

Group field theory with non-commutative metric variables

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joint work with D. Oriti

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Introduction

Why metric variables for GFT's?

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1. Exact duality spin foam models/simplicial gravity path integrals

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Tracking simplicial data is critical

- ▶ to impose simplicity constraints in a geometrically consistent way
- ▶ to understand relation with simplicial geometry, tackle issue of semi-classical limit

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- ▶ Similar question in LQG: Relation LQG/simplicial geometry
Dittrich-Ryan, Dittrich-Speziale.

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2. Indications of a role played by non-commutative geometry in spin foams and GFTs:

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2. Indications of a role played by **non-commutative geometry** in spin foams and GFTs:
 - Emergence of effective **non-commutative space-time** in 3d gravity
Freidel, Livine '05.
 - Matter dynamics as a phase of GFT: effective theories with deformed Poincaré symmetry *Fairbairn, Girelli, Livine, Oriti.*

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Suggests non commutative geometry lie hidden in GFT formalism

- Similar hint in LQG: metric data encoded in electric flux variables which **do not commute**.

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 - ▶ Dual description of LQG kin. Hilbert space in terms of cylindrical functions on Lie algebras A.B, B. Dittrich, D. Oriti J. Tambornino
Dual variables x_j interpreted as elementary flux variables
(see Johannes' talk)

Outline

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Simplicial representation of 3d GFT

Towards 4d gravity models

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- Higher D. **generalization of matrix models**, so successful for 2d gravity
- Feynman diagrams are D -stranded graphs dual to **simplicial complexes**, Feynman amplitudes are **spin foam models**
 - ▶ GFT for 3D gravity **Boulatov '92**
 - ▶ Generalized to 4D lattice BF theories **Ooguri '92**.
 - ▶ GFT for the Barrett-Crane spin foam model: **De Pietri, Freidel, Krasnov, Rovelli '00**.

- **Universal structure** behind spin foam framework:

Any local spin foam model can be viewed as Feynman graph of a GFT

Rovelli and Reisenberger '00

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- For the spin foam perspective, GFTs tackle issue of **triangulation dependence**
- But they do much more:
 - ▶ Provides a framework to compute quantum gravity amplitudes including the **sum over all topologies**.
 - ▶ This is a **field theory**: lots of tools at our disposal!
Symmetry, renormalization...

Group field theory in a nutshell

Building up space-time

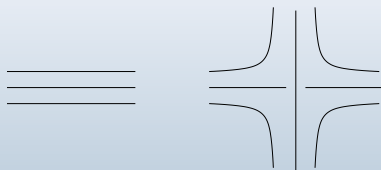
GFT for 3d Riemannian gravity: $D = 3$, $G = \text{SO}(3)$.

- Field $\varphi_{123} := \varphi(g_1, g_2, g_3)$, with invariance $\varphi(hg_i) = \varphi(g_i) \ \forall h \in G$.

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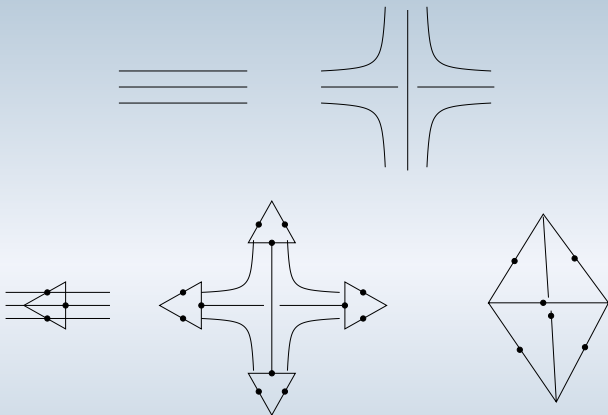
- Field $\varphi_{123} := \varphi(g_1, g_2, g_3)$, with invariance $\varphi(hg_i) = \varphi(g_i) \ \forall h \in G$.
- Dynamics governed by the action: :

$$S = \frac{1}{2} \int [dg]^3 \varphi_{123} \varphi_{123} - \frac{\lambda}{4!} \int [dg]^6 \varphi_{123} \varphi_{345} \varphi_{526} \varphi_{641}$$



Group field theory in a nutshell

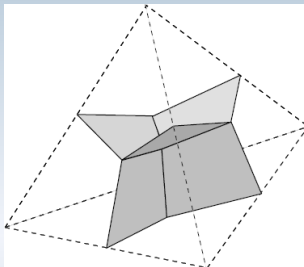
Building up space-time



Feynman diagrams as 2-complexes dual to simplicial complexes
(triangulated spaces)

Group field theory in a nutshell

GFT and lattice gauge theory



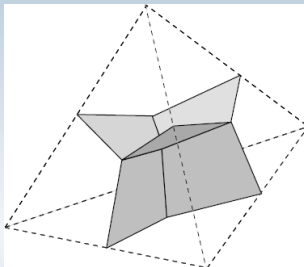
- **bulk** variables $h_{t\tau}$: holonomy from the triangle t to the tetrahedron τ
- **boundary** variables g_{et} : holonomy from the edge e to the triangle t .

