

Exercice 1.

(A) $|l, m\rangle$ $\hat{L}^2$ op $\hat{L}_x$ $l=1$

$l=1$ $m=1$ $m=0$ $m=-1$

$\hat{L}^2 |l, m\rangle = \hbar^2 l(l+1) |l, m\rangle$

$\hat{L}^2 |1, 1\rangle = \hbar^2 \cdot 1(2) |1, 1\rangle$

$\hat{L}^2 |1, 1\rangle = 2\hbar^2 |1, 1\rangle$

$\hat{L}^2 |1, 0\rangle = 2\hbar^2 |1, 0\rangle$

$\hat{L}^2 |1, -1\rangle = 2\hbar^2 |1, -1\rangle$

$\langle l', m' | \hat{L}_x | l, m \rangle$

$\langle 1, 1 | \hat{L}_x | 1, 1 \rangle = \hbar$

$\langle 1, 1 | \hat{L}_x | 1, 0 \rangle = 0$

$\langle 1, 1 | \hat{L}_x | 1, -1 \rangle = 0$

$\begin{pmatrix} \hbar & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & -\hbar \end{pmatrix} = \hbar \begin{pmatrix} 1 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & -1 \end{pmatrix}$

(B) $\hat{L}_+ |l, m\rangle = \hbar \cdot \sqrt{l(l+1) - m(m+1)} |l, m+1\rangle$

$\hat{L}_- |l, m\rangle = \hbar \cdot \sqrt{l(l+1) - m(m-1)} |l, m-1\rangle$

$\hat{L}_+ |1, 1\rangle = \hbar \cdot \sqrt{1 \cdot 2 - 1 \cdot 2} |1, 2\rangle = 0$

$\hat{L}_+ |1, 0\rangle = \hbar \cdot \sqrt{1 \cdot 2 - 0} |1, 1\rangle = \sqrt{2} \hbar |1, 1\rangle$

$\hat{L}_+ |1, -1\rangle = \hbar \cdot \sqrt{1 \cdot 2 + 1(-1)} |1, 0\rangle = \sqrt{2} \hbar |1, 0\rangle$

$\hat{L}_- |1, 1\rangle = \hbar \cdot \sqrt{1 \cdot 2 - 1(0)} |1, 0\rangle = \sqrt{2} \hbar |1, 0\rangle$

$\hat{L}_- |1, 0\rangle = \hbar \cdot \sqrt{1 \cdot 2 - 0} |1, -1\rangle = \sqrt{2} \hbar |1, -1\rangle$

$\hat{L}_- |1, -1\rangle = \hbar \cdot \sqrt{1 \cdot 2 + 1(-2)} |1, -2\rangle = 0$

$\hat{L}_+ = \hat{L}_x + i\hat{L}_y$

$\hat{L}_- = \hat{L}_x - i\hat{L}_y$

$2\hat{L}_x = \hat{L}_+ + \hat{L}_-$

$\hat{L}_x = \frac{1}{2}(\hat{L}_+ + \hat{L}_-)$

$\hat{L}_x |l, m\rangle = \left(\frac{1}{2}(\hat{L}_+ + \hat{L}_-) \right) |l, m\rangle = \frac{1}{2} \hat{L}_+ |l, m\rangle + \frac{1}{2} \hat{L}_- |l, m\rangle$

$\hat{L}_x |1, m\rangle = \frac{1}{2} (\hat{L}_+ + \hat{L}_-) |1, m\rangle = \frac{1}{2} \hat{L}_+ |1, m\rangle - \frac{1}{2} \hat{L}_- |1, m\rangle$

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Anna Socha

resolution 1.

