Examples of explicit Fefferman-Graham metrics

Paweł Nurowski (joint work with Ian Anderson)

Instytut Fizyki Teoretycznej Uniwersytet Warszawski

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Plan

Ambient metrics and distributions

- Fefferman-Graham construction
- Conformal structures and Cartan's paper

The main theorem

- An ansatz
- The theorem



Examples of explicit ambient metrics

- Solutions analytic in ρ
- Nonanalytic in ρ solutions
- Poincaré-Einstein picture

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Conformal structure

A conformal structure $(M^n, [g])$ on an $n = n_+ + n_-$ dimensional manifold M^n is an equivalence [g] class of (n_+, n_-) -signature metrics on M^n , such that two metrics g and \hat{g} are in the same class [g] if and only if there exists a function ϕ on M^n , such that

 $\hat{g} = \mathrm{e}^{2\phi}g.$

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Ambient metric

- Consider a conformal structure (*Mⁿ*, [g]) as defined on the previous slide.
- An *ambient space* \tilde{M} for $(M^n, [g])$ is locally a product

 $\tilde{M} =]0, +\infty[\times M^n \times] - \epsilon, \epsilon[, \quad \epsilon > 0,$

with respective coordinates (t, x^i, ρ) , and the *ambient metric* \tilde{g} for $(M^n, [g])$ is an $(n_+ + 1, n_- + 1)$ -signature *Ricci flat* metric on \tilde{M} given by:

$$\tilde{g} = 2\mathrm{d}t\mathrm{d}(\rho t) + t^2 g(\mathbf{x}^i, \rho)$$

such that

$$g(\mathbf{x}^i,\rho)_{|\rho=0}=g(\mathbf{x}^i),$$

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Explicit ambient metrics?

Assuming that the metric \tilde{g} admits a power series expansion with integer powers in ρ one can see that:

• If [g] contains the *flat* metric g_0 than

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and a 5-manifold M^5 parametrized by (x, y, p, q, z), and equipped with a conformal structure [g] represented by

 $g = 2\omega^{1}\omega^{5} - 2\omega^{2}\omega^{4} + (\omega^{3})^{2},$ $\omega^{1} = dy - pdx, \quad \omega^{2} = dz - (q^{2} + f + bz)dx - \frac{\sqrt{2}}{2}q\omega^{3}, \quad \omega^{3} = 2\sqrt{2}(dp - qdx), \quad \omega^{4} = 3dx,$ $\omega^{5} = \frac{\sqrt{2b}}{2}\omega^{3} - 6dq + 3(2bq + f_{p})dx + \frac{1}{10}(9f_{pp} + 4b^{2})\omega^{1},$

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$$\begin{aligned} A &= \frac{27}{8} f_6 \rho^2 - \frac{9}{5} (f_4 + 5\rho f_5 + 15\rho^2 f_6) \rho, \\ B &= \frac{1}{16} (f_5 + 6\rho f_6) \rho^2 - \frac{3}{20} (f_3 + 4\rho f_4 + 10\rho^2 f_5 + 20\rho^3 f_6) \rho, \\ C &= \frac{1}{360} (f_4 + 5\rho f_5 + 15\rho^2 f_6) \rho^2 - \frac{1}{45} (f_2 + 3\rho f_3 + 6\rho^2 f_4 + 10\rho^3 f_5 + 15\rho^4 f_6) \rho. \\ &= 4 \Box \nu + 4 \Box \mu +$$

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Distributions associated with z' = F(x, y, y', y'', z)

Associated with a differential equation

 $\mathbf{z}' = \mathbf{F}(\mathbf{x}, \mathbf{y}, \mathbf{p}, \mathbf{q}, \mathbf{z}),$

where p = y', q = y'', there is a 5-manifold M^5 parametrized by (x, y, p, q, z), and a distribution

$$\mathcal{D} = \operatorname{Span}\left(\partial_q, \partial_x + p\partial_y + q\partial_p + F\partial_z\right).$$

whose differential invariants, when $F_{qq} \neq 0$, are in one-to-one correspondance with *conformal* invariants of a certain conformal class $[g_D]$ of metrics of signature (3, 2) on M^5 .

Fefferman-Graham construction Conformal structures and Cartan's paper

Distributions associated with z' = F(x, y, y', y'', z)

Associated with a differential equation

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Fefferman-Graham construction Conformal structures and Cartan's paper

The conformal class for $F = q^2 + f(x, p) + bz$

If $F = q^2 + f(x, p) + bz$, where *b* is a real constant, the conformal class may be represented by a metric g_{D_f} in a relatively simple form:

$$g_{D_f} = 8 \left(dp - q dx \right)^2 - 6 \left(dz - 2q dp + (q^2 - f - bz) dx \right) dx - 2 \left(dy - p dx \right) \left(6 dq - 2b dp - (\frac{2}{5}b^2 + \frac{9}{10}f_{\rho\rho})(dy - p dx) - (4bq + 3f_{\rho})dx \right).$$

QUESTION: Can we find explicit formulae for Fefferman-Graham ambient metrics for the conformal class $(M^5, [g_{D_f}])$?

