

Local Phenomenological Models in M-Theory and F-Theory

Jacob L. Bourjaily
Princeton University & IAS

[arXiv:0905.0142],
[arXiv:0901.3785],
[arXiv:0706.3364], [arXiv:0804.1132],
[arXiv:0704.0445], and [arXiv:0704.0444]

String Phenomenology 2009, Warsaw

Outline

- 1 Building Blocks and Global Architecture
 - An Engineer's Guide to Model Building in F/M-Theory
 - Local Structural Engineering: A Group-Theoretic Classification of Locally-Engineered Effective Theories
 - Global Architecture: The Ubiquity and Uniqueness of E_8
- 2 Unfolding The Standard Model Out of E_8
 - Geometric Analogues to Grand Unification
 - Physics from Geometry:
 Novel Approaches to Model Building
 - Examples with Monodromies:
 The Diamond Ring of F-Theory
- 3 Conclusions and Future Directions

An Engineer's Toolbox: the Basic Building Blocks

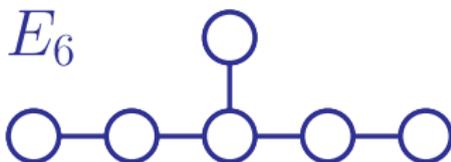
- Gauge symmetries arise via co-dim 4 'ADE'-type orbifold singularities

An Engineer's Toolbox: the Basic Building Blocks

- Gauge symmetries arise via co-dim 4 'ADE'-type orbifold singularities
 - These singularities are named according to the resulting gauge group, and are of the same structure in F-theory as in M-theory

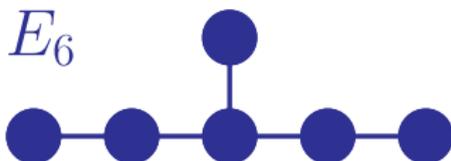
An Engineer's Toolbox: the Basic Building Blocks

- Gauge symmetries arise via co-dim 4 'ADE'-type orbifold singularities
 - These singularities are named according to the resulting gauge group, and are of the same structure in F-theory as in M-theory



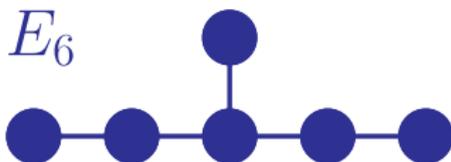
An Engineer's Toolbox: the Basic Building Blocks

- Gauge symmetries arise via co-dim 4 'ADE'-type orbifold singularities
 - These singularities are named according to the resulting gauge group, and are of the same structure in F-theory as in M-theory



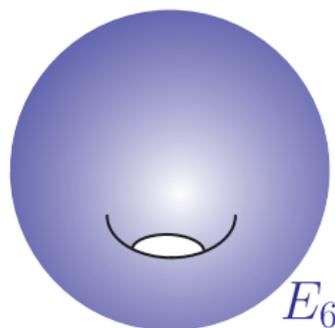
An Engineer's Toolbox: the Basic Building Blocks

- **Gauge symmetries** arise via co-dim 4 '**ADE**'-type orbifold singularities
 - These singularities are named according to the resulting gauge group, and are of the same structure in F-theory as in M-theory



- F-theory: these are **complex two-cycles** in an elliptic Calabi-Yau four-fold
- M-theory: these are **real three-folds** in a G_2 -manifold

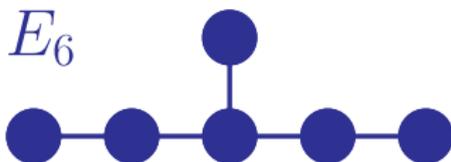
F-Theory



E_6 -gauge theory
via ' E_6 '-singularities
along a del-Pezzo 8

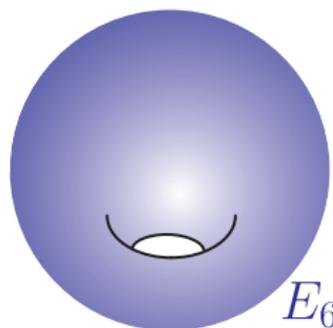
An Engineer's Toolbox: the Basic Building Blocks

