

**Continuous and discrete Neumann systems  
on Stiefel varieties**

Yuri Fëdorov

In collaboration with Bozidar Jovanovic (Belgrade, Serbia)  
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- **Classical Neumann system** on  $T^* S^{n-1} = \left\{ \langle q, q \rangle = 1, \langle q, p \rangle = 0 \right\} \subset \mathbb{R}^{2n}$ .

$$H = \frac{1}{2}\langle p, p \rangle + \frac{1}{2}\langle Aq, q \rangle, \quad A = \text{diag}(a_1, \dots, a_n),$$

Hamilton equations (with respect to the Dirac bracket on  $\mathbb{R}^{2n}$ )

$$\dot{q} = p, \quad \dot{p} = -Aq + \nu q, \quad \nu = \langle p, p \rangle - \langle q, Aq \rangle$$

Big ( $n \times n$ ) Lax pair by Moser, Small ( $2 \times 2$ ) Lax pair by Mumford:

$$\begin{aligned} \dot{L}(\lambda) &= [L(\lambda), A(\lambda)], \\ L(\lambda) &= \begin{pmatrix} \sum_{i=1}^n \frac{q_i p_i}{\lambda - a_i} & \sum_{i=1}^n \frac{q_i^2}{\lambda - a_i} \\ -1 - \sum_{i=1}^n \frac{p_i^2}{\lambda - a_i} & -\sum_{i=1}^n \frac{q_i p_i}{\lambda - a_i} \end{pmatrix} = Y + \sum_{i=1}^n \frac{N_i}{\lambda - a_i}, \quad A(\lambda) = \begin{pmatrix} 0 & 1 \\ \lambda + \nu(p, q) & 0 \end{pmatrix} \end{aligned}$$

First integrals

$$F_i = q_i^2 + \sum_{j \neq i} \frac{(p_i q_j - p_j q_i)^2}{a_j - a_i}, \quad i = 1, \dots, n$$

Real generic tori  $\mathbb{T}^{n-1} \mapsto$  complex tori  $\mathbb{T}_{\mathbb{C}}^{n-1} =$  open subsets of  $\text{Jac}(\Gamma)$ ,  
 $\Gamma$  being the hyperelliptic genus  $n - 1$  spectral curve of  $L(\lambda)$ .

- **Generalizations to higher-order potentials**

(S. Wojciechowski, O. Bogoyavlenski, P. Saksida)

- **Discrete Neumann system** (Bäcklund transformation)  $\mathcal{B}_{\lambda^*} : (p, q) \mapsto (\tilde{p}, \tilde{q})$   
(Veselov, Kuznetsov, Vanhaecke, Suris)

$$\begin{aligned}\tilde{q} &= A^{-1/2}(\lambda^*)(\beta q + p), \quad \tilde{p} = -A^{1/2}(\lambda^*)q + A^{-1/2}(\lambda^*)(\beta^2 q + \beta p), \\ A(\lambda^*) &= \lambda^* \mathbf{I} - A \quad \beta = \langle \tilde{q}, A^{1/2}(\lambda^*)q \rangle,\end{aligned}$$

$\lambda^* \in \mathbb{C}$  being arbitrary step parameter.  $\lambda^* \rightarrow \infty$  gives the continuous limit.

Alternative form

$$p = A^{1/2}(\lambda^*)\tilde{q} - \beta q, \quad \tilde{p} = -A^{1/2}(\lambda^*)q + \beta\tilde{q}.$$

**Theorem 1.** 1). Up to the action of the group of reflections  $(p_i, q_i) \rightarrow (-p_i, -q_i)$ , the map  $\mathcal{B}_{\lambda^*}$  is equivalent to the discrete Lax pair

$$\tilde{L}(\lambda)M(\lambda|\lambda^*) = M(\lambda|\lambda^*)L(\lambda), \quad M(\lambda|\lambda^*) = \begin{pmatrix} -\beta & 1 \\ \lambda - \lambda^* + \beta^2 & -\beta \end{pmatrix}.$$

2) (Veselov)  $\mathcal{B}_{\lambda^*}$  is given by a shift on  $\text{Jac}(\Gamma)$  by  $\mathcal{A}(\lambda^*) - \mathcal{A}(\infty)$