$|1, 1\rangle$
 $|1, 0\rangle$
 $|1, -1\rangle$

$\hat{L}_x |l, m\rangle = m \hbar |l, m\rangle$

$\hat{L}_x |1, 1\rangle = \hbar |1, 1\rangle$

$\hat{L}_x |1, 0\rangle = 0 \cdot \hbar |1, 0\rangle$

$\hat{L}_x |1, -1\rangle = -\hbar |1, -1\rangle$

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$$\begin{aligned}\hat{L}_x |1,1\rangle &= \frac{1}{2} \hat{L}_+ |1,1\rangle + \frac{1}{2} \hat{L}_- |1,1\rangle = 0 + \frac{\sqrt{2} \hbar}{2} |1,0\rangle \\ \hat{L}_x |1,0\rangle &= \frac{1}{2} \hat{L}_+ |1,0\rangle + \frac{1}{2} \hat{L}_- |1,0\rangle = \frac{\sqrt{2} \hbar}{2} |1,1\rangle + \frac{\sqrt{2} \hbar}{2} |1,-1\rangle \\ \hat{L}_x |1,-1\rangle &= \frac{1}{2} \hat{L}_+ |1,-1\rangle + \frac{1}{2} \hat{L}_- |1,-1\rangle = \frac{\sqrt{2} \hbar}{2} |1,0\rangle + 0 \\ \hat{L}_y |1,1\rangle &= \frac{1}{2i} \hat{L}_+ |1,1\rangle + \frac{1}{2i} \hat{L}_- |1,1\rangle = 0 - \frac{1}{2i} \sqrt{2} \hbar |1,0\rangle \\ \hat{L}_y |1,0\rangle &= \frac{1}{2i} \hat{L}_+ |1,0\rangle - \frac{1}{2i} \hat{L}_- |1,0\rangle = \frac{\sqrt{2} \hbar}{2i} |1,1\rangle - \frac{\sqrt{2} \hbar}{2i} |1,-1\rangle \\ \hat{L}_y |1,-1\rangle &= \frac{1}{2i} \hat{L}_+ |1,-1\rangle - \frac{1}{2i} \hat{L}_- |1,-1\rangle = \frac{\sqrt{2} \hbar}{2i} |1,0\rangle + 0\end{aligned}$$

$$\begin{aligned}\langle 1,1 | \hat{L}_x | 1,1 \rangle &= \frac{\sqrt{2} \hbar}{2} \langle 1,1 | 1,0 \rangle = 0 & \langle 1,1 | 1,1 \rangle &= 1 \\ \langle 1,1 | \hat{L}_x | 1,0 \rangle &= \frac{\sqrt{2} \hbar}{2} & \langle 1,1 | 1,-1 \rangle &= 0 \\ \langle 1,1 | \hat{L}_x | 1,-1 \rangle &= 0 & \langle 1,1 | 1,0 \rangle &= 0\end{aligned}$$

$$\begin{aligned}\langle 1,0 | \hat{L}_x | 1,1 \rangle &= \frac{\sqrt{2} \hbar}{2} \\ \langle 1,0 | \hat{L}_x | 1,0 \rangle &= 0 \\ \langle 1,0 | \hat{L}_x | 1,-1 \rangle &= \frac{\sqrt{2} \hbar}{2} \\ \langle 1,-1 | \hat{L}_x | 1,1 \rangle &= 0 \\ \langle 1,-1 | \hat{L}_x | 1,0 \rangle &= \frac{\sqrt{2} \hbar}{2}\end{aligned}$$

$$\hat{L}_x = \begin{pmatrix} 0 & \frac{\sqrt{2} \hbar}{2} & 0 \\ \frac{\sqrt{2} \hbar}{2} & 0 & \frac{\sqrt{2} \hbar}{2} \\ 0 & \frac{\sqrt{2} \hbar}{2} & 0 \end{pmatrix} = \frac{\sqrt{2} \hbar}{2} \begin{pmatrix} 0 & 1 & 0 \\ 1 & 0 & 1 \\ 0 & 1 & 0 \end{pmatrix}$$

$$\hat{L}_y = \begin{pmatrix} 0 & -\frac{\sqrt{2} \hbar}{2i} & 0 \\ \frac{\sqrt{2} \hbar}{2i} & 0 & -\frac{\sqrt{2} \hbar}{2i} \\ 0 & \frac{\sqrt{2} \hbar}{2i} & 0 \end{pmatrix} = \frac{\sqrt{2} \hbar}{2i} \begin{pmatrix} 0 & -1 & 0 \\ 1 & 0 & -1 \\ 0 & 1 & 0 \end{pmatrix}$$

