_1 = I_2 = I_3$$

Equations of motion:

$$\dot{M} = M \times \Omega$$

$$\dot{u} = -rB(\gamma \times \Omega) + BV + W$$

$$\dot{B} = B\hat{\Omega}$$

$$M = I\Omega + mr^2\gamma \times (\Omega \times \gamma) + mr\gamma \times (V + B^{-1}W)$$

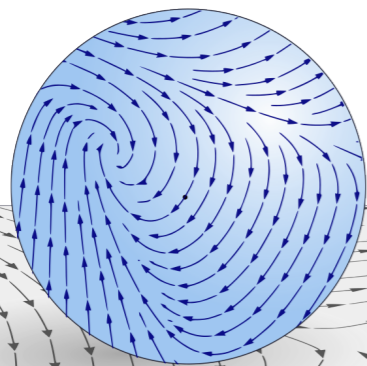
First integrals:

$$\|M\|^2, \quad \langle M, \gamma \rangle, \quad \|\gamma\|^2$$

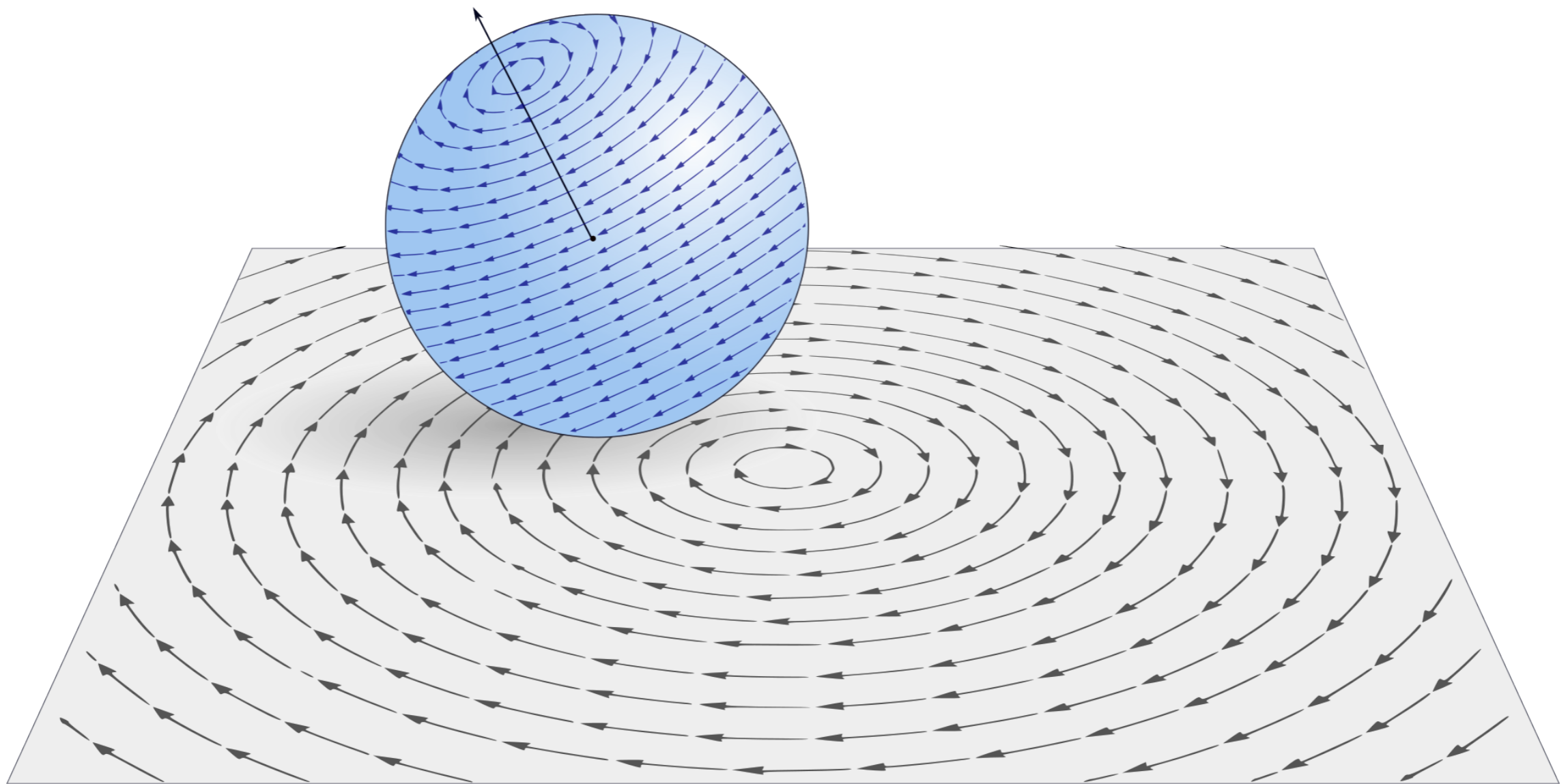
Theorem (C., García Naranjo, 2023)