Fefferman-Graham construction Conformal structures and Cartan's paper

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An ansatz The theorem

Plan

Ambient metrics and distributions Fefferman-Graham construction Conformal structures and Cartan's paper

- The main theorem
 - An ansatz
 - The theorem
- 3 Examples of explicit ambient metrics
 - Solutions analytic in ρ
 - Nonanalytic in ρ solutions
 - Poincaré-Einstein picture

An ansatz

- Observation: The Schouten tensor for the class [g_{D_f}] has the form: P = α ⋅ (ω¹)² + 2β ⋅ ω¹ω⁴ + γ ⋅ (ω⁴)², with ω¹ = dy pdx and ω⁴ = 3dx, and α, β, γ functions depending on f and its derivatives.
- Idea: Make an *ansatz* for the ambient metric \tilde{g}_{D_f} in which $g_{D_f}(x^i, \rho)$ assumes a similar form.
- Explicitly, make the following ansatz for $\tilde{g}_{\mathcal{D}_f}$:

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with *unknown* functions $A = A(x, p, \rho)$, $B = B(x, p, \rho)$ and $C = C(x, p, \rho)$.

An ansatz The theorem

Plan

Ambient metrics and distributions

- Fefferman-Graham construction
- Conformal structures and Cartan's paper

The main theorem

- An ansatz
- The theorem

3 Examples of explicit ambient metrics

- Solutions analytic in ρ
- Nonanalytic in ρ solutions
- Poincaré-Einstein picture

An ansatz The theorem

Theorem (Ian Anderson + PN)

The metric $\tilde{g}_{\mathcal{D}_f}$, as above, is an ambient metric for the conformal class $(M^5, [g_{\mathcal{D}_f}])$, if and only if the unknown functions $A = A(x, p, \rho)$, $B = B(x, p, \rho)$ and $C = C(x, p, \rho)$, satisfy the initial conditions $A_{|\rho=0} \equiv 0$, $B_{|\rho=0} \equiv 0$, $C_{|\rho=0} \equiv 0$ and the following system of PDEs:

$$LA = \frac{9}{40} f_{pppp}$$

$$LB = -\frac{1}{36} A_p + \frac{3}{40} f_{ppp}$$

$$LC = -\frac{1}{18} B_p + \frac{1}{324} A + \frac{1}{30} f_{pp} - \frac{2}{15} b^2,$$

with the linear operator L given by

$$L = 2\rho \frac{\partial^2}{\partial \rho^2} - 3\frac{\partial}{\partial \rho} - \frac{1}{8}\frac{\partial^2}{\partial \rho^2}.$$

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Solutions analytic in ρ Nonanalytic in ρ solutions Poincaré-Einstein picture

Plan

Ambient metrics and distributions Fefferman-Graham construction

- Conformal structures and Cartan's paper
- 2 The main theorem
 - An ansatz
 - The theorem
- 3

Examples of explicit ambient metrics

- Solutions analytic in ρ
- Nonanalytic in ho solutions
- Poincaré-Einstein picture

Solutions analytic in ρ Nonanalytic in ρ solutions Poincaré-Einstein picture

Power series expansion in ρ

One can solve the above equations, assuming power series expansion in ρ :

$$A = \sum_{k=1}^{\infty} a_k(x,p) \rho^k, \quad B = \sum_{k=1}^{\infty} b_k(x,p) \rho^k, \quad C = \sum_{k=1}^{\infty} c_k(x,p) \rho^k,$$

obtaining:

$$\begin{split} A &= \sum_{k=1}^{\infty} \frac{3}{5} \cdot \frac{(2k-1)(2k-3)}{2^{2k}(2k)!} \cdot \frac{\partial^{(2k+2)}f}{\partial \rho^{(2k+2)}} \cdot \rho^k, \\ B &= -\sum_{k=1}^{\infty} \frac{1}{15} \cdot \frac{(2k-1)(2k-3)(2k-5)}{2^{2k}(2k)!} \cdot \frac{\partial^{(2k+1)}f}{\partial \rho^{(2k+1)}} \cdot \rho^k, \\ C &= \sum_{k=1}^{\infty} \left(\frac{2}{135} \cdot \frac{(k-3)(2k-1)(2k-3)(2k-5)}{2^{2k}(2k)!} \cdot \frac{\partial^{2k}f}{\partial \rho^{2k}} + \frac{2}{45}b^2\delta_{1k}\right) \cdot \rho^k. \end{split}$$

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Solutions analytic in ρ Nonanalytic in ρ solutions Poincaré-Einstein picture

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Solutions analytic in ρ Nonanalytic in ρ solutions Poincaré-Einstein picture

Solutions being ploynomials in ρ

An important feature of the analytic solutions is that their coefficients behave as:

$$a_k(x,p)\sim rac{\partial^{(2k+2)}f}{\partial p^{(2k+2)}}, \quad b_k(x,p)\sim rac{\partial^{(2k+1)}f}{\partial p^{(2k+1)}}, \quad c_k(x,p)\sim rac{\partial^{(2k)}f}{\partial p^{(2k)}}.$$

Thus, if we want to have an example of an ambient metric that does not involve powers in ρ higher than k_0 we need to have $\frac{\partial^{(2k_0+2)}f}{\partial p^{(2k_0+2)}} \equiv 0$, i.e. the function f = f(x, p) defining the distribution must be a polynomial of order no higher than $2k_0 + 1$. Because of $c_3(x, p) \equiv 0$, this statement can be improved, if we want to have ambient metrics truncated at order $k_0 = 2$. Here *f* must be a polynomial of order no higher than $2k_0 + 2 = 6$, which is the case of examples of Leistner and PN.

