Kinks in the dispersion of strongly correlated electrons

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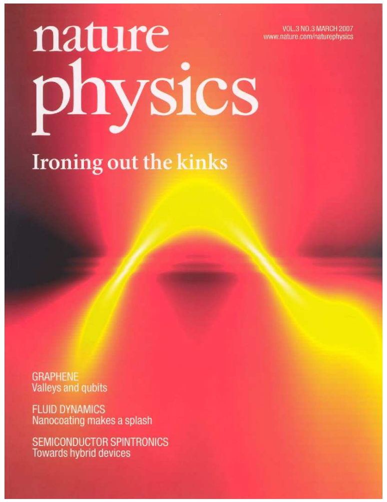


K. Byczuk, M. Kollar, K. Held, Y.-F. Yang, I.A. Nekrasov, Th. Pruschke, D. Vollhardt

Nature Physics **3**, 168 (2007)

Collaboration

- M. Kollar, D. Vollhardt, Augsburg, Germany
- K. Held, Y.-F. Yang, Stuttgart, Germany
- I. Nekrasov, Ekaterinburg, Russia
- T. Pruschke, Göttingen, Germany

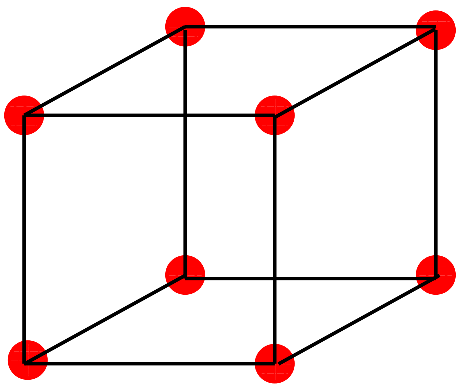
Support from SFB 484

Aim of this talk

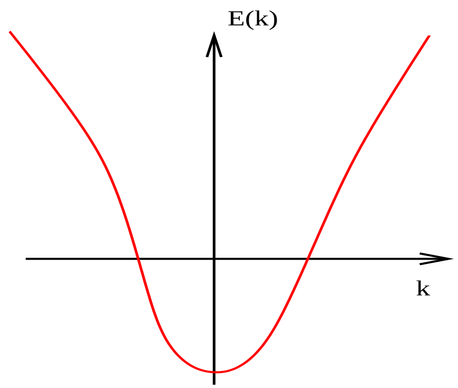
New Purely Electronic Mechanism for Kinks in Electronic Dispersion Relations

- in strongly correlated electron systems
- characteristic energy scale
- range of validity for Fermi liquid theory

Standard model of quantum many-body system



emergent particles
quasiparticle
quasihole
holon
spinon
plasmon
magnon
phonon
polariton
exciton
anyon
g-on
...



(i) well defined dispersion relation $E(\mathbf{k})$

(ii) long (infinite) life-time τ

(iii) proper set of quantum numbers

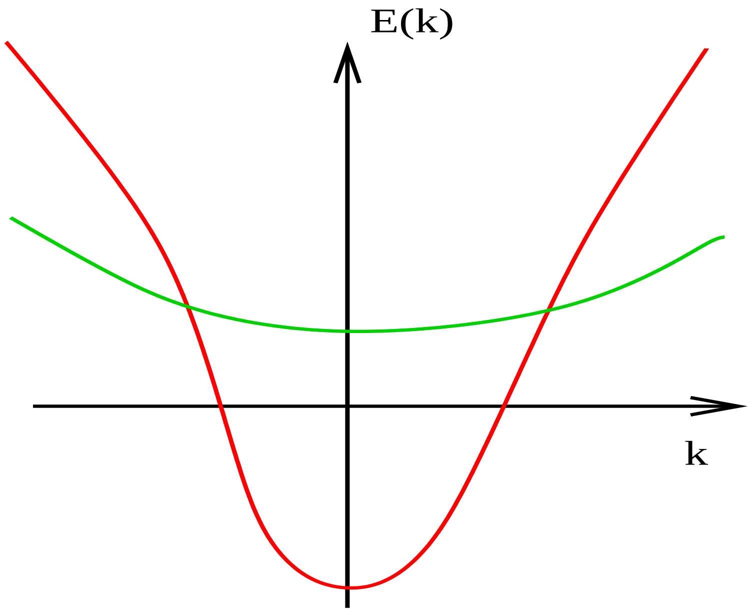
(iv) statistics

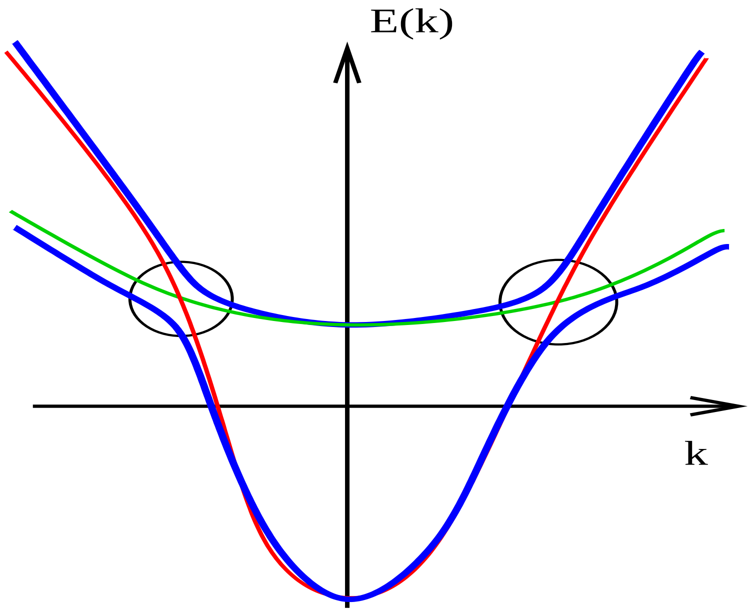
Dispersions and kinks

Coupling/hybridization $\hat{V}$ between different particles/modes

$$\langle \Psi | \hat{V} | \Phi \rangle \neq 0$$

anticrossing, lifting degeneracy, ...

