

Muon $g - 2$ theory: SM and beyond the SM

Dominik Stöckinger, TU Dresden

Warsaw Physics Colloquium, 26th April 2021

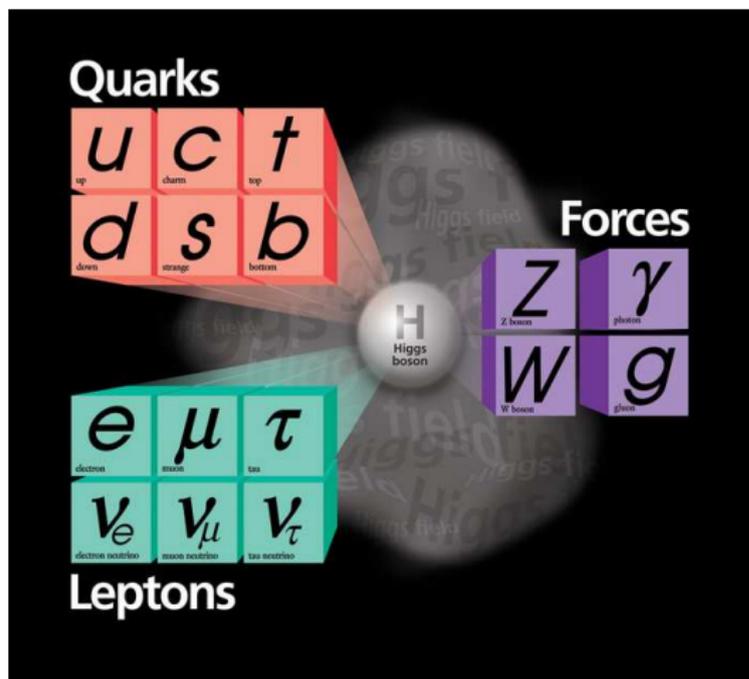
Collaborations:

Muon $g - 2$ collaboration

SM Theory Initiative

BSM collaborators: Peter Athron, Csaba Balasz, Douglas Jacob,
Wojciech Kotlarski, Hyejung Stöckinger-Kim

Standard Model of particle physics (est. 1967...1973))

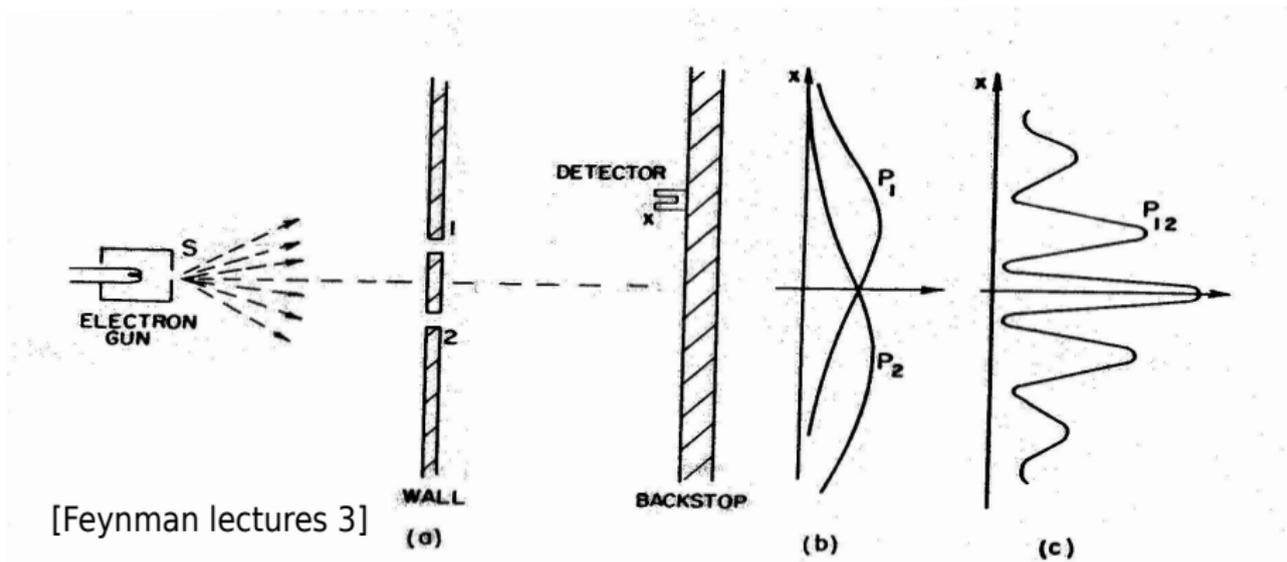


SM very well confirmed!

- All known interactions (\neq gravity)
- relativistic QFT
 \rightsquigarrow renormalizable
- gauge invariance
 \rightsquigarrow specific interactions
- spontaneous EWSB
 \rightsquigarrow Higgs

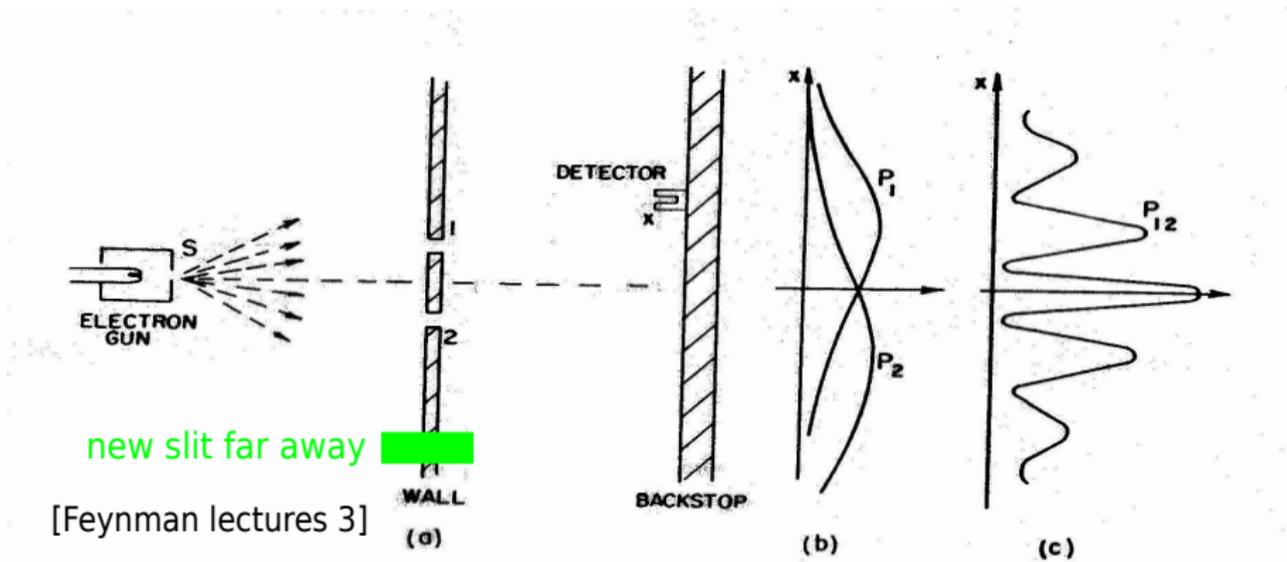
a_μ sensitive to all particles and forces via quantum fluctuations!

Quantum fluctuations: double slit experiment



- essence of quantum mechanics
- All possible paths contribute, probability amplitudes interfere

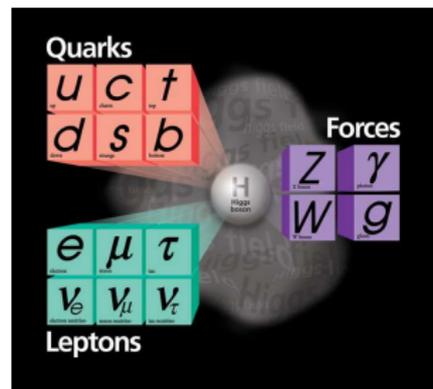
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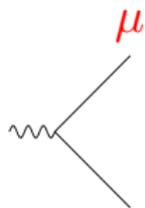
Our “paths” for $g - 2$: Feynman diagrams

μ couples to B -field directly or via virtual particles



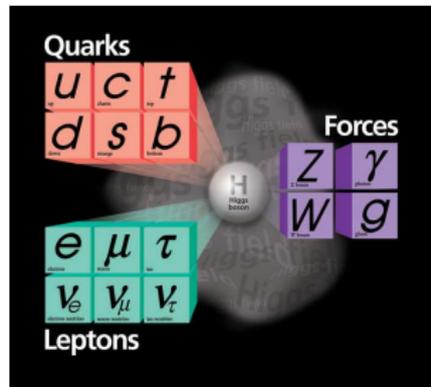
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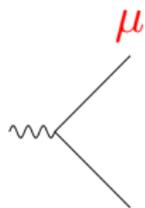
Dirac equation/direct \rightsquigarrow “pointlike”

$$g = 2$$



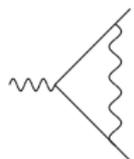
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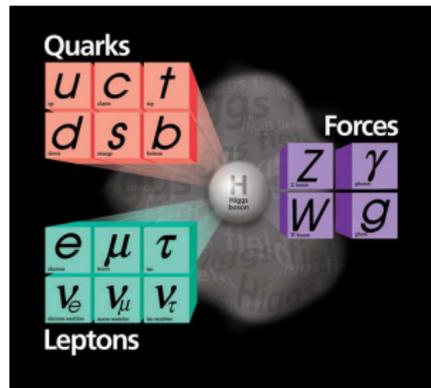
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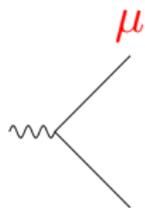
Schwinger (1948): quantum fluctuations \rightsquigarrow “non-pointlike”

