

## Feynman diagrams

- To each perturbative contribution generated by Wick's theorem, one can associate a graph called a Feynman diagram.
- All contributions can be derived from a subclass that contains only connected terms, represented by connected diagrams.

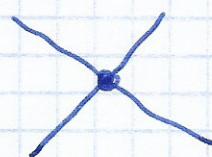
$$\begin{aligned}
 Z(\lambda) &= \int d^n x e^{-A(x)} + \lambda V(x) = & A(x) = \frac{1}{2} \sum_{i,j} x_i A_{ij} x_j \\
 &= \sum_{k=0}^{\infty} \frac{(-\lambda)^k}{k!} \int d^n x V(x)^k e^{-A(x)} = & V(x) - \text{polynomial} \\
 &= Z(0) \sum_{k=0}^{\infty} \frac{(-\lambda)^k}{k!} \langle V(x)^k \rangle_0
 \end{aligned}$$

where

$$\langle V(x)^k \rangle_0 = \frac{1}{Z(0)} \int d^n x V(x)^k e^{-A(x)}$$

## Feynman diagrams

Each monomial contributing to a perturbator  $V(I)$  is represented by a point (a vertex) from which originates a number of lines equal to the degree of the monomial



$$x^4 \text{ vertex in } V(x) = \frac{1}{4!} \sum_{i=1}^n x_i^4$$

Each pairing is represented by a line joining the vertices to which belongs the corresponding variable.

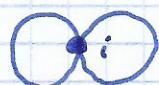
(1)

Example

$$V(\lambda) = \frac{1}{4!} \sum_{i=1}^n x_i^4$$

$$\begin{aligned}\frac{Z(\lambda)}{Z(0)} &= 1 - \frac{1}{4!} \lambda \sum_i \langle x_i^4 \rangle_0 + \frac{1}{2!(4!)^2} \lambda^2 \sum_{ij} \langle x_i^4 x_j^4 \rangle_0 + O(\lambda^3) = \\ &= 1 - \frac{1}{8} \lambda \sum_i \Delta_{ii}^2 + \frac{1}{128} \lambda^2 \sum_i \Delta_{ii}^2 \sum_j \Delta_{jj}^2 + \\ &\quad + \lambda^2 \sum_{ii} \left( \frac{1}{16} \Delta_{ii} \Delta_{jj} \Delta_{ij}^2 + \frac{1}{48} \Delta_{ij}^4 \right) + O(\lambda^3)\end{aligned}$$

diagrams of order  $\lambda^1$ :



$$\langle x_i^4 \rangle_0$$

diagrams of order  $\lambda^2$ :



- non-connected

contributions to  $\langle x_i^4 x_j^4 \rangle_0$   
(product of sums)



- connected

contributions to  
 $\langle x_i^4 x_j^4 \rangle_0$

Latter we will not indicate indices of summation

(Σ)