Solutions analytic in ρ Nonanalytic in ρ solutions Poincaré-Einstein picture

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Polynomial solutions have G₂ holonomy

It is a matter of checking that the so obtained analytic in ρ Fefferman-Graham metrics generically have full G_2 holonomy. As an example we give a formula for a Fefferman-Graham full G_2 holonomy metric that truncates at order 4 in ρ :

$$\begin{split} &f = f_0 + f_1 \rho + f_2 \rho^2 + f_3 \rho^3 + f_4 \rho^4 + f_5 \rho^5 + f_6 \rho^6 + f_7 \rho^7 + f_8 \rho^8 + f_9 \rho^9, \\ &\omega^1 = dy - \rho dx, \quad \omega^2 = dz - (q^2 + t + bz) dx - \frac{\sqrt{2}}{2} q \omega^3, \quad \omega^3 = 2\sqrt{2} (dp - q dx), \quad \omega^4 = 3 dx, \\ &\omega^5 = \frac{\sqrt{2} b}{2} \omega^3 - 6 dq + 3(2bq + f_p) dx + \frac{4}{10} (9 f_{pp} + 4b^2) \omega^1, \\ & \bar{g}_{D_f} = 2 dt d(\rho t) + t^2 \left(2 \omega^1 \omega^5 - 2 \omega^2 \omega^4 + (\omega^3)^2 + A \cdot (\omega^1)^2 + 2B \cdot \omega^1 \omega^4 + C \cdot (\omega^4)^2 \right), \\ &A = \frac{63}{8} (f_8 + 9 \rho f_9) \rho^3 + \frac{27}{8} (f_6 + 7 \rho f_7 + 28 \rho^2 f_8 + 84 \rho^3 f_9) \rho^2 - \frac{9}{8} (f_4 + 5 \rho f_5 + 15 \rho^2 f_6 + 35 \rho^3 f_7 + 70 \rho^4 f_8 + 126 \rho^5 f_9) \rho \\ &B = -\frac{63}{256} f_9 \rho^4 - \frac{7}{64} (f_7 + 8 \rho f_8 + 36 \rho^2 f_9) \rho^3 + \frac{1}{16} (f_5 + 6 \rho f_6 + 21 \rho^2 f_7 + 56 \rho^3 f_8 + 126 \rho^4 f_9) \rho^2 - \frac{3}{20} (f_3 + 4 \rho f_4 + 10 \rho^2 f_5 + 20 \rho^3 f_6 + 35 \rho^4 f_7 + 56 \rho^5 f_8 + 84 \rho^6 f_9) \rho, \\ &C = \frac{7}{1152} (f_8 + 9 \rho f_9) \rho^4 + \frac{1}{360} (f_4 + 5 \rho f_5 + 15 \rho^2 f_6 + 35 \rho^3 f_7 + 70 \rho^4 f_8 + 126 \rho^5 f_9) \rho^2 + \frac{1}{45} (2b^2 - f_2 - 3\rho f_3 - 6\rho^2 f_4 - 10 \rho^3 f_5 - 15 \rho^4 f_6 - 21 \rho^5 f_7 - 28 \rho^6 f_8 - 36 \rho^7 f_9) \rho. \end{split}$$

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$$\begin{split} & f = f_0 + f_1 \rho + f_2 \rho^2 + f_3 \rho^3 + f_4 \rho^4 + f_5 \rho^5 + f_6 \rho^6 + f_7 \rho^7 + f_8 \rho^8 + f_9 \rho^9, \\ & \omega^1 = dy - \rho dx, \quad \omega^2 = dz - (q^2 + f + bz) dx - \frac{\sqrt{2}}{2} q \omega^3, \quad \omega^3 = 2\sqrt{2}(d\rho - q dx), \quad \omega^4 = 3dx, \\ & \omega^5 = \frac{\sqrt{2}b}{2} \omega^3 - 6dq + 3(2bq + f_\rho) dx + \frac{1}{10} (9f_{\rho\rho} + 4b^2) \omega^1, \\ & \tilde{g}_{D_f} = 2dtd(\rho t) + t^2 \left(2\omega^1 \omega^5 - 2\omega^2 \omega^4 + (\omega^3)^2 + A \cdot (\omega^1)^2 + 2B \cdot \omega^1 \omega^4 + C \cdot (\omega^4)^2 \right), \\ & A = \frac{63}{8} (f_8 + 9pf_9) \rho^3 + \frac{27}{8} (f_6 + 7pf_7 + 28\rho^2 f_8 + 84\rho^3 f_9) \rho^2 - \frac{9}{5} (f_4 + 5pf_5 + 15\rho^2 f_6 + 35\rho^3 f_7 + 70\rho^4 f_8 + 126\rho^5 f_9) \rho, \\ & B = -\frac{63}{256} f_9 \rho^4 - \frac{7}{64} (f_7 + 8pf_8 + 36\rho^2 f_9) \rho^3 + \frac{1}{16} (f_5 + 6pf_6 + 21\rho^2 f_7 + 56\rho^3 f_8 + 126\rho^4 f_9) \rho^2 - \frac{3}{20} (f_3 + 4pf_4 + 10\rho^2 f_5 + 20\rho^3 f_6 + 35\rho^4 f_7 + 56\rho^5 f_8 + 84\rho^6 f_9) \rho, \\ & C = \frac{7}{1152} (f_8 + 9pf_9) \rho^4 + \frac{1}{360} (f_4 + 5pf_5 + 15\rho^2 f_6 + 35\rho^3 f_7 + 70\rho^4 f_8 + 126\rho^5 f_9) \rho^2 + \frac{1}{45} (2b^2 - f_2 - 3pf_3 - 6\rho^2 f_4 - 10\rho^3 f_5 - 15\rho^4 f_6 - 21\rho^5 f_7 - 28\rho^6 f_8 - 36\rho^7 f_9) \rho. \end{split}$$