Invariant measure:

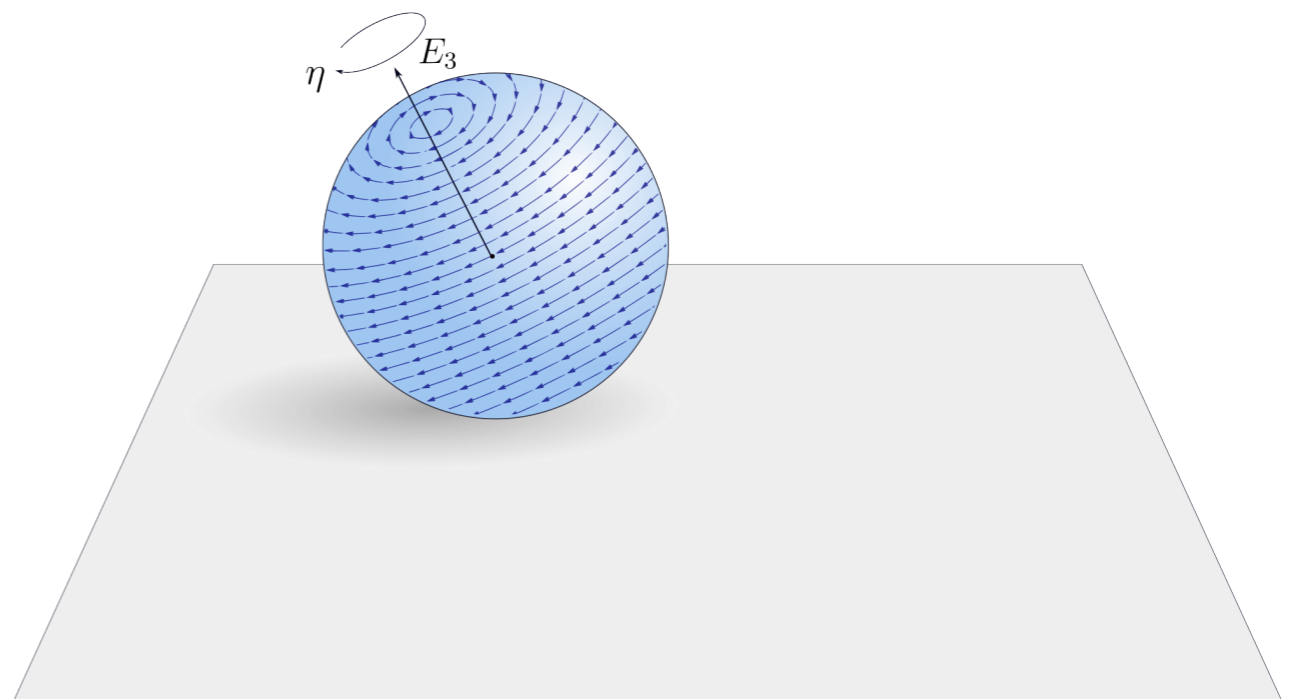
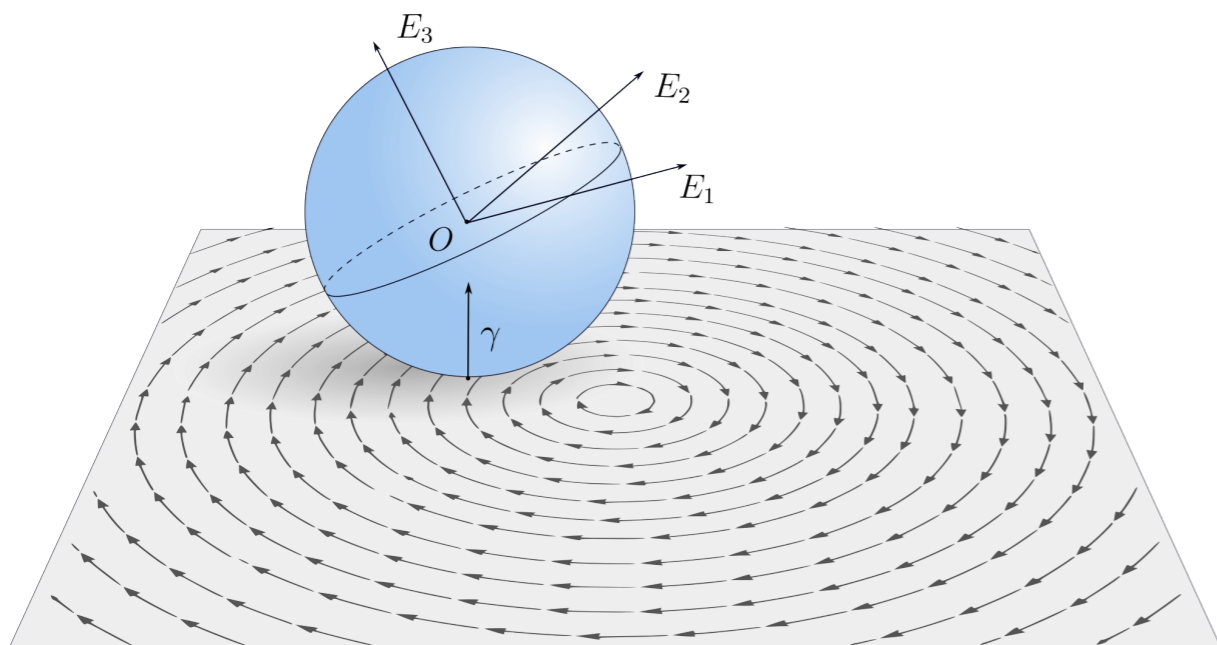
$$dMdx dy d\alpha d\beta d\gamma \quad \text{if } \operatorname{div}_{\mathbb{R}^2} W = 0 \text{ and } \operatorname{div}_{S^2} V = 0$$



