

Conformal actions, Kummer tables and hypergeometric-type functions

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How the action of the symmetry group $SO(6, \mathbb{C})$ and the choice of a nice set of parameters help one to understand and present logically a whole bunch of special functions.

History: early beginnings

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WikipediaTM says

The term „hypergeometric series” was first used by **John Wallis** in his 1655 book *Arithmetica Infinitorum*. Hypergeometric series were studied by **Leonhard Euler**, but the first full systematic treatment was given by **Carl Friedrich Gauss (1813)**. Studies in the nineteenth century included those of **Ernst Kummer (1836)**, and the fundamental characterisation by **Bernhard Riemann** of the hypergeometric function by means of the differential equation it satisfies. Riemann showed that the second-order differential equation for ${}_2F_1(z)$, examined in the complex plane, could be characterised (on the Riemann sphere) by its three regular singularities.

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Beware

I will try to follow the historical path on which the subject was being discovered.

Euler Gamma function

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- Euler integral of the first kind

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$$\frac{1}{\Gamma(z)} = \frac{\sin \pi z}{\pi} \Gamma(1-z) \quad (4)$$

Hypergeometric series

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Hypergeometric series of type ${}_2F_1(a, b; c; z)$

$${}_2F_1(a, b; c; z) = \sum_{n=0}^{\infty} \frac{(a)_n (b)_n z^n}{(c)_n n!} \quad (5)$$

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$$(a)_n = \prod_{k=0}^{n-1} (a+k) = a(a+1) \cdot \dots \cdot (a+n-1) \quad (6)$$

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Examples

$$(1)_n = n! \quad (-k)_n = 0, \quad n \geq k \quad (c)_n = \frac{\Gamma(c+n)}{\Gamma(c)} \quad (7)$$

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- Due to problems with the denominator one should introduce

$${}_2\mathbf{F}_1(a, b; c; z) = \frac{{}_2F_1(a, b; c; z)}{\Gamma(c)} = \sum_{n=0}^{\infty} \frac{(a)_n (b)_n z^n}{\Gamma(c+n)n!} \quad (6)$$

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- In fact, it is a subclass of the confluent function ${}_1F_1$
- It is classically known as the modified Bessel function,

$$I_\alpha(w) = \left(\frac{w}{2}\right)^\alpha {}_0F_1\left(\alpha + 1; \frac{w^2}{4}\right) \quad (9)$$

Hypergeometric equation

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Hypergeometric operator $\mathcal{F}(a, b; c; z, \partial_z)$

$$\mathcal{F}(a, b; c; z, \partial_z) = z(1-z)\partial_z^2 + (c - (a+b+1)z)\partial_z - ab \quad (10)$$

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- The operator has three regular singular points on the Riemann sphere 0, 1 and ∞
- Parameters a , b solve the index equation for $z = \infty$. The indices at 0 and 1 are respectively $1 - c$ and $c - a - b$

A nightmarish zoo from Abramowitz & Stegun

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15.2.10

$$(c-a)F(a-1, b; c; z) + (2a-c-az+bz)F(a, b; c; z) + a(z-1)F(a+1, b; c; z) = 0$$

15.2.11

$$(c-b)F(a, b-1; c; z) + (2b-c-bz+az)F(a, b; c; z) + b(z-1)F(a, b+1; c; z) = 0$$

15.2.12

$$c(c-1)(z-1)F(a, b; c-1; z) + c[c(c-1)-(2c-a-b-1)z]F(a, b; c; z) + (c-a)(c-b)zF(a, b; c+1; z) = 0$$

15.2.13

$$[c-2a-(b-a)z]F(a, b; c; z) + a(1-z)F(a+1, b; c; z) - (c-a)F(a-1, b; c; z) = 0$$

15.2.14

$$(b-a)F(a, b; c; z) + aF(a+1, b; c; z) - bF(a, b+1; c; z) = 0$$

15.2.15

$$(c-a-b)F(a, b; c; z) + a(1-z)F(a+1, b; c; z) - (c-b)F(a, b-1; c; z) = 0$$

15.2.16

$$[c-2b+(b-a)z]F(a, b; c; z) + b(1-z)F(a, b+1; c; z) - (c-b)F(a, b-1; c; z) = 0$$

15.2.23

$$c[b-(c-a)z]F(a, b; c; z) - bc(1-z)F(a, b+1; c; z) + (c-a)(c-b)zF(a, b; c+1; z) = 0$$