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$$\begin{split} f &= f_0 + f_1 \rho + f_2 \rho^2 + f_3 \rho^3 + f_4 \rho^4 + f_5 \rho^5 + f_6 \rho^6 + f_7 \rho^7 + f_8 \rho^8 + f_9 \rho^9, \\ \omega^1 &= dy - \rho dx, \quad \omega^2 &= dz - (q^2 + f + bz) dx - \frac{\sqrt{2}}{2} q \omega^3, \quad \omega^3 &= 2\sqrt{2} (dp - q dx), \quad \omega^4 &= 3dx, \\ \omega^5 &= \frac{\sqrt{2}b}{2} \omega^3 - 6dq + 3(2bq + f_p) dx + \frac{1}{10} (9f_{pp} + 4b^2) \omega^1, \\ \tilde{g}_{D_f} &= 2dtd(\rho t) + t^2 \left(2\omega^1 \omega^5 - 2\omega^2 \omega^4 + (\omega^3)^2 + A \cdot (\omega^1)^2 + 2B \cdot \omega^1 \omega^4 + C \cdot (\omega^4)^2 \right), \\ A &= \frac{63}{8} (f_8 + 9\rho f_9) \rho^3 + \frac{27}{8} (f_6 + 7\rho f_7 + 28\rho^2 f_8 + 84\rho^3 f_9) \rho^2 - \frac{9}{5} (f_4 + 5\rho f_5 + 15\rho^2 f_6 + 35\rho^3 f_7 + 70\rho^4 f_8 + 126\rho^5 f_9) \rho, \\ B &= -\frac{63}{256} f_9 \rho^4 - \frac{7}{64} (f_7 + 8\rho f_8 + 36\rho^2 f_9) \rho^3 + \frac{1}{16} (f_5 + 6\rho f_6 + 21\rho^2 f_7 + 56\rho^3 f_8 + 126\rho^4 f_9) \rho^2 - \frac{3}{20} (f_3 + 4\rho f_4 + 10\rho^2 f_5 + 20\rho^3 f_6 + 35\rho^4 f_7 + 56\rho^5 f_8 + 84\rho^6 f_9) \rho, \\ C &= \frac{7}{1152} (f_8 + 9\rho f_9) \rho^4 + \frac{1}{360} (f_4 + 5\rho f_5 + 15\rho^2 f_6 - 35\rho^3 f_7 + 70\rho^4 f_8 + 126\rho^5 f_9) \rho^2 + \frac{1}{45} (2b^2 - f_2 - 3\rho f_3 - 6\rho^2 f_4 - 10\rho^3 f_5 - 15\rho^4 f_6 - 21\rho^5 f_7 - 28\rho^6 f_8 - 36\rho^7 f_9) \rho. \end{split}$$

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$$\begin{split} f &= f_0 + f_1 p + f_2 p^2 + f_3 p^3 + f_4 p^4 + f_5 p^5 + f_6 p^6 + f_7 p^7 + f_8 p^8 + f_9 p^9, \\ \omega^1 &= dy - p dx, \quad \omega^2 &= dz - (q^2 + f + bz) dx - \frac{\sqrt{2}}{2} q \omega^3, \quad \omega^3 &= 2\sqrt{2} (dp - q dx), \quad \omega^4 &= 3 dx, \\ \omega^5 &= \frac{\sqrt{2b}}{2} \omega^3 - 6 dq + 3(2bq + f_p) dx + \frac{1}{10} (9f_{pp} + 4b^2) \omega^1, \\ \tilde{g}_{D_f} &= 2dtd(\rho t) + t^2 \left(2\omega^1 \omega^5 - 2\omega^2 \omega^4 + (\omega^3)^2 + A \cdot (\omega^1)^2 + 2B \cdot \omega^1 \omega^4 + C \cdot (\omega^4)^2 \right), \\ A &= \frac{63}{8} (f_8 + 9pf_9) \rho^3 + \frac{27}{8} (f_6 + 7pf_7 + 28p^2 f_8 + 84p^3 f_9) \rho^2 - \frac{9}{5} (f_4 + 5pf_5 + 15p^2 f_6 + 35p^3 f_7 + 70p^4 f_8 + 126p^5 f_9) \rho, \\ B &= -\frac{63}{256} f_9 \rho^4 - \frac{7}{764} (f_7 + 8pf_8 + 36p^2 f_9) \rho^3 + \frac{1}{16} (f_5 + 6pf_6 + 21p^2 f_7 + 56p^3 f_8 + 126p^4 f_9) \rho^2 - \frac{3}{20} (f_3 + 4pf_4 + 10p^2 f_5 + 20p^3 f_6 + 35p^4 f_7 + 56p^5 f_8 + 84p^6 f_9) \rho, \\ C &= \frac{7}{1152} (f_8 + 9pf_9) \rho^4 + \frac{1}{360} (f_4 + 5pf_5 + 15p^2 f_6 + 35p^3 f_7 + 70p^4 f_8 + 126p^5 f_9) \rho^2 + \frac{1}{45} (2b^2 - f_2 - 3pf_3 - 6p^2 f_4 - 10p^3 f_5 - 15p^4 f_6 - 21p^5 f_7 - 28p^6 f_8 - 36p^7 f_9) \rho. \end{split}$$