$$g = 2 \left(1 + \frac{\alpha}{2\pi} \right)$$



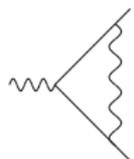
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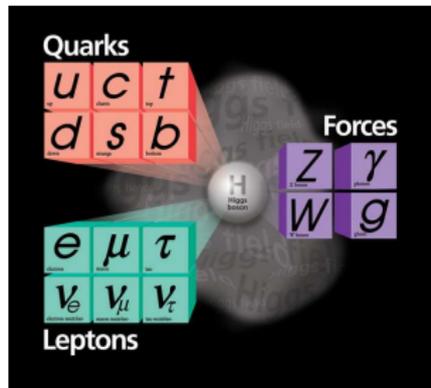
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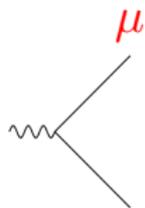
Technically: effective Lagrangian/Hamiltonian

$$\mathcal{L}_{\text{eff}} = -\frac{Q_e}{4m_\mu} a_\mu \times \bar{\psi}_L \sigma_{\mu\nu} \psi_R F^{\mu\nu}$$

$$H_{\text{eff}} = -2(1 + a_\mu) \frac{e}{2m_\mu} \vec{B} \cdot \vec{S}$$

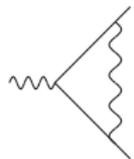
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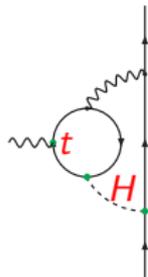
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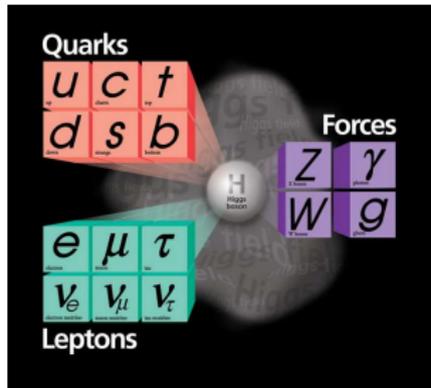
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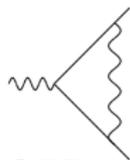
All SM particles contribute, even Higgs and top!

$$g = 2 \left(1 + \dots - 1.5 \times 10^{-11} \right)$$

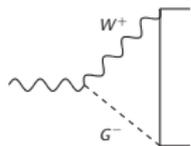


Theory Initiative prediction $a_\mu^{\text{SM}} = (11\,659\,181.0 \ (4.3)) [10^{-10}]$

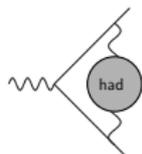
since 2017, 6 workshops, White Paper (2020), 132 authors, ongoing effort



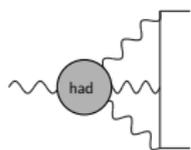
QED: 11 658 471.9 (0.0)



Weak: 15.36 (0.1)



Hadronic vac.pol.: 684.5 (4.0)



Hadr. light-by-light: 9.0 (1.7)

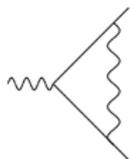
number = conservative
combination of approaches

very precise particle theory
prediction

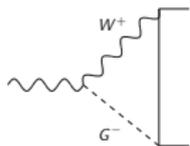
evaluation uses all conceivable
QFT methods & tricks

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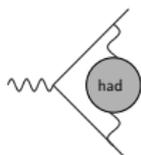
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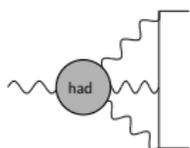
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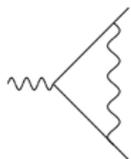
- 4-/5-loop Feynman integrals analytical & numerical
- most precise QED test: $(g - 2)_e$!

Weak:

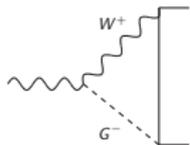
- first EW 2-loop calculation, also use renormalization group methods

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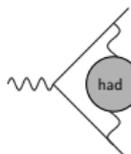
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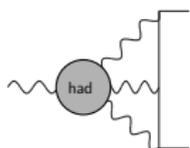
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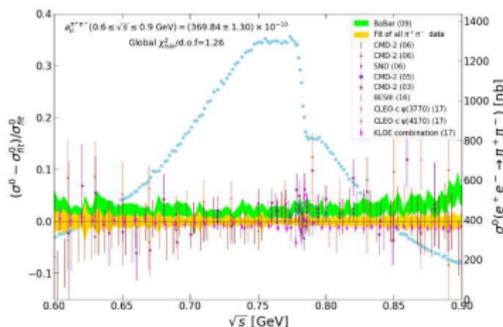


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Hadronic vacuum polarization:

- unitarity+causality \rightsquigarrow
exact dispersion relation

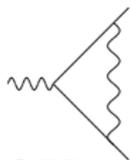
$$2 \text{Im} \left[\text{had.} \right] = \sum_{\text{had.}} \int d\Phi \left| \text{had.} \right|^2$$



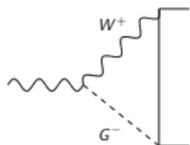
- lattice QCD impressive progress
(not yet used in TI value)

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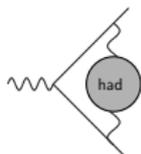
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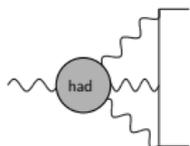
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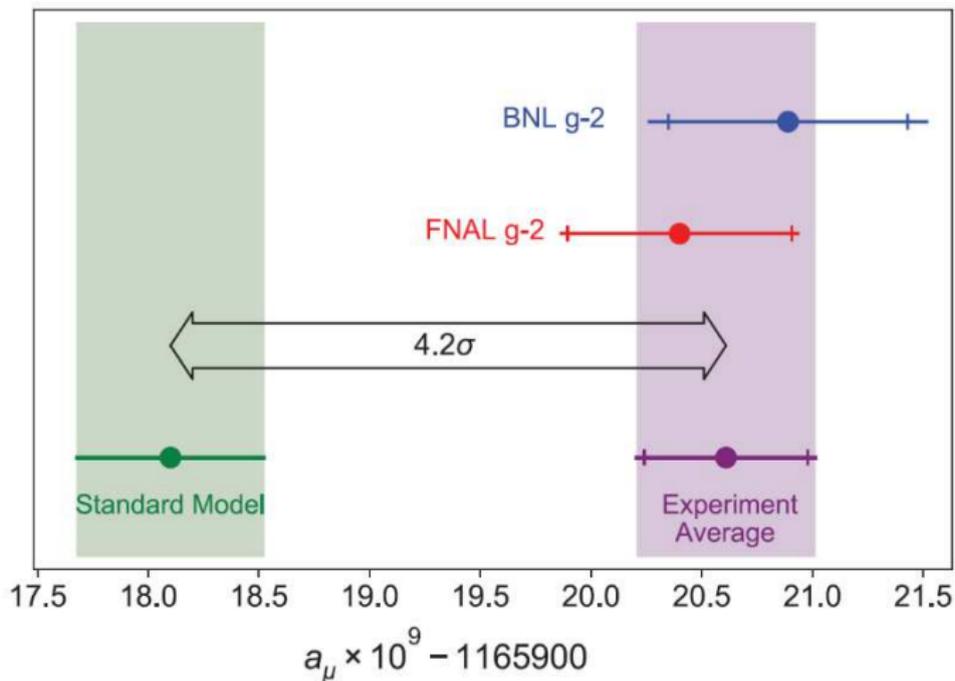


Hadr. light-by-light: 9.0 (1.7)

Hadronic light-by-light:

- difficult QFT problem
- Traditionally: low-energy models
- Recently: data-driven (dispersion relations) & lattice QCD results
- consistent results
- uncertainty better under control

Finally: Fermilab Run 1 versus Theory Initiative SM value



Discrepancy

SM prediction too low by $\approx (25 \pm 6) \times 10^{-10}$

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Large!

discrepancy $\approx 2 \times a_{\mu}^{\text{SM,weak}}$

how to explain without conflict to LHC etc?

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Large!

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how to explain without conflict to LHC etc?

Questions: Which models can(not) explain it?

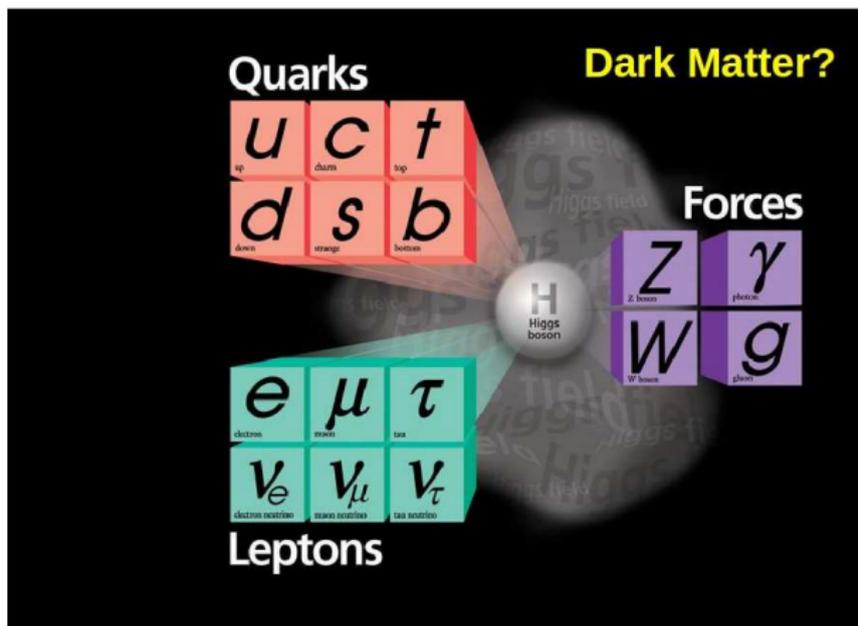
Why is a single number so interesting?

“Why are you happy about a discrepancy?”

