

# Relativity of superluminal observers in $1 + 3$ spacetime

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Student's problem: find all linear transformations preserving the speed of light.

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$$c^2 dt^2 - dx^2 = -c^2 dt'^2 + dx'^2$$

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# Quantum principle of relativity

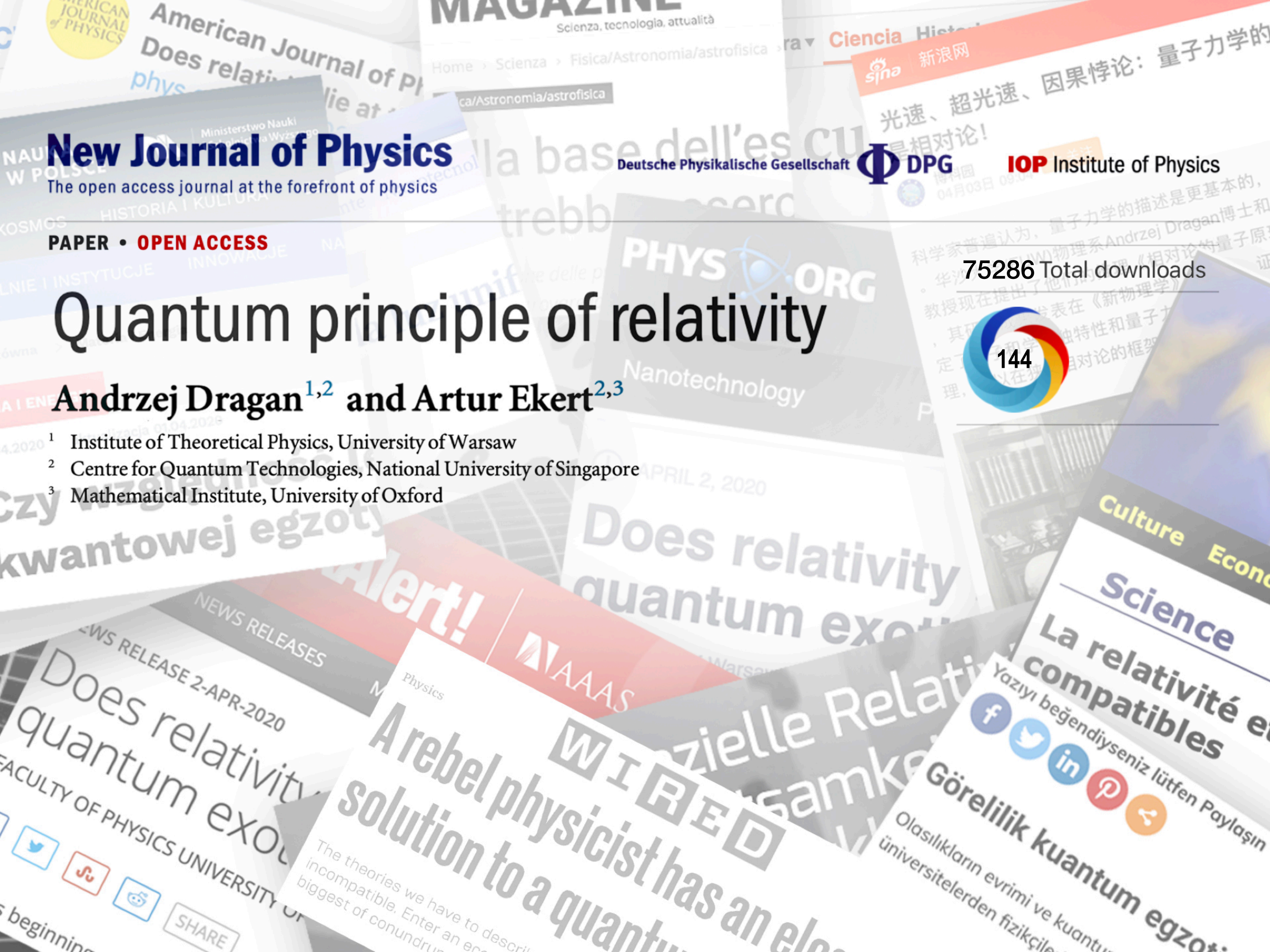
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# Special relativity in 1+3 dimensional spacetime