Solutions analytic in ρ Nonanalytic in ρ solutions Poincaré-Einstein picture

Plan

Ambient metrics and distributions Fefferman-Graham construction Conformal structures and Contan's and Contan'

- Conformal structures and Cartan's paper
- 2 The main theorem
 - An ansatz
 - The theorem
- Examples of explicit ambient metrics
 - Solutions analytic in ρ
 - Nonanalytic in ρ solutions
 - Poincaré-Einstein picture

Solutions analytic in ρ Nonanalytic in ρ solutions Poincaré-Einstein picture

Indicial exponents

To find *all*, and in particular nonanalytic in ρ , solutions to the system

$$LA = \frac{9}{40} f_{pppp}, \qquad LB = -\frac{1}{36} A_p + \frac{3}{40} f_{ppp},$$
$$LC = -\frac{1}{18} B_p + \frac{1}{324} A + \frac{1}{30} f_{pp} - \frac{2}{15} b^2,$$

we first observe that the two independent solutions to $L(\rho^k) = 0$ are ρ^0 and $\rho^{5/2}$. Thus, the most general solution to the above system can be obtained by the following series:

$$A = \sum_{k=1}^{\infty} a_k(x, p) \rho^k + \rho^{5/2} \sum_{k=0}^{\infty} \alpha_k(x, p) \rho^k,$$

$$B = \sum_{k=1}^{\infty} b_k(x, p) \rho^k + \rho^{5/2} \sum_{k=0}^{\infty} \beta_k(x, p) \rho^k,$$

$$C = \sum_{k=1}^{\infty} c_k(x, p) \rho^k + \rho^{5/2} \sum_{k=0}^{\infty} \gamma_k(x, p) \rho^k.$$

Solutions analytic in ρ Nonanalytic in ρ solutions Poincaré-Einstein picture

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Solutions analytic in ρ Nonanalytic in ρ solutions Poincaré-Einstein picture

General solutions

$$A = \sum_{k=1}^{\infty} \frac{3}{5} \cdot \frac{(2k-1)(2k-3)}{2^{2k}(2k)!} \cdot \frac{\partial^{(2k+2)}f}{\partial p^{(2k+2)}} \cdot \rho^k$$

 Ambient metrics and distributions
 Solutions analytic in p

 The main theorem
 Nonanalytic in p solutions

 Examples of explicit ambient metrics
 Poincaré-Einstein picture

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Ambient metrics and distributions The main theorem Examples of explicit ambient metrics Solutions analytic in solutions Poincaré-Einstein picture

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Ambient metrics and distributions The main theorem Examples of explicit ambient metrics Solutions analytic in polytocaré-Einstein picture

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Note that the analytic solutions are totally determined by the distribution, i.e. by the function f and the constant b.

 Ambient metrics and distributions
 Solutions analytic in ρ

 The main theorem
 Nonanalytic in ρ solutions

 Examples of explicit ambient metrics
 Poincaré-Einstein picture

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 Ambient metrics and distributions
 Solutions analytic in ρ

 The main theorem
 Nonanalytic in ρ solutions

 Examples of explicit ambient metrics
 Poincaré-Einstein picture

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 Ambient metrics and distributions
 Solutions analytic in ρ

 The main theorem
 Nonanalytic in ρ solutions

 Examples of explicit ambient metrics
 Poincaré-Einstein picture

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Ambient metrics and distributions
 Solutions analytic in ρ

 The main theorem
 Nonanalytic in ρ solutions

 Examples of explicit ambient metrics
 Poincaré-Einstein picture

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 Ambient metrics and distributions The main theorem
 Solutions analytic in ρ

 Examples of explicit ambient metrics
 Poincaré-Einstein picture

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Note that the analytic solutions are totally determined by the distribution, i.e. by the function *f* and the constant *b*. On the other hand, the nonalytic sloutions do not depend on a distribution at all!! They depend on α_0 , β_0 and γ_0 , which can be *arbitrary* functions of the variables *x* and *p*.

Solutions analytic in ρ Nonanalytic in ρ solutions Poincaré-Einstein picture

- If α₀ ≡ β₀ ≡ γ₀ ≡ 0 and for a randomly chosen *f* the holonomy of the corresponding ambient metric is equal to G₂ (Graham-Willse result).
- What about the holonomy of FG metrics corresponding to the solutions with nontrivial $\rho^{5/2+k}$ terms?
- Problems:
 - These solutions are only defined for $\rho \ge 0$.
 - They are only twice differentiable at $\rho = 0$.
 - Holonomy on a manifold with a boundary?
 - First calculate holonomy in the points where $\rho > 0, ...$

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Solutions analytic in ρ Nonanalytic in ρ solutions Poincaré-Einstein picture

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Solutions with $f \equiv 0$, b = 0

In case of a flat distribution, i.e. when $f \equiv 0$ and b = 0 the solutions are:

$$\begin{split} A &= \rho^{5/2} \sum_{k=0}^{\infty} 60 \cdot \frac{(k+2)(k+1)}{2^{2k}(2k+5)!} \cdot \frac{\partial^{2k} \alpha_0}{\partial \rho^{2k}} \cdot \rho^k, \\ B &= \rho^{5/2} \sum_{k=0}^{\infty} \frac{20}{3} \cdot \frac{(k+1)(k+2)}{2^{2k}(2k+5)!} \cdot \left(9 \frac{\partial^{2k} \beta_0}{\partial \rho^{2k}} - 2k \frac{\partial^{(2k-1)} \alpha_0}{\partial \rho^{(2k-1)}}\right) \cdot \rho^k, \\ C &= \rho^{5/2} \sum_{k=0}^{\infty} \frac{20}{27} \cdot \frac{(k+1)(k+2)}{2^{2k}(2k+5)!} \cdot \left(81 \frac{\partial^{2k} \gamma_0}{\partial \rho^{2k}} - 36k \frac{\partial^{(2k-1)} \beta_0}{\partial \rho^{(2k-1)}} + 2k(2k-1) \frac{\partial^{(2k-2)} \alpha_0}{\partial \rho^{(2k-2)}}\right) \cdot \rho^k, \end{split}$$

Solutions analytic in ρ Nonanalytic in ρ solutions Poincaré-Einstein picture

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Solutions analytic in ρ Nonanalytic in ρ solutions Poincaré-Einstein picture

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Solutions with $f \equiv 0$, b = 0

In case of a flat distribution, i.e. when $f \equiv 0$ and b = 0 the solutions are:

$$\begin{split} A &= \rho^{5/2} \sum_{k=0}^{\infty} 60 \cdot \frac{(k+2)(k+1)}{2^{2k}(2k+5)!} \cdot \frac{\partial^{2k} \alpha_0}{\partial \rho^{2k}} \cdot \rho^k, \\ B &= \rho^{5/2} \sum_{k=0}^{\infty} \frac{20}{3} \cdot \frac{(k+1)(k+2)}{2^{2k}(2k+5)!} \cdot \left(9 \frac{\partial^{2k} \beta_0}{\partial \rho^{2k}} - 2k \frac{\partial^{(2k-1)} \alpha_0}{\partial \rho^{(2k-1)}}\right) \cdot \rho^k, \\ C &= \rho^{5/2} \sum_{k=0}^{\infty} \frac{20}{27} \cdot \frac{(k+1)(k+2)}{2^{2k}(2k+5)!} \cdot \left(81 \frac{\partial^{2k} \gamma_0}{\partial \rho^{2k}} - 36k \frac{\partial^{(2k-1)} \beta_0}{\partial \rho^{(2k-1)}} + 2k(2k-1) \frac{\partial^{(2k-2)} \alpha_0}{\partial \rho^{(2k-2)}}\right) \cdot \rho^k, \end{split}$$

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and, as in the general case, they depend on three arbitrary functions $\alpha_0, \beta_0, \gamma_0$ of variables *x* and *p*.

As an illustration we discuss holonomy properties of the corresponding ambient metrics on a very simple example, in which we have made a particular choice of these 3 functionsons. We believe that the discussed behaviour is a typical one.

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Take

 $\begin{aligned} \alpha_0 = \beta(x) + p\alpha(x), \quad \beta_0 = f_0(x) + pf_1(x) + 252cp^2\alpha(x), \\ \gamma_0 = f_3(x) + pf_4(x) + \frac{1}{81}p^2(2268f_2(x) - \beta(x) + 18f_1(x)), \end{aligned}$

with c a real constant.

This gives the following solution:

$$\begin{split} A &= \left(252\rho\alpha(x) + \beta(x)\right)\rho^{5/2} \\ B &= (9c - 1)\alpha(x)\rho^{7/2} + \left(f_0(x) + \rho f_1(x) + 252c\rho^2\alpha(x)\right)\rho^{5/2} \\ C &= \left(f_2(x) + \left(\frac{1}{9} - 4c\right)\rho\alpha(x)\right)\rho^{7/2} + \\ &\left(f_3(x) + \rho f_4(x) + \frac{1}{81}\rho^2\left(18f_1(x) + 2268f_2(x) - \beta(x)\right)\right)\rho^{5/2} \end{split}$$

corresponding to the following Fefferman-Graham *family* of metrics for the *flat* conformal structure:

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- The family depends on *seven* arbitrary functions $\alpha = \alpha(x), \beta = \beta(x), f_0 = f_0(x), f_1 = f_1(x), \dots, f_4 = f_4(x)$ and a real constant *c*.
- \tilde{g} is an ambient metric for the flat conformal structure represented by a flat metric

 $g = 8(\mathrm{d}p - q\mathrm{d}x)^2 - 6(\mathrm{d}z - 2q\mathrm{d}p + q^2\mathrm{d}x)\mathrm{d}x - 12(\mathrm{d}y - p\mathrm{d}x)\mathrm{d}q.$

Note that ğ̃ is only two times differentiable at ρ = 0; the third derivative at ρ = 0 does not exist.

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- In general these metrics have full SO(4,3) holonomy!!!!
- Even if we put: $\beta(x) \equiv f_0(x) \equiv f_1(x) \equiv \cdots \equiv f_4(x) \equiv 0$, and $\alpha(x) \equiv 1$, the holonomy algebra behaves as this:
 - the curvature defines 6 independent components of the holonomy algebra
 - the first covariant derivative of the curvature produces next
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- Interestingly these metrics include, as special cases, metrics with full G₂ holonomy, which can not be extended to anything larger!!!
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Solutions analytic in ρ Nonanalytic in ρ solutions Poincaré-Einstein picture

- If we put: $\alpha(x) \equiv \beta(x) \equiv f_0(x) \equiv f_1(x) \equiv f_2(x) \equiv f_4(x) \equiv 0$, and c = 1/9. Then the holonomy algebra behaves as this:
 - the curvature defines 4 independent components of the holonomy algebra
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Solutions analytic in ρ Nonanalytic in ρ solutions Poincaré-Einstein picture

Plan

Ambient metrics and distributions Fefferman-Graham construction Conformal structures and Cartan's pairs

- 2 The main theorem
 - An ansatz
 - The theorem
- 3 Examples of explicit ambient metrics
 - Solutions analytic in ρ
 - Nonanalytic in ρ solutions
 - Poincaré-Einstein picture

Solutions analytic in ρ Nonanalytic in ρ solutions Poincaré-Einstein picture

Passing from ρ to r such that $\rho = r^2$

- The nonanalytic in ρ solutions have troubles at $\rho \leq 0$ because they are expessible in odd powers of $\sqrt{\rho}$.
- One can try to remedy the situation by passing to the coordinate *r* such that *ρ* = *r*².
- On doing this we first assume that r > 0, and bring the metric g̃ to the form

$$\begin{split} \tilde{g} &= 2dtd(r^2t) + t^2(8(dp - qdx)^2 - 6(dz - 2qdp + q^2dx)dx - 12(dy - pdx)dq + \\ (252p\alpha + \beta)r^5 \cdot (dy - pdx)^2 + \\ 6((9c - 1)\alpha r^7 + (f_0 + pf_1 + 252cp^2\alpha)r^5) \cdot (dy - pdx)dx + \\ 9((f_2 + (\frac{1}{9} - 4c)p\alpha)r^7 + (f_3 + pf_4 + \frac{1}{91}p^2(18f_1 + 2268f_2 - \beta))r^5) \cdot dx^2). \end{split}$$

This metric is regular and *Ricci flat* for all r ≠ 0, but because 2dtd(r²t) = 4rdrdt + 2r²dt², it is degenerate at r = 0.

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Solutions analytic in ρ Nonanalytic in ρ solutions Poincaré-Einstein picture

Poincaré-Einstein metrics in general

Given a normal form of a (4, 3)-signature ambient metric \tilde{g}_{D_f}

$$\begin{split} ilde{g}_{\mathcal{D}_f} =& 2 \mathrm{d} t \mathrm{d}(
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one can associate with it a (3,3) signature metric g_{PE} , called *a Poincaré-Einstein metric*, obtained in the following way:

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Poincaré-Einstein metrics in general

- Let ℝ⁶ be coordinatized by (*r*, *x*, *y*, *p*, *q*, *z*), and consider an open neigbourhood U₆ around a point with *r* ≠ 0 there.
- Imbedd \mathcal{U}_6 in \tilde{M} by $\iota : \mathcal{U}_6 \to \tilde{M}$, where ι is given by:

$$\iota(r, x, y, p, q, z) := \left(t = \frac{1}{r}, \rho = r^2, x, y, p, q, z\right).$$

• Pullback $ilde{g}_{\mathcal{D}_f}$ from $ilde{M}$ to \mathcal{U}_6 obtaining

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$$egin{aligned} g_{\mathsf{PE}} &:= \iota^*(ilde{g}) = \ & rac{1}{r^2} \Big(-2\mathrm{d} r^2 + g_{\mathcal{D}_f} + A\cdot(\omega^1)^2 + 2B\cdot\omega^1\omega^4 + C\cdot(\omega^4)^2 \Big). \end{aligned}$$

• g_{PE} is a (3,3)-signature metric everywhere except r = 0.

Solutions analytic in ρ Nonanalytic in ρ solutions Poincaré-Einstein picture

Poincaré-Einstein metrics in general

- Let ℝ⁶ be coordinatized by (*r*, *x*, *y*, *p*, *q*, *z*), and consider an open neigbourhood U₆ around a point with *r* ≠ 0 there.
- Imbedd \mathcal{U}_6 in \tilde{M} by $\iota : \mathcal{U}_6 \to \tilde{M}$, where ι is given by:

$$\iota(\mathbf{r},\mathbf{x},\mathbf{y},\mathbf{p},\mathbf{q},\mathbf{z}) := \left(t = \frac{1}{r}, \rho = r^2, \mathbf{x}, \mathbf{y}, \mathbf{p}, \mathbf{q}, \mathbf{z}\right).$$

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Solutions analytic in ρ Nonanalytic in ρ solutions Poincaré-Einstein picture

Poincaré-Einstein metrics in general

• The metric g_{PE} is Einstein,

 $Ric(g_{PE}) = \frac{5}{2}g_{PE},$

if and only if the functions A, B, C satisfy

$$LA = \frac{9}{40} f_{pppp}, \qquad LB = -\frac{1}{36} A_p + \frac{3}{40} f_{ppp}$$
$$LC = -\frac{1}{18} B_p + \frac{1}{324} A + \frac{1}{30} f_{pp} - \frac{2}{15} b^2,$$

i.e. iff they correspond to the Fefferman-Graham (Ricci flat) metric $\tilde{g}_{\mathcal{D}_f}$.

Note that in coordinate r the linear operator is

$$L = \frac{1}{2} \frac{\partial^2}{\partial r^2} - \frac{2}{r} \frac{\partial}{\partial r} - \frac{1}{8} \frac{\partial^2}{\partial p^2},$$

so now the indicial exponents are 0 and 5, and we have no troubles with r < 0 range.