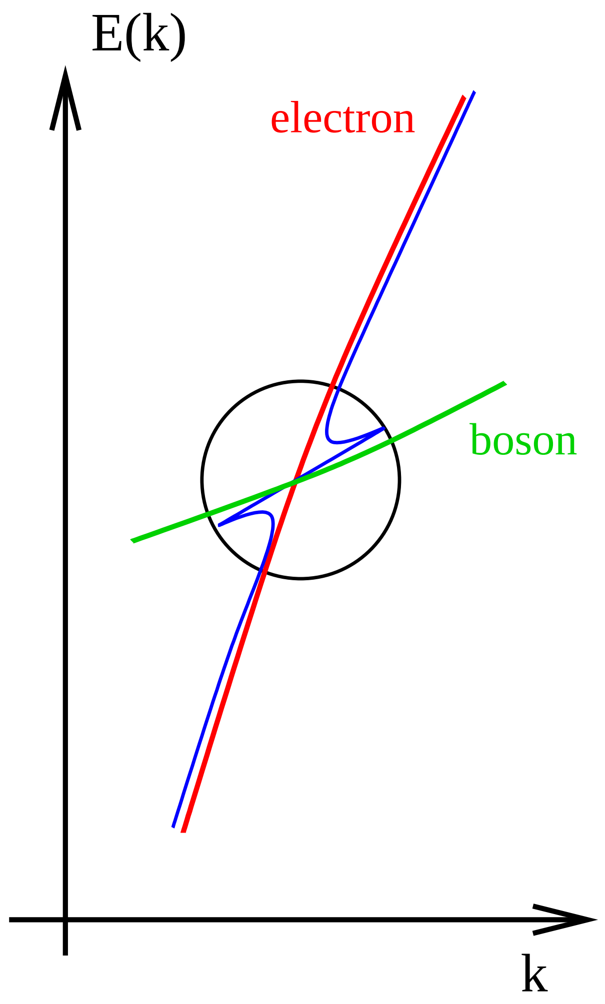




Df. Kinks are abrupt slope changes in the dispersion relations

They provide information on coupling between modes

Dispersions and kinks - coupling to bosons



energy of a kink is related to energy of a bosonic fluctuation

Dispersion of correlated electrons

One-particle spectral function - excitations at $\mathbf{k}$ and ω

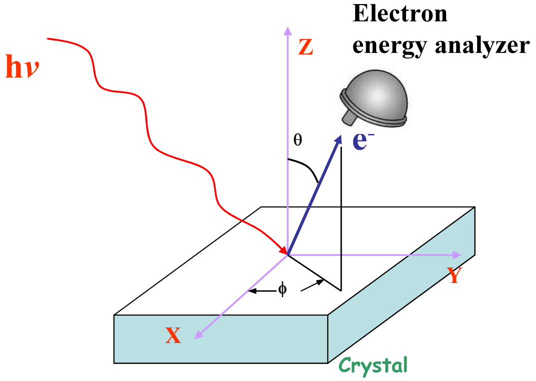
$$A(\mathbf{k},\omega) = -\frac{1}{\pi}\mathrm{Im}\frac{1}{\omega + \mu - \epsilon_{\mathbf{k}} - \Sigma(\mathbf{k},\omega)}$$

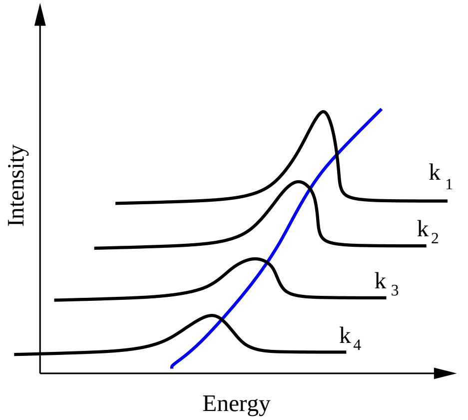
Dispersion relation $E_{\mathbf{k}}$

$$E_{\mathbf{k}} = \{\omega \text{ where } A(\mathbf{k},\omega) = \max\}$$

Dispersion relation is experimentally measured

Angular Resolved Photoemission Spectroscopy



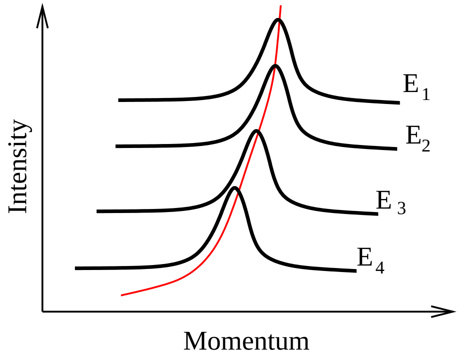


energy distribution curve (EDC)

$$k_x = k \cos \phi$$
$$k_y = k \sin \phi$$

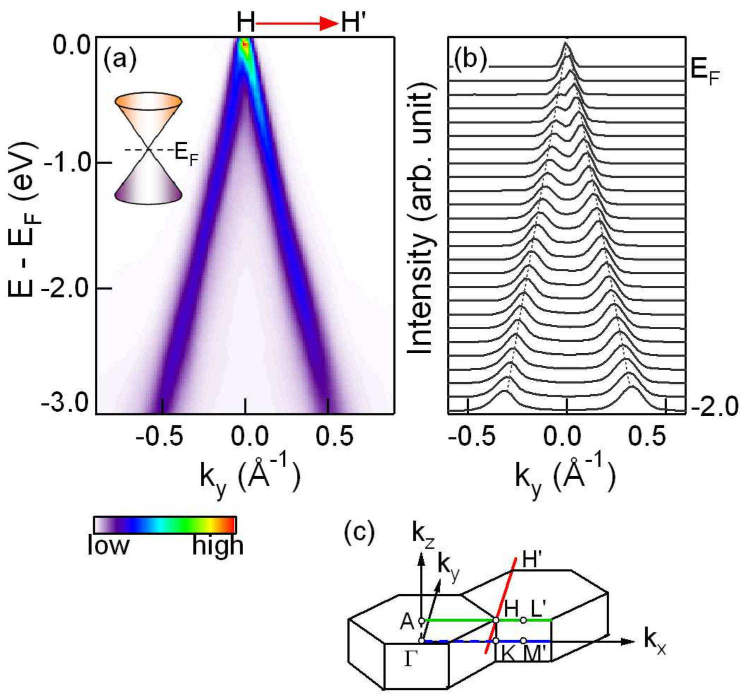
$$E = k^2/2m$$

energy resolution 1meV



momentum distribution curve (MDC)