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$$\boldsymbol{r}' - \frac{\boldsymbol{r}' \cdot \boldsymbol{V}}{V^2} \boldsymbol{V} = \boldsymbol{r} - \frac{\boldsymbol{r} \cdot \boldsymbol{V}}{V^2} \boldsymbol{V}$$

# Special relativity in 1+3 dimensional spacetime

$$\mathbf{r}' - \frac{\mathbf{r}' \cdot \mathbf{V}}{V^2} \mathbf{V} = \mathbf{r} - \frac{\mathbf{r} \cdot \mathbf{V}}{V^2} \mathbf{V}$$

$$\frac{\mathbf{r}' \cdot \mathbf{V}}{V^2} \mathbf{V} = \frac{\frac{\mathbf{r} \cdot \mathbf{V}}{V^2} - t}{\sqrt{1 - \frac{V^2}{c^2}}} \mathbf{V}$$

# Special relativity in 1+3 dimensional spacetime

$$\left\{ \begin{array}{l} \mathbf{r}' = \mathbf{r} - \frac{\mathbf{r} \cdot \mathbf{V}}{V^2} \mathbf{V} + \frac{\frac{\mathbf{r} \cdot \mathbf{V}}{V^2} - t}{\sqrt{1 - \frac{V^2}{c^2}}} \mathbf{V}, \\ ct' = \frac{ct - \frac{\mathbf{r} \cdot \mathbf{V}}{c}}{\sqrt{1 - \frac{V^2}{c^2}}}. \end{array} \right.$$

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$$c^2 dt^2 - d\mathbf{r} \cdot d\mathbf{r} = c^2 dt'^2 - d\mathbf{r}' \cdot d\mathbf{r}'.$$

# Superluminal branch of solutions



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$$c^2 dt^2 - d\mathbf{r} \cdot d\mathbf{r} = -c^2 dt' \cdot dt' + dr'^2$$

# Superluminal branch of solutions

$$\left\{ \begin{array}{l} r' = \frac{Vt - \frac{\mathbf{r} \cdot \mathbf{V}}{V}}{\sqrt{\frac{V^2}{c^2} - 1}}, \\ ct' = \mathbf{r} - \frac{\mathbf{r} \cdot \mathbf{V}}{V^2} \mathbf{V} + \frac{\frac{\mathbf{r} \cdot \mathbf{V}}{Vc} - \frac{ct}{V}}{\sqrt{\frac{V^2}{c^2} - 1}} \mathbf{V} \end{array} \right.$$

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$$r' \frac{\mathbf{V}}{V^2} = -t'$$

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$$\left\{ \begin{array}{l} r' = \frac{Vt - \frac{r \cdot V}{V}}{\sqrt{\frac{V^2}{c^2} - 1}}, \\ ct' = r - \frac{r \cdot V}{V^2} V + \frac{\frac{r \cdot V}{Vc} - \frac{ct}{V}}{\sqrt{\frac{V^2}{c^2} - 1}} V \end{array} \right.$$

$$\begin{aligned} r' &= ct, \\ ct' &= r, \end{aligned}$$

$$r' \frac{V}{V^2} = -t'$$

$$v' = \frac{dr'}{dt'} \frac{dt'}{dt'}$$

# Superluminal branch of solutions

$$\boldsymbol{v}' = \frac{\sqrt{1 - \frac{c^2}{V^2}} \left( \boldsymbol{v} - \frac{\boldsymbol{v} \cdot \boldsymbol{V}}{V^2} \boldsymbol{V} \right) - \left( \frac{c^2}{V^2} \boldsymbol{V} - \frac{\boldsymbol{v} \cdot \boldsymbol{V}}{V^2} \boldsymbol{V} \right)}{1 - \frac{\boldsymbol{v} \cdot \boldsymbol{V}}{V^2}} \left( 1 - \frac{\left( 1 - \frac{c^2}{V^2} \right) \left( 1 - \frac{v^2}{c^2} \right)}{\left( 1 - \frac{\boldsymbol{v} \cdot \boldsymbol{V}}{V^2} \right)^2} \right)^{-1}$$