Homogeneous sphere

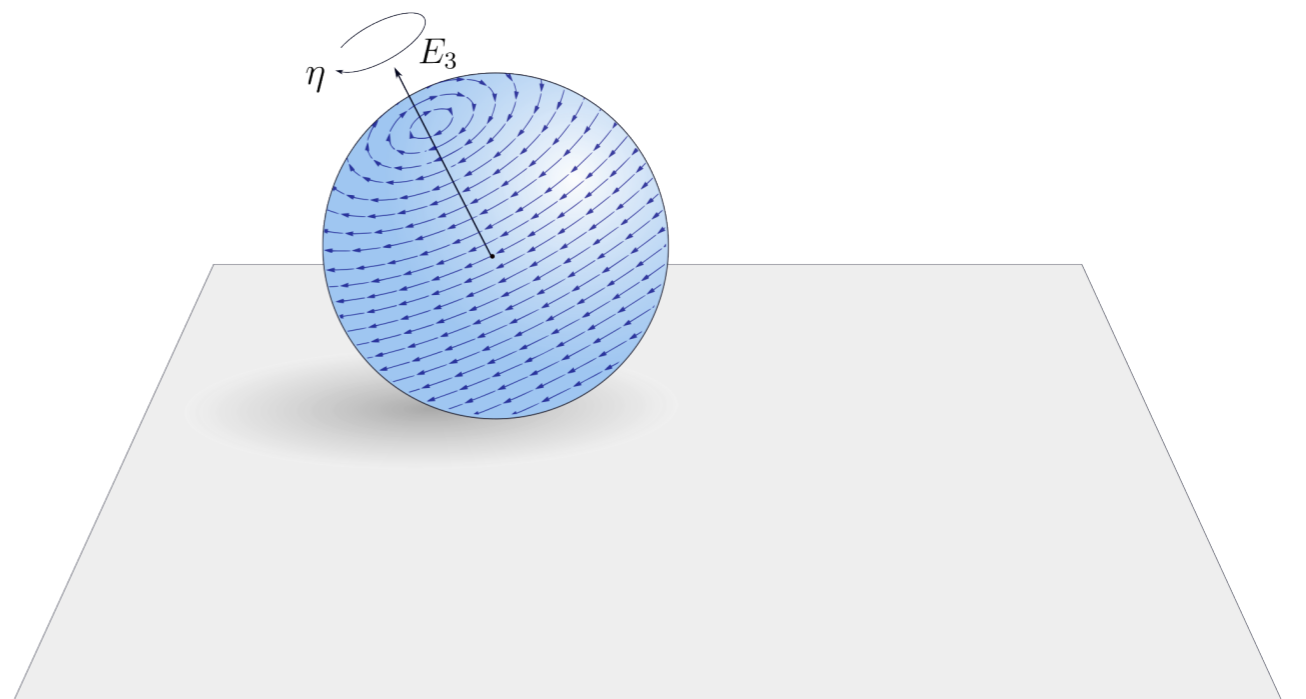
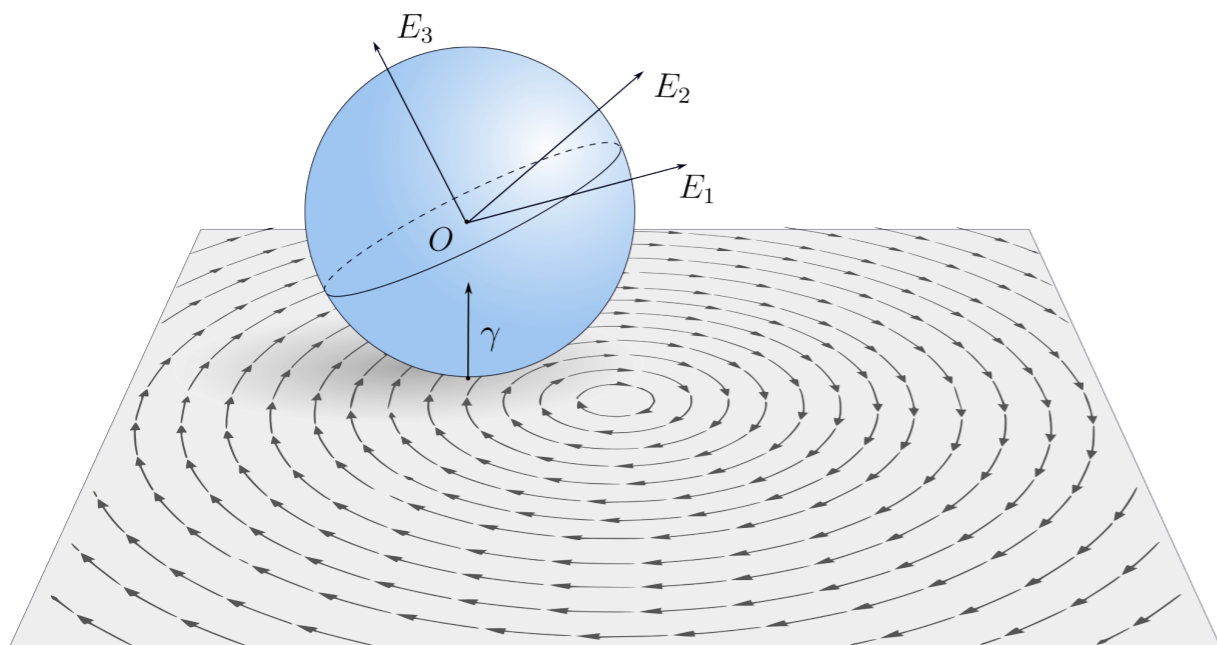