- Very active area (> 70 papers)
- Here: general remarks and examples from survey [2104.03691](#)

[Peter Athron, Csaba Balasz, Douglas Jacob, Wojciech Kotlarski, DS, Hyejung Stöckinger-Kim]

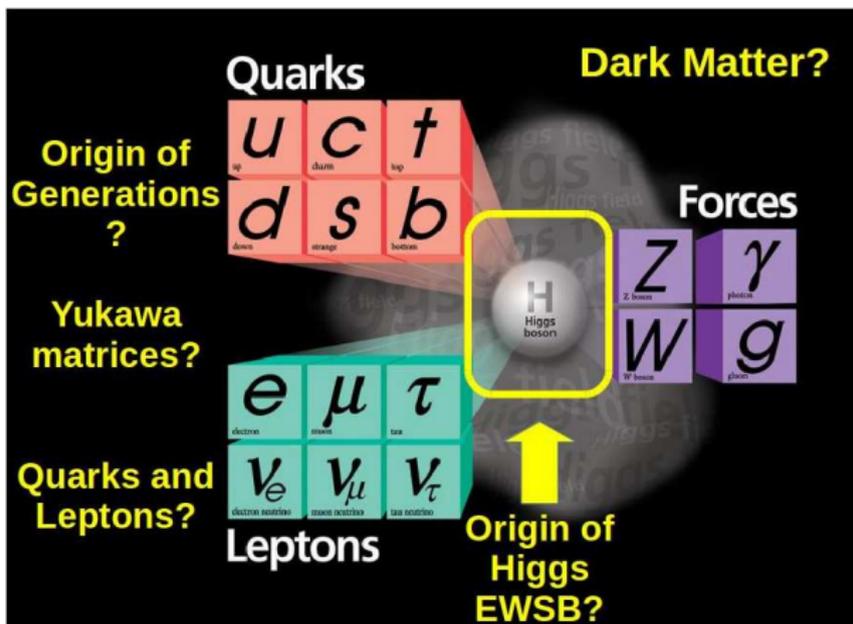
Open questions require Beyond the Standard Model (BSM) physics



Open questions!

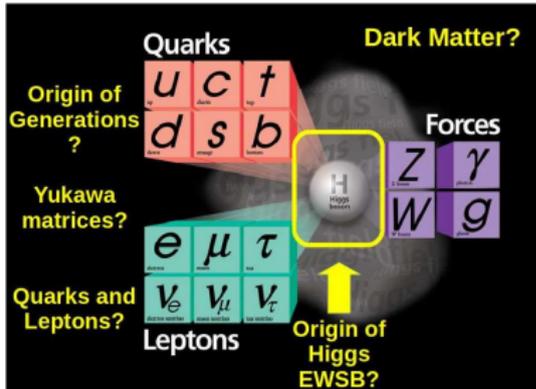
- experimental clues needed! $\rightsquigarrow g - 2!$
not easy to explain!
- relevant and deep questions may be related to $g - 2$

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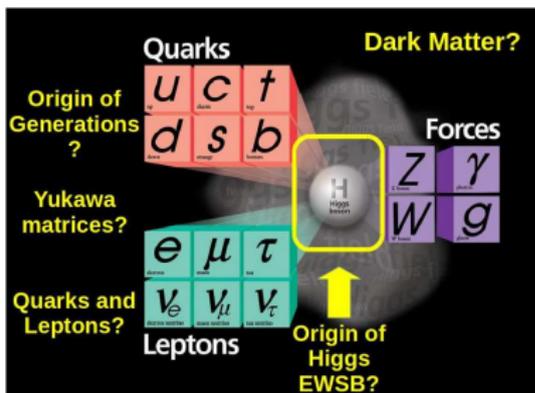
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Window to muon mass generation mechanism?

Dark Matter? Hard to see in detectors
but could couple to muon \rightsquigarrow large effects possible!

≥ 15 explanations of $g - 2$ via dark matter



Window to muon mass generation mechanism?

papers: 16 SUSY, 7 2HDM, 3 mass-generation, 4 leptoquark

(+B-anomalies), 4-field model (+B-anomalies and DM), 1 vector-like lepton, 1 GUT(331)

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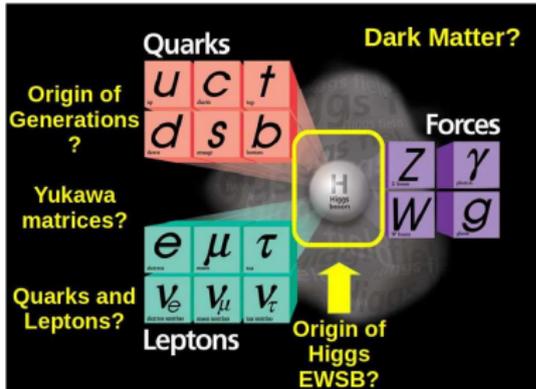
≥ 15 explanations of $g - 2$ via dark matter

Window to the muon mass generation mechanism (Higgs/Yukawa sectors)

Technically: QFT operators for m_μ and a_μ are **chirality flipping** and break gauge invariance:

$$m_\mu \bar{\psi}_L \psi_R$$

$$\frac{a_\mu}{m_\mu} \bar{\psi}_L \sigma_{\mu\nu} \psi_R F^{\mu\nu}$$



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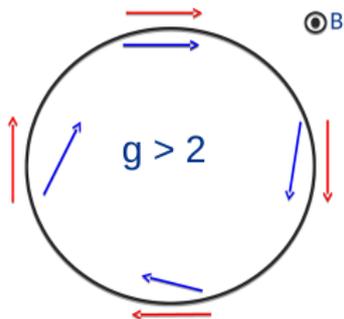
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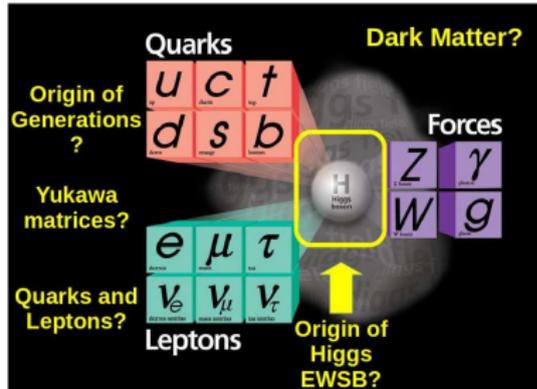
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(continuous spin rotation requires rest mass!)





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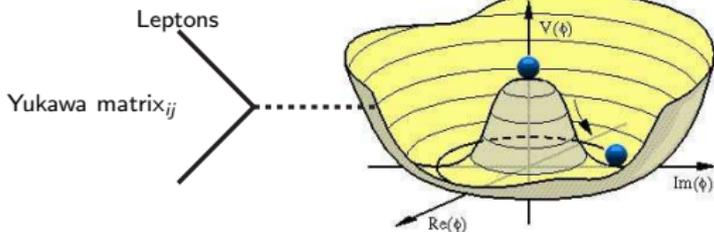
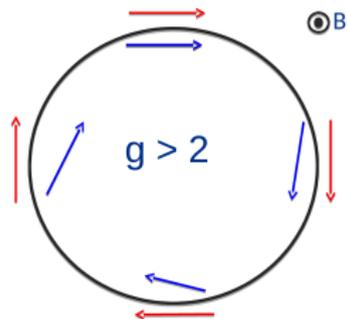
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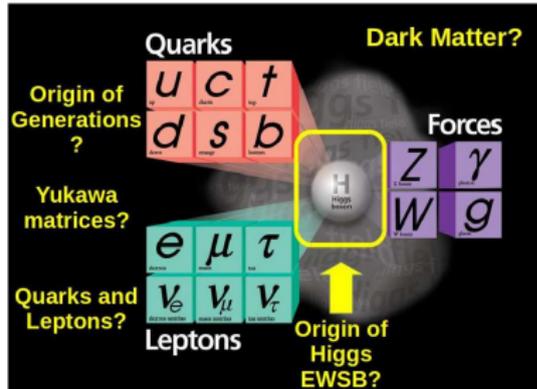
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(changed by new physics?)



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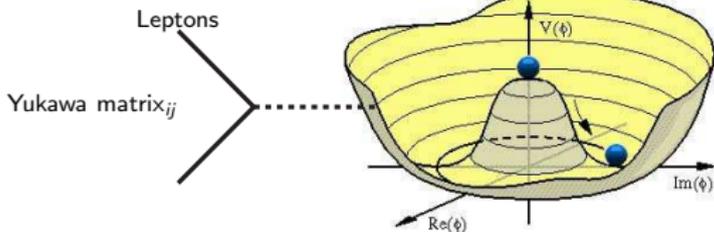
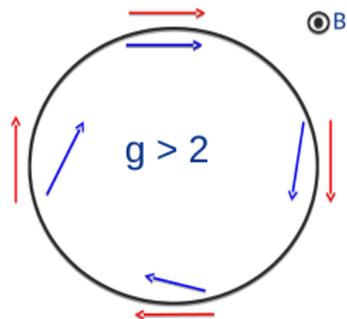
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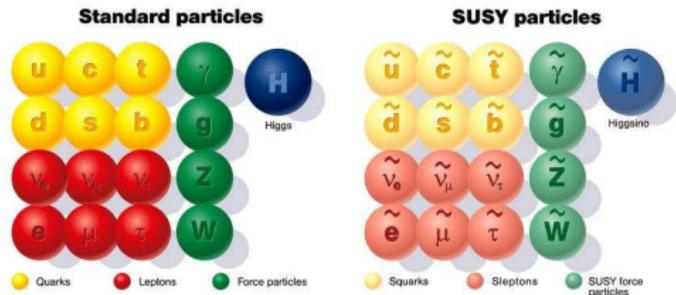
(changed by new physics?)