# Superluminal branch of solutions

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$$\left( 1 - \frac{c^2}{v'^2} \right) = \frac{\left( 1 - \frac{c^2}{V^2} \right) \left( 1 - \frac{v^2}{c^2} \right)}{\left( 1 - \frac{\boldsymbol{v} \cdot \boldsymbol{V}}{V^2} \right)^2}$$

# Energy and momentum of superluminal particles

$$E \equiv \frac{\sigma m c^2}{\sqrt{\frac{v^2}{c^2} - 1}}$$

$$\sigma = \pm 1$$

$$p \equiv \frac{\sigma m v}{\sqrt{\frac{v^2}{c^2} - 1}}$$

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$$\sigma' = \sigma \operatorname{sgn} \left( 1 - \frac{\mathbf{v} \cdot \mathbf{V}}{c^2} \right)$$

„Anti-particles”  
are relative

# Covariant dynamics

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$$\begin{array}{ccc} S \equiv \int L \, dt & & S' \equiv \int L' \, d^3t' \\ \psi & \psi \rightarrow \psi' & S = S' \end{array}$$

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$$\psi(t, \boldsymbol{r})$$

## Example: Maxwell's theory

$$\nabla_{t'} \cdot \mathbf{E}' = -\frac{1}{\varepsilon_0 c} j'$$

$$\nabla_{t'} \cdot \mathbf{B}' = 0$$

$$\nabla_{t'} \times \mathbf{E}' = -\partial_{r'} \mathbf{B}'$$

$$\nabla_{t'} \times \mathbf{B}' = -\mu_0 c \boldsymbol{\varrho}' + \frac{1}{c} \partial_{r'} \mathbf{E}'$$

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$$\partial_{r'} \mathbf{E}' = \frac{\rho'}{\varepsilon_0}$$

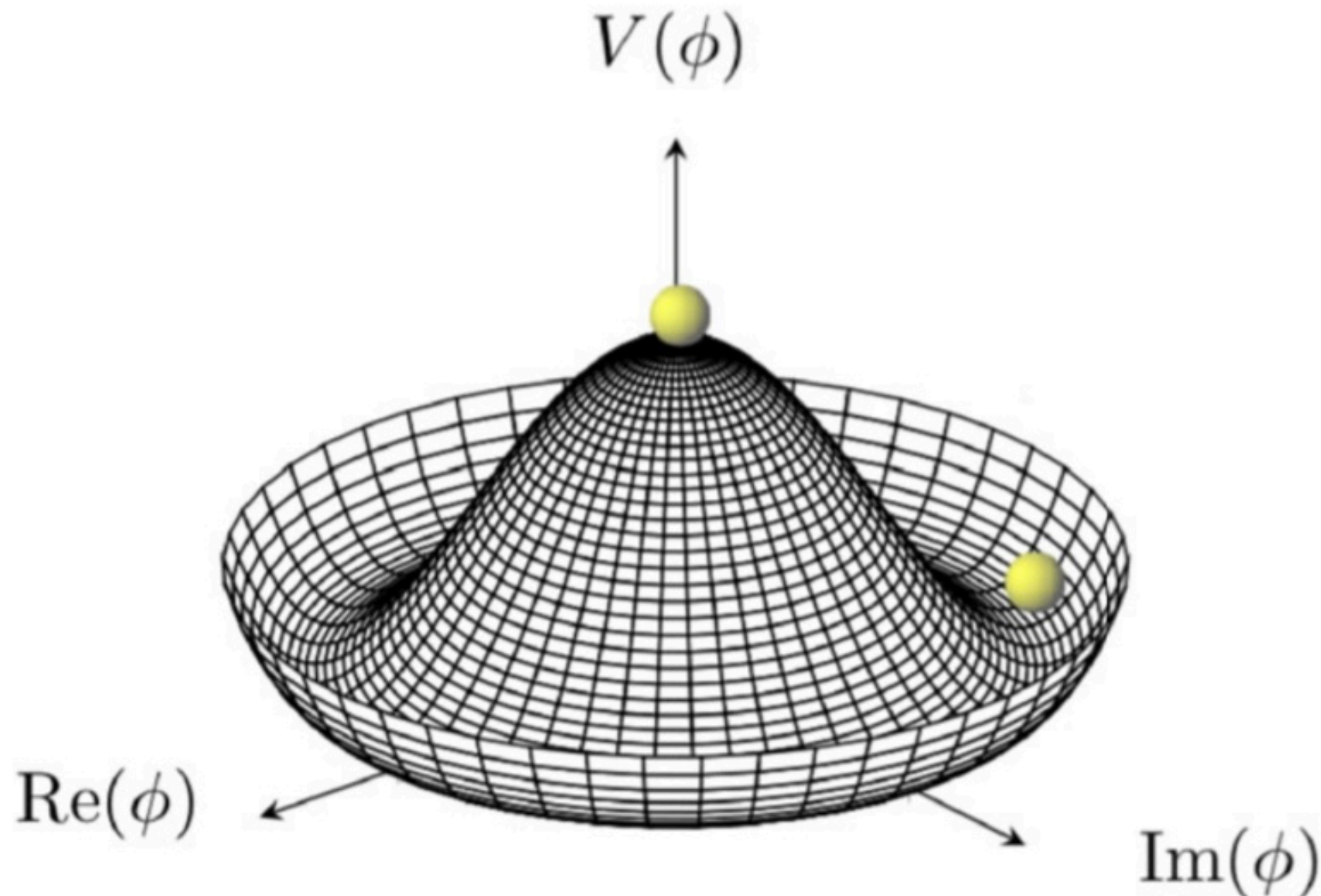
“Gauss law”