GFT Propagator and vertex:

$$\int dh_t \prod_{e \in t} \delta(g_{et} h_t \tilde{g}_{et}^{-1}), \quad \int \prod_t dh_{t\tau} \prod_{\langle et \rangle} \delta(g_{et} h_{t\tau} \tilde{g}_{et}^{-1})$$

Group field theory in a nutshell

GFT and lattice gauge theory



- Feynman amplitude: integral over **discrete flat connections**

$$I(\Gamma) = \int \prod_t dh_t \prod_e \delta\left(\prod_{t \supset e} h_t\right)$$

Group field theory in a nutshell

From GFT to spin foam models

- **Spin representation** of GFT using harmonic analysis on the gauge group

$$\varphi_{123} = \sum_{j_1, j_2, j_3} \phi_{m_1, m_2, m_3}^{j_1, j_2, j_3} D_{m_1 n_1}^{j_1}(g_1) D_{m_2 n_2}^{j_2}(g_2) D_{m_3 n_3}^{j_3}(g_3) C_{n_1, n_2, n_3}^{j_1, j_2, j_3}$$

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- Field pictured as a 3-valent **spin network vertex**, interpreted as a quantized triangle.

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- Feynman amplitudes

$$I(\Gamma) = \sum_{\{j_t\}} \prod_t (2j_t + 1) \prod_{\tau} \left\{ \begin{array}{ccc} j_1 & j_2 & j_3 \\ j_4 & j_5 & j_6 \end{array} \right\}$$

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- Boundary observables described in terms of spin networks

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From GFT to simplicial path integrals?

- Lattice gauge theory picture:

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- Simplicial path integral picture :

$$I(\Gamma) = \int \prod_t dh_t \prod_e dX_e \prod_e e^{i\text{Tr} X_e H_e(h)}, \quad H(e) = \delta\left(\prod_{t \supset e} h_t\right)$$

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Simplicial representation of GFT?

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Simplicial representation of 3d GFT

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 \vec{p}_g coordinates on the group manifold
- Define **Fourier transform** $\hat{f}(x) = \int dg f(g) e_g(x)$
- Algebra structure on $\widehat{\text{Im}}$: $e_{g_1} \star e_{g_2} = e_{g_1 g_2}$
inherited from the convolution product on the group

Simplicial representation of 3d GFT

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- For $G = \text{SU}(2)$, we choose $\vec{p}_g = \text{Tr}|g|\vec{\tau}$, $|g| := \text{sign}(\text{Tr}g)g$

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- Fourier transform **invertible** on functions $f(g) = f(-g)$ of $\text{SO}(3) \sim \text{SU}(2)/\mathbb{Z}_2$:

$$f(g) = \int d^3x (\widehat{f} \star e_{g^{-1}})(x)$$

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- With more work full Fourier transform on $\mathrm{SU}(2)$

[Freidel Majid 05; Joung, Mourad, Noui 08]

Simplicial representation of 3d GFT

Going to metric variables: Fourier transform

$$\widehat{f}(x) = \int dg f(g) \mathbf{e}_g(x)$$

What functions of \mathbb{R}^3 does the Fourier transform hit?

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★-product \leftrightarrow deformed addition of momenta preserving the bound.

Simplicial representation of 3d GFT

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★-product \leftrightarrow deformed addition of momenta preserving the bound.
- \mathbb{R}^3 seen with a **finite resolution**: Functions that can be **sampled by discrete values** f_{mn}^j without loss of information:

$$\widehat{f}(x) = \sum_{j,m,n} f_{mn}^j \widehat{D}_{mn}^j(x)$$

Simplicial representation of 3d GFT

Dual field, action, Feynman rules

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Dual field, action, Feynman rules

- Fourier transform of the Boulatov field φ_{123} :

$$\widehat{\varphi}_{123} := \widehat{\varphi}(x_1, x_2, x_3) = \int [dg]^3 \varphi_{123} \mathbf{e}_{g_1}(x_1) \mathbf{e}_{g_2}(x_2) \mathbf{e}_{g_3}(x_3)$$

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- Gauge invariance dual to a **closure constraint**:

$$\widehat{P}\widehat{\varphi} = \widehat{C} \star \widehat{\varphi}, \quad \widehat{C}(x_1, x_2, x_3) = \delta_0(x_1 + x_2 + x_3)$$

where $\delta_0(x) := \int dg e_g(x)$ plays the role of a Dirac distribution:

$$\int d^3x (\delta_0 \star f)(x) = f(0)$$

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- ▶ Dual field as a **(non-commutative) triangle**

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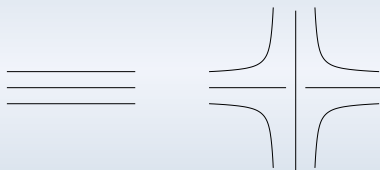
- ▶ Dual field as a **(non-commutative) triangle**
- ▶ Field variables as **metric variables** associated to the edges.