Example BSM idea

- fundamental new QFT symmetry
- predicts Higgs potential/mass
- dark matter candidate
- **chirality flip enhancement** $\rightsquigarrow g - 2$
- **viable (LHC)?**

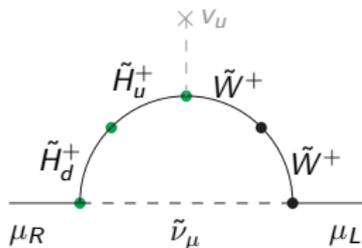
Example BSM idea Minimal SUSY Standard Model

- fundamental new QFT symmetry
- predicts Higgs potential/mass
- dark matter candidate
- **chirality flip enhancement** $\rightsquigarrow g - 2$
- **viable (LHC)?**



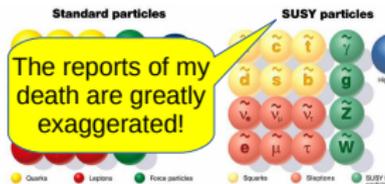
Superpartners and SUSY Higgs sector $\rightsquigarrow \tan \beta = \frac{v_u}{v_d}$, Higgsino mass μ

MSSM can explain $g - 2$ and dark matter

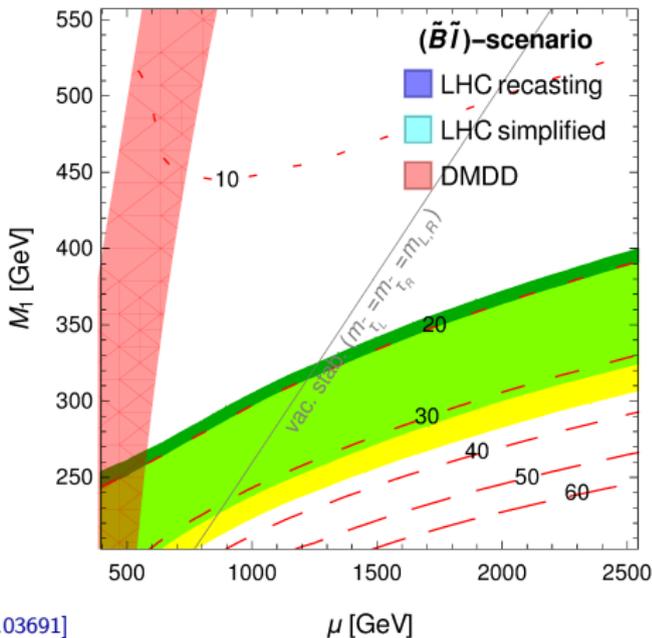


$$a_{\mu}^{\text{SUSY}} \approx 25 \times 10^{-10} \frac{\tan \beta}{50} \frac{\mu}{M_{\text{SUSY}}} \left(\frac{500 \text{ GeV}}{M_{\text{SUSY}}} \right)^2$$

- “Dark matter mass” versus μ
- explains $g - 2$ in large region (expands for $\tan \beta \neq 40$)
- DM explained by stau/slepton-coannihilation
- this automatically evades (current) LHC limits

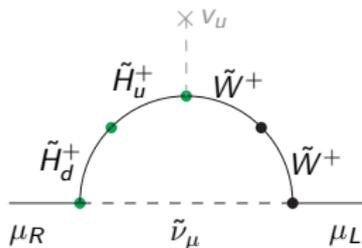


$m_{L,R} = M_1 + 50 \text{ GeV}, M_2 = 1200 \text{ GeV}, \tan \beta = 40$



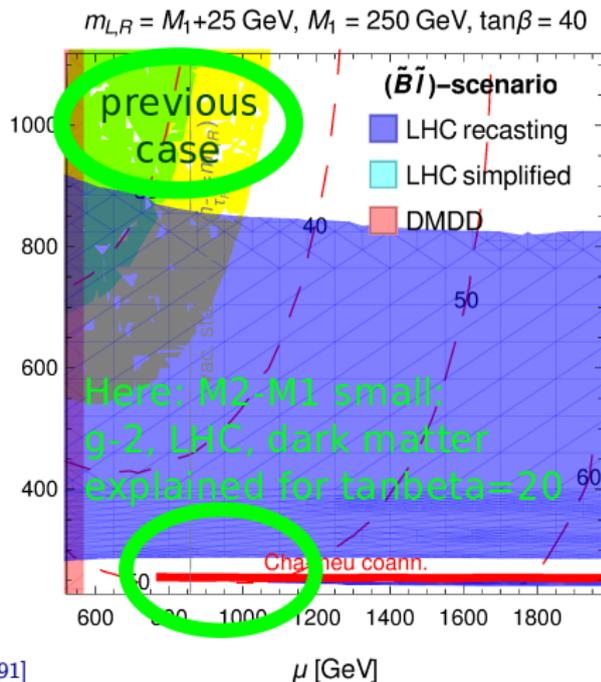
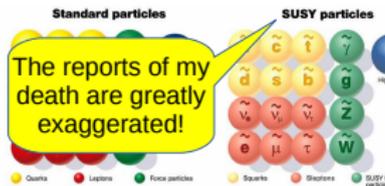
[2104.03691]

MSSM can explain $g - 2$ and dark matter



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- Strong LHC limits on M_2
- DM also explained by Wino-coannihilation
- again evades (current) LHC limits



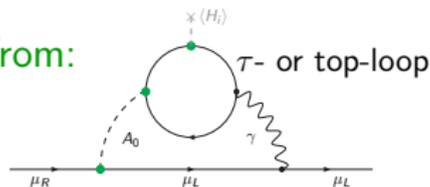
[2104.03691]

BSM with smaller masses, hidden from colliders?

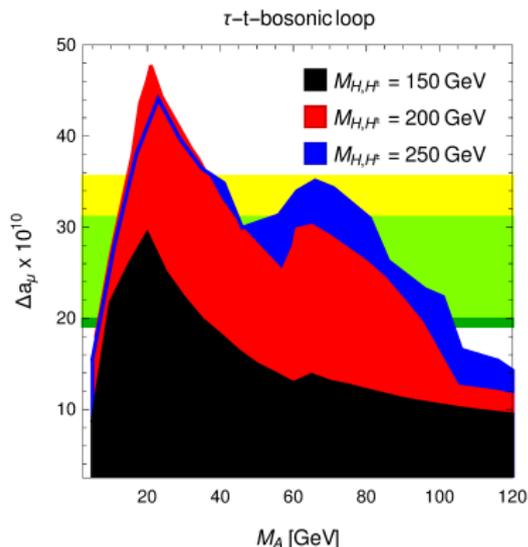
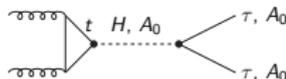
- Aligned 2-Higgs doublet model, rich new Higgs/Yukawa sectors

[compare: Type 2 pioneered by Maria Krawczyk, excluded for $g - 2$ e.g. by Mikolaj Misiak et al]

a_μ from:



LHC constraints:

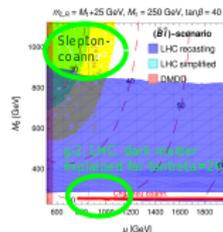
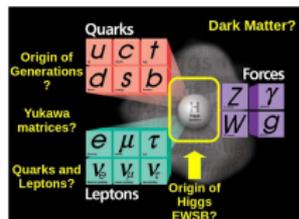
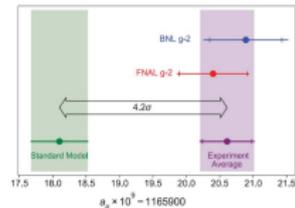


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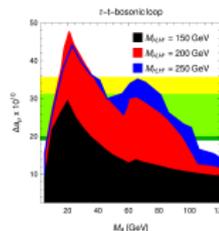
- can explain $g - 2$
- need large new Yukawa couplings
- under pressure, testable at LHC, lepton colliders, B-physics

Conclusions

- **SM prediction for $g - 2$:**
 - ▶ All known particles relevant (and all QFT tricks)
 - ▶ Theory Initiative: worldwide (ongoing!) effort, agreed & conservative value
- **BSM contributions to $g - 2$:**
 - ▶ large effect needed
 - ▶ Connections to deep questions
 - ▶ many viable models ... but
 - ▶ constraints from LHC, DM ...



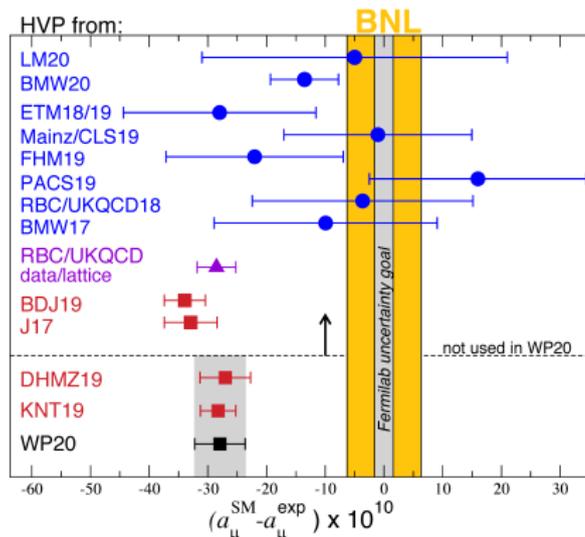
- **Outlook**
 - ▶ Fermilab Run-2,3,...
 - ▶ SM: scrutinize lattice, more e^+e^- -data, MUonE
 - ▶ Exp. tests of BSM models: Higgs couplings, B -physics, CLFV, EDM, light-particle searches, e^+e^-/μ on collider



20 years after BNL... deviation confirmed ... very promising future!

Details on hadronic vacuum polarization

a_μ^{HVP} : Status of Hadronic Vacuum Polarisation contributions



Lattice QCD + QED

- impressive progress, but...
- large spread between results
- tensions when looking at 'Euclidean time window' comparisons
- large systematic uncertainties (e.g. from non-trivial extrapolation to continuum limit, finite size)

Dispersive/lattice hybrid ('window' method)

For WP20: **Dispersive data-driven** from DHMZ and KNT

TI White Paper 2020 value:

$$a_\mu^{\text{HVP}} = 6845 (40) \times 10^{-11}$$

- TI WP2020 prediction uses **dispersive data-driven** evaluations with **minimal model dependence**
- a_μ^{HVP} **value and error** obtained by **merging** procedure \Rightarrow accounts for tensions in input data and differences in data treatment & combination (going beyond usual χ^2_{min} inflation)

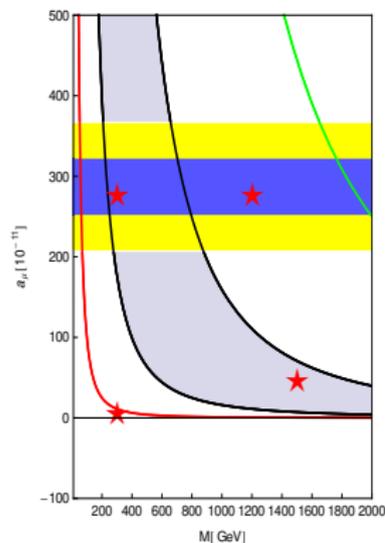
Thomas Teubner

10

Typical behaviour: \sim chirality flip (\rightsquigarrow Higgs!) and masses

$$a_\mu \sim \frac{m_\mu \times (\text{some VEV}) \times (\mu_{L \leftrightarrow R}\text{-flipping parameter})}{M_{\text{typical}}^2} \left[\lesssim \frac{m_\mu^2}{M_{\text{typical}}^2} \text{ (no finetuning)} \right]$$

$$\Delta m_\mu \sim (\text{some VEV}) \times (\mu_{L \leftrightarrow R}\text{-flipping parameter})$$