Homogeneous sphere



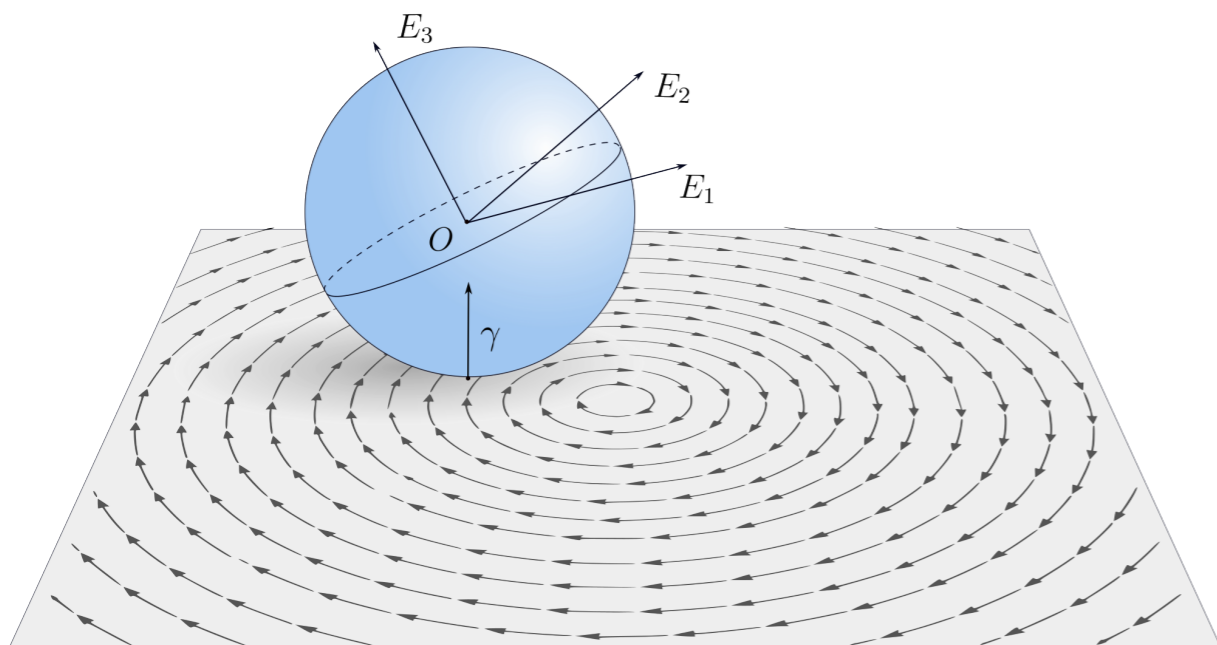
Homogeneous sphere

Integrable

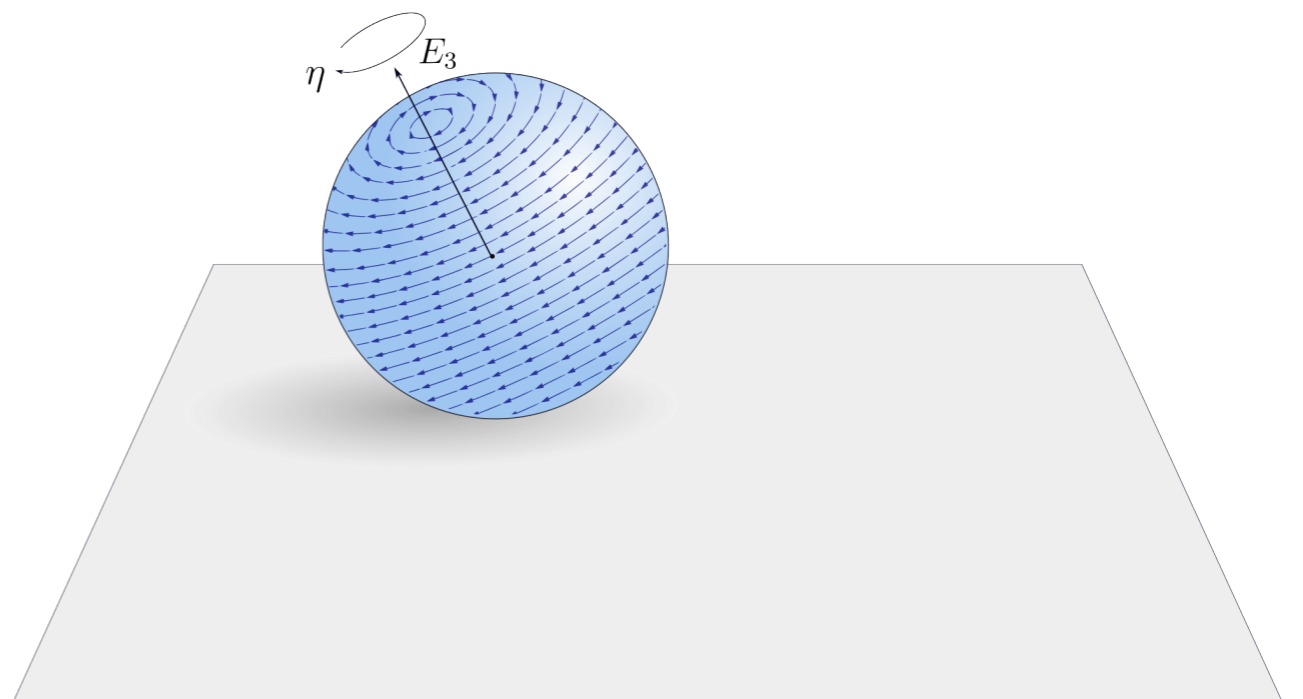


Homogeneous sphere

Integrable



Integrable



Homogeneous sphere

Equations of motion:

$$\dot{M} = M \times \Omega$$