Simplicial representation of 3d GFT

Dual field, action, Feynman rules

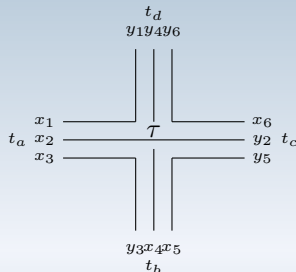
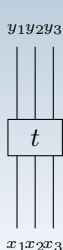
Action

$$S = \frac{1}{2} \int [dx]^3 \hat{\varphi}_{123} \star \hat{\varphi}_{321} - \frac{\lambda}{4!} \int [dx]^6 \hat{\varphi}_{123} \star \hat{\varphi}_{345} \star \hat{\varphi}_{526} \star \hat{\varphi}_{641}$$



Simplicial representation of 3d GFT

Dual field, action, Feynman rules



$$\int dh_t \prod_{i=1}^3 (\delta_{-x_i} \star e_{h_t})(y_i), \quad \int \prod_t dh_{t\tau} \prod_{i=1}^6 (\delta_{-x_i} \star e_{h_{t\tau}})(y_i)$$

- h_t : parallel transport through the triangle t .
- $h_{t\tau}$: parallel transport from the tetrahedron τ to triangle t .
- $h_{tt'} := h_{t\tau} h_{\tau t'}$

Simplicial representation of 3d GFT

Feynman amplitudes

- Join strands using the \star -product, keeping track of ordering.
- Each loop of strands bound a **face** of the 2-complex, dual to an **edge** of the triangulation.
- Under integration over holonomies h , product of **face amplitudes** $A_f[h]$.

Simplicial representation of 3d GFT

Feynman amplitudes

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Feynman amplitudes

- Let $\{\tau_j\}_{0 \leq j \leq N}$ ordered sequence of tetrahedra around the edge.

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$A_f[h]$ is the **cyclic \star -product**:

$$A_f[h] = \int \prod_{j=0}^N dx_j \star_{j=0}^{\vec{N}+1} (\delta_{x_j} \star e_{h_{jj+1}})(x_{j+1})$$

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Simplicial representation of 3d GFT

Feynman amplitudes

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Simplicial representation of 3d GFT

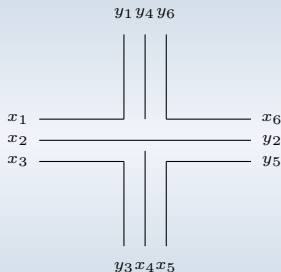
Feynman amplitudes

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Simplicial representation of 3d GFT

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Towards 4d gravity models

Going up dimensions: Ooguri model

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$$S = \frac{1}{2} \int \varphi_{1234}^2 - \frac{\lambda}{5!} \int \varphi_{1234} \varphi_{4567} \varphi_{7389} \varphi_{96210} \varphi_{10851}.$$

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Towards 4d gravity models

Imposing simplicity: non-commutative tetrahedron

- $S_{\text{grav}}[e, A] = \int \text{tr } B \wedge F(A), \quad B = \star e \wedge e$
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$$S_k \varphi(g_j) = \prod_{j=1}^4 \int_{\text{SO}(3)} du_j \varphi(k^{-1} u_j k g_j^-, u_j g_j^+)$$

Projector onto fields on $\text{SO}(4)/\text{SO}(3)_k$. $k = 1$: **Barrett-Crane projector**

Towards 4d gravity models

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- Combining the simplicity projector $\widehat{S} := \widehat{S}_1$ with closure: $\widehat{S} \star \widehat{C} \star \widehat{\varphi}$
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imposes simplicity of bivectors x_t **in each of the frames** associated to the simplices j around t .

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Beyond Barrett-Crane

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Beyond Barrett-Crane: imposing simplicity covariantly

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- ▶ Study of the model in progress...

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Explicit GFT duality spin foams/ simplicial path integrals

- Path integral representation of the BC model as a non-commutative B observable in BF
- New proposal for GFT model: study under way.

Conclusion

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 - ▶ deepen links between GFT formalism and non-commutative geometry