Typical behaviour: \sim chirality flip (\rightsquigarrow Higgs!) and masses

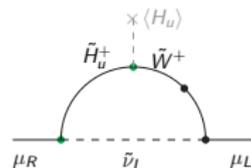
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• EWSM: $\alpha \frac{m_\mu^2}{M_W^2}$

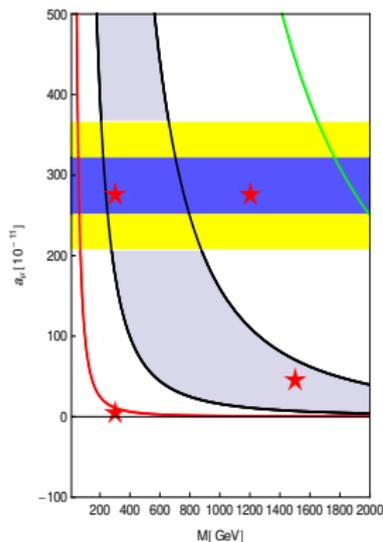
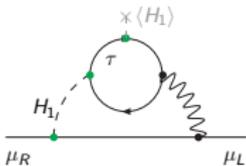


• SUSY: $\alpha \frac{m_\mu^2 \tan \beta}{M_{\text{SUSY}}^2}$

Well-motivated theory. Many other advantages



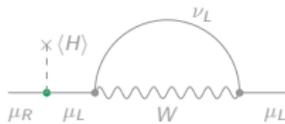
• 2HDM: $\alpha^2 \tan^2 \beta \frac{m_\mu^2}{M_H^2}$



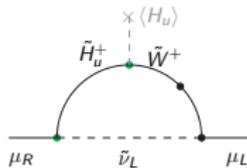
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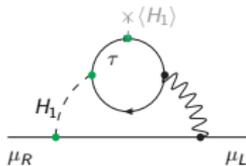
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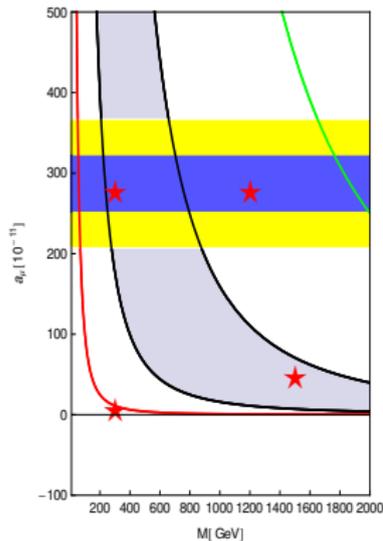
• SUSY: $\alpha \frac{m_\mu^2 \tan \beta}{M_{\text{SUSY}}^2}$



• 2HDM: $\alpha^2 \tan^2 \beta \frac{m_\mu^2}{M_H^2}$



Well motivated; many variants; many constraints



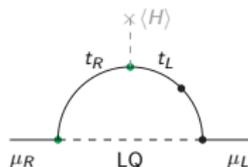
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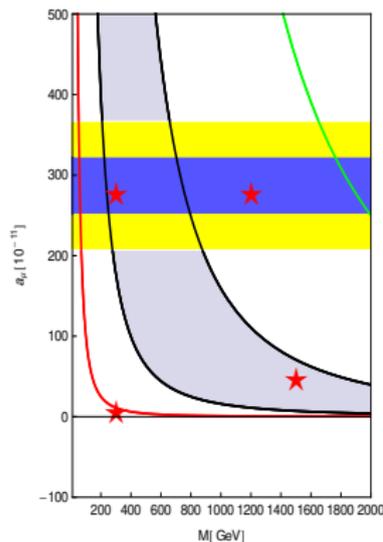
- EWSM: $\propto \frac{m_\mu^2}{M_W^2}$



- LQ: $gLGR \frac{m_\mu m_t}{M_{LQ}^2}$



- rad. $m_\mu \sim \frac{m_\mu^2}{M_{NP}^2}$



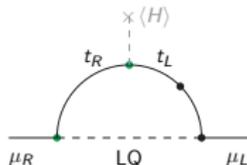
Typical behaviour: \sim chirality flip (\rightsquigarrow Higgs!) and masses

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- EWSM: $\alpha \frac{m_\mu^2}{M_W^2}$



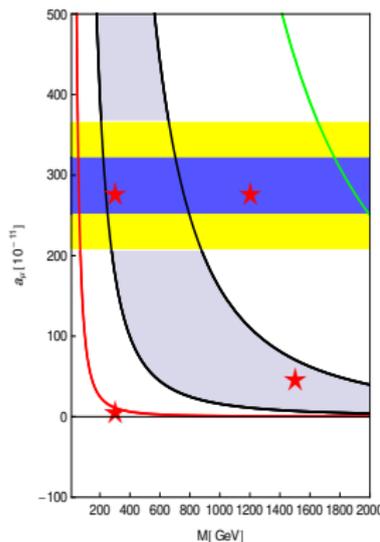
- LQ: $g_{LGR} \frac{m_\mu m_t}{M_{LQ}^2}$



Can also involve Higgs couplings to b , c or new particles.

Beware: $\Delta m_\mu / m_\mu \sim g_{LGR} m_t / m_\mu$ restricts couplings

- rad. $m_\mu \sim \frac{m_\mu^2}{M_{NP}^2}$



Take-home message 2:

Many models involve **enhancement mechanisms**

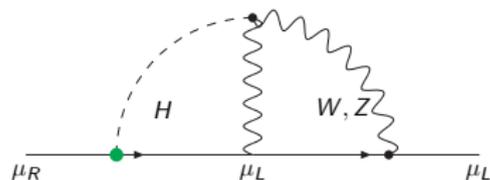
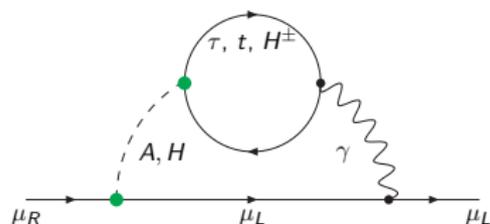
but: **experimental constraints!**

- LHC, LEP (particle mass limits; Higgs-data; EW-precision)
- Dark matter (direct detection limits!)
- Quark and lepton flavour, EDMs

Interlude — Role of BSM loops — examples

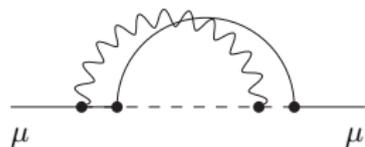
- EWSM: 2-loop = -20% of 1-loop — dominated by $\log(M_W/m_\mu)$
- 2HDM: 2-loop = leading order \Rightarrow full 2-loop prediction motivated

[Cherchiglia, Kneschke, DS, Stöckinger-Kim '16, '17]

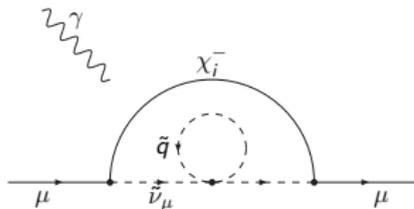


Role of BSM loops — examples

- SUSY: several 2-loop effects $\mathcal{O}(10\%)$ or more



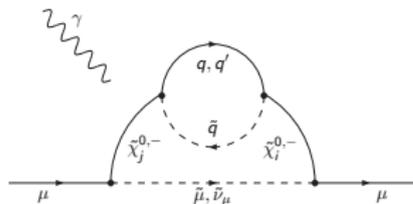
$$\rightarrow a_{\mu}^{1L} \times \frac{\alpha}{4\pi} \log\left(\frac{M_{\text{SUSY}}}{m_{\mu}}\right) \approx -9\%$$



$$\rightarrow a_{\mu}^{1L} \times \frac{\alpha}{4\pi} \frac{M_{\tilde{q}}^2}{M_{\mu}^2}$$

can* be $\mathcal{O}(100\%)$

* artifact of $\overline{\text{MS}}$ scheme
not present in on-shell scheme!



$$\rightarrow a_{\mu}^{1L} \times \frac{\alpha}{4\pi} \log\left(\frac{M_{\tilde{q}}}{M_{\text{SUSY}}}\right) \text{ up to } \mathcal{O}(10\%)$$

non-decoupling effect

Outline

1 Concrete models

2 Conclusions

Outline

1 Concrete models

- SUSY/MSSM, Two-Higgs doublet model, Leptoquarks, Vector-like leptons ...
- Overview of old results + [Athron,Balazs,Jacob,Kotlarski,DS,Stöckinger-Kim, preliminary]

2 Conclusions

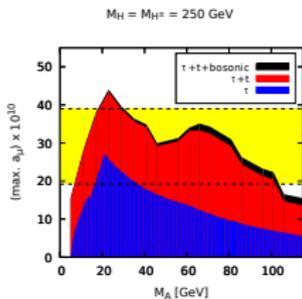
Outline

1 Concrete models

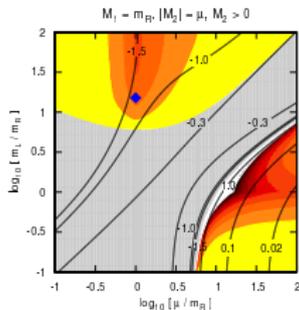
- SUSY/MSSM, Two-Higgs doublet model, Leptoquarks, Vector-like leptons . . .
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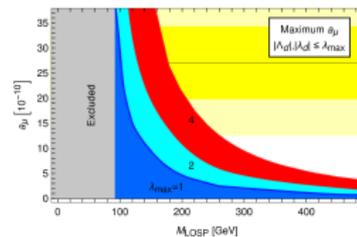
Example models and a_μ : “Largest results”