Consider the theory (introduced in section 22) of a complex scalar field  $\phi$  with

# Do tachyons exist?

$$\mathcal{L} = -\partial^\mu \phi^\dagger \partial_\mu \phi - m^2 \phi^\dagger \phi - \frac{1}{4} \lambda (\phi^\dagger \phi)^2. \quad (32.1)$$



and... Thus we have a continuous family of minima of the potential, parameterized by  $\theta$ . Under the  $U(1)$  transformation of eq. (32.2),  $\theta$  changes to  $\theta + \alpha$ ; thus the different minimum-energy field configurations are all related



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$$\mathcal{L} =$$

(32.1)

This lagrangian is ob

transformation

(32.2)

where  $\alpha$  is a real nu

Now suppose th  
eq. (32.1) is achieve

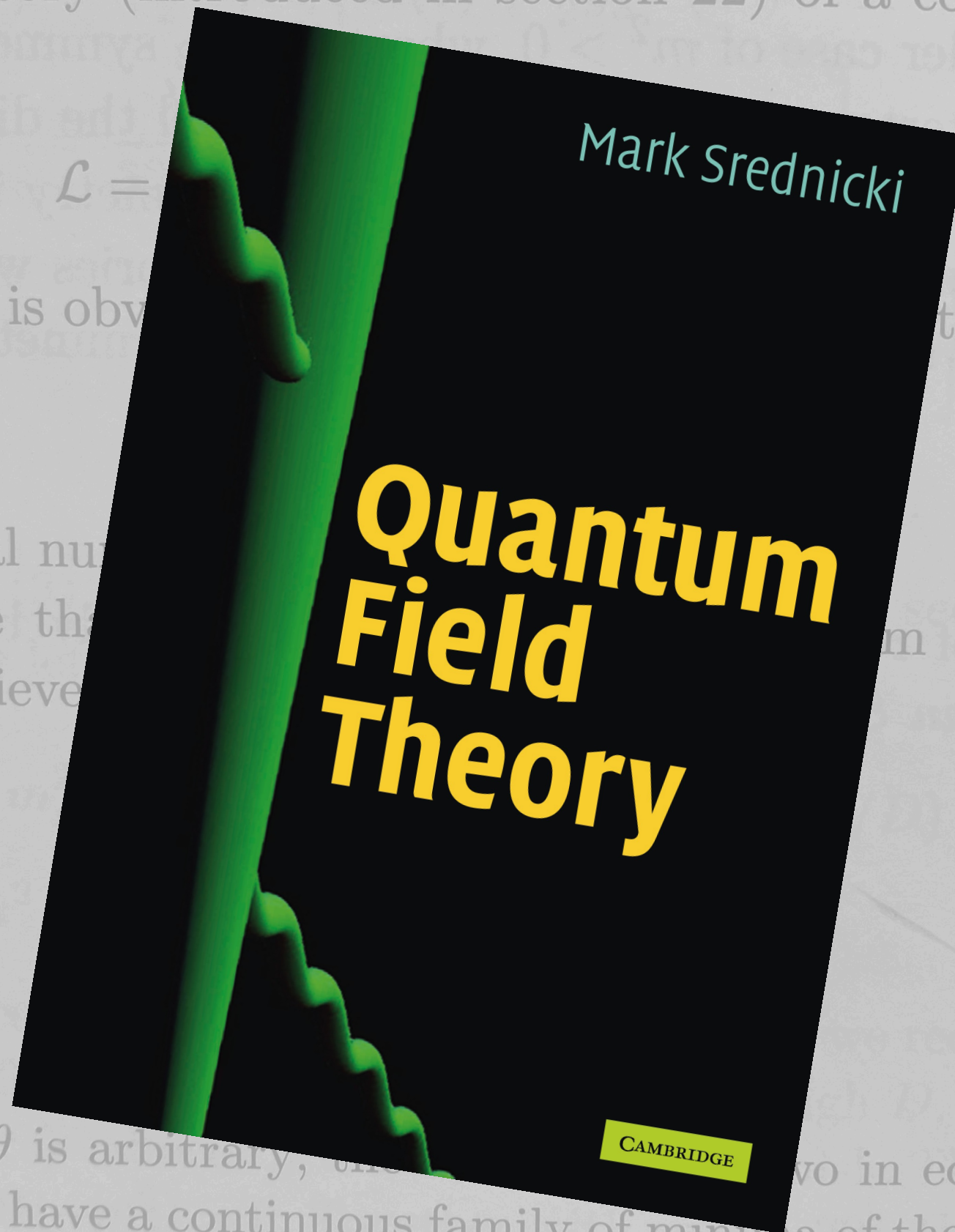
m of the potential of

(32.3)

where

(32.4)

and the phase  $\theta$  is arbitrary, the minimum in eq. (32.3) is conventional. Thus we have a continuous family of minima of the potential, parameterized by  $\theta$ . Under the U(1) transformation of eq. (32.2),  $\theta$  changes to  $\theta + \alpha$ ; thus the different minimum-energy field configurations are all related





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$$\mathcal{L} = -\partial^\mu \varphi^\dagger \partial_\mu \varphi - m^2 \varphi^\dagger \varphi - \frac{1}{4} \lambda (\varphi^\dagger \varphi)^2. \quad (32.1)$$

This lagrangian is obviously invariant under the U(1) transformation

$$\varphi(x) \rightarrow e^{-i\alpha} \varphi(x), \quad (32.2)$$

where  $\alpha$  is a real number.

Now suppose that  $m^2$  is negative. The minimum of the potential of eq. (32.1) is achieved for


$$\varphi(x) = \frac{1}{\sqrt{2}} v e^{-i\theta}, \quad (32.3)$$

where

$$v = (4|m^2|/\lambda)^{1/2}, \quad (32.4)$$

and the phase  $\theta$  is arbitrary; the factor of root-two in eq. (32.3) is conventional. Thus we have a continuous family of minima of the potential, parameterized by  $\theta$ . Under the U(1) transformation of eq. (32.2),  $\theta$  changes to  $\theta + \alpha$ ; thus the different minimum-energy field configurations are all related



Thanks for listening