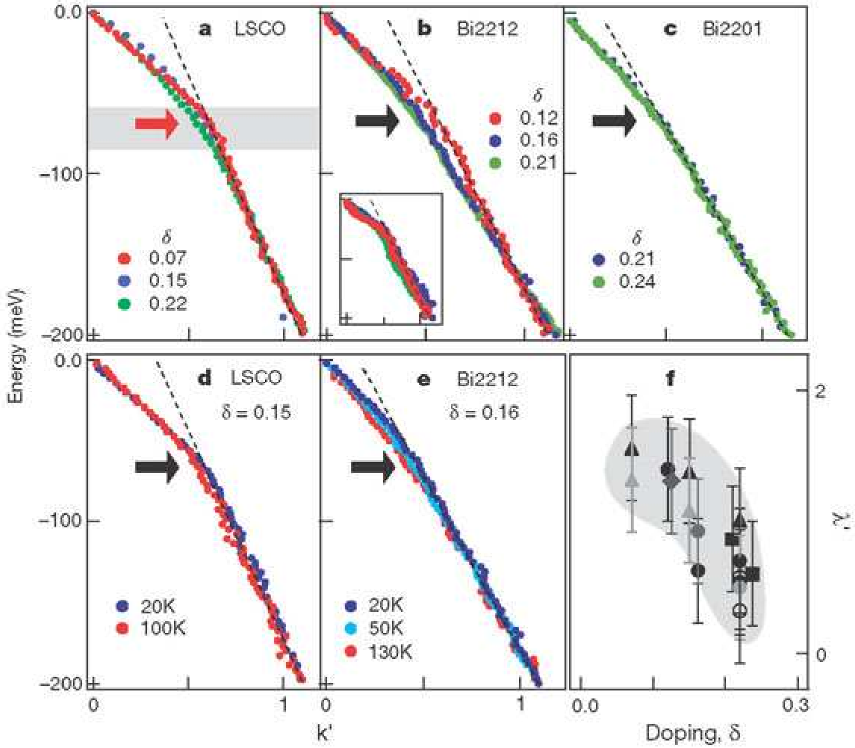
ARPES and graphene



Dirac linear dispersion relation for graphene

cond-mat/0608069

Kinks in HTC

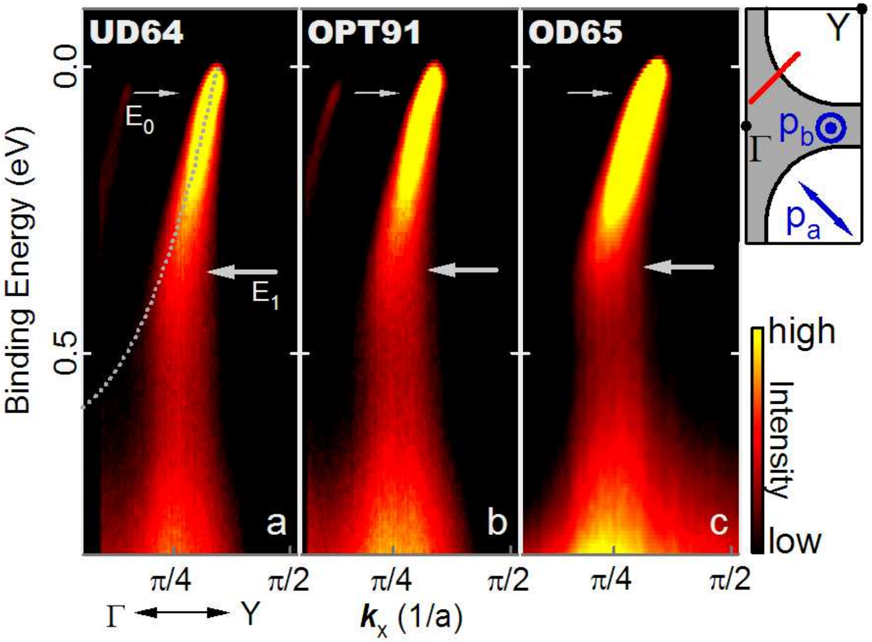


cond-mat/0604284

Kinks at $40 - 70$meV

electron-phonon or electron-spin fluctuations coupling

“Waterfalls” in HTC

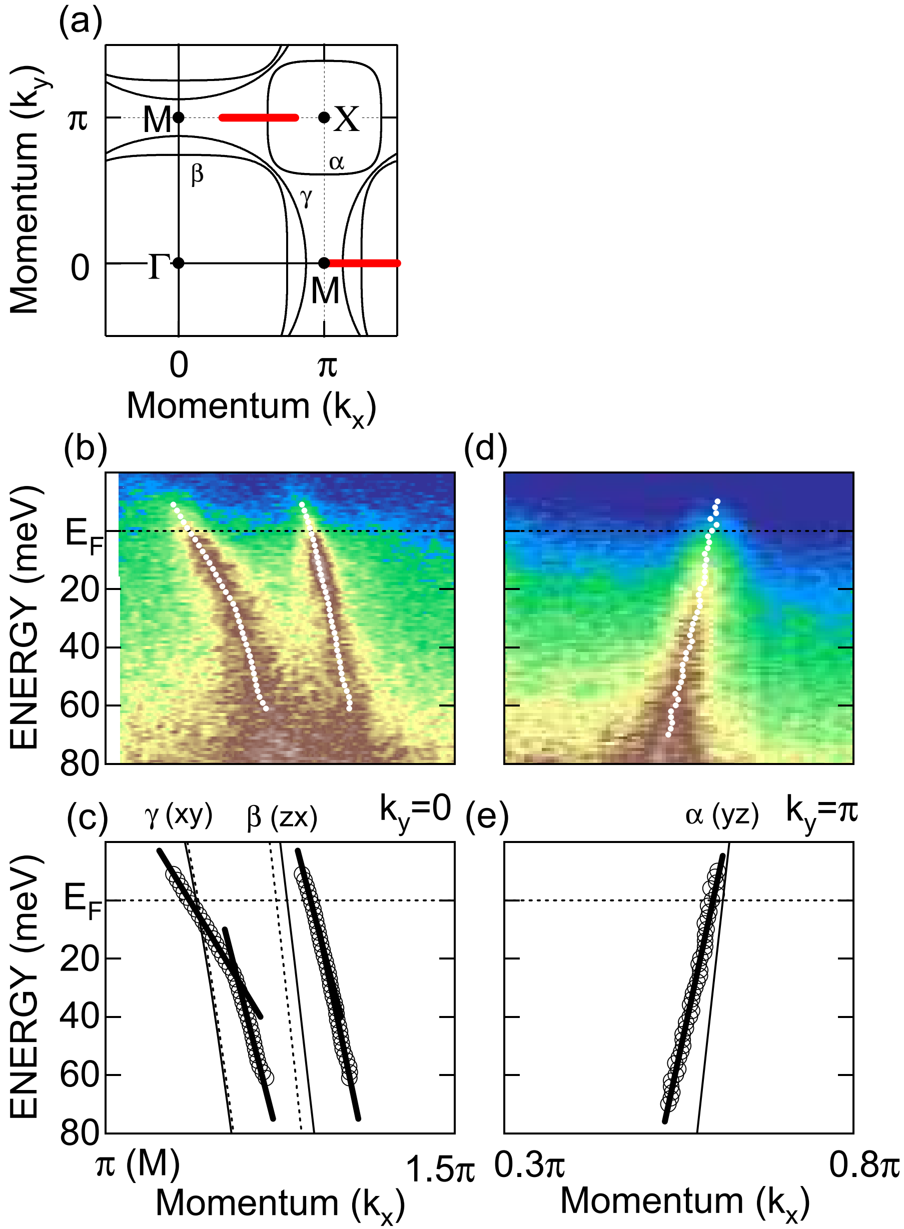


different HTC systems, cond-mat/0607319

Kinks seen experimentally between 300-800 meV

Origin: phonons, spin fluctuations, not known yet

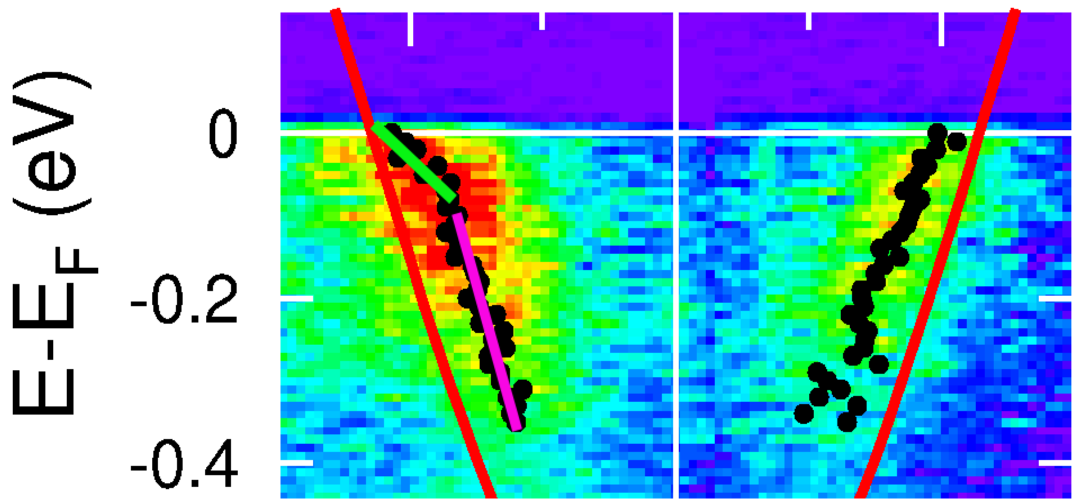
Kinks orbital selective



Kink at 30meV in γ-band only

Sr_2RuO_4, cond-mat/0508312

More examples of kinks in ARPES



$SrVO_3$, cond-mat/0504075

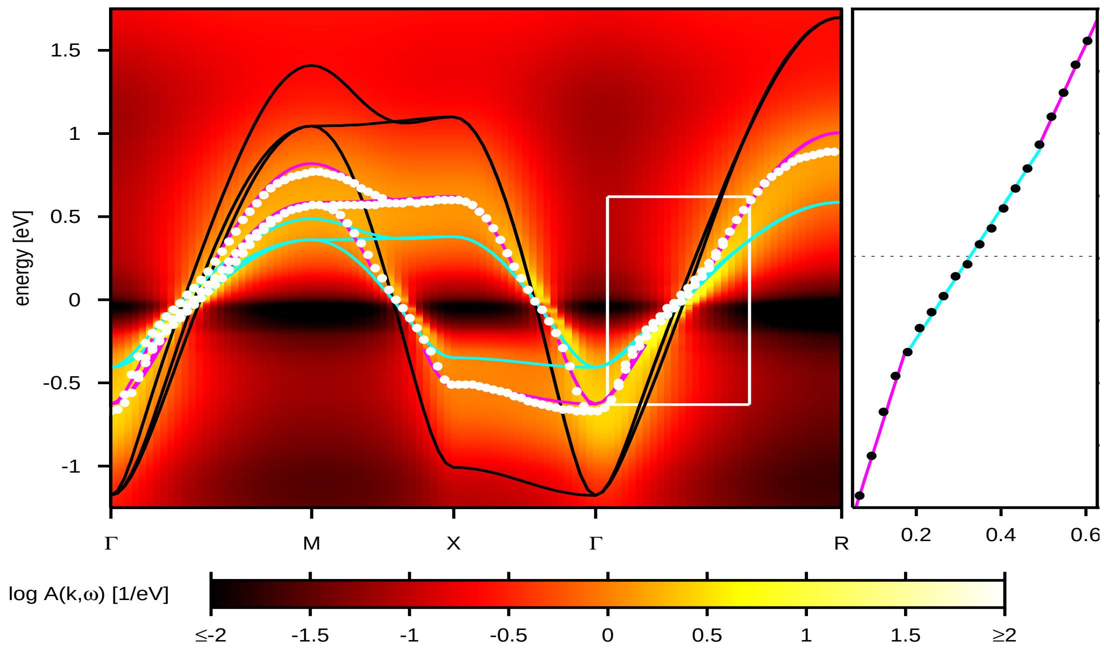
Kinks seen experimentally at 150 meV