~~$$[\hat{L}_x, \hat{L}_y] = \hat{L}_x \hat{L}_y - \hat{L}_y \hat{L}_x =$$~~

~~$$\frac{\sqrt{2} \hbar}{2} \begin{pmatrix} 0 & 1 & 0 \\ 1 & 0 & 1 \\ 0 & 1 & 0 \end{pmatrix} \cdot \frac{\sqrt{2} \hbar}{2i} \begin{pmatrix} 0 & -1 & 0 \\ 1 & 0 & -1 \\ 0 & 1 & 0 \end{pmatrix} -$$~~

~~$$\frac{\sqrt{2} \hbar}{2} \begin{pmatrix} 0 & -1 & 0 \\ 1 & 0 & -1 \\ 0 & 1 & 0 \end{pmatrix} \cdot \frac{\sqrt{2} \hbar}{2i} \begin{pmatrix} 0 & 1 & 0 \\ 1 & 0 & 1 \\ 0 & 1 & 0 \end{pmatrix} =$$~~

~~$$\frac{1}{2} \hbar^2 \begin{pmatrix} 1 & 0 & -1 \\ 0 & 0 & 0 \\ 1 & 0 & -1 \end{pmatrix} - \frac{1}{2} \hbar^2 \begin{pmatrix} -1 & 0 & 1 \\ 0 & 0 & 0 \\ 1 & 0 & 1 \end{pmatrix} =$$~~

~~$$\frac{1}{2} \hbar^2 \begin{pmatrix} 2 & 0 & 2 \\ 0 & 0 & 0 \\ 0 & 0 & -2 \end{pmatrix} = \frac{1}{2} \hbar^2 \begin{pmatrix} 1 & 0 & 1 \\ 0 & 0 & 0 \\ 0 & 0 & -1 \end{pmatrix}$$~~

~~$$\begin{pmatrix} 0 & 1 & 0 & | & 1 & 0 & -1 \\ 1 & 0 & 1 & | & 0 & 0 & 0 \\ 0 & 1 & 0 & | & 1 & 0 & -1 \\ \hline 0 & 1 & 0 & | & 1 & 0 & -1 \\ 1 & 0 & 1 & | & 0 & 0 & 0 \\ 0 & 1 & 0 & | & 1 & 0 & -1 \end{pmatrix}$$~~

$$|\psi\rangle = \frac{1}{\sqrt{3}}|\phi_0\rangle + \sqrt{\frac{2}{3}}|\phi_1\rangle$$

$$\downarrow$$

$$\frac{\hbar\omega}{2} \quad \frac{3\hbar\omega}{2}$$

$$E_0 = \frac{\hbar\omega}{2} \quad E_1 = \frac{3\hbar\omega}{2}$$

Przedpostawienie:

~~Korzystamy z~~

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$$A) |\langle\phi_0|\psi\rangle|^2 = \left|\frac{1}{\sqrt{3}}\langle\phi_0|\phi_0\rangle + \sqrt{\frac{2}{3}}\langle\phi_0|\phi_1\rangle\right|^2 = \frac{1}{3}$$

Przedpostawienie, że energia kol w stanie $|\psi\rangle$ jest w stanie podstaw = $\frac{1}{3}$

b) ~~Korzystamy z~~

$$\langle\psi|\hat{H}|\psi\rangle$$

Wzrostem operatora harmonicznego

duży jest użycie

$$\hat{H} = \frac{\hat{p}^2}{2m} + V(x) = \frac{-\hbar^2}{2m} \frac{\partial^2}{\partial x^2} + \frac{1}{2} m \omega^2 x^2$$

~~Korzystamy z~~

$$\langle\psi|\hat{H}|\psi\rangle = \langle\psi|\left(-\frac{\hbar^2}{2m} \frac{\partial^2}{\partial x^2} + \frac{1}{2} m \omega^2 x^2\right)|\psi\rangle$$

Równanie Schrödingera

$$\hat{H}|\psi\rangle = E|\psi\rangle$$

gdzie E - wartość własna

ty energia w określonym stanie

~~Korzystamy z~~

$$\hat{H}|\phi_0\rangle = E_0|\phi_0\rangle$$

$$\hat{H}|\phi_0\rangle = \frac{\hbar\omega}{2}|\phi_0\rangle$$

$$\hat{H}|\phi_1\rangle = E_1|\phi_1\rangle$$

$$\hat{H}|\phi_1\rangle = \frac{3\hbar\omega}{2}|\phi_1\rangle$$

$$\langle\psi|\hat{H}|\psi\rangle = \langle\psi|\left(\frac{1}{\sqrt{3}}\hat{H}|\phi_0\rangle + \sqrt{\frac{2}{3}}\hat{H}|\phi_1\rangle\right) =$$