15.2.24

$$(c-b-1)F(a, b; c; z) + bF(a, b+1; c; z) - (c-1)F(a, b; c-1; z) = 0$$

15.2.25

$$c(1-z)F(a, b; c; z) - cF(a, b-1; c; z) + (c-a)zF(a, b; c+1; z) = 0$$

15.2.26

$$[b-1-(c-a-1)z]F(a, b; c; z) + (c-b)F(a, b-1; c; z) - (c-1)(1-z)F(a, b; c-1; z) = 0$$

15.2.27

$$c[c-1-(2c-a-b-1)z]F(a, b; c; z) + (c-a)(c-b)zF(a, b; c+1; z) - c(c-1)(1-z)F(a, b; c-1; z) = 0$$

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- $\mu = b - a$, the difference between the indices at infinity
- Inverse relations are

$$a = \frac{1}{2}(1 + \alpha + \beta + \mu) \quad b = \frac{1}{2}(1 + \alpha + \beta - \mu) \quad c = 1 + \alpha \quad (11)$$

Hypergeometric operator $\mathcal{F}_{\alpha,\beta,\mu}(z, \partial_z)$

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Once we employ Lie-algebraic parameters into action we get

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$$\mathcal{F}_{\alpha, \beta, \mu} = z(1-z)\partial_z^2 + ((1+\alpha)(1-z) - (1+\beta)z)\partial_z - \frac{(1+\alpha+\beta)^2 - \mu^2}{4} \quad (12)$$

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- Its symmetries are becoming a lot more visible
- We will see that the parameters really do have Lie-algebraic interpretation

$\mathcal{F}_{\alpha, \beta, \mu}(z, \partial_z)$ in balanced form

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Balanced form and Schrödinger form

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It is known that any ODE of our type can be shown in the form of a Schrödinger operator by simple substitutions. Similarly a so-called balanced form can always be obtained

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Canonical form of the hypergeometric case

$$\begin{aligned} \mathcal{F}_{\alpha,\beta,\mu}(z, \partial_z) &= \\ &= z^{-\alpha}(1-z)^{-\beta} \partial_z z^{\alpha+1}(1-z)^{\beta+1} \partial_z - \frac{(1+\alpha+\beta)^2 - \mu^2}{4} \end{aligned} \quad (13)$$

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Balanced form of the hypergeometric operator

$$\begin{aligned} z^{\frac{\alpha}{2}}(1-z)^{\frac{\beta}{2}} \mathcal{F}_{\alpha,\beta,\mu}(z, \partial_z) z^{-\frac{\alpha}{2}}(1-z)^{-\frac{\beta}{2}} &= \\ &= \partial_z z(1-z) \partial_z - \frac{\alpha^2}{4z} - \frac{\beta^2}{4(1-z)} - \frac{1-\mu^2}{4} \end{aligned} \quad (13)$$

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Notice!

The symmetries are becoming **even** more v-i-s-i-b-l-e.