$$\dot{U} = -r(\gamma \times \Omega) + U \times \Omega - r\eta_1(\gamma \times E_3) - \eta_2\gamma \times U$$

$$\dot{\gamma} = \gamma \times \Omega$$

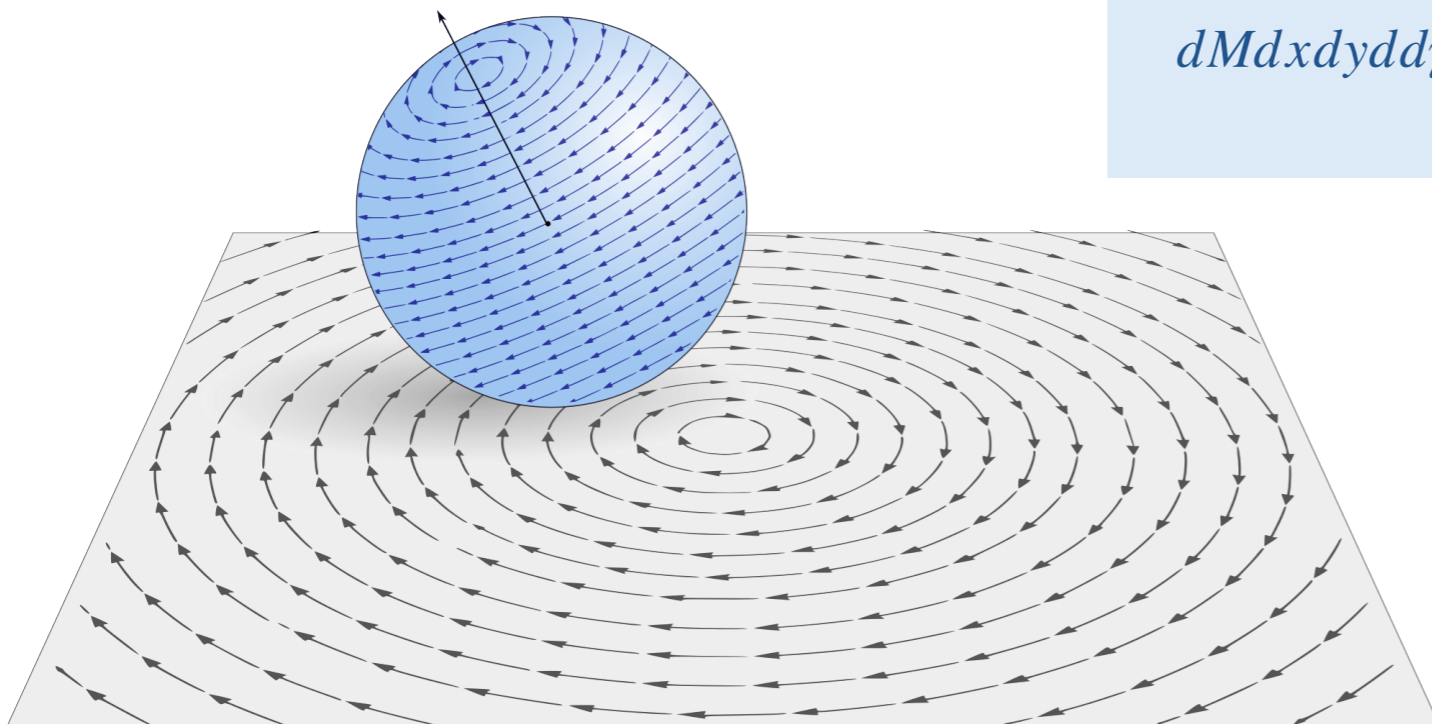
$$M = I\Omega + mr^2\gamma \times (\Omega \times \gamma) - mr^2\eta_1\gamma \times (\gamma \times E_3) + \eta_2\gamma \times (\gamma \times U)$$