~~Korzystamy z~~

$$\langle\psi|\frac{1}{\sqrt{3}}\hat{H}|\phi_0\rangle + \sqrt{\frac{2}{3}}\langle\psi|\hat{H}|\phi_1\rangle =$$

$$\frac{1}{\sqrt{3}}\langle\psi|\frac{\hbar\omega}{2}|\phi_0\rangle + \sqrt{\frac{2}{3}}\langle\psi|\frac{3\hbar\omega}{2}|\phi_1\rangle = \frac{\hbar\omega}{2\sqrt{3}}\langle\phi_0|\phi_0\rangle +$$

$$\sqrt{\frac{2}{3}}\langle\phi_1|\phi_0\rangle + \sqrt{\frac{2}{3}}\frac{3\hbar\omega}{2}\langle\phi_1|\phi_1\rangle = \frac{\hbar\omega}{2\sqrt{3}} + \sqrt{\frac{2}{3}}\langle\phi_1|\phi_1\rangle = \frac{\hbar\omega}{2\sqrt{3}} + \sqrt{\frac{2}{3}}\hbar\omega$$

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c) ψ inexakte Potenzreihe

$$\langle x' | x \rangle = x \delta(x - x')$$

$$|\phi_0\rangle = \sqrt{\frac{1}{2^0 \cdot 0!}} \left(\frac{m\omega}{\hbar \pi} \right)^{\frac{1}{4}} e^{-\frac{m\omega x^2}{2\hbar}} \mathcal{H}_0\left(\sqrt{\frac{m\omega}{\hbar}} x\right) = \left(\frac{m\omega}{\hbar \pi} \right)^{\frac{1}{4}} e^{-\frac{m\omega x^2}{2\hbar}} \cdot 1$$

$$|\phi_1\rangle = \sqrt{\frac{1}{2 \cdot 1!}} \left(\frac{m\omega}{\hbar \pi} \right)^{\frac{1}{4}} e^{-\frac{m\omega x^2}{2\hbar}} \mathcal{H}_1\left(\sqrt{\frac{m\omega}{\hbar}} x\right) = \frac{1}{\sqrt{2}} \left(\frac{m\omega}{\hbar \pi} \right)^{\frac{1}{4}} e^{-\frac{m\omega x^2}{2\hbar}}$$

$$2 \cdot \sqrt{\frac{m\omega}{\hbar}} \cdot x = \frac{2}{\sqrt{2}} \left(\frac{m\omega}{\hbar \pi} \right)^{\frac{1}{4}} e^{-\frac{m\omega x^2}{2\hbar}} \cdot \sqrt{\frac{m\omega}{\hbar}} \cdot x$$

~~Erwartungswert~~

$$\langle x \rangle = \int_{-\infty}^{+\infty} \psi^* x \psi dx = \int_{-\infty}^{+\infty} \left(\frac{1}{\sqrt{2}} \cdot \left(\frac{m\omega}{\hbar \pi} \right)^{\frac{1}{4}} e^{-\frac{m\omega x^2}{2\hbar}} \right) x dx$$

$$+ \sqrt{\frac{2}{3}} \cdot \frac{2}{\sqrt{2}} \left(\frac{m\omega}{\hbar \pi} \right)^{\frac{1}{4}} e^{-\frac{m\omega x^2}{2\hbar}} \cdot \sqrt{\frac{m\omega}{\hbar}} \cdot x \cdot \left(\frac{1}{\sqrt{2}} \left(\frac{m\omega}{\hbar \pi} \right)^{\frac{1}{4}} e^{-\frac{m\omega x^2}{2\hbar}} \right) dx =$$

$$+ x \cdot \sqrt{\frac{2}{3}} \cdot \frac{2}{\sqrt{2}} \left(\frac{m\omega}{\hbar \pi} \right)^{\frac{1}{4}} e^{-\frac{m\omega x^2}{2\hbar}} dx =$$