A less nightmarish zoo from JD – Kummer table

$w = z :$	$(-z)^{-\alpha}(z-1)^{-\beta}$ $(z-1)^{-\beta}$ $(-z)^{-\alpha}$	$\mathcal{F}_{\alpha,\beta,\mu}(z, \partial_z),$ $\mathcal{F}_{-\alpha,-\beta,\mu}(z, \partial_z)$ $\mathcal{F}_{\alpha,-\beta,-\mu}(z, \partial_z)$ $\mathcal{F}_{-\alpha,\beta,-\mu}(z, \partial_z)$	$(-z)^\alpha(z-1)^\beta$ $(z-1)^\beta,$ $(-z)^\alpha;$
$w = 1 - z :$	$(z-1)^{-\alpha}(-z)^{-\beta}$ $(z-1)^{-\alpha}$ $(-z)^{-\beta}$	$\mathcal{F}_{\beta,\alpha,\mu}(z, \partial_z),$ $\mathcal{F}_{-\beta,-\alpha,\mu}(z, \partial_z)$ $\mathcal{F}_{\beta,-\alpha,-\mu}(z, \partial_z)$ $\mathcal{F}_{-\beta,\alpha,-\mu}(z, \partial_z)$	$(z-1)^\alpha(-z)^\beta,$ $(z-1)^\alpha,$ $(-z)^\beta;$
$w = \frac{1}{z} :$	$(-z)^{\frac{1}{2}(\alpha+\beta+\mu+1)}$ $(-z)^{\frac{1}{2}(\alpha+\beta-\mu+1)}(z-1)^{-\beta}$ $(-z)^{\frac{1}{2}(\alpha+\beta+\mu+1)}(z-1)^{-\beta}$ $(-z)^{\frac{1}{2}(\alpha+\beta-\mu+1)}$	$(-z)\mathcal{F}_{\mu,\beta,\alpha}(z, \partial_z)$ $(-z)\mathcal{F}_{-\mu,-\beta,\alpha}(z, \partial_z)$ $(-z)\mathcal{F}_{\mu,-\beta,-\alpha}(z, \partial_z)$ $(-z)\mathcal{F}_{-\mu,\beta,-\alpha}(z, \partial_z)$	$(-z)^{\frac{1}{2}(-\alpha-\beta-\mu-1)},$ $(-z)^{\frac{1}{2}(-\alpha-\beta+\mu-1)}(z-1)^\beta,$ $(-z)^{\frac{1}{2}(-\alpha-\beta-\mu-1)}(z-1)^\beta,$ $(-z)^{\frac{1}{2}(-\alpha-\beta+\mu-1)};$
$w = 1 - \frac{1}{z} :$	$(-z)^{\frac{1}{2}(\alpha+\beta+\mu+1)}$ $(-z)^{\frac{1}{2}(\alpha+\beta-\mu+1)}(z-1)^{-\alpha}$ $(-z)^{\frac{1}{2}(\alpha+\beta+\mu+1)}(z-1)^{-\alpha}$ $(-z)^{\frac{1}{2}(\alpha+\beta-\mu+1)}$	$(-z)\mathcal{F}_{\mu,\alpha,\beta}(z, \partial_z)$ $(-z)\mathcal{F}_{-\mu,-\alpha,\beta}(z, \partial_z)$ $(-z)\mathcal{F}_{\mu,-\alpha,-\beta}(z, \partial_z)$ $(-z)\mathcal{F}_{-\mu,\alpha,-\beta}(z, \partial_z)$	$(-z)^{\frac{1}{2}(-\alpha-\beta-\mu-1)},$ $(-z)^{\frac{1}{2}(-\alpha-\beta+\mu-1)}(z-1)^\alpha,$ $(-z)^{\frac{1}{2}(-\alpha-\beta-\mu-1)}(z-1)^\alpha,$ $(-z)^{\frac{1}{2}(-\alpha-\beta+\mu-1)};$
$w = \frac{1}{1-z} :$	$(z-1)^{\frac{1}{2}(\alpha+\beta+\mu+1)}$ $(-z)^{-\beta}(z-1)^{\frac{1}{2}(\alpha+\beta-\mu+1)}$ $(z-1)^{\frac{1}{2}(\alpha+\beta-\mu+1)}$ $(-z)^{-\beta}(z-1)^{\frac{1}{2}(\alpha+\beta+\mu+1)}$	$(z-1)\mathcal{F}_{\beta,\mu,\alpha}(z, \partial_z)$ $(z-1)\mathcal{F}_{-\beta,-\mu,\alpha}(z, \partial_z)$ $(z-1)\mathcal{F}_{\beta,-\mu,-\alpha}(z, \partial_z)$ $(z-1)\mathcal{F}_{-\beta,\mu,-\alpha}(z, \partial_z)$	$(z-1)^{\frac{1}{2}(-\alpha-\beta-\mu-1)},$ $(-z)^\beta(z-1)^{\frac{1}{2}(-\alpha-\beta+\mu-1)},$ $(z-1)^{\frac{1}{2}(-\alpha-\beta+\mu-1)},$ $(-z)^\beta(z-1)^{\frac{1}{2}(-\alpha-\beta-\mu-1)};$
$w = \frac{z}{z-1} :$	$(z-1)^{\frac{1}{2}(\alpha+\beta+\mu+1)}$ $(-z)^{-\alpha}(z-1)^{\frac{1}{2}(\alpha+\beta-\mu+1)}$ $(z-1)^{\frac{1}{2}(\alpha+\beta-\mu+1)}$ $(-z)^{-\alpha}(z-1)^{\frac{1}{2}(\alpha+\beta+\mu+1)}$	$(z-1)\mathcal{F}_{\alpha,\mu,\beta}(z, \partial_z)$ $(z-1)\mathcal{F}_{-\alpha,-\mu,\beta}(z, \partial_z)$ $(z-1)\mathcal{F}_{\alpha,-\mu,-\beta}(z, \partial_z)$ $(z-1)\mathcal{F}_{-\alpha,\mu,-\beta}(z, \partial_z)$	$(z-1)^{\frac{1}{2}(-\alpha-\beta-\mu-1)},$ $(-z)^\alpha(z-1)^{\frac{1}{2}(-\alpha-\beta+\mu-1)},$ $(z-1)^{\frac{1}{2}(-\alpha-\beta+\mu-1)},$ $(-z)^\alpha(z-1)^{\frac{1}{2}(-\alpha-\beta-\mu-1)}.$

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These are enough to generate all possible permutations of the three singular points! There are six of those.