- **Gauge symmetries** arise via co-dim 4 '**ADE**'-type orbifold singularities
 - These singularities are named according to the resulting gauge group, and are of the same structure in F-theory as in M-theory



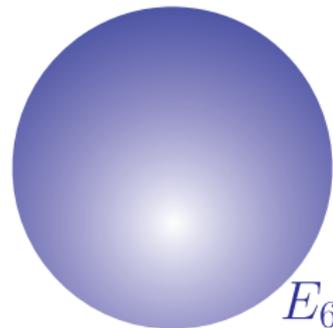
- F-theory: these are **complex two-cycles** in an elliptic Calabi-Yau four-fold
- M-theory: these are **real three-folds** in a G_2 -manifold

F-Theory



E_6 -gauge theory
 via ' E_6 '-singularities
 along a del-Pezzo 8

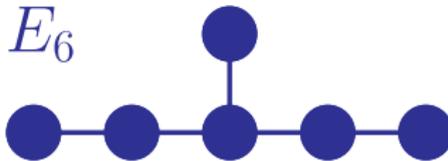
M-Theory



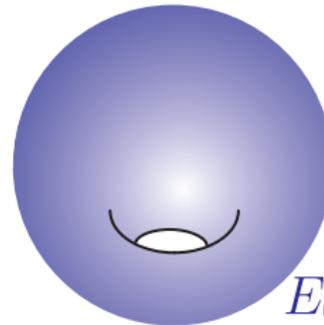
E_6 -gauge theory
 via ' E_6 '-singularities
 along a Poincaré-' S^3 '

An Engineer's Toolbox: the Basic Building Blocks

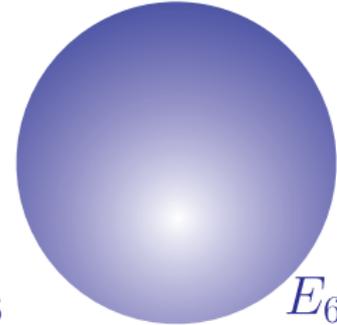
- Gauge symmetries arise via co-dim 4 'ADE'-type orbifold singularities
- Charged, chiral matter from local **enhancements** of ADE singularities



F-Theory

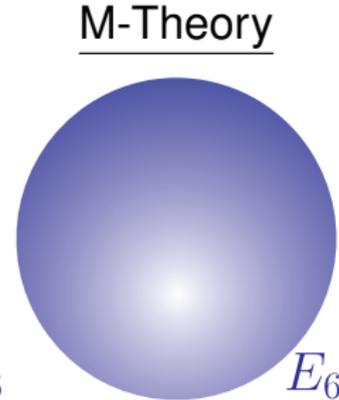
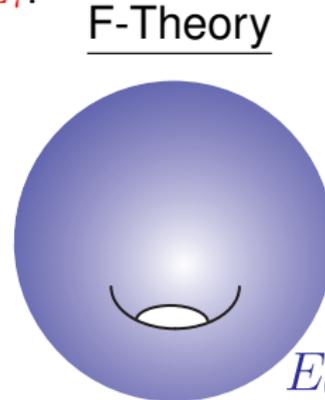
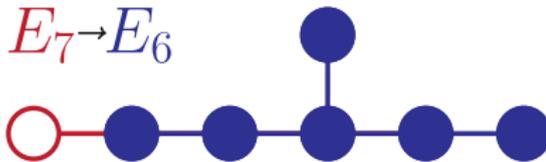