**The Stiefel variety**  $V(n, r) = SO(n)/SO(n - r)$ , ( $r < n$ ) the set of  $n \times r$  matrices

$$X = (e_1 \cdots e_r), \quad e_s \in \mathbb{R}^n, \quad X^T X = \mathbf{I}_r,$$

The cotangent bundle  $T^*V(n, r)$ , the set of  $n \times r$  pairs  $(X, P)$ ,  $P = (p_1 \cdots p_r)$ ,  $p \in \mathbb{R}^n$

$$X^T X = \mathbf{I}_r, \quad X^T P + P^T X = 0,$$

giving  $r(r + 1)$  independent scalar constraints.

- Reduction to the oriented Grassmannian  $G(r, n) = V(n, r)/SO(r)$

The canonical symplectic structure  $\omega$  on  $T^*V(n, r)$  is the restriction of the canonical 2-form in the ambient space  $T^*M_{n,r}(\mathbb{R})$ ,

$$\omega_0 = \sum_{i=1}^n \sum_{s=1}^r dp_s^i \wedge de_s^i$$

Given  $H(X, P)$ , the Hamiltonian equations with respect to the Dirac Poisson structure or involve  $r \times r$  symmetric matrix Lagrange multipliers  $\Lambda, \Pi$ ,

$$\dot{X} = \frac{\partial H}{\partial P} - X\Pi, \quad \dot{P} = -\frac{\partial H}{\partial X} + X\Lambda + P\Pi$$

- **Two natural  $SO(n)$ -invariant metrics** on  $V(n, r)$ :

The Euclidean metric

$$L_E(X, \dot{X}) = \frac{1}{2} \text{Tr}(\dot{X}^T \dot{X}) \mapsto H_E(X, P) = \frac{1}{2} \text{Tr}(P^T P)$$

The normal metric

$$H_0(X, P) = \frac{1}{2} \langle \Phi, \Phi \rangle = \frac{1}{2} \text{Tr}(P^T P) - \frac{1}{2} \text{Tr}((X^T P)^2)$$

*There are other  $SO(n)$ -invariant metrics (....)*

- The potential

$$V = \frac{1}{2} \text{Tr}(X^T A X) = \frac{1}{2} \sum_{i=1}^r (e_i, A e_i)$$

- The Neumann system with the normal metric

$$\dot{X} = P - X P^T X, \quad \dot{P} = -A X - X P^T P + P X^T P + X X^T A X,$$

- The Neumann system with the Euclidean metric

$$\dot{X} = P, \quad \dot{P} = -A X - X P^T P + X X^T A X$$

**Theorem 2.** *The above equations admit the same  $n \times n$  matrix Lax representation (by Reyman and Semonov-Tian-Shanski)*

$$\frac{d}{dt} \mathcal{L}(\lambda) = [\mathcal{A}(\lambda), \mathcal{L}(\lambda)],$$

$$\mathcal{L}(\lambda) = \lambda M + XX^T - \lambda^2 A, \quad M = PX^T - X P^T, \quad \mathcal{A}(\lambda) = \Phi - \lambda A,$$

This yields commuting (with respect to  $\omega$ ) integrals

$$\mathfrak{F} = \left\{ \text{Tr} (\lambda(PX^T - X P^T) + XX^T - \lambda^2 A)^k \mid k = 1, \dots, n, \lambda \in \mathbb{R} \right\}$$

**Theorem 3.** *If all the eigenvalues of  $A$  are distinct, then the Neumann systems are completely integrable in the non-commutative sense with the set of integrals  $\mathfrak{F}$  and by the  $r(r-1)/2$  components of the  $SO(r)$ -momentum mapping  $\boxed{\Psi(X, P) = X^T P - P^T X}$ . The generic motions of the system are quasi-periodic over the isotropic tori of dimension*

$$d = \frac{1}{2} \left( 2r(n-r) + \frac{r(r-1)}{2} - \left[ \frac{r}{2} \right] \right) + \left[ \frac{r}{2} \right], \quad d < \frac{1}{2} \dim T^*V(n, r), \quad (r > 1)$$