Largest THDM

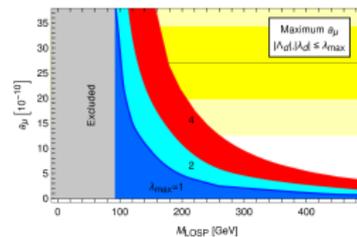
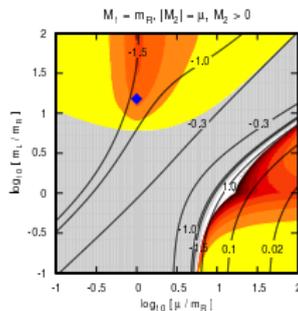
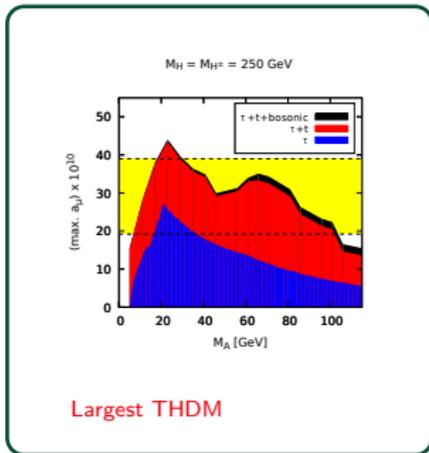


Largest SUSY ($\tan \beta \rightarrow \infty$)

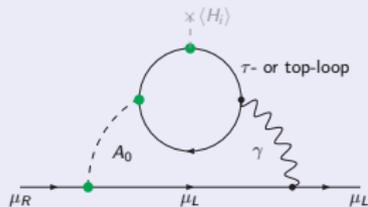


Largest MRSSM

Example models and a_μ : “Largest results”

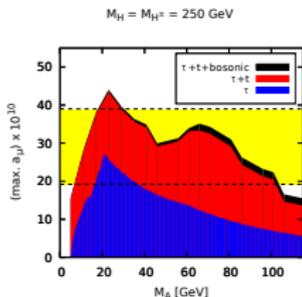


a_μ in the 2-Higgs doublet model? [Cherchiglia,DS,Stöckinger-Kim '17]

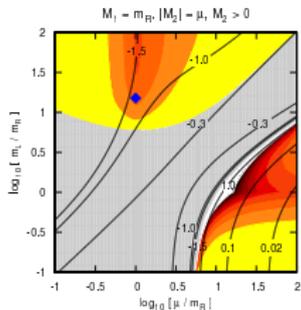


Results: a_μ explained in tightly constrained parameter space;
testable by many observables: $Z \rightarrow \tau\tau$, τ^- and b -decays, LHC $gg \rightarrow A, H \rightarrow \tau\tau$, future ILC?

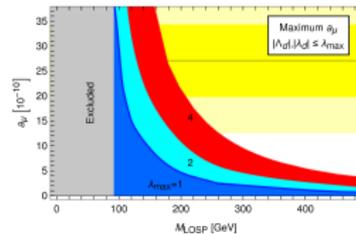
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Largest THDM



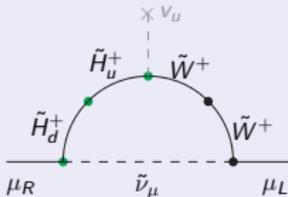
Largest SUSY ($\tan \beta \rightarrow \infty$)



Largest MRSSM

a_μ in special SUSY: $\tan \beta \rightarrow \infty$, MRSSM

[Bach,Park,DS,Stöckinger-Kim '15]
[Kotlarski,DS,Stöckinger-Kim '19]

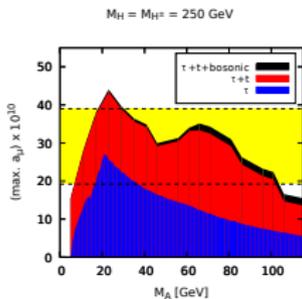


MSSM \neq SUSY! SUSY can be realized differently!

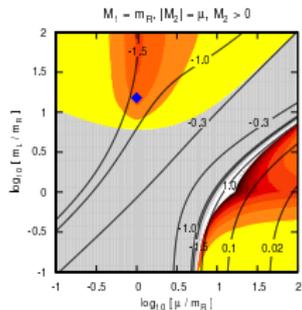
$\tan \beta \rightarrow \infty$: radiative m_μ

Result: a_μ explained even if $M_{LSP} > 1$ TeV

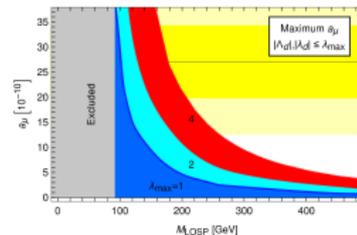
Example models and a_μ : “Largest results”



Largest THDM



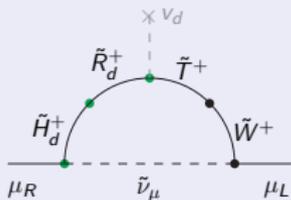
Largest SUSY ($\tan \beta \rightarrow \infty$)



Largest MRSSM

a_μ in special SUSY: $\tan \beta \rightarrow \infty$, MRSSM

[Bach,Park,DS,Stöckinger-Kim '15]
[Kotlarski,DS,Stöckinger-Kim '19]



MRSSM has more symmetry but
no $\tan \beta$ -enhancement!

Result: a_μ explained for $M_{\text{SUSY}} \sim 100\text{GeV}$,
compressed spectra;
testable by LHC/ILC, $\mu \rightarrow e/\mu \rightarrow e\gamma$

Which models are now promising in view of

a_{μ}^{BNL} , LHC and dark matter?

[Peter Athron, Csaba Balasz, Douglas Jacob, Wojciech Kotlarski, DS, Hyejung Stöckinger-Kim]

Which models can still accommodate large deviation?

SUSY: MSSM, MRSSM

- MSugra... many other generic scenarios
- Bino-dark matter+some coannihil.+mass splittings
- Wino-LSP+specific mass patterns

Two-Higgs doublet model

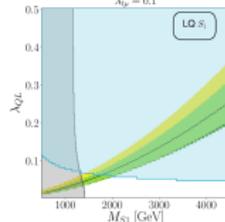
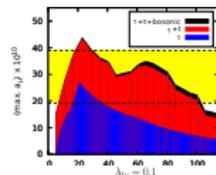
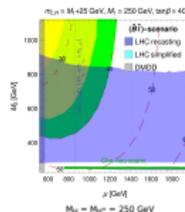
- Type I, II, Y, Type X(lepton-specific), flavour-aligned

Lepto-quarks, vector-like leptons

- scenarios with muon-specific couplings to μ_L and μ_R

Simple models (one or two new fields)

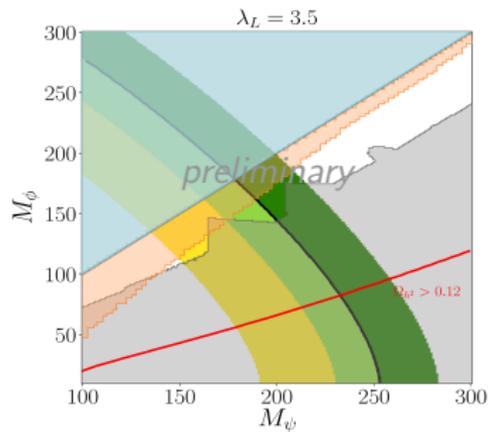
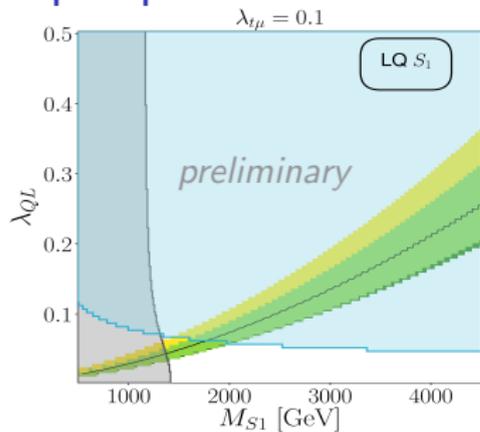
- Mostly excluded
- light N.P. (ALPs, Dark Photon, Light $L_\mu - L_\tau$)



Model	Mass	Exclusion
1	1.0 (1, 1)	Excluded (LQ X)
2	1.0 (1, 1)	Excluded (LQ X)
3	1.0 (1, 1)	Excluded (LQ X)
4	1.0 (1, 1)	Excluded (LQ X)
5	1.0 (1, 1)	Excluded (LQ X)
6	1.0 (1, 1)	Excluded (LQ X)
7	1.0 (1, 1)	Excluded (LQ X)
8	1.0 (1, 1)	Excluded (LQ X)
9	1.0 (1, 1)	Excluded (LQ X)
10	1.0 (1, 1)	Excluded (LQ X)
11	1.0 (1, 1)	Excluded (LQ X)
12	1.0 (1, 1)	Excluded (LQ X)
13	1.0 (1, 1)	Excluded (LQ X)
14	1.0 (1, 1)	Excluded (LQ X)
15	1.0 (1, 1)	Excluded (LQ X)
16	1.0 (1, 1)	Excluded (LQ X)
17	1.0 (1, 1)	Excluded (LQ X)
18	1.0 (1, 1)	Excluded (LQ X)
19	1.0 (1, 1)	Excluded (LQ X)
20	1.0 (1, 1)	Excluded (LQ X)