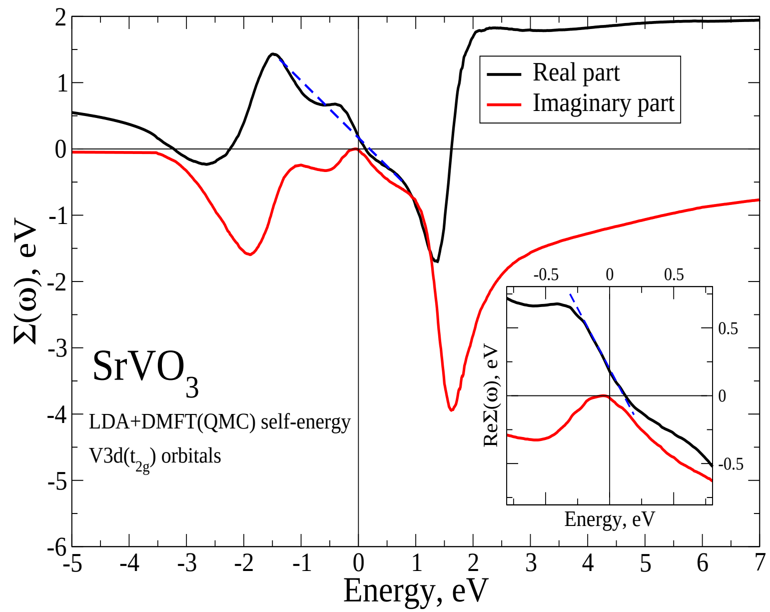
Pure electronic origin?

Kinks in LDA+DMFT study of $SrVO_3$

plain band model with local correlations, no other bosons, ... but kinks!

I.A. Nekrasov *et al.*, cond-mat/0508313, PRB (2006)





$$G_{\mathbf{k}}(\omega) = \frac{1}{\omega+\mu-\epsilon_{\mathbf{k}}-\Sigma(\omega)} \qquad \rightarrow E_{\mathbf{k}}+\mu-\epsilon_{\mathbf{k}}-\mathrm{Re}\Sigma(E_{\mathbf{k}})=0$$

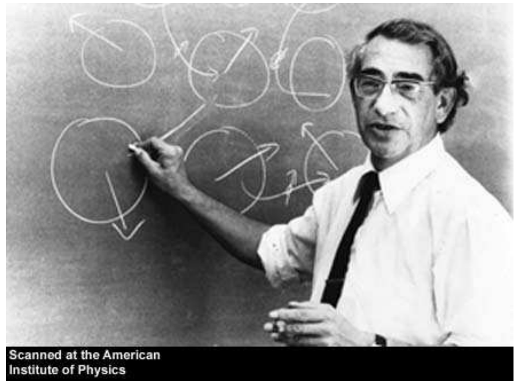
Not found in SIAM with simple hybridization function! → DMFT self-consistency effect

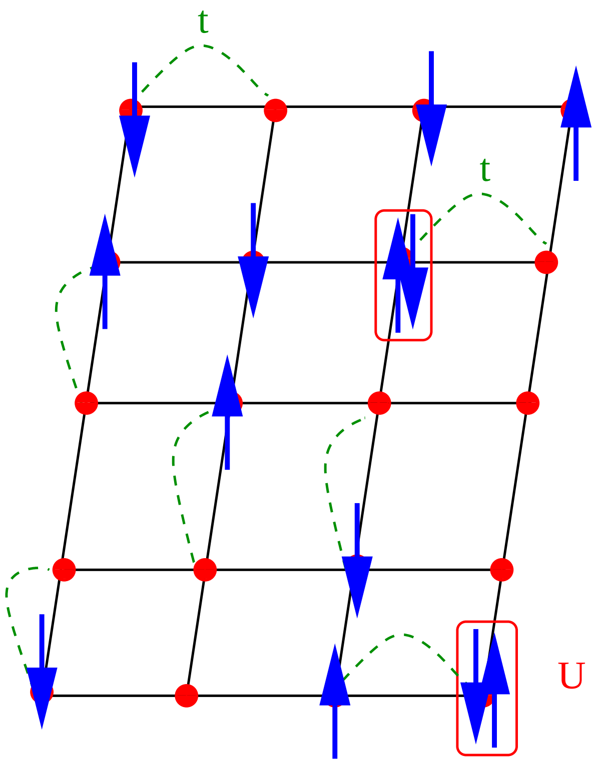
New purely electronic mechanism

- in strongly correlated systems
- characteristic energy scale
- range of validity for Fermi liquid theory