$$\int_{-\infty}^{+\infty} \left(\frac{1}{\sqrt{2}} \left(\frac{m\omega}{\hbar \pi} \right)^{\frac{1}{4}} e^{-\frac{m\omega x^2}{2\hbar}} + 2x \sqrt{\frac{m\omega}{\hbar}} \cdot x \cdot e^{-\frac{m\omega x^2}{2\hbar}} \right) \cdot x \left(\frac{1}{\sqrt{2}} \left(\frac{m\omega}{\hbar \pi} \right)^{\frac{1}{4}} e^{-\frac{m\omega x^2}{2\hbar}} \right) dx +$$

$$\left(e^{-\frac{m\omega x^2}{2\hbar}} \cdot 2x \cdot \sqrt{\frac{m\omega}{\hbar}} \right) = \int_{-\infty}^{+\infty} \frac{1}{3} \left(\frac{m\omega}{\hbar \pi} \right)^{\frac{1}{2}} \int_{-\infty}^{+\infty} \left(e^{-\frac{m\omega x^2}{2\hbar}} + 2x \sqrt{\frac{m\omega}{\hbar}} \cdot x \cdot e^{-\frac{m\omega x^2}{2\hbar}} \right) dx$$

$$\left(x \cdot e^{-\frac{m\omega x^2}{2\hbar}} + 2x^2 \sqrt{\frac{m\omega}{\hbar}} e^{-\frac{m\omega x^2}{2\hbar}} \right) = \frac{1}{3} \left(\frac{m\omega}{\hbar \pi} \right)^{\frac{1}{2}} \int_{-\infty}^{+\infty} \left(e^{-\frac{m\omega x^2}{2\hbar}} + 2x^2 \sqrt{\frac{m\omega}{\hbar}} e^{-\frac{m\omega x^2}{2\hbar}} \right) dx$$

$$= \int_{-\infty}^{+\infty} \frac{m\omega}{\hbar} 2x^2 e^{-\frac{m\omega x^2}{2\hbar}} dx + \int_{-\infty}^{+\infty} \frac{m\omega x^2}{\hbar} dx = \frac{m\omega}{\hbar} \int_{-\infty}^{+\infty} 2x^2 e^{-\frac{m\omega x^2}{2\hbar}} dx + \frac{m\omega}{\hbar} \int_{-\infty}^{+\infty} x^2 dx =$$

$$\left\{ \begin{array}{l} \frac{m\omega}{\hbar} x = u \\ dx = \frac{\hbar}{m\omega} du \end{array} \right.$$

oder wiederum Hermitesche $\mathcal{H}_1(x) = 2x$

Potenzreihe

Chcemy zobaczyć
rozwiązanie

$$\begin{aligned} & \frac{1}{3} \left(\frac{m\omega}{\hbar \cdot \pi} \right)^{\frac{1}{2}} \left[\int_{-\infty}^{+\infty} \frac{\hbar^2}{m^2 \omega^2} \cdot u \cdot e^{-u^2} du + 4 \int_{-\infty}^{+\infty} \frac{\hbar^3}{m^3 \omega^3} \cdot u^2 \cdot e^{-u} du \right] \\ & + 4 \int_{-\infty}^{+\infty} \frac{m\omega}{\hbar} \cdot \frac{\hbar^4}{m^4 \omega^4} u^3 \cdot e^{-u} du = \frac{1}{3} \left(\frac{m\omega}{\hbar \cdot \pi} \right)^{\frac{1}{2}} \cdot \left(\frac{\hbar^2}{m^2 \omega^2} \right) \int_{-\infty}^{+\infty} u \cdot e^{-u^2} du + \frac{4}{3} \left(\frac{m\omega}{\hbar \cdot \pi} \right)^{\frac{1}{2}} \int_{-\infty}^{+\infty} u^2 \cdot e^{-u} du \\ & + \frac{4}{3} \left(\frac{m\omega}{\hbar \cdot \pi} \right)^{\frac{1}{2}} \int_{-\infty}^{+\infty} u^3 \cdot e^{-u} du = \frac{1}{3} \left(\frac{m\omega}{\hbar \cdot \pi} \right)^{\frac{1}{2}} \cdot \left(\frac{\hbar^2}{m^2 \omega^2} \right) \cdot \frac{5}{2} \cdot \frac{\sqrt{\pi}}{2} \cdot \frac{2}{2} \\ & = \frac{1}{3} \cdot \left(\frac{m\omega}{\hbar} \right)^{\frac{1}{2}} \cdot \left(\frac{\hbar}{m\omega} \right)^{\frac{5}{2}} \cdot \frac{2\hbar^2}{3m^2\omega^2} \end{aligned}$$