## Bi-Hamiltonian description

$$\begin{aligned}\Theta : T^*V(n, r) &\mapsto gl(n, \mathbb{R}) = so(n) + Symm(n) \\ (X, P) &\mapsto (XP^T - P^T X) + XX^T\end{aligned}$$

On  $gl^*(n)$ , the pair of compatible Poisson brackets given by Poisson tensors (Bolsinov)

$$\begin{aligned}\Lambda_1(\xi + \eta, \zeta + \theta)|_x &= \langle x, [\xi, \zeta] + [\xi, \theta] + [\eta, \zeta] \rangle, \\ \Lambda_2(\xi + \eta, \zeta + \theta)|_x &= \langle x - A, [\xi + \eta, \zeta + \theta] \rangle,\end{aligned}$$

$x \in gl^*(n)$ ,  $\xi, \zeta \in so(n)$ ,  $\eta, \theta \in Symm(n)$

The Casimirs of  $\Lambda_1 + \lambda\Lambda_2$ ,  $\lambda \in \mathbb{P}^1$  appear as integrals of the Lax matrix  $P + \lambda M - \lambda^2 A$  (integrable tops)

Dim of generic  $\mathcal{S}$  in  $(gl^*(n), \Lambda_1) = n^2 - n$ , which is less than

$$\text{Dim of generic } \mathcal{S} \text{ in } \Theta^*T^*V(n, r) = 2r(n - r) + \frac{r(r - 1)}{2} - \left[ \frac{r}{2} \right]$$

The image of  $\Theta(T^*V(n, r))$  is the union of *singular* symplectic leaves of  $(gl^*(n), \Lambda_1)$ .

**Theorem 4.** *The set of functions  $\mathfrak{F}$  is complete on generic symplectic leaves in  $\Theta(T^*V(n, r)) \subset (gl^*(n), \Lambda_1)$ .*

Marsden–Weinstein reduction to Grassmannian  $G(n, r)$ ,  
the quotient space of  $V(n, r)$  by the right  $SO(r)$ -action.

Assume  $\Psi = X^T P - P^T X = 0$ .

The reduced phase space  $\Psi^{-1}(0)/SO(r)$  is symplectomorphic to the cotangent bundle  $T^*G(n, r)$  with a canonical symplectic structure.

**Theorem 5.** *If all eigenvalues of  $A$  are distinct, then the reduced Neumann system on  $T^*G(n, r)$  is completely integrable in the Liouville sense by means of the integrals  $\mathfrak{F}$ .  
The dimension of generic invariant tori =  $\frac{1}{2}$  dimension of  $T^*G(n, r)$ .*

- "Small"  $(2r \times 2r)$  matrix Lax representation generalizing the Mumford Lax pair (a modification of that in an unpublished manuscript by S. Kapustin)

**Theorem 6.** *Up to the action of the discrete group generated by reflections  $(X, P) \longmapsto (\pm X, \pm P)$ , the Neumann flows are equivalent to the Lax pair with*

$$\frac{d}{dt}L(\lambda) = [L(\lambda), A(\lambda)], \quad \lambda \in \mathbb{R},$$

$$L(\lambda) = \begin{pmatrix} -X^T(\lambda\mathbf{I}_n - A)^{-1}P & -X^T(\lambda\mathbf{I}_n - A)^{-1}X \\ \mathbf{I}_r + P^T(\lambda\mathbf{I}_n - A)^{-1}P & P^T(\lambda\mathbf{I}_n - A)^{-1}X \end{pmatrix} = \begin{pmatrix} 0 & 0 \\ \mathbf{I}_r & 0 \end{pmatrix} + \sum_i \frac{\mathcal{N}_i}{\lambda - a_i},$$

where for the systems with the normal, respectively Euclidean, metric,

$$A(\lambda) = \begin{pmatrix} X^T P & \mathbf{I}_r \\ \Lambda - \lambda\mathbf{I}_r & -P^T X \end{pmatrix}, \quad \text{respectively,} \quad A(\lambda) = \begin{pmatrix} 0 & \mathbf{I}_r \\ \Lambda - \lambda\mathbf{I}_r & 0 \end{pmatrix},$$

$$\Lambda = X^T A X - P^T P.$$

( compare with  $2 \times 2$  Lax matrix for the classical Neumann system

$$L(\lambda) = \begin{pmatrix} -q^T(\lambda\mathbf{I}_n - A)^{-1}p & -q^T(\lambda\mathbf{I}_n - A)^{-1}q \\ 1 + p^T(\lambda\mathbf{I}_n - A)^{-1}p & p^T(\lambda\mathbf{I}_n - A)^{-1}q \end{pmatrix}$$

