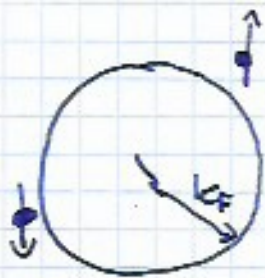


Cooper pair

A simplified model showing formation of a bound state - a Cooper pair



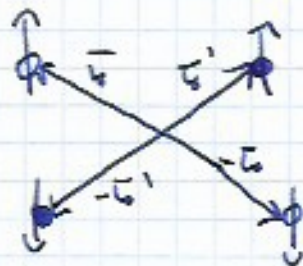
Variational (trial) wave function/state

$$|\Psi\rangle = \sum_{\vec{k} \neq 0} d_{\vec{k}} c_{-\vec{k}\downarrow}^{\dagger} c_{\vec{k}\uparrow}^{\dagger} |FS\rangle$$

Model Hamiltonian with interaction

$$\hat{H} = \sum_{\vec{k}\sigma} \epsilon_{\vec{k}} c_{\vec{k}\sigma}^{\dagger} c_{\vec{k}\sigma} + \sum_{\vec{k}\vec{k}'} V_{\vec{k}\vec{k}'} c_{\vec{k}\uparrow}^{\dagger} c_{-\vec{k}\downarrow}^{\dagger} c_{-\vec{k}'\downarrow} c_{\vec{k}'\uparrow}$$

$V_{\vec{k}\vec{k}'}$ - scattering potential, if $V_{\vec{k}\vec{k}'} < 0$
attractive



Normalization

$$c_{k\sigma} c_{k'\sigma'}^\dagger + c_{k'\sigma'}^\dagger c_{k\sigma} = \delta_{kk'} \delta_{\sigma\sigma'}$$

$$1 = \langle \Psi | \Psi \rangle = \sum_{k\sigma} \alpha_k^\dagger \alpha_k \langle FS | c_{k\uparrow} c_{-k\downarrow} c_{-k\downarrow}^\dagger c_{k\uparrow}^\dagger | FS \rangle =$$

$\delta_{kk} - c_{k\downarrow}^\dagger c_{k\downarrow}$

$$= \sum_{k\sigma} \alpha_k^\dagger \alpha_k \langle FS | c_{k\uparrow} c_{k\uparrow}^\dagger \delta_{kk} - c_{k\uparrow} c_{-k\downarrow} c_{-k\downarrow} c_{k\uparrow}^\dagger | FS \rangle =$$

$\delta_{kk} - c_{k\uparrow}^\dagger c_{k\uparrow}$ $= 0$ $(-)$

$$= \sum_{k\sigma} \alpha_k^\dagger \alpha_k \langle FS | \delta_{kk} - c_{k\uparrow}^\dagger c_{k\uparrow} + c_{k\uparrow} c_{-k\downarrow} c_{k\uparrow}^\dagger c_{-k\downarrow} | FS \rangle =$$

δ_{kk}

$$= \sum_{k\sigma} \delta_{kk} \alpha_k^\dagger \alpha_k \langle FS | FS \rangle = \sum_k |\alpha_k|^2 = 1$$

Matrix elements

$$\sum_{k\sigma} \epsilon_{k\sigma} c_{k\sigma}^\dagger c_{k\sigma} \sum_p \alpha_p c_{-p\downarrow}^\dagger c_{p\uparrow}^\dagger | FS \rangle =$$

$$= \sum_{k\sigma} \sum_p \epsilon_{k\sigma} \alpha_p c_{k\sigma}^\dagger (\delta_{k-p} \delta_{\sigma\downarrow} - c_{-p\downarrow}^\dagger c_{k\sigma}) c_{p\uparrow}^\dagger | FS \rangle =$$

$$= \sum_{k\sigma} \sum_p \epsilon_{k\sigma} \alpha_p (c_{k\sigma}^\dagger c_{p\uparrow}^\dagger \delta_{k-p} \delta_{\sigma\downarrow} - c_{k\sigma}^\dagger c_{-p\downarrow}^\dagger c_{k\sigma} c_{p\uparrow}^\dagger) | FS \rangle =$$

$$= \sum_{k\sigma} \sum_p \epsilon_{k\sigma} \alpha_p \delta_{k-p} \delta_{\sigma\downarrow} c_{k\sigma}^\dagger c_{p\uparrow}^\dagger | FS \rangle -$$

$$- \sum_{k\sigma} \sum_p \epsilon_{k\sigma} \alpha_p c_{k\sigma}^\dagger c_{-p\downarrow}^\dagger (\delta_{k\uparrow} \delta_{\sigma\uparrow} - c_{p\uparrow}^\dagger c_{k\sigma}) | FS \rangle =$$