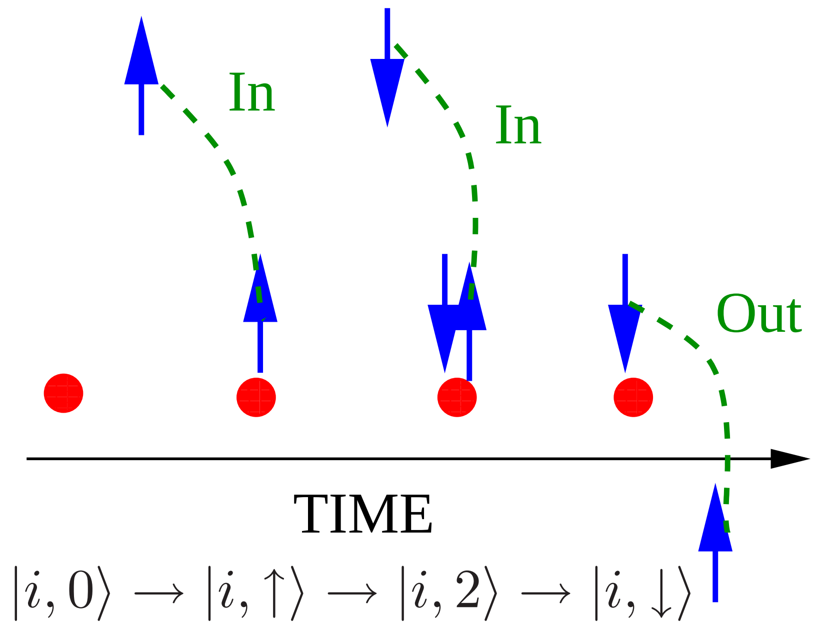
Hubbard model for strongly correlated electrons

$$H = -\sum_{ij\sigma} t_{ij} c^{\dagger}_{i\sigma} c_{j\sigma} + U \sum_{i} n_{i\uparrow} n_{i\downarrow}$$





Local Hubbard physics



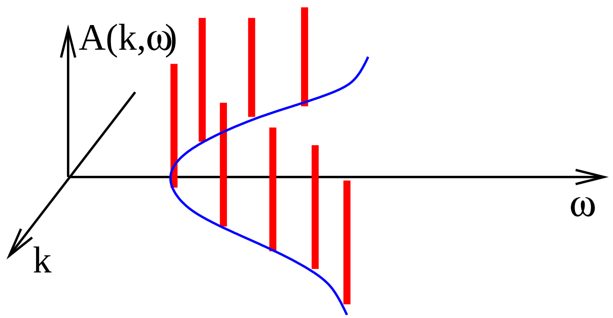
All what we know about Hubbard model

Solved in $U = 0$ limit (non-interacting limit)

$$G_\sigma(\mathbf{k},\omega) = \frac{1}{\omega + \mu - \epsilon_\mathbf{k}}$$

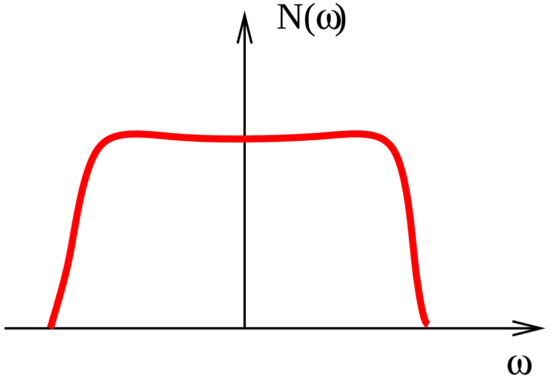
Dispersion relation

$$\epsilon_\mathbf{k} = \sum_{j(i)} t_{ij} e^{i\mathbf{k}(\mathbf{R}_i - \mathbf{R}_j)}$$



Spectral function - one-particle excitations

$$A_\sigma(\mathbf{k},\omega) \equiv -\frac{1}{\pi}\mathrm{Im}G(\mathbf{k},\omega) = \delta(\omega + \mu - \epsilon_\mathbf{k})$$



Density of states (DOS) - thermodynamics

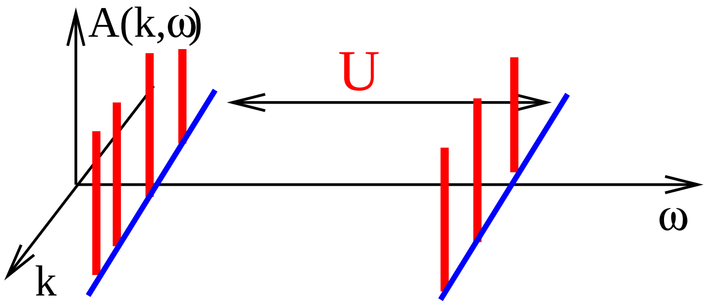
$$N_\sigma(\omega) \equiv \sum_\mathbf{k} A(\mathbf{k},\omega) = \sum_\mathbf{k} \delta(\omega + \mu - \epsilon_\mathbf{k})$$

All what we know about Hubbard model

Solved in $t = 0$ limit (atomic limit)

$$G_\sigma(\mathbf{k},\omega) = \frac{1-n_{-\sigma}}{\omega+\mu} + \frac{n_{-\sigma}}{\omega+\mu-U} = \frac{1}{\omega+\mu-\Sigma_\sigma(\omega)}$$

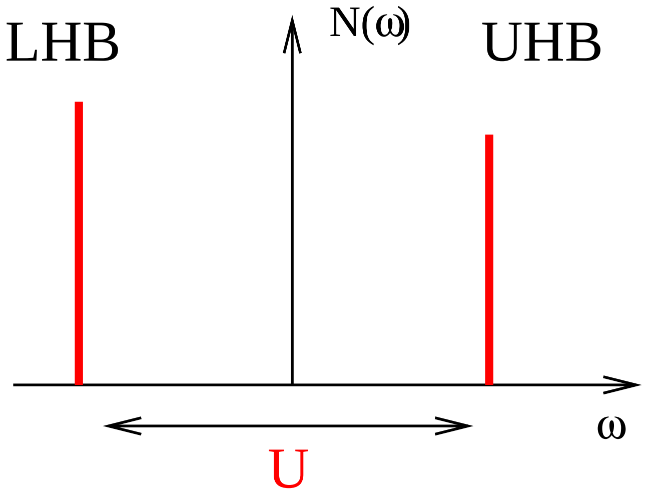
Real self-energy



$$\Sigma_\sigma(\omega) = n_\sigma U + \frac{n_{-\sigma}(1-n_{-\sigma})U^2}{\omega+\mu-(1-n_{-\sigma})U}$$