M-Theory



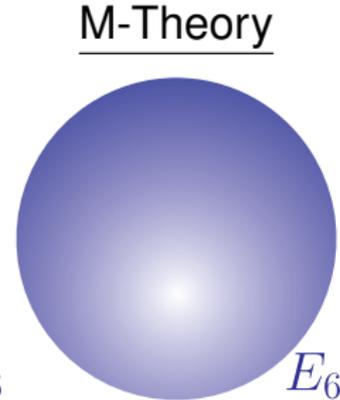
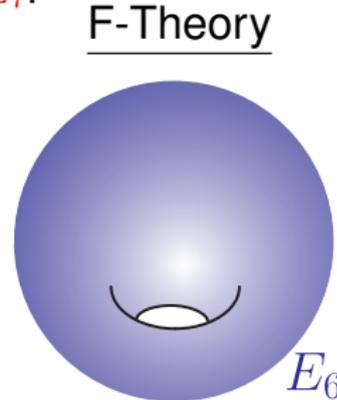
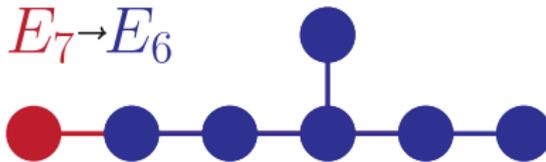
An Engineer's Toolbox: the Basic Building Blocks

- Gauge symmetries arise via co-dim 4 'ADE'-type orbifold singularities
- Charged, chiral matter from local **enhancements** of ADE singularities
 - For example, a **27** of E_6 can be supported wherever a co-dim 4 E_6 -singularity is enhanced to E_7 .



An Engineer's Toolbox: the Basic Building Blocks

- **Gauge symmetries** arise via co-dim 4 'ADE'-type orbifold singularities
- **Charged, chiral matter** from local **enhancements** of ADE singularities
 - For example, a **27** of E_6 can be supported wherever a co-dim 4 E_6 -singularity is enhanced to E_7 .



An Engineer's Toolbox: the Basic Building Blocks

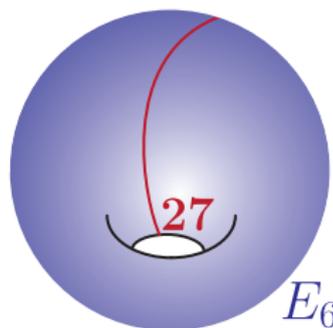
- **Gauge symmetries** arise via co-dim 4 'ADE'-type orbifold singularities
- **Charged, chiral matter** from local **enhancements** of ADE singularities
 - For example, a **27** of E_6 can be supported wherever a co-dim 4 E_6 -singularity is enhanced to E_7 .

- F-theory: these are 'curves'
(Riemann surfaces) in the co-dim 4
singular surface (think D7-branes)

- only curves with non-vanishing flux
generate chiral matter, making
'exotic' matter curves relatively
easy to ignore.

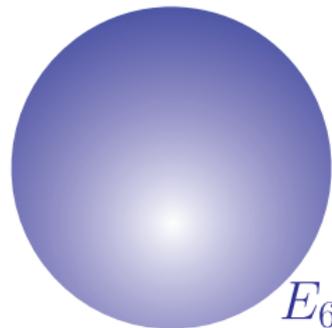
- allows a single curve to support
multiple 'generations'
- allows for a nice solution to
doublet-triplet splitting

F-Theory



an E_7 'matter-curve'
can support a
massless **27** of E_6

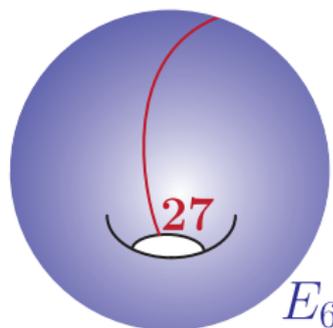
M-Theory



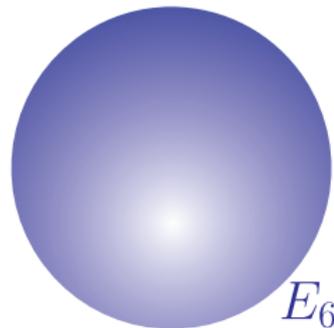
An Engineer's Toolbox: the Basic Building Blocks

- **Gauge symmetries** arise via co-dim 4 'ADE'-type orbifold singularities
- **Charged, chiral matter** from local **enhancements** of ADE singularities
 - For example, a **27** of E_6 can be supported wherever a co-dim 4 E_6 -singularity is enhanced to E_7 .
- F-theory: these are 'curves' (Riemann surfaces) in the co-dim 4 singular surface (think D7-branes)
- only curves with non-vanishing flux generate chiral matter, making 'exotic' matter curves relatively easy to ignore.
 - allows a single curve to support multiple 'generations'
 - allows for a nice solution to doublet-triplet splitting