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ten wzór może mieć
dobry, ale nie pierwszy
wzór
 $\psi(t) = \psi \cdot e^{-\frac{i\hbar\omega}{2}t}$
ten wzór ma być
i to w
długości

(0) $t=0 \rightarrow |\psi\rangle$ ewolucja operatora $|\psi(t)\rangle$

$$|\psi(t)\rangle = \frac{1}{\sqrt{3}} |\phi_0(t)\rangle + \sqrt{\frac{2}{3}} |\phi_1(t)\rangle$$

$$\begin{aligned} |\psi(t)\rangle &= \frac{1}{\sqrt{3}} \left(\frac{m\omega}{\hbar \cdot \pi} \right)^{\frac{1}{4}} \cdot e^{-\frac{m\omega x^2}{2\hbar}} \cdot e^{-\frac{i\hbar\omega t}{2\hbar \cdot \omega}} - \frac{1}{\sqrt{3}} \left(\frac{m\omega}{\hbar \cdot \pi} \right)^{\frac{1}{4}} \cdot e^{-\frac{m\omega x^2}{2\hbar}} \cdot e^{-\frac{i\hbar\omega t}{2\hbar \cdot \omega}} \\ &= \frac{1}{\sqrt{3}} \left(\frac{m\omega}{\hbar \cdot \pi} \right)^{\frac{1}{4}} \cdot e^{-\frac{m\omega x^2}{2\hbar}} \cdot e^{-\frac{i\hbar\omega t}{2\hbar \cdot \omega}} \left(e^{-\frac{i\hbar\omega t}{2\hbar \cdot \omega}} + 2e^{-\frac{3i\hbar\omega t}{2\hbar \cdot \omega}} \right) \end{aligned}$$

~~$$-\frac{1}{\sqrt{3}} \left(\frac{m\omega}{\hbar \cdot \pi} \right)^{\frac{1}{4}} \cdot e^{-\frac{m\omega x^2}{2\hbar}} \cdot e^{-\frac{i\hbar\omega t}{2\hbar \cdot \omega}} \cdot e^{-\frac{i\hbar\omega t}{2\hbar \cdot \omega}}$$~~

$$\langle \psi(t) | \psi(t) \rangle = \int_{-\infty}^{+\infty} \frac{1}{3} \left(\frac{m\omega}{\hbar \cdot \pi} \right)^{\frac{1}{2}} \cdot e^{-\frac{m\omega x^2}{\hbar}} \cdot \left(e^{\frac{i\hbar\omega t}{2\hbar \cdot \omega}} + 2e^{-\frac{3i\hbar\omega t}{2\hbar \cdot \omega}} \right) \cdot \left(e^{-\frac{i\hbar\omega t}{2\hbar \cdot \omega}} + 2e^{-\frac{3i\hbar\omega t}{2\hbar \cdot \omega}} \right) dx$$

~~$$-\frac{1}{3} \left(\frac{m\omega}{\hbar \cdot \pi} \right)^{\frac{1}{2}} \cdot e^{-\frac{m\omega x^2}{\hbar}} \cdot e^{-\frac{i\hbar\omega t}{2\hbar \cdot \omega}} \cdot e^{-\frac{i\hbar\omega t}{2\hbar \cdot \omega}}$$~~

$$= \frac{1}{3} \left(\frac{m\omega}{\hbar \cdot \pi} \right)^{\frac{1}{2}} \cdot \int_{-\infty}^{+\infty} e^{-\frac{m\omega x^2}{\hbar}} \cdot \left(x + 2x \cdot e^{-\frac{i\hbar\omega t}{2\hbar \cdot \omega}} + 2x \cdot e^{-\frac{3i\hbar\omega t}{2\hbar \cdot \omega}} \right) dx =$$