δ_{kk}

$$= \sum_k \sum_p \underbrace{\epsilon_{-p} \alpha_p}_{\beta_p} c_{-p\downarrow}^\dagger c_{p\uparrow}^\dagger | FS \rangle - \sum_k \sum_p \epsilon_p \alpha_p \underbrace{c_{p\uparrow}^\dagger c_{-p\downarrow}^\dagger}_{\beta_p} | FS \rangle =$$

$$= 2 \sum_p \beta_p \alpha_p c_{-p\downarrow}^\dagger c_{p\uparrow}^\dagger | FS \rangle$$

$\epsilon_{\vec{k}} = \epsilon_{-\vec{k}}$ Kramer's
Symmetry
time-reversal
Symmetry (2)

$$\sum_{k \neq l} V_{kl} c_k^\dagger c_l^\dagger c_{-k} c_{-l} \underbrace{\sum_p d_p c_{-p} c_p^\dagger}_{=0} |FS\rangle =$$

$$= - \sum_{k \neq l} \sum_p V_{kl} d_p c_k^\dagger c_{-k} c_{-l} c_l^\dagger (\delta_{k-p} - c_p^\dagger c_p) |FS\rangle$$

$$= - \sum_{k \neq l} \sum_p V_{kl} d_p c_k^\dagger c_{-k} d_{k-p} (\delta_{-k-p} - c_p^\dagger c_p) |FS\rangle$$

$$= - \sum_k \sum_p V_{kp} d_p c_k^\dagger c_{-k} |FS\rangle = + \sum_{k \neq p} d_p V_{kp} c_{-k} c_k^\dagger |FS\rangle$$

$$\langle \Psi | \hat{H} | \Psi \rangle = \sum_{k \neq l} 2 \epsilon_q d_k d_l \underbrace{c_k^\dagger c_{-k} c_{-l} c_l^\dagger}_{\text{only } k=l} |FS\rangle$$

$$+ \sum_{k \neq l} d_p d_l V_{kl} \underbrace{c_k^\dagger c_{-k} c_{-l} c_l^\dagger}_{\text{only } k=l} |FS\rangle = \text{energy functional}$$

$$= \sum_k 2 \epsilon_k |d_k|^2 + \sum_{k \neq l} V_{kl} d_k^\dagger d_l \equiv E(\{d_k\})$$

$$\frac{\partial}{\partial d_k} \left[E(\{d_k\}) - \lambda \left(\sum_k |d_k|^2 - 1 \right) \right] = 0 \quad \text{variational principle}$$

↳ Lagrange multiplier

$$2 \epsilon_k d_k + \sum_{k'} V_{kk'} d_{k'} - \lambda d_k = 0$$

$$\alpha_{\vec{k}} = - \frac{\sum_{\vec{k}'} V(\vec{k}-\vec{k}') \alpha_{\vec{k}'}}{2 \epsilon_{\vec{k}} - \lambda}$$

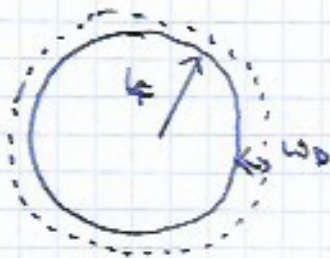
integral
equation on
 $\alpha_{\vec{k}}$ with eigenvalue λ

Model potential

$$V(\vec{k}-\vec{k}') = \begin{cases} -V_0 & \text{attraction} \\ 0 & \text{otherwise} \end{cases}$$