F-Theory



M-Theory

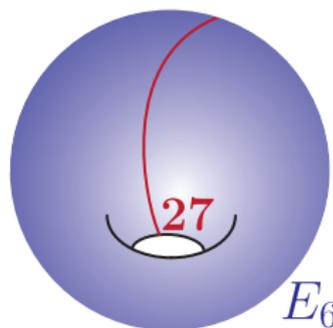


an E_7 'matter-curve' can support a massless **27** of E_6

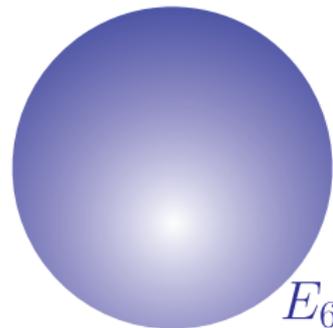
An Engineer's Toolbox: the Basic Building Blocks

- **Gauge symmetries** arise via co-dim 4 'ADE'-type orbifold singularities
- **Charged, chiral matter** from local **enhancements** of ADE singularities
 - For example, a **27** of E_6 can be supported wherever a co-dim 4 **E_6 -singularity** is enhanced to E_7 .
- F-theory: these are 'curves' (Riemann surfaces) in the co-dim 4 singular surface (think D7-branes)
- only curves with non-vanishing flux generate chiral matter, making 'exotic' matter curves relatively easy to ignore.
 - allows a single curve to support multiple 'generations'
 - allows for a nice solution to doublet-triplet splitting

F-Theory



M-Theory

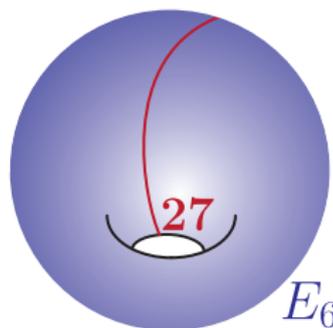


an E_7 'matter-curve' can support a massless **27** of E_6

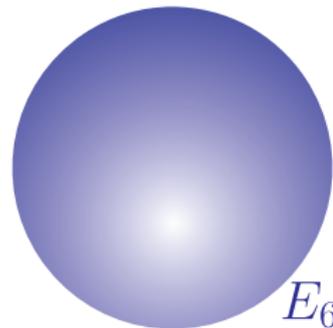
An Engineer's Toolbox: the Basic Building Blocks

- **Gauge symmetries** arise via co-dim 4 'ADE'-type orbifold singularities
- **Charged, chiral matter** from local **enhancements** of ADE singularities
 - For example, a **27** of E_6 can be supported wherever a co-dim 4 E_6 -singularity is enhanced to E_7 .
- F-theory: these are 'curves' (Riemann surfaces) in the co-dim 4 singular surface (think D7-branes)
- only curves with non-vanishing flux generate chiral matter, making 'exotic' matter curves relatively easy to ignore.
 - allows a single curve to support multiple 'generations'
 - allows for a nice solution to doublet-triplet splitting

F-Theory



M-Theory

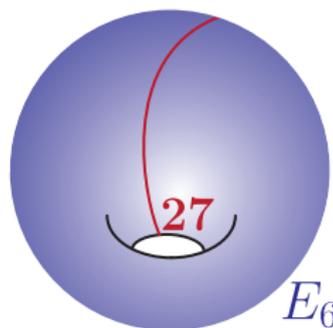


an E_7 'matter-curve' can support a massless **27** of E_6

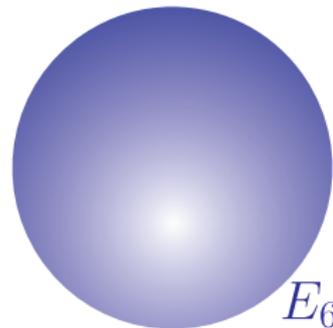
An Engineer's Toolbox: the Basic Building Blocks