First integrals

$$\|M\|^2, \quad \langle M, \gamma \rangle, \quad \|\gamma\|^2, \quad \langle U, \gamma \rangle = 0, \quad E_{mov}$$

Invariant measure

$$dMdx dy d\gamma$$



Homogeneous sphere

Current work

Equations of motion:

$$\dot{M} = M \times \Omega$$

$$\dot{U} = -r(\gamma \times \Omega) + U \times \Omega - r\eta_1(\gamma \times E_3) - \eta_2\gamma \times U$$

$$\dot{\gamma} = \gamma \times \Omega$$

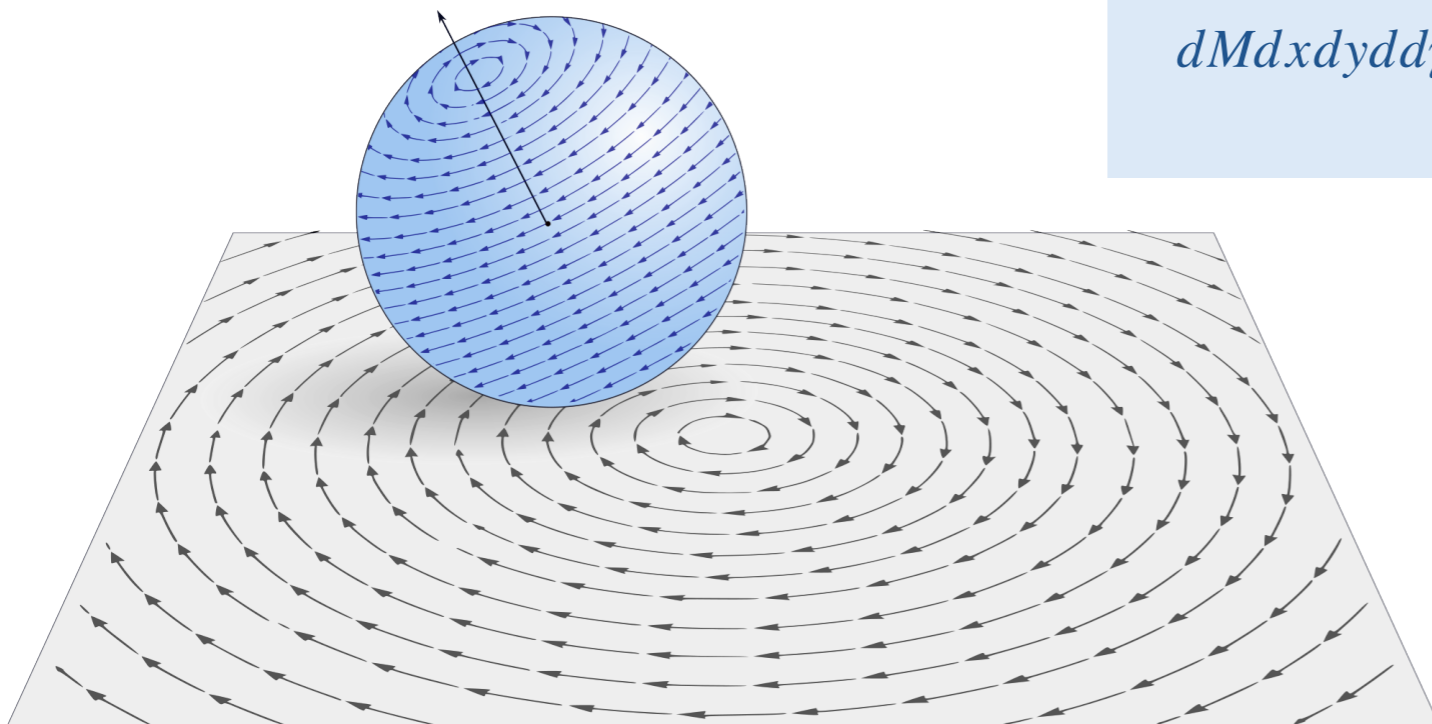
$$M = I\Omega + mr^2\gamma \times (\Omega \times \gamma) - mr^2\eta_1\gamma \times (\gamma \times E_3) + \eta_2\gamma \times (\gamma \times U)$$

