

# 2-point correlation function

$$\hat{H} = \frac{\hat{p}^2}{2m} + \frac{1}{2} m \omega^2 \hat{x}^2$$

$$G(\tau) = \frac{1}{Z} \text{Tr} \left[ e^{-\beta \hat{H}} \hat{x}(\tau) \hat{x}(0) \right]$$

$$\hat{x}(\tau) = e^{\frac{\hat{H}\tau}{\hbar}} \hat{x} e^{-\frac{\hat{H}\tau}{\hbar}}, \quad 0 < \tau < \beta \hbar$$

Show, using standard QM that

$$G(\tau) = \frac{\hbar}{2m\omega} \frac{\cosh\left[\left(\frac{\beta\hbar}{2} - \tau\right)\omega\right]}{\sinh\left[\frac{\beta\hbar\omega}{2}\right]}$$

## Path integral

$$G(\tau) = \frac{\int_{x(\beta\hbar)=x(0)} \mathcal{D}[x(\tau)] x(\tau) x(0) e^{-\frac{SE}{\hbar}}}{\int_{x(\beta\hbar)=x(0)} \mathcal{D}[x(\tau)] e^{-\frac{SE}{\hbar}}}$$

The  $c'$  factor drops out!

$$x(\tau) = \frac{1}{\beta} \left[ a_0 + \sum_{\ell=1}^{\infty} \left[ (a_{\ell} + i b_{\ell}) e^{i\omega_{\ell}\tau} + (a_{\ell} - i b_{\ell}) e^{-i\omega_{\ell}\tau} \right] \right]$$

$$x(0) = \frac{1}{\beta} \left[ a_0 + \sum_{\ell=1}^{\infty} 2a_{\ell} \right]$$

$$G(\tau) = \langle x(\tau) x(0) \rangle = \frac{\int da_0 \int \prod_{n \geq 1} da_n db_n x(\tau) x(0) e^{-\frac{SE}{\hbar}}}{\int da_0 \int \prod_{n \geq 1} da_n db_n e^{-\frac{SE}{\hbar}}}$$

since,  $a_0, a_n, b_n \in \mathbb{R}$  and  $SE$  is Gaussian

$$\langle a_0 a_n \rangle = \langle a_0 b_n \rangle = \langle a_n b_n \rangle = 0$$

$$\langle a_n a_{\ell} \rangle \approx \langle b_n b_{\ell} \rangle \approx \delta_{n\ell}$$

$\langle \dots \rangle$  in sense of  $(*)$



$$\bullet G(z) = \frac{1}{\beta^2} \left\langle a_0^2 + \sum_{\ell=1}^{\infty} 2a_{\ell}^2 (e^{i\omega_{\ell}z} + e^{-i\omega_{\ell}z}) \right\rangle$$

where

$$\begin{aligned} \langle a_0^2 \rangle &= \frac{\int da_0 a_0^2 e^{-\frac{m}{2\beta} \omega^2 a_0^2}}{\int da_0 e^{-\frac{m}{2\beta} \omega^2 a_0^2}} = \\ &= - \frac{2}{m\omega^2} \frac{\partial}{\partial (\frac{\beta}{2})} \ln \int da_0 e^{-\frac{m}{2\beta} \omega^2 a_0^2} = \\ &= - \frac{2}{m\omega^2} \frac{\partial}{\partial (\frac{\beta}{2})} \ln \sqrt{\frac{2\pi\beta}{m\omega^2}} = \frac{\beta}{m\omega^2} \end{aligned}$$

$$\langle a_{\ell}^2 \rangle = \frac{\int da_{\ell} a_{\ell}^2 e^{-\frac{m}{\beta} (\omega_{\ell}^2 + \omega^2) a_{\ell}^2}}{\int da_{\ell} e^{-\frac{m}{\beta} (\omega_{\ell}^2 + \omega^2) a_{\ell}^2}} = \frac{\beta}{2m(\omega_{\ell}^2 + \omega^2)}$$

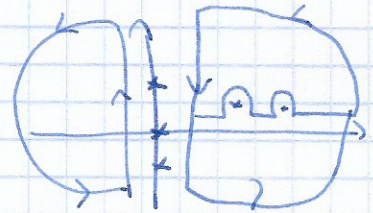
Thus we get

$$G(z) = \frac{1}{m\beta} \left( \frac{1}{\omega^2} + \sum_{\ell=1}^{\infty} \frac{e^{i\omega_{\ell}z} + e^{-i\omega_{\ell}z}}{\omega_{\ell}^2 + \omega^2} \right) = \frac{1}{m\beta} \sum_{k=-\infty}^{\infty} \frac{e^{i\omega_k z}}{\omega_k^2 + \omega^2}$$

Using a standard sum over Matsubara frequency (lecture)

$$\sum_n f(i\omega_n) = - \frac{\beta}{2\pi i} \oint_C \frac{f(z)}{e^{\beta z} - 1} dz$$

$$\oint_C f(z) dz = 2\pi i \sum_i \text{Res} f(z_i)$$



We get

$$\bullet G(z) = \frac{1}{2m\omega} \frac{\text{ch} \left[ \left( \frac{\beta t}{2} - z \right) \omega \right]}{\text{ch} \left[ \frac{\beta t \omega}{2} \right]}$$