- **Gauge symmetries** arise via co-dim 4 'ADE'-type orbifold singularities
- **Charged, chiral matter** from local **enhancements** of ADE singularities
 - For example, a **27** of E_6 can be supported wherever a co-dim 4 E_6 -singularity is enhanced to E_7 .
- F-theory: these are 'curves' (Riemann surfaces) in the co-dim 4 singular surface (think D7-branes)
- only curves with non-vanishing flux generate chiral matter, making 'exotic' matter curves relatively easy to ignore.
 - allows a single curve to support multiple 'generations'
 - allows for a nice solution to doublet-triplet splitting

F-Theory



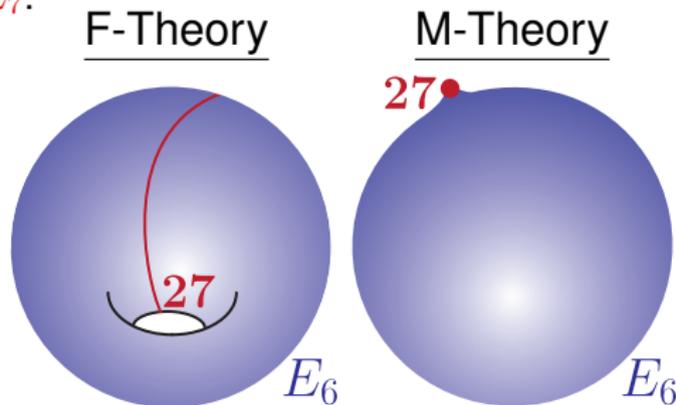
M-Theory



an E_7 'matter-curve' can support a massless **27** of E_6

An Engineer's Toolbox: the Basic Building Blocks

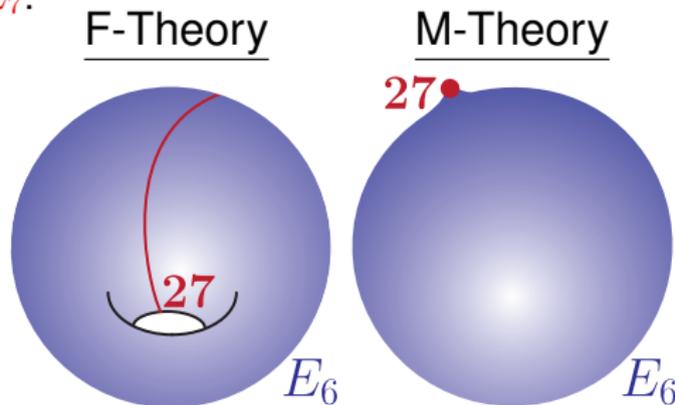
- **Gauge symmetries** arise via co-dim 4 'ADE'-type orbifold singularities
- **Charged, chiral matter** from local **enhancements** of ADE singularities
 - For example, a **27** of E_6 can be supported wherever a co-dim 4 **E_6 -singularity** is enhanced to E_7 .
- M-theory: these are isolated points along the co-dim 4 singular surface (think D6-branes)
- **every** such conical singularity gives rise to ± 1 chiral multiplet making 'exotic' matter much harder to ignore
 - makes M-theory models relatively more constrained (and hence predictive)
 - softens flavour problems



a conical E_7 -singularity can support a massless **27** of E_6

An Engineer's Toolbox: the Basic Building Blocks

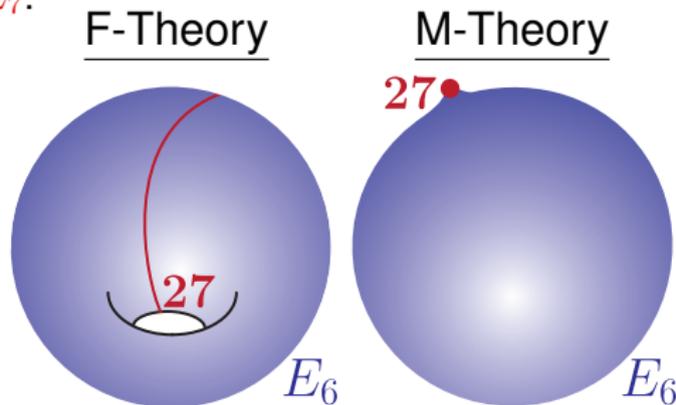
- **Gauge symmetries** arise via co-dim 4 'ADE'-type orbifold singularities
- **Charged, chiral matter** from local **enhancements** of ADE singularities
 - For example, a **27** of E_6 can be supported wherever a co-dim 4 **E_6 -singularity** is enhanced to E_7 .
- M-theory: these are isolated points along the co-dim 4 singular surface (think D6-branes)
- **every** such conical singularity gives rise to ± 1 chiral multiplet, making 'exotic' matter much harder to ignore
 - makes M-theory models relatively more constrained (and hence predictive)
 - softens flavour problems



