

Asymptotics of Spin Foams: What we learned

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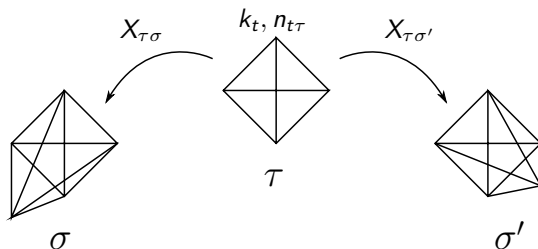
The Statesum

The statesum of EPRL, FK is given in terms of a triangulation of a manifold and can be written in terms of variables living on triangles t , tetrahedra τ and 4-simplices σ :

$$k_t \in \mathbb{N}/2 = \text{Irrep}(\text{SU}(2))$$

$$n_{t\tau} \in S^2$$

$$X_{\tau\sigma} \in \text{SL}(2, \mathbb{C}) \text{ or } \text{Spin}(4)$$



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It is then given by:

$$Z(\mathcal{T}) = \sum_{k_t} \int_{n_{t\tau}} \int_{X_{\tau\sigma}} \prod_t f_2(k_t, n_{t\tau}) \prod_{\tau} f_3(k_t, n_{t\tau}) \prod_{\sigma} f_4(k_t, n_{t\tau}, X_{\tau\sigma})$$

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Heuristics

$X_{\tau\sigma}$ integration is local per 4-simplex. Heuristically: most spins in the sum are large, so approximately evaluate this integral using stationary phase ($S_4 = \ln f_4$ is linear in k):

$$\int_{X_{\tau\sigma}} f_4(\lambda k_t, n_{t\tau}, X_{\tau\sigma}) \sim \sum_{\delta S_4=0, \text{Re}(S_4)=0} \frac{e^{S_4(\lambda k_t, n_{t\tau}, X_{\tau\sigma})}}{\sqrt{\det H}}$$

(same spirit as in Conrady, Freidel '08, works very well in the PR model as seen in Dowdall, Gomes, FH '09)

Geometry

$\delta S_4 = 0$ and $\text{Re}(S_4) = 0$ (and $n_{t\tau}$ 3d non-degeneracy) imply:

- ▶ $b_{t\tau}$ from $(0, b_{t\tau}) = k_t X_{\tau\sigma}(0, n_{t\tau})$ is a constant $\mathfrak{su}(2)$ valued 2-form on σ (aka vector geometries), or
- ▶ $B_{t\tau}(\sigma_g) = \mu_\sigma \star k_t X_{\tau\sigma}(1, 0, 0, 0) \wedge (0, n_{t\tau})$ are the bivectors of a geometric 4-simplex (σ_g) ,

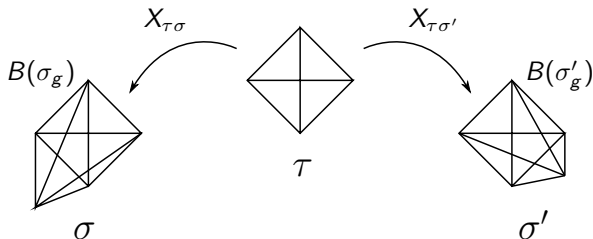
$$S^4 = \mu_\sigma S_{\text{Regge}}(\sigma_g).$$

The $\mathfrak{su}(2)$ sector is much larger than the geometric sector (dominant scaling contribution in BC (Barrett, Steele 02)).

The geometric sector can be distinguished in various ways: 4d non-degeneracy (Conrady, Freidel) gluing constraints (Dittrich, Ryan, Speziale), Regge likeness (Barrett, Dowdall, Fairbairn, Gomes, FH, Pereira)

Geometry gluing

The geometry of the 4-simplex σ_g agrees with the boundary geometry induced on the tetrahedra τ by $k_t n_{t\tau}$. Thus we get a consistent geometry throughout the manifold. X provide the connection between the frames in different 4-simplices.

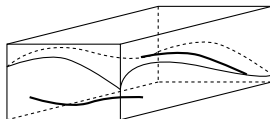


Two questions:

- ▶ Do the $\mathfrak{su}(2)$ solutions spoil the geometric sector?
- ▶ What is the meaning of μ ?

$su(2)$ where are you?

The remaining k, n integrations should peak us on solutions to the variation of the action S_4 with respect to the solution space.



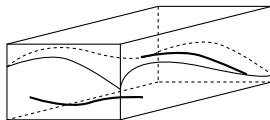
Do the $su(2)$ solutions induce spurious equation of motion? That is, is the geometric space a subspace of the $su(2)$ space?

Euclidean:

$(1 \pm \star)B_{t\tau}(\sigma_g)$ are constant $su(2)$ valued 2-form on σ . If we add an arbitrary constant $su(2)$ solution $b_{t\tau}$ to it we obtain good boundary data k, n but are out of the geometric sector (15 possible perturbations out of which only 10 are geometric) \Rightarrow possibly spurious equations of motion for the geometric sector.

$su(2)$ where are you?

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Do the $su(2)$ solutions induce spurious equation of motion? That is, is the geometric space a subspace of the $su(2)$ space?

Lorentzian:

The solutions require hyperbolic space to glue up, no way to perturb them into the $su(2)$ solutions. \Rightarrow Regge equations of motion.

The meaning of μ

Bivectors reconstruct geometry and orientation. This orientation does not need to be constant throughout the manifold μ measures the orientation relative to a fiducial reference one, e.g. standard orientation in \mathbb{R}^4 . Using $N_\tau = X_{\tau\sigma}(1, 0, 0, 0)$.

$$B_{t\tau}(\sigma_g) = \star k_t \frac{N_\tau \wedge N_{\tau'}}{|N_\tau \wedge N_{\tau'}|} = \mu_\sigma \star k_t X_{\tau\sigma}(1, 0, 0, 0) \wedge (0, n_{t\tau})$$

Then:

$$\frac{N_\tau \wedge N_{\tau'}}{|N_\tau \wedge N_{\tau'}|} = \mu_\sigma N_\tau \wedge X_{\tau\sigma}(0, n_{t\tau})$$

and thus $\mu X_{\tau\sigma}(0, n_{t\tau}) \cdot N_{\tau'} > 0$ and $\mu X_{\tau\sigma}(0, n_{t\tau})$ are the outward pointing triangle normals.

The meaning of μ (2)

Bivectors reconstruct geometry and orientation. This orientation does not need to be constant throughout the manifold μ measures the orientation relative to a fiducial reference one, e.g. standard orientation in \mathbb{R}^4 . So what?

The action on the manifold is

$$\sum_{\sigma} \mu_{\sigma} S_{\text{Regge}}(\sigma_g).$$

These actions have very different solutions for different set of signs μ . In Ponzano Regge we get non-flat manifolds and divergences.

We need to interpret these solutions (anti space?) or make them go away (modify boundary state space?).

Conclusions

Caveats

- ▶ $n_{t\tau}$ required to be 3d non-degenerate.
- ▶ Uncertain status of inserting the asymptotics into the full statesum.
- ▶ No knowledge of the phase of the Hessian.
- ▶ Unclear if Regge likeness propagates from the boundary (maybe unlikely).

but if....

- ▶ Full geometric understanding of the non-metric solutions now available.
- ▶ Larger triangulations look vastly more feasible than in BC.
- ▶ Possibility of spurious EOMs in the Euclidean case, Lorentzian model looks better, maybe "good enough".
- ▶ Orientations problematic. Possibly indicates problem with the boundary state space (half the holomorphic states?).