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$$\mathcal{F}_{\alpha, \beta, \mu}(z, \partial_z) = z^{-\alpha} \mathcal{F}_{-\alpha, \beta, -\mu}(z, \partial_z) z^{\alpha} \quad (16)$$

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$$\mathcal{F}_{\alpha, \beta, \mu}(z, \partial_z) = (1-z)^{-\beta} \mathcal{F}_{\alpha, -\beta, -\mu}(z, \partial_z) (1-z)^{\beta} \quad (17)$$

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- Three simple building blocks are enough for all of that!

A less nightmarish zoo from JD – Kummer table – again

$w = z :$	$(-z)^{-\alpha}(z-1)^{-\beta}$ $(z-1)^{-\beta}$ $(-z)^{-\alpha}$	$\mathcal{F}_{\alpha,\beta,\mu}(z, \partial_z),$ $\mathcal{F}_{-\alpha,-\beta,\mu}(z, \partial_z)$ $\mathcal{F}_{\alpha,-\beta,-\mu}(z, \partial_z)$ $\mathcal{F}_{-\alpha,\beta,-\mu}(z, \partial_z)$	$(-z)^\alpha(z-1)^\beta$ $(z-1)^\beta,$ $(-z)^\alpha;$
$w = 1 - z :$	$(z-1)^{-\alpha}(-z)^{-\beta}$ $(z-1)^{-\alpha}$ $(-z)^{-\beta}$	$\mathcal{F}_{\beta,\alpha,\mu}(z, \partial_z),$ $\mathcal{F}_{-\beta,-\alpha,\mu}(z, \partial_z)$ $\mathcal{F}_{\beta,-\alpha,-\mu}(z, \partial_z)$ $\mathcal{F}_{-\beta,\alpha,-\mu}(z, \partial_z)$	$(z-1)^\alpha(-z)^\beta,$ $(z-1)^\alpha,$ $(-z)^\beta;$
$w = \frac{1}{z} :$	$(-z)^{\frac{1}{2}(\alpha+\beta+\mu+1)}$ $(-z)^{\frac{1}{2}(\alpha+\beta-\mu+1)}(z-1)^{-\beta}$ $(-z)^{\frac{1}{2}(\alpha+\beta+\mu+1)}(z-1)^{-\beta}$ $(-z)^{\frac{1}{2}(\alpha+\beta-\mu+1)}$	$(-z)\mathcal{F}_{\mu,\beta,\alpha}(z, \partial_z)$ $(-z)\mathcal{F}_{-\mu,-\beta,\alpha}(z, \partial_z)$ $(-z)\mathcal{F}_{\mu,-\beta,-\alpha}(z, \partial_z)$ $(-z)\mathcal{F}_{-\mu,\beta,-\alpha}(z, \partial_z)$	$(-z)^{\frac{1}{2}(-\alpha-\beta-\mu-1)},$ $(-z)^{\frac{1}{2}(-\alpha-\beta-\mu-1)}(z-1)^\beta,$ $(-z)^{\frac{1}{2}(-\alpha-\beta-\mu-1)}(z-1)^\beta,$ $(-z)^{\frac{1}{2}(-\alpha-\beta-\mu-1)};$
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Try to multiply

It is now clear that $2^3 \cdot 3! = 8 \cdot 2 \cdot 3 = 2 \cdot 24$ is the order of the Weyl group for $SO(6, \mathbb{C})$. We have briefly described the action of the discrete Weyl group on hypergeometric operators. Let us discover quickly how the whole group acts.

Introduction

The hypergeometric equation

$SO(6, \mathbb{C})$ conformal action on hypergeometric functions

Endnotes

Conformal reduction from $n + 2$ to n complex dimensions

$so(6, \mathbb{C})$ Lie algebra

Coordinates

Root operators in the hand picked coordinates

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It is not obvious that one can also „push forward” the Laplace operator. In fact it descends to the quotient space when one considers homogeneous functions of order $1 - \frac{n}{2}$ on the quadric.