Debye cutoff
 \downarrow
 $\epsilon_{\vec{k}} < \omega_D$
 otherwise

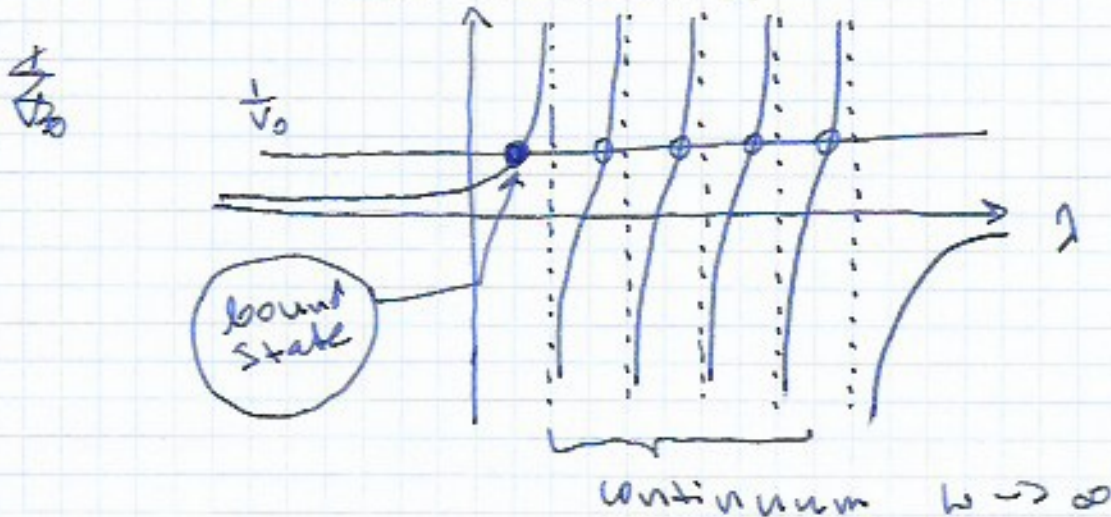
$V_0 > 0$



$$\alpha_{\vec{k}} = V_0 \frac{\sum_{\vec{k}'} \alpha_{\vec{k}'}}{2 \epsilon_{\vec{k}} - \lambda} \quad / \quad \sum_{\vec{k}'} \alpha_{\vec{k}'}$$

$$\sum_{\vec{k}'} \alpha_{\vec{k}'} = V_0 \sum_{\vec{k}'} \frac{1}{2 \epsilon_{\vec{k}} - \lambda} \sum_{\vec{k}'} \alpha_{\vec{k}'}$$

$$\frac{1}{V_0} = \sum_{\vec{k}'} \frac{1}{2 \epsilon_{\vec{k}} - \lambda}$$

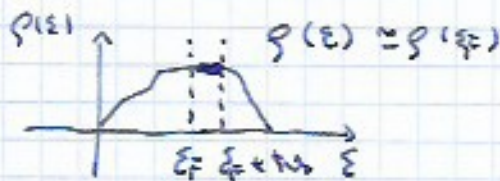


Solution

$$\frac{1}{V_0} = \int_{E_F}^{E_F + \hbar \omega_D} d\varepsilon$$

DOS

$$\frac{\rho(\varepsilon)}{2\varepsilon - \lambda} \approx \rho(E_F) \int_{E_F}^{E_F + \hbar \omega_D} d\varepsilon \frac{1}{2\varepsilon - \lambda}$$



$$= \rho(E_F) \frac{1}{2} \ln \left| \frac{2(E_F + \hbar \omega_D) - \lambda}{2E_F - \lambda} \right|$$

$$\lambda = 2E_F - 2\hbar \omega_D \frac{e^{-\frac{2}{\rho(E_F)V_0}}}{1 - e^{-\frac{2}{\rho(E_F)V_0}}} \approx 2E_F - 2\hbar \omega_D e^{-\frac{2}{V_0 \rho(E_F)}}$$

What is λ ?

$$E(\{d_k\}) \rightarrow \sum_k (d_k)^2 + \lambda =$$

$$= \sum_k 2\varepsilon_k |d_k|^2 + \sum_{k,k'} V_{kk'} d_k d_{k'} - \lambda \sum_k |d_k|^2 + \lambda =$$

$$= \sum_k d_k^* \left(2\varepsilon_k d_k + \sum_{k'} V_{kk'} d_{k'} - \lambda d_k \right) + \lambda = \lambda$$

λ is the energy at the minimum

$$E = 2E_F - 2\hbar \omega_D e^{-\frac{2}{V_0 \rho(E_F)}}$$

two electrons
at the Fermi surface

binding
energy

(*) non-perturbative result, $\rho V_0 \ll 1$ not expandable

Schrödinger equation

$$\hat{H} |\Psi\rangle = E |\Psi\rangle$$

$$\left| \sum_{k \neq p} \epsilon_k c_{k\uparrow} c_{k\downarrow} + \sum_{k \neq p} V_{pk} c_{k\uparrow} c_{k\downarrow} c_{p\downarrow} c_{p\uparrow} \right] \sum_p \alpha_p c_{p\downarrow} c_{p\uparrow} |\Psi\rangle = E \sum_p \alpha_p c_{p\downarrow} c_{p\uparrow} |\Psi\rangle$$

$$\sum_p \left[2 \epsilon_p \alpha_p c_{p\downarrow} c_{p\uparrow} |\Psi\rangle + \sum_{k \neq p} \alpha_k V_{kp} c_{k\downarrow} c_{k\uparrow} |\Psi\rangle \right] = E \sum_p \alpha_p c_{p\downarrow} c_{p\uparrow} |\Psi\rangle$$

$$\sum_p \left[\underbrace{2 \epsilon_p \alpha_p + \sum_k V_{kp} \alpha_k}_{=0} - E \alpha_p \right] c_{p\downarrow} c_{p\uparrow} |\Psi\rangle = 0$$

$$\alpha_p = - \frac{\sum_k V_{kp} \alpha_k}{2 \epsilon_p - E}$$

the same integral equation as before

The total wave function $|\Psi\rangle$ is apparently an exact state of \hat{H} .