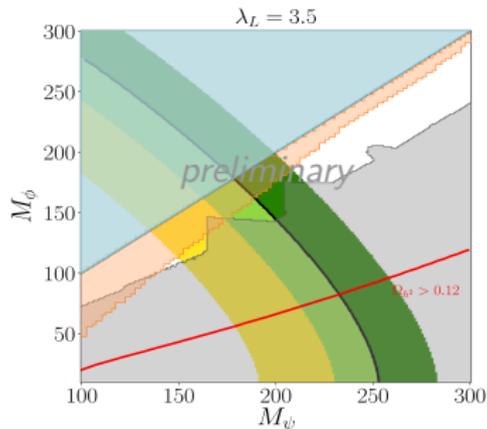
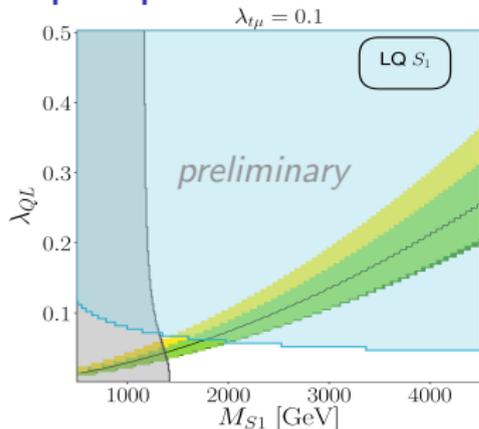
[Athron, Balazs, Jacob, Kotlarski, DS, Stöckinger-Kim, preliminary]

Leptoquarks and Model L with 2 fields



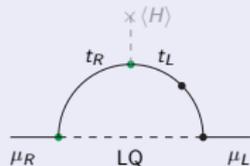
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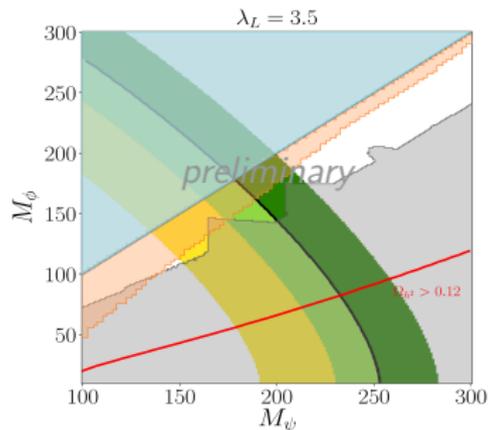
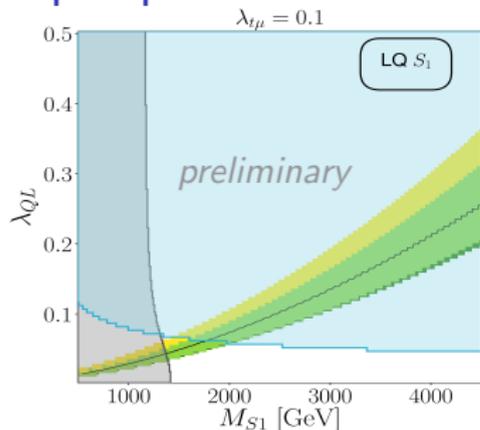
[Athron, Balazs, Jacob, Kotlarski, DS, Stöckinger-Kim, preliminary]

a_μ from LQ (or VLL) (motivation: simple/by GUTs, extra dim, ...)



- Chiral enhancement $\sim y_{\text{top}}, y_{\text{VLL}}$ versus y_μ
- LHC: lower mass limits
- Flavour constraints \rightsquigarrow
assume **only couplings to muons**
- Viable window above LHC (without m_μ -finetuning)

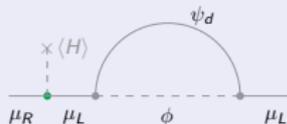
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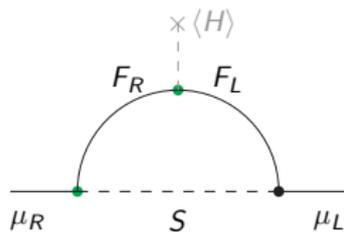
[Athron, Balazs, Jacob, Kotlarski, DS, Stöckinger-Kim, preliminary]

a_μ from 2-field model L

- No chiral enhancement, need very large couplings
- LHC: lower mass limits
- Dark matter candidate, but incompatible with large a_μ



Three-field models



- many models: viable, large chirality enhancements
- can explain a_μ^{BNL} and LHC and dark matter

MSSM fits very well: Chirality flip and $\tan \beta$ -enhancement

Charginos=mixtures $\tilde{H}-\tilde{B}-\tilde{W}$:

couplings to μ_L and to μ_R

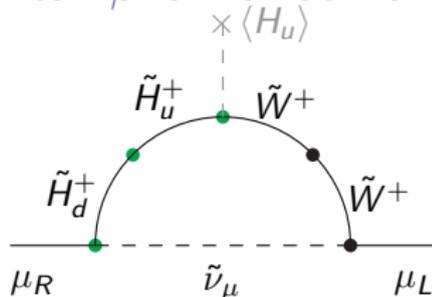


Diagram enhanced by Yukawa and large “other” vev, $\tan \beta = \langle H_u \rangle / \langle H_d \rangle$

$$\propto y_\mu \langle H_u \rangle \mu = m_\mu \tan \beta \mu \quad \rightarrow \quad a_\mu^{\text{SUSY}} \propto \tan \beta \text{sign}(\mu) \frac{m_\mu^2}{M_{\text{SUSY}}^2}$$

numerically: fits well if $\tan \beta = 10 \dots 50$ and $M_{\text{SUSY}} < \sim 500 \text{ GeV}$

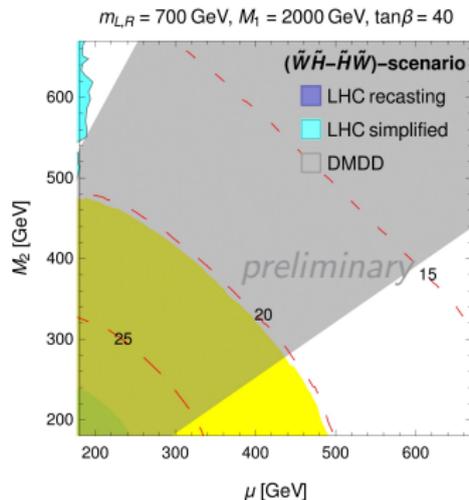
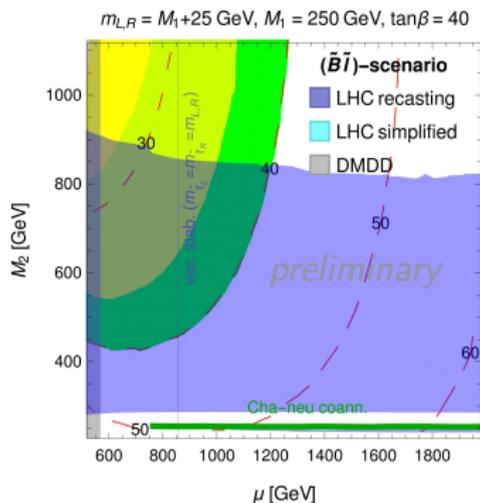
$$a_\mu^{\text{SUSY}} \sim 12 \times 10^{-10} \times \tan \beta \left(\frac{100 \text{ GeV}}{M_{\text{SUSY}}} \right)^2 \text{sign}(\mu)$$

Status of SUSY and a_μ

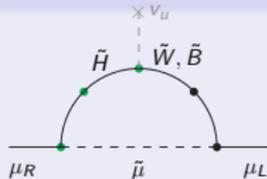
- SUSY and MSSM well-motivated!
- In general: large a_μ possible, precision computations available [GM2Calc]
- scenarios with large a_μ require only 3 light sparticles $\lesssim 500$ GeV
- constraints pull in different directions:

LHC \leftrightarrow M_h \leftrightarrow dark matter \leftrightarrow finetuning \leftrightarrow a_μ

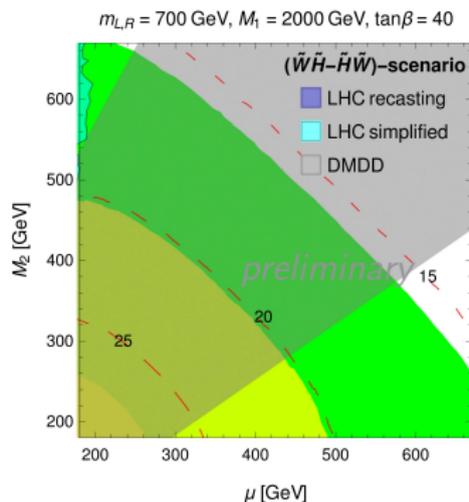
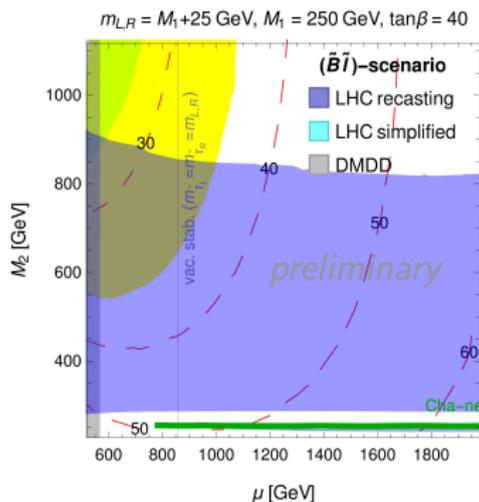
- “Constrained MSSM”: Higgs+LHC \Rightarrow stops, squarks very heavy \Rightarrow sleptons heavy \Rightarrow a_μ tiny!



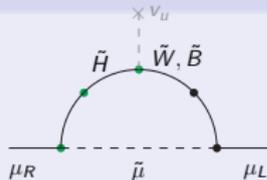
MSSM: well motivated, can explain large deviation (but...)



- LHC + Dark Matter \Rightarrow mass patterns!
- Co-annihil. regions; large $\mu \equiv m_{\tilde{H}}$; Wino-LSP; ...
- Excludes many simple scenarios (MSugra, ...)



MSSM: well motivated, can explain large deviation (but...)



- LHC + Dark Matter \Rightarrow mass patterns!
- Co-annihil. regions; large $\mu \equiv m_{\tilde{H}}$; Wino-LSP; ...
- Excludes many simple scenarios (MSugra, ...)

Outline

1 Concrete models

2 Conclusions

Summary of main points

discrepancy $\approx 2 \times a_\mu^{\text{SM,weak}}$

but: expect $a_\mu^{\text{NP}} \sim a_\mu^{\text{SM,weak}} \times \left(\frac{M_W}{M_{\text{NP}}}\right)^2 \times \text{couplings}$

Many models involve **enhancement mechanisms**

but: **experimental constraints!**

Take-home message 3:

Which models can still accommodate large deviation?