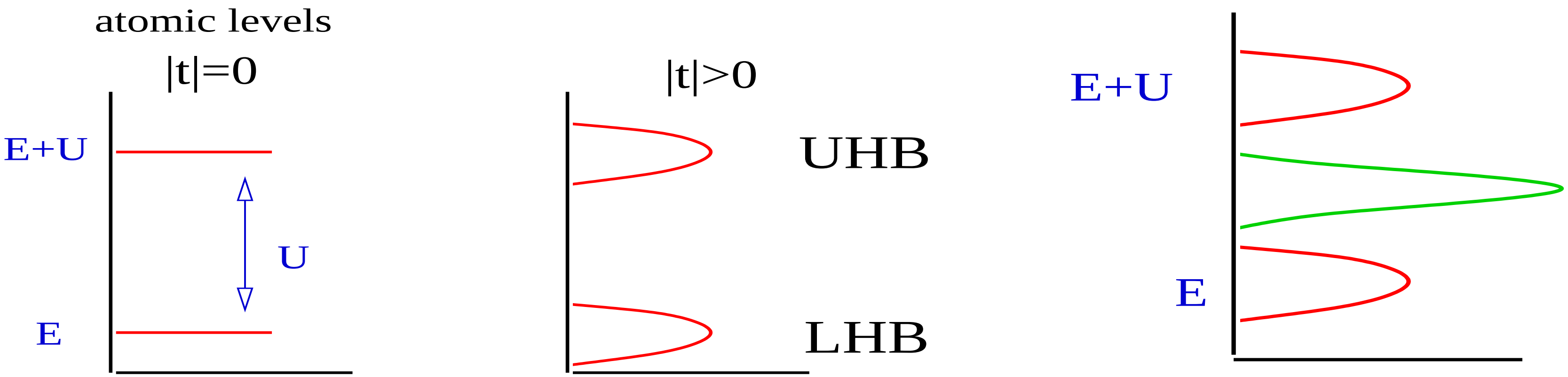
Spectral function

$$A_\sigma(\mathbf{k},\omega) = (1-n_{-\sigma})\delta(\omega+\mu) + n_{-\sigma}\delta(\omega+\mu-U)$$



Green function and self-energy are local,
i.e. $\mathbf{k}$ independent

Hubbard subbands and quasiparticle peak



spin flip on central site

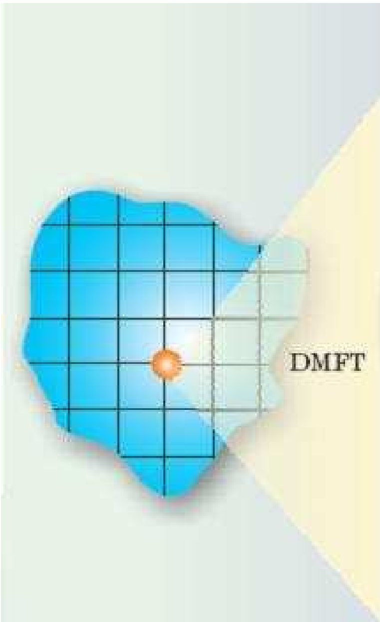
at $U = U_c$ resonance disapears

gaped insulator

dynamical processes with spin-flips inject states into correlation the gap
giving rise to a quasiparticle resonance peak

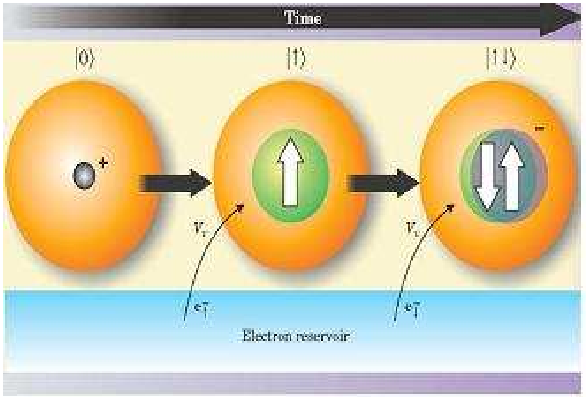
DMFT for lattice fermions

Replace (map) full many-body lattice problem by a single-site coupled to dynamical reservoir and solve such problem self-consistently



Metzner, Vollhardt 89; Georges at al. 96

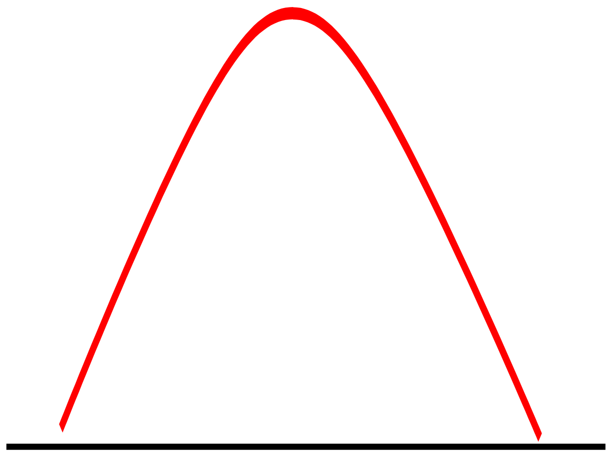
Kotliar, Vollhardt, Physics Today 57 No.3, 53 (2004)