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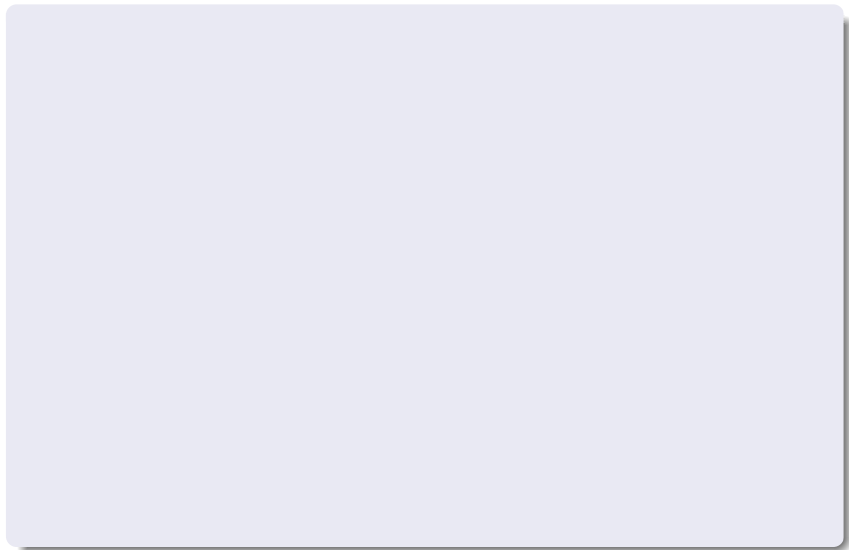
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- ... and three Cartan operators

$$N_i := B_{-i i} = -B_{i -i} = z_i\partial_{z_i} - z_{-i}\partial_{z_{-i}} \quad (22)$$

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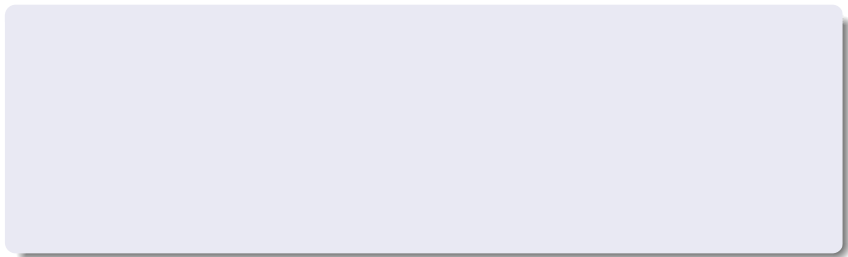
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In this coordinates the bilinear form is $Q(\vec{z}) = r^2 + p^2$. Therefore the reduction to the quadric will be given by $p = ir$.

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The four dimensional part with respect to coordinates w, u_1, u_2, u_3 gives the hypergeometric equation provided one makes a certain ansatz (roughly – the details of the reduction, even though most interesting, have been skipped for simplicity of this presentation)!

$$F(w, u_1, u_2, u_3) = u_1^\alpha u_2^\beta u_3^\mu F(w) \quad (23)$$

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Further we will frequently use those operators as the hypergeometric functions are their eigenvectors.

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Note how obvious and elegant the action of the Weyl group of $SO(6, \mathbb{C})$ looks using the above shown forms!

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 B_{32} &= \frac{i}{2} \frac{1}{u_2 u_3} \sqrt{1-w} \left[\lambda - 2w \partial_w + \frac{N_2}{1-w} + N_3 \right] \\
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The symmetries and relations for hypergeometric-type operators are visible here in the simplest way. Also this algebra can easily generate all possible formulae which arise through that construction! Please remember that the change of 1 into 2 involves the change of w into $1 - w$ and that is all!

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- The solution was involving hypergeometric-type functions
- We looked at the geometric interpretation of the six exactly solvable Schrödinger potentials (by means of hypergeometric functions)

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- We have studied Gell-Mann – Low adiabatic limit approach to scattering, and calculated **rigorously** a time-ordered exponential for a **non-commutative** family of operators
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- Jan introduced me to the global picture of the Lie group of symmetries for the hypergeometric equation which I liked very much and which finally helped me to understand the real beauty of the great plurality of possibilities.

A little about literature

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There is a great lot of literature on this subject. It would not be wise to present a long list here. For a condensed, full presentation, involving Lie groups of symmetries, I will show references to works of Jan Dereziński and recently also PM.

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- *M. Abramowitz, I. Stegun*, Handbook of Mathematical Functions, multiple editions, i.e. tenth printing, December 1972

Thank you for your attention!

It was a pleasure to give a talk to such an audience. 😊