Many models! General ideas still viable (SUSY, THDM, LQ, VLL, ...)

but: **restricted parameter space!** Specific scenarios excluded!

What can now happen? Deviation confirmed? Back to SM?

Which models can still accommodate large deviation?

SUSY: **MSSM**, **MRSSM**

- **MSugra**. . . many other generic scenarios
- **Bino-dark matter**+some coannihil.+mass splittings
- **Wino-LSP**+specific mass patterns

Two-Higgs doublet model

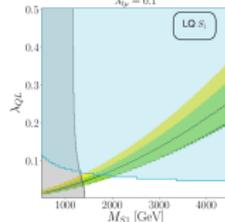
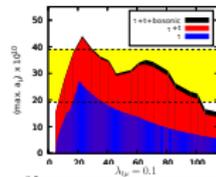
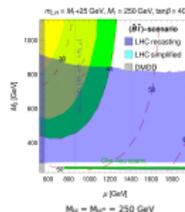
- **Type I, II, Y**, **Type X**(lepton-specific), flavour-aligned

Lepto-quarks, vector-like leptons

- scenarios with muon-specific couplings to μ_L and μ_R

Simple models (one or two new fields)

- **Mostly excluded**
- light N.P. (**ALPs**, **Dark Photon**, **Light $L_\mu - L_\tau$**)



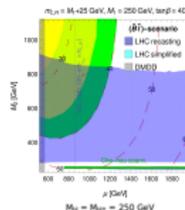
Model	Mass	Spin	CP	Notes
1	1.0	0	+	Scalar
2	1.0	0	-	Scalar
3	1.0	1	-	Vector
4	1.0	1	+	Vector
5	1.0	1	0	Vector
6	1.0	1	0	Vector
7	1.0	1	0	Vector
8	1.0	1	0	Vector
9	1.0	1	0	Vector
10	1.0	1	0	Vector
11	1.0	1	0	Vector
12	1.0	1	0	Vector
13	1.0	1	0	Vector
14	1.0	1	0	Vector
15	1.0	1	0	Vector
16	1.0	1	0	Vector
17	1.0	1	0	Vector
18	1.0	1	0	Vector
19	1.0	1	0	Vector
20	1.0	1	0	Vector
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25	1.0	1	0	Vector
26	1.0	1	0	Vector
27	1.0	1	0	Vector
28	1.0	1	0	Vector
29	1.0	1	0	Vector
30	1.0	1	0	Vector
31	1.0	1	0	Vector
32	1.0	1	0	Vector
33	1.0	1	0	Vector
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36	1.0	1	0	Vector
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38	1.0	1	0	Vector
39	1.0	1	0	Vector
40	1.0	1	0	Vector
41	1.0	1	0	Vector
42	1.0	1	0	Vector
43	1.0	1	0	Vector
44	1.0	1	0	Vector
45	1.0	1	0	Vector
46	1.0	1	0	Vector
47	1.0	1	0	Vector
48	1.0	1	0	Vector
49	1.0	1	0	Vector
50	1.0	1	0	Vector

[Athron,Balazs,Jacob,Kottarski,DS,Stöckinger-Kim, preliminary]

Which models can still accommodate large deviation?

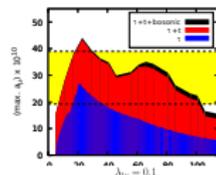
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- **Wino-LSP**+specific mass patterns



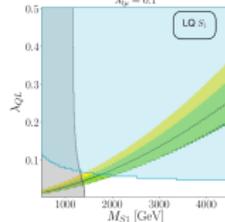
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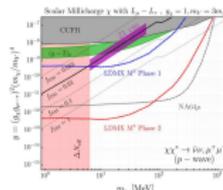
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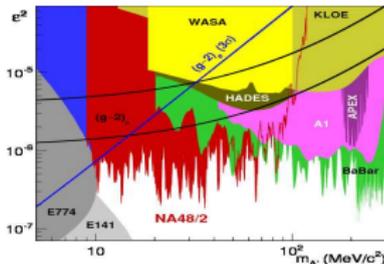
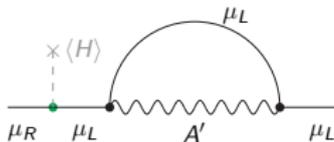
[Berlin, Blinov, Krnjaic, Schuster, Toro '18]

Light/dark sectors — compatible with large a_μ ?

Very light, very weakly interacting new particles

- “dark photon” **NO**

$$\mathcal{L} = -\frac{\epsilon}{2 \cos \theta_W} F^{\mu\nu} B_{\mu\nu} \quad a_\mu \sim \frac{\alpha}{2\pi} \epsilon^2$$

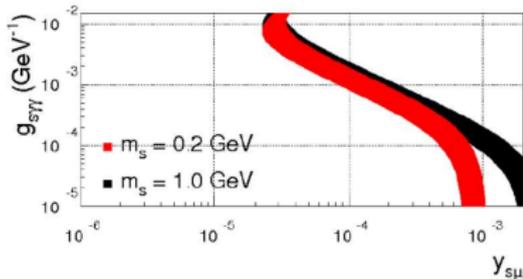
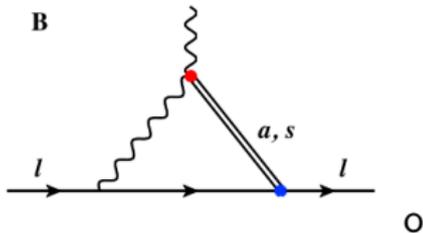


[NA48: 1504.00607]
excludes minimal dark photon for a_μ

- “ALPs” **YES**

$$\mathcal{L} = \frac{1}{4} g_{s\gamma\gamma} s F^{\mu\nu} F_{\mu\nu} + y_s s \bar{\mu} \mu$$

$$a_\mu^{\text{BZ}} \sim \frac{m_\mu}{4\pi^2} g_{s\gamma\gamma} y_s \ln(\Lambda/m_s)$$



[Marciano, Masiero, Paradisi, Passera '16]

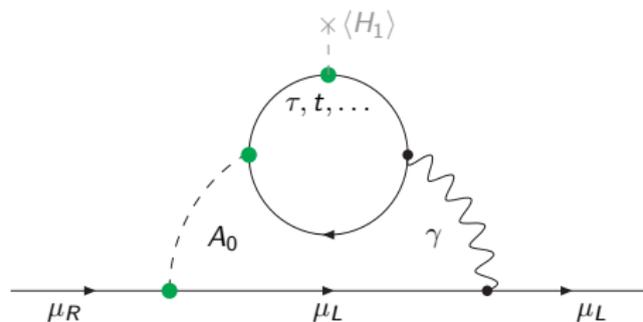
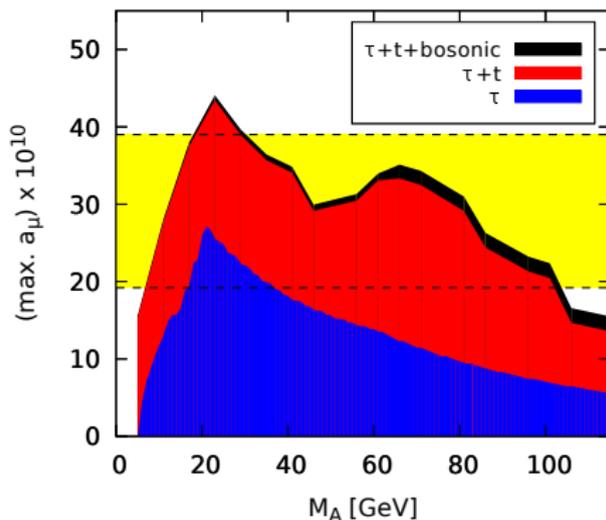
Two-Higgs doublet model — can it explain Δa_μ ?

[Cherchiglia, DS, Stöckinger-Kim '17]

Answer: YES (in small par. space)!

- light A_0 -boson, large couplings to leptons (and top-quarks)

$$M_H = M_{H^\pm} = 250 \text{ GeV}$$



• full computation

[Cherchiglia, Kneschke, DS, Stöckinger-Kim '16]

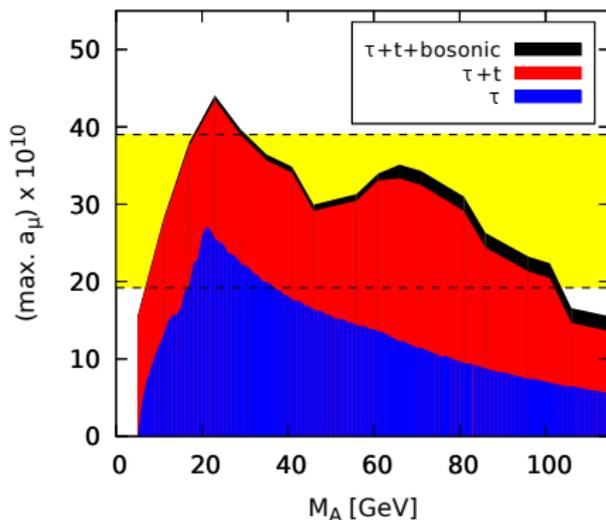
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constrained/testable by

- τ -, Z -decays, LEP
- b -decays, LHC

\Rightarrow maximum Yukawa couplings

- lepton Yukawa $\lesssim 100$
- quark Yukawas $\lesssim 0.5$
- (for $M_A = 20 \dots 100$ GeV, else even stronger)

More on bounds on dark photons, ALPs

- beam dump: dark bremsstrahlung (works for specific coupling range)
- electron fixed target (APEX,A1)
- Babar, KLOE, WASA: meson decays
- often assumed: $A' \rightarrow e^+e^-$ dominant
- if not: $K \rightarrow \pi A'$, $A' \rightarrow$ invisible and $e^+e^- \rightarrow \gamma +$ invisible lead to bounds

[Davoudiasl, Lee, Marciano '14][Izaguirre et al '13]

- generalization: also mass mixing “dark Z” with more general couplings, also strongly constrained
- for ALPs: [Marciano et al '16]
 $e^+e^- \rightarrow \gamma a$, $e^+e^- a$ at future low-E experiments