It provides integrals  $\mathfrak{F}$ , as well as Casimirs  $\text{Tr } \Psi^k$ .

- What are generic complex tori ?

The spectral curve  $\mathcal{S}$  of  $L(\lambda)$

$$F(\lambda, w) = |(\lambda - a_1) \cdots (\lambda - a_n) \mathcal{L}(\lambda) - w \mathbf{I}_n| = 0$$

- 1). The regularized spectral curve  $\mathcal{S}'$  has genus  $g > d$ ,  
infinite points  $\infty_1, \dots, \infty_r$  and admits the involution  $\sigma : (\lambda, w) \rightarrow (\lambda, -w)$ ;  
Let  $\mathcal{C} = \mathcal{S}'/\sigma$ , "small" curve of genus  $g_0$ .
- 2).  $\sigma$  extends to  $\text{Jac}(\mathcal{S}') = \mathbb{C}^g/\Lambda \cong \text{Jac}(\mathcal{C}) \oplus \text{Prym}(\mathcal{S}'/\sigma)$   
Dym.  $(\text{Prym}(\mathcal{S}'/\sigma)) = g - g_0 < d$ ;
- 3).  $\exists$  precisely  $[r/2]$  meromorphic differentials  $\Omega_i$  with pairs of simple poles at  $\infty$   
such that  $\sigma^* \Omega_i = -\Omega_i$

$$\text{Generalized Jacobian } \widetilde{\text{Jac}}(\mathcal{S}', \Omega_i) = \mathbb{C}^{g+[r/2]}/\tilde{\Lambda} \cong \text{Jac}(\mathcal{S}') \times \underbrace{\mathbb{C}^* \times \cdots \times \mathbb{C}^*}_{[r/2]}.$$

$\sigma$  extends also to  $\widetilde{\text{Jac}}(\mathcal{S}', \Omega_i)$ .

- Theorem 7.** 1). Generic complex tori of the Neumann system on  $T^*V(n, r)$   
are open subsets of  $\widetilde{\text{Prym}}(\mathcal{S}'/\sigma, \Omega_i)$ .
- 2). Generic complex tori of the Neumann system on  $T^*G(n, r)$   
are open subsets of  $\text{Prym}(\mathcal{S}'/\sigma)$ .

## Some integrable generalizations.

- The Neumann system on the Grassmannian  $G(n, r)$  with *quartic* potential:

$$H(X, P) = \frac{1}{2} \operatorname{Tr}(P^T P) + \operatorname{Tr}(X^T A^2 X) - \operatorname{Tr}(X^T A X X^T A X),$$

which for  $r = 1$ , takes the form

$$H = \frac{1}{2}(p, p) + \sum_{i=1}^n a_i^2 e_i^2 - \left( \sum_{i=1}^n a_i e_i^2 \right)^2$$

(S. Wojciechowski, P. Saksida).

- Neumann Systems on Complex Stiefel Manifolds  $W(n, r) \cong U(n)/U(n - r)$ ,  
the set of  $n \times r$  matrices  $Z \in M_{n,r}(\mathbb{C})$  satisfying  $\bar{Z}^T Z = \mathbf{I}_r$ .

The real Stiefel manifold  $V(n, r)$  is a submanifold of  $W(n, r)$  given by  $Z = \bar{Z}$ .