First integrals

$$\|M\|^2, \quad \langle M, \gamma \rangle, \quad \|\gamma\|^2, \quad \langle U, \gamma \rangle = 0, \quad E_{mov}$$

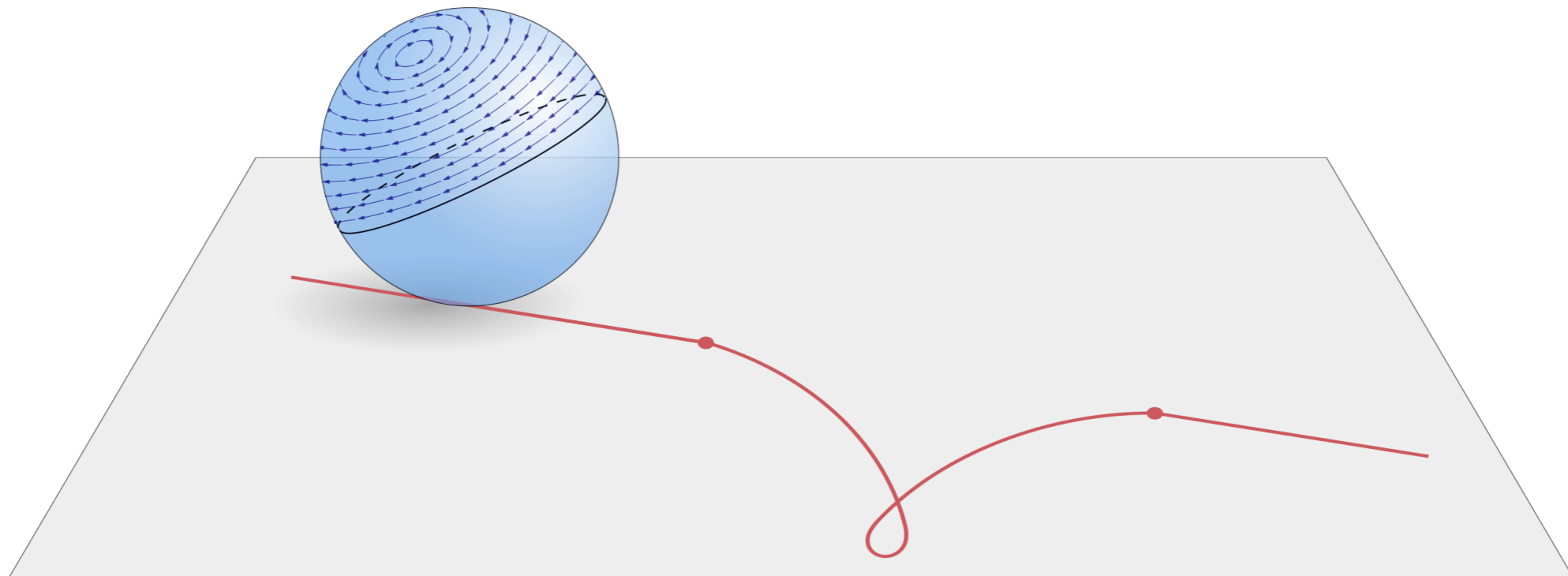
Invariant measure

$$dMdx dy d\gamma$$



Homogeneous sphere with rotating shell

We proved an analogous to the
Anais-billiard phenomenon



Conclusions

Conclusions

- Identified a rich class of examples to analyze dynamical properties of affine nonholonomic systems
- Recognized a robust mechanism leading to existence of first integrals (momenta)
- Found new instances of moving energy
- Observed integrable and chaotic behavior and transition as function of problem parameters

Future work

Future work

- Role of reversibility in affine nonholonomic systems
- Momentum first integrals for affine nonholonomic systems
- Develop a theory to frame these examples:
Integrability, perturbation theory and chaos of nonholonomic systems

Thank you!