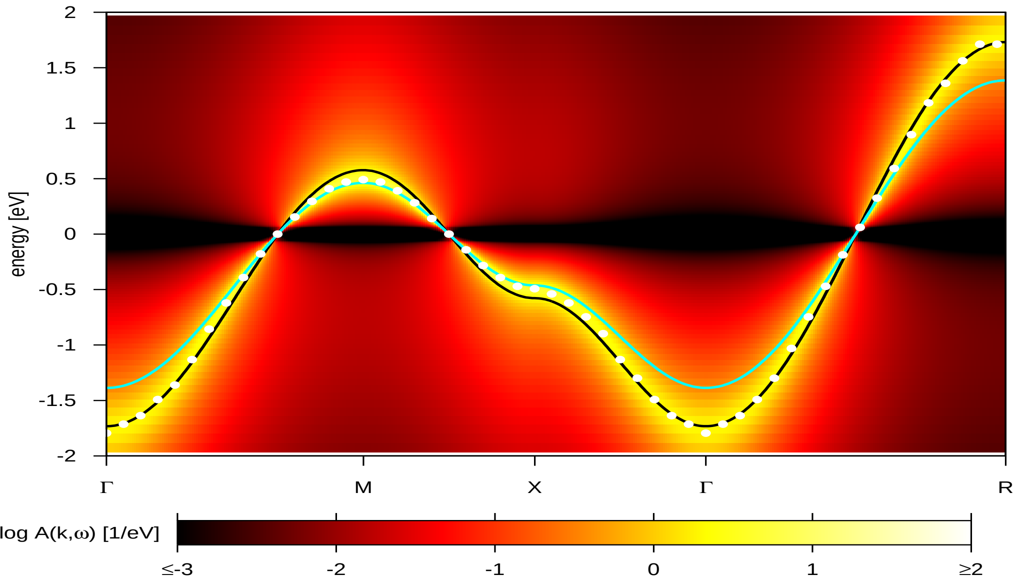


All local dynamical correlations included exactly

Space correlations neglected - mean-field approximation

Weakly correlated system

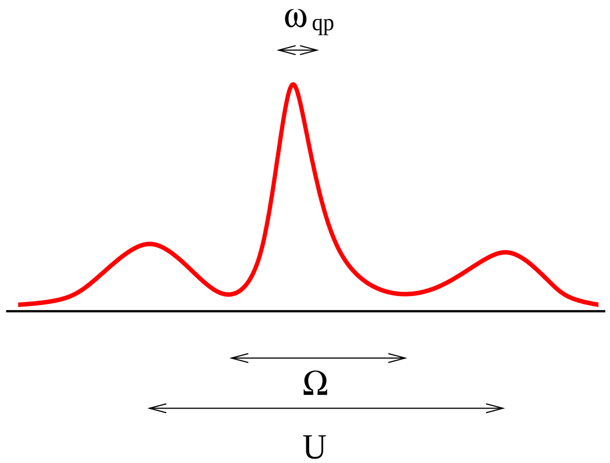


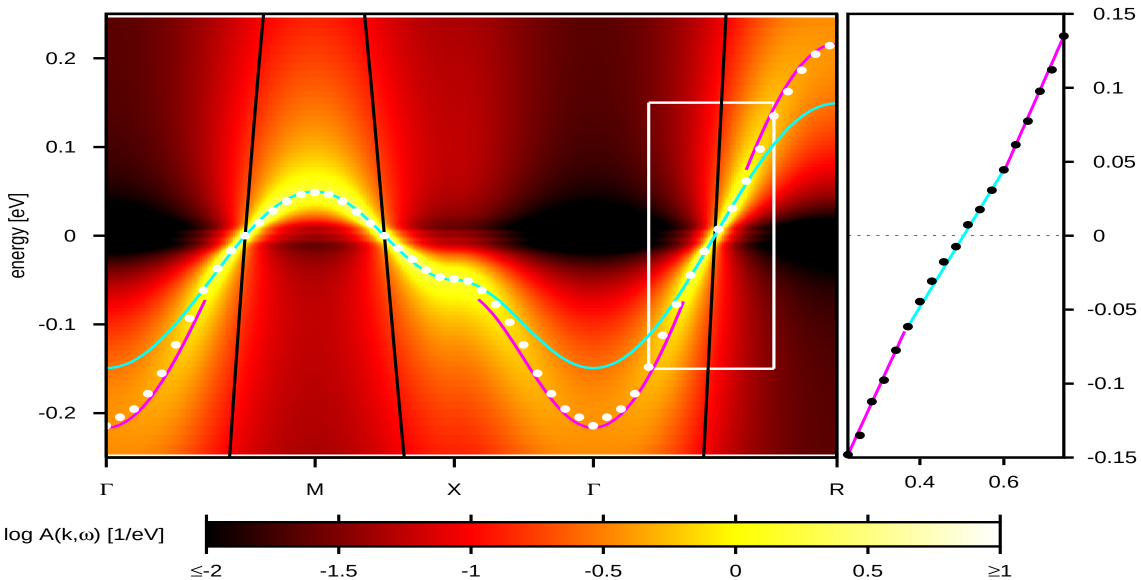


Fermi liquid $Z_{FL} \lesssim 1$: $E_{\mathbf{k}} = Z_{FL}\epsilon_{\mathbf{k}}$ for $|E_{\mathbf{k}}| < \omega_*$

$E_{\mathbf{k}} = \epsilon_{\mathbf{k}}$ for $|E_{\mathbf{k}}| > \omega_*$

Kinks due to strong correlations





Fermi liquid $Z_{FL} \ll 1$: $E_{\mathbf{k}} = Z_{FL}\epsilon_{\mathbf{k}}$ for $|E_{\mathbf{k}}| < \omega_*$

Different renormalization $Z_{CP} \ll 1$: $E_{\mathbf{k}} = Z_{CP}\epsilon_{\mathbf{k}} \pm c$ for $|E_{\mathbf{k}}| > \omega_*$

Mathematical explanation of kinks within DMFT

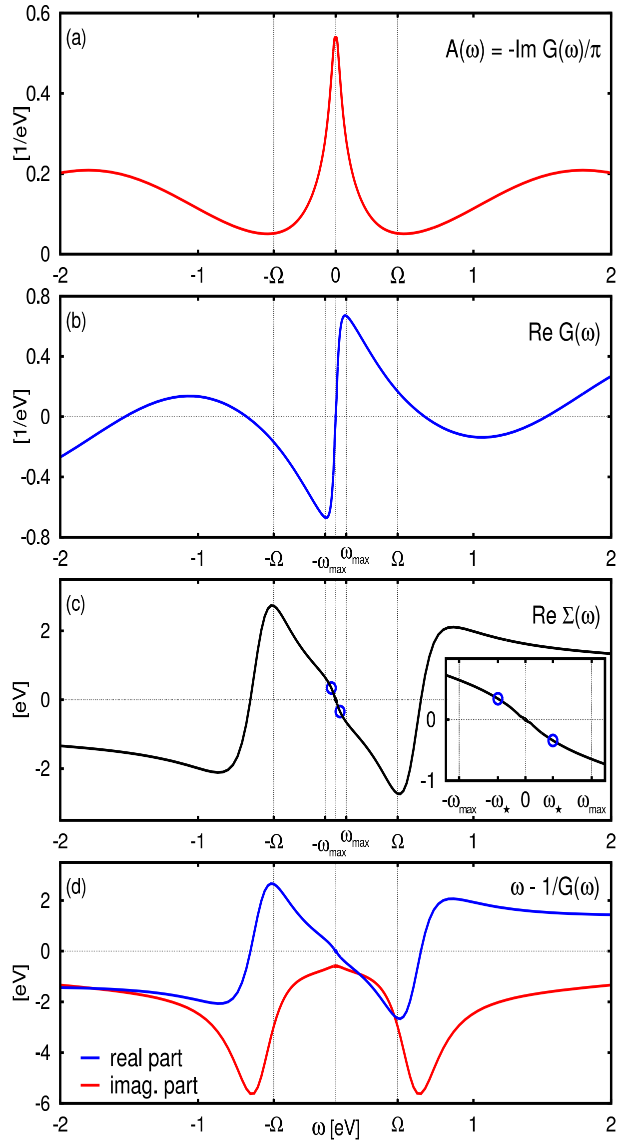
DMFT self-consistency condition

$$\Sigma(\omega) = \omega - 1/G(\omega) - \Delta(G(\omega))$$

$$\Delta(G(\omega)) \approx (m_2 - m_1^2)G(\omega) + \ldots$$

Three-peak structure sufficient condition

Fermi-liquid for $|\omega| < \omega_* \sim Z_{FL}$



Microscopic predictions

Starting from:

- $\epsilon_{\mathbf{k}}$ - bare dispersion relation

$$G_0(\omega) = \sum_{\mathbf{k}} \frac{1}{\omega - \epsilon_{\mathbf{k}}}$$

- Z_{FL}

we predict that:

Microscopic predictions

- Kink position

$$\omega_* = 0.41 Z_{FL} \frac{\mathrm{Im} 1/G_0}{\mathrm{Re} G_0'/G_0^2}$$

- Intermediate energy regime

$$Z_{CP} = Z_{FL} \frac{1}{\mathrm{Re} G_0'/G_0^2}$$

- Change in the slope Z_{FL}/Z_{CP} interaction independent
- Curvature of the kink $\sim Z_{FL}^2$
- Sharpness of the kink $\sim 1/Z_{FL}^2$
- Sharper for stronger U

Outlook: possible origin of the "waterfalls"