~~$$\frac{1}{3} \left(\frac{m\omega}{\hbar \cdot \pi} \right)^{\frac{1}{2}} \cdot \int_{-\infty}^{+\infty} e^{-\frac{m\omega x^2}{\hbar}} \cdot x dx = \frac{3}{3} \left(\frac{m\omega}{\hbar \cdot \pi} \right)^{\frac{1}{2}} \cdot \int_{-\infty}^{+\infty} x \cdot e^{-\frac{m\omega x^2}{\hbar}} dx$$~~

3. $\psi(r, \theta, \phi) = \frac{1}{\sqrt{\pi}} \left(\frac{1}{a_0} \right)^{3/2} \cdot e^{-r/a_0}$

(A) $r = a_0$

Dany jest stan $\psi(r, \theta, \phi)$ dla $\theta(r, \theta, \phi)$ taki, że $r = a_0$

$\theta(a_0, \theta, \phi) = \frac{1}{\sqrt{\pi}} \left(\frac{1}{a_0} \right)^{3/2} \cdot e^{-1} = \frac{1}{e\sqrt{\pi} a_0^{3/2}}$

Przebadajemy, że elektron znajduje się w strefie θ wynosi

$|\langle \theta | \psi \rangle|^2 = \frac{16}{a_0^2 \cdot e^4}$

$\langle \theta | \psi \rangle = \int d^3r \psi^*(r, \theta, \phi) \cdot \psi(r, \theta, \phi) =$

$\int_0^{2\pi} \int_0^\pi \int_0^\infty d\phi \cdot \sin\theta \cdot \frac{1}{\sqrt{\pi}} \left(\frac{1}{a_0} \right)^{3/2} \cdot e^{-r/a_0} \cdot \frac{1}{\sqrt{\pi}} \left(\frac{1}{a_0} \right)^{3/2} \cdot e^{-r/a_0} \cdot a_0^2 \cdot dr =$

$2\pi \cdot \frac{1}{\pi} \left(\frac{1}{a_0} \right)^3 \cdot a_0^2 \int_0^\infty \frac{1}{e} \cdot e^{-r/a_0} \cdot \sin\theta d\theta = \frac{2}{e a_0} \cdot e^{-r/a_0} \cdot (\cos\theta) \Big|_0^\pi =$

$\frac{2}{e a_0} \cdot e^{-r/a_0} \cdot 2 = \frac{4}{a_0 e} \cdot \frac{1}{a_0 e^2} = \frac{4}{a_0^2 e^3}$

(B) Przechodzimy do tego, że elektron znajduje się w obs. nie większej niż jeden promień Bohra

Wtedy założymy, że chodzi o to, że elektron znajduje się w strefie $(0, a_0)$

$\theta(r_0, \theta, \phi) = \frac{1}{\sqrt{\pi}} \left(\frac{1}{a_0} \right)^{3/2} \cdot e^{-r/a_0}$

$|\langle \theta | \psi \rangle|^2 = \int d^3r \psi^* \psi = \int_0^{2\pi} \int_0^\pi \int_0^{a_0} d\phi \cdot \sin\theta \cdot \frac{1}{\sqrt{\pi}} \left(\frac{1}{a_0} \right)^3 \cdot e^{-r/a_0} \cdot e^{-r/a_0} \cdot r^2 \cdot dr =$

$= \frac{2\pi}{\pi} \cdot \left(\frac{1}{a_0} \right)^3 \cdot (-\cos\theta) \Big|_0^\pi \cdot \int_0^{a_0} r^2 \cdot dr = 4 \cdot \left(\frac{1}{a_0} \right)^3 \cdot \int_0^{a_0} r^2 \cdot dr = 4 \cdot \frac{1}{a_0^3} \cdot \frac{r^3}{3} \Big|_0^{a_0} = \frac{4}{3} \cdot \frac{a_0^3}{a_0^3} = \frac{4}{3}$