Solutions analytic in ρ Nonanalytic in ρ solutions Poincaré-Einstein picture

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Solutions analytic in ρ Nonanalytic in ρ solutions Poincaré-Einstein picture

Nontrivial Poincare-Einstein metrics associated with flat conformal structure

For example: a Poincaré-Einstein metric corresponding to the Fefferman-Graham metric associated with the flat conformal structure discussed few slides ago is then given by:

$$\begin{split} \tilde{g} &= r^{-2} \Big(-2dr^2 + 8(dp - qdx)^2 - 6(dz - 2qdp + q^2dx)dx - 12(dy - pdx)dq + \\ (252p\alpha + \beta)r^5 \cdot (dy - pdx)^2 + 6((9c - 1)\alpha r^7 + (f_0 + pf_1 + 252cp^2\alpha)r^5) \cdot (dy - pdx)dx + \\ 9((f_2 + (\frac{1}{9} - 4c)p\alpha)r^7 + (f_3 + pf_4 + \frac{1}{81}p^2(18f_1 + 2268f_2 - \beta))r^5) \cdot dx^2) \Big). \end{split}$$

Solutions analytic in ρ Nonanalytic in ρ solutions Poincaré-Einstein picture

Nonanalytic PE metrics and flat structure

The most general one associated with $f \equiv 0$ and b = 0 is:

$$\tilde{g} = r^{-2} \Big(-2dr^2 + 8(dp - qdx)^2 - 6(dz - 2qdp + q^2dx)dx - 12(dy - pdx)dq + A(dy - pdx)^2 + 6B(dy - pdx)dx + 9Cdx^2 \Big)$$

$$\begin{split} &A = \rho^{5/2} \sum_{k=0}^{\infty} 60 \cdot \frac{(k+2)(k+1)}{2^{2k}(2k+5)!} \cdot \frac{\partial^{2k} \alpha_{\mathbf{0}}}{\partial \rho^{2k}} \cdot \rho^{k}, \\ &B = \rho^{5/2} \sum_{k=0}^{\infty} \frac{20}{3} \cdot \frac{(k+1)(k+2)}{2^{2k}(2k+5)!} \cdot \left(9 \frac{\partial^{2k} \beta_{\mathbf{0}}}{\partial \rho^{2k}} - 2k \frac{\partial^{(2k-1)} \alpha_{\mathbf{0}}}{\partial \rho^{(2k-1)}}\right) \cdot \rho^{k}, \\ &C = \rho^{5/2} \sum_{k=0}^{\infty} \frac{20}{27} \cdot \frac{(k+1)(k+2)}{2^{2k}(2k+5)!} \cdot \left(81 \frac{\partial^{2k} \gamma_{\mathbf{0}}}{\partial \rho^{2k}} - 36k \frac{\partial^{(2k-1)} \beta_{\mathbf{0}}}{\partial \rho^{(2k-1)}} + 2k(2k-1) \frac{\partial^{(2k-2)} \alpha_{\mathbf{0}}}{\partial \rho^{(2k-2)}}\right) \cdot \rho^{k}. \end{split}$$

This can be truncated at *any* half-integer order of ρ , starting at $\rho^{5/2}$ by choosing $\alpha_0, \beta_0, \gamma_0$ as polynomials in *p* of an apropriate order.

Solutions analytic in ρ Nonanalytic in ρ solutions Poincaré-Einstein picture

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TT tensors

Nonanalytic in ρ solutions Poincaré-Einstein picture

The explicit solutions for the PE metrics can be used to calculate the trace-free, divergence-free tensors for each of the conformal structure $[g_{D_f}]$:

$$TT = \alpha_0(x,p)(\mathrm{d}y - p\mathrm{d}x)^2 + 6\beta_0(x,p)(\mathrm{d}y - p\mathrm{d}x)\mathrm{d}x + 9\gamma_0(x,p)\mathrm{d}x^2.$$

For all choices of the free functions α_0, β_0 and γ_0 they are trace-free and divergence-free in the metric

$$g_{\mathcal{D}_{f}} = 8 \left(dp - q dx \right)^{2} - 6 \left(dz - 2q dp + (q^{2} - f - bz) dx \right) dx - 2 \left(dy - p dx \right) \left(6 dq - 2b dp - (\frac{2}{5}b^{2} + \frac{9}{10}f_{pp})(dy - p dx) - (4bq + 3f_{p}) dx \right),$$

associated with $z' = (y'')^2 + f(x, y') + bz$.

Solutions analytic in ρ Nonanalytic in ρ solutions Poincaré-Einstein picture

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For all choices of the free functions α_0, β_0 and γ_0 they are trace-free and divergence-free in the metric

$$g_{\mathcal{D}_f} = 8\left(\mathrm{d}p - q\mathrm{d}x\right)^2 - 6\left(\mathrm{d}z - 2q\mathrm{d}p + (q^2 - f - bz)\mathrm{d}x\right)\mathrm{d}x - 2\left(\mathrm{d}y - p\mathrm{d}x\right)\left(6\mathrm{d}q - 2b\mathrm{d}p - (\frac{2}{5}b^2 + \frac{9}{10}f_{\rho\rho})(\mathrm{d}y - p\mathrm{d}x) - (4bq + 3f_{\rho})\mathrm{d}x\right),$$

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Solutions analytic in ρ Nonanalytic in ρ solutions Poincaré-Einstein picture

THANK YOU!

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37/37