"Waterfalls": kinks at $\omega_\star \approx 300\text{-}400$ meV in cuprates

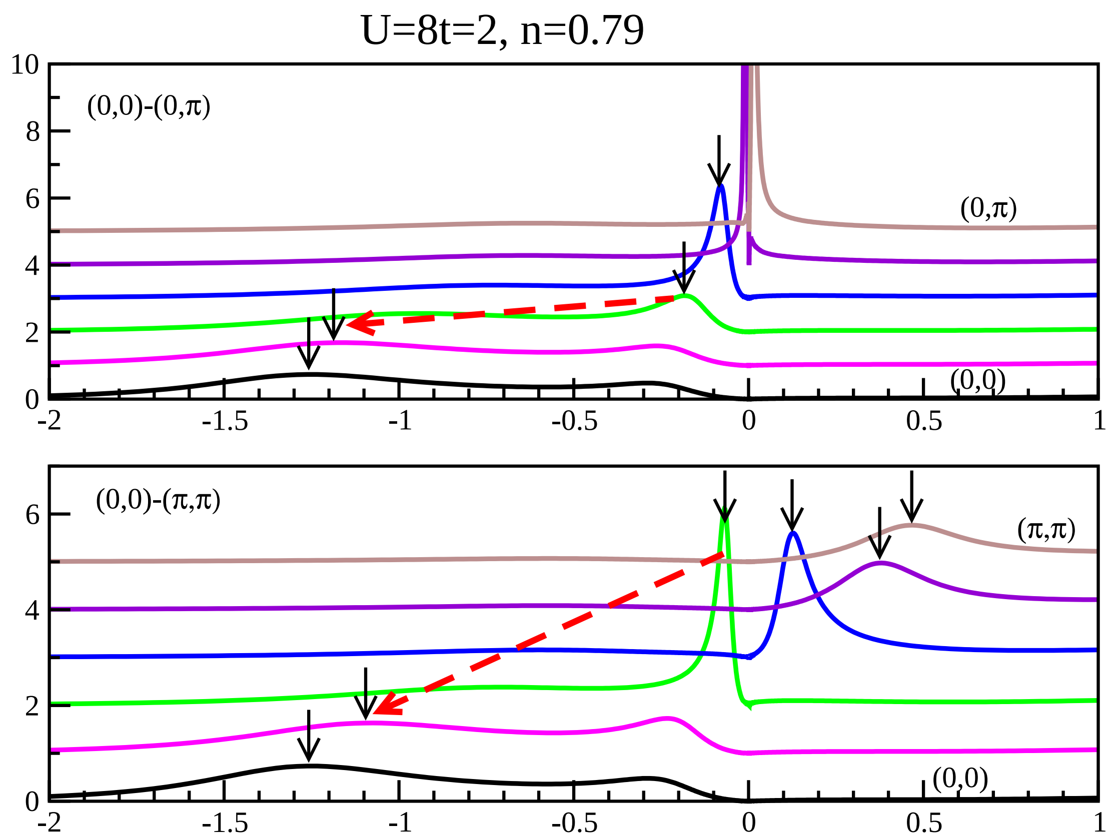
- crossover to Hubbard bands? Wang et al. (2006)
- $U \gg t \Rightarrow$ dispersion goes from central peak to Hubbard band

K. Byczuk, M. Kollar (unpublished)

$$\Sigma(\omega) = \Sigma_0 + \frac{\Sigma_1}{\omega} + O\left(\frac{1}{\omega^2}\right)$$

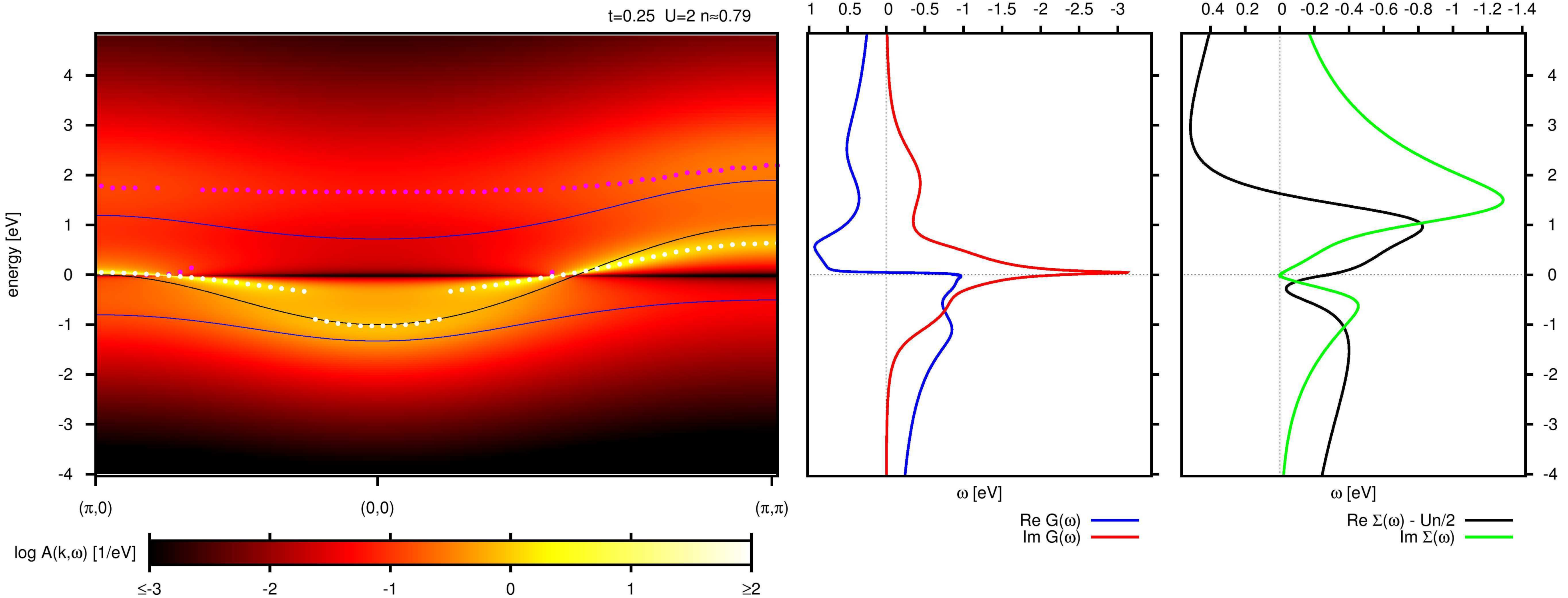
$$\Downarrow$$

$$E_{\boldsymbol{k}}^{\text{UHB,LHB}} \approx \frac{1}{2}\left[\epsilon_{\boldsymbol{k}} \pm \sqrt{\epsilon_{\boldsymbol{k}}^2 + cU^2}\right]$$



Crossover to Hubbard bands

Hubbard model, square lattice, DMFT(NRG), $U = 8t$, $n = 0.79$



- ImΣ decays faster than ReΣ
- for large energies: $E_{\boldsymbol{k}}$ approaches $E_{\boldsymbol{k}}^{\text{UHB,LHB}}$
- waterfalls from central peak to LHB

K. Byczuk, M. Kollar (unpublished)
Y.-F. Yang, K. Held (unpublished)

Conclusions

- Strong correlations (three peak spectral function) a sufficient condition for electronic kinks
- Energy scale for electronic kinks $\omega_* = Z_{FL} D$ determined by Fermi-liquid renormalization and bare (LDA) density of states
- ω_* sets the energy scale for Fermi-liquid regime where $E_{\mathbf{k}} = Z_{FL}\epsilon_{\mathbf{k}}$ for $|E_{\mathbf{k}}| < \omega_*$
- Beyond Fermi-liquid regime the dispersion is still renormalized and useful $E_{\mathbf{k}} = Z_{CP}\epsilon_{\mathbf{k}} \pm c$ for $|E_{\mathbf{k}}| > \omega_*$ where the offset c and Z_{CP} determined by Z_{FL} and D
- Electronic kinks are within cluster extension of DMFT (DCA) $\Sigma_{\mathbf{K}}(\omega) = \omega - \frac{1}{G_{\mathbf{K}}(\omega)} - \Delta(G_{\mathbf{K}}(\omega))$
- Electronic kinks and waterfalls are generic feature of strongly correlated systems