$\left\{ \begin{array}{l} \frac{2\pi}{a_0} = \pi \\ \frac{dr}{dr} = \frac{a_0}{a} \end{array} \right\} = \frac{4}{3} \cdot \left(\frac{1}{a_0} \right)^3 \cdot \int_0^{a_0} r^2 \cdot dr = \frac{4}{3} \cdot \frac{1}{a_0^3} \cdot \frac{r^3}{3} \Big|_0^{a_0} = \frac{4}{3} \cdot \frac{a_0^3}{a_0^3} = \frac{4}{3}$

$$\frac{1}{2} \int_0^2 e^{-u} \cdot u \, du = \frac{1}{2} \cdot e^{-u} \left(\frac{u^2}{(-1)} - \frac{2u}{(-1)^2} + \frac{2}{(-1)^3} \right) \Big|_0^2 =$$

$$\frac{1}{2} [e^{-2}(-4-4-2) - \cancel{0} (0-0-2)] =$$

(1b)

$$\frac{1}{2} (e^{-2}(-10)+2) = \underline{1-5 \cdot e^{-2}} \quad | \text{Prawdopodobieństwo } |\langle 0|\psi \rangle|^2 = \underline{(1-5 \cdot e^{-2})^2}$$

$$\begin{aligned} \text{(c)} \quad \langle \psi | \hat{r} | \psi \rangle &= \int d^3r \, \psi^* \cdot r \cdot \psi = \int_0^{2a_0} \int_0^\pi \int_0^{2\pi} d\varphi \cdot \int_0^\pi d\theta \sin\theta \cdot \int_0^{2a_0} r^2 \cdot \left(\frac{1}{\pi} \cdot \left(\frac{1}{a_0} \right)^{3/2} \cdot e^{-\frac{2r}{a_0}} \right) \cdot r \cdot dr \\ &= \pi \cdot \frac{1}{\pi} \cdot \left(\frac{1}{a_0} \right)^{3/2} \cdot e^{-\frac{2r}{a_0}} \cdot \int_0^{2a_0} r^3 \cdot dr = \int_0^{2a_0} \frac{1}{\pi} \cdot \left(\frac{1}{a_0} \right)^3 \cdot r^3 \cdot e^{-\frac{2r}{a_0}} \cdot dr \\ &= 4 \left(\frac{1}{a_0} \right)^3 \cdot \int_0^{2a_0} r^3 \cdot e^{-\frac{2r}{a_0}} \cdot dr \quad \left\{ \begin{array}{l} \frac{2}{a_0} r = u \\ dr = \frac{a_0}{2} du \end{array} \right\} = \end{aligned}$$

(2/2)

$$4 \left(\frac{1}{a_0} \right)^3 \cdot \int_0^{2a_0} \frac{a_0^4}{16} \cdot u^3 \cdot e^{-u} \cdot du = \frac{a_0}{4} \cdot \int_0^{2a_0} u^3 \cdot e^{-u} \cdot du = \underline{\frac{6a_0}{4} - \frac{3}{2}a_0}$$

(d) Jak najbardziej prawdopodobny całkowity max. wartość energia $\langle 0|\psi \rangle$

$$\langle 0|\psi \rangle = \int d^3r \, \psi^* \cdot \psi = \int_0^{2a_0} \int_0^\pi \int_0^{2\pi} d\varphi \cdot \int_0^\pi d\theta \sin\theta \cdot d\theta \cdot \frac{1}{\pi} \cdot \left(\frac{1}{a_0} \right)^3 \cdot e^{-\frac{2r}{a_0}} =$$

$$4 \cdot \left(\frac{1}{a_0} \right)^3 \cdot \int_0^{2a_0} r^2 \cdot dr = \int_0^{2a_0} \frac{2r}{a_0} \cdot r \cdot dr = \int_0^{2a_0} \frac{2r}{a_0} \cdot r \cdot dr = \int_0^{2a_0} \frac{2r}{a_0} \cdot r \cdot dr =$$

$$4 \left(\frac{1}{a_0} \right)^3 \cdot \int_0^{2a_0} \frac{2r}{a_0} \cdot r \cdot dr = \frac{1}{2} \cdot \int_0^{2a_0} u^2 \cdot e^{-u} \cdot du = \int_0^{2a_0} \frac{1}{2} \cdot u^2 \cdot e^{-u} \cdot du =$$

zgodnie ze wzorem

$$\int x^2 e^{-x} = e^{-x} \left(-\frac{x^2}{1} - \frac{2x}{1} + \frac{2}{1} \right) = -\frac{1}{e^x} (x^2 - 2x + 2)$$

to jest prawdopodobieństwo ze danego elektron znajduje się gdzieś

dominującym wkładem będzie $\frac{1}{e^x} \rightarrow$ ~~nie ma sensu~~

~~jest to wartość~~