a conical E_7 -singularity can support a massless **27** of E_6

An Engineer's Toolbox: the Basic Building Blocks

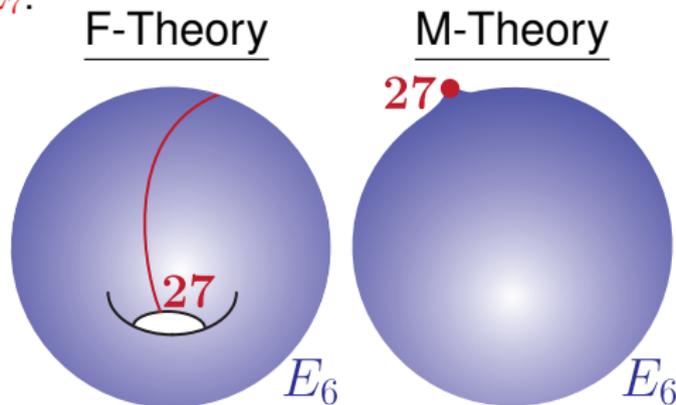
- **Gauge symmetries** arise via co-dim 4 'ADE'-type orbifold singularities
- **Charged, chiral matter** from local **enhancements** of ADE singularities
 - For example, a **27** of E_6 can be supported wherever a co-dim 4 **E_6 -singularity** is enhanced to E_7 .
- M-theory: these are isolated points along the co-dim 4 singular surface (think D6-branes)
- **every** such conical singularity gives rise to ± 1 chiral multiplet, making 'exotic' matter much harder to ignore
 - makes M-theory models relatively more constrained (and hence predictive)
 - softens flavour problems



a conical E_7 -singularity can support a massless **27** of E_6

An Engineer's Toolbox: the Basic Building Blocks

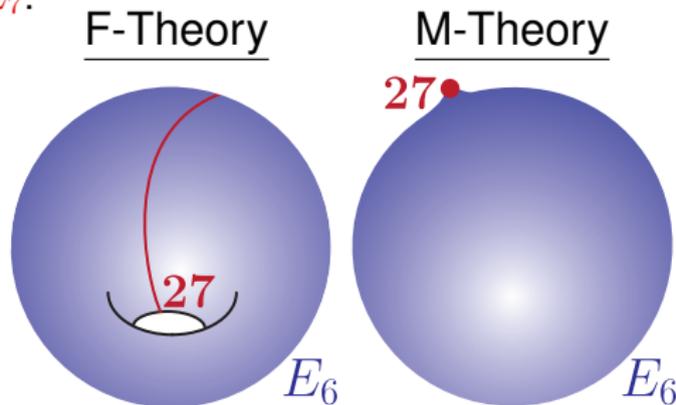
- **Gauge symmetries** arise via co-dim 4 'ADE'-type orbifold singularities
- **Charged, chiral matter** from local **enhancements** of ADE singularities
 - For example, a **27** of E_6 can be supported wherever a co-dim 4 **E_6 -singularity** is enhanced to E_7 .
- M-theory: these are isolated points along the co-dim 4 singular surface (think D6-branes)
- **every** such conical singularity gives rise to ± 1 chiral multiplet, making 'exotic' matter much harder to ignore
 - makes M-theory models relatively more constrained (and hence predictive)
 - softens flavour problems



a conical E_7 -singularity can support a massless **27** of E_6

An Engineer's Toolbox: the Basic Building Blocks