The Euler-Lagrange equations with multipliers

$$\begin{aligned}\ddot{Z} &= -AZ + Z\Lambda, & \ddot{\bar{Z}} &= -A\bar{Z} + \bar{Z}\bar{\Lambda}, \\ \Lambda &= \bar{Z}^T AZ - \dot{\bar{Z}}^T \dot{Z} = \bar{\Lambda}^T.\end{aligned}$$

Matrix  $2r \times 2r$  Lax representation

$$\begin{aligned}\frac{d}{dt} \mathcal{L}^*(\lambda) &= [\mathcal{L}^*(\lambda), \mathcal{A}^*(\lambda)], \\ \mathcal{L}^*(\lambda) &= \begin{pmatrix} -Z^T(\lambda\mathbf{I}_n - A)^{-1}\dot{\bar{Z}} & -Z^T(\mathbf{I}_n - \lambda A)^{-1}\bar{Z} \\ \mathbf{I}_r + \dot{Z}^T(\mathbf{I}_n - \lambda A)^{-1}\dot{\bar{Z}} & \dot{Z}^T(\lambda\mathbf{I}_n - A)^{-1}\bar{Z} \end{pmatrix}, \quad \mathcal{A}^*(\lambda) = \begin{pmatrix} 0 & \mathbf{I}_r \\ \bar{\Lambda} - \lambda\mathbf{I}_r & 0 \end{pmatrix},\end{aligned}$$

- **Integrable discretization of the Neumann systems on  $V(n, r)$ ,**  
a family of Bäcklund transformations  $\mathcal{B}_{\lambda^*} : (X, P) \mapsto (\tilde{X}, \tilde{P})$

$$\begin{aligned} P &= A^{1/2}(\lambda^*) \tilde{X} - X \Gamma(\lambda^*), \\ \tilde{P} &= -A^{1/2}(\lambda^*) X + \tilde{X} \Gamma(\lambda^*), \\ A(\lambda) &= \lambda \mathbf{I}_n - A, \quad \Gamma(\lambda^*) = \frac{1}{2} \left( \tilde{X}^T A^{1/2}(\lambda^*) X + X^T A^{1/2}(\lambda^*) \tilde{X} \right). \end{aligned}$$

The alternative form (discrete Lagrange equations on  $V(n, r)$ ) ( $\lambda^* = 0$ , Veselov–Moser)

$$X + \tilde{\tilde{X}} = A^{-1/2}(\lambda^*) \tilde{X} B, \quad B = \frac{1}{2} \left( \tilde{X}^T A^{1/2}(\lambda^*) (X + \tilde{\tilde{X}}) + (X + \tilde{\tilde{X}})^T A^{1/2}(\lambda^*) \tilde{X} \right)$$

For generic  $\lambda^*$ , the *complex* map  $\mathcal{B}_{\lambda^*}$  is  $2^r$ -valued.

*Why the integrals of the continuous system are preserved?*

**Theorem 8.** *The discrete Neumann system is equivalent to the intertwining  $2r \times 2r$  matrix relation*

$$\begin{aligned} \tilde{L}(\lambda) M(\lambda|\lambda^*) &= M(\lambda|\lambda^*) L(\lambda), \\ L(\lambda) &= \begin{pmatrix} X^T (\lambda \mathbf{I}_n - A)^{-1} P & X^T (\lambda \mathbf{I}_n - A)^{-1} X \\ \mathbf{I}_r - P^T (\lambda \mathbf{I}_n - A)^{-1} P & -P^T (\lambda \mathbf{I}_n - A)^{-1} X \end{pmatrix} \text{ (as in the cont. case)}, \\ M(\lambda|\lambda^*) &= \begin{pmatrix} -\Gamma(\lambda^*) & \mathbf{I}_r \\ (\lambda - \lambda^*) \mathbf{I}_r + \Gamma^2(\lambda^*) & -\Gamma(\lambda^*) \end{pmatrix}, \end{aligned}$$

*How the map is described on the invariant tori ?*

**Theorem 9.** *The  $2^r$ -valued map  $\mathcal{B}_{\lambda^*}$  is given by translations by one of the vectors in  $\text{Prym}(\mathcal{S}'/\sigma, \Omega)$*

$$\tilde{\mathcal{A}}(\lambda_i^*) - \tilde{\mathcal{A}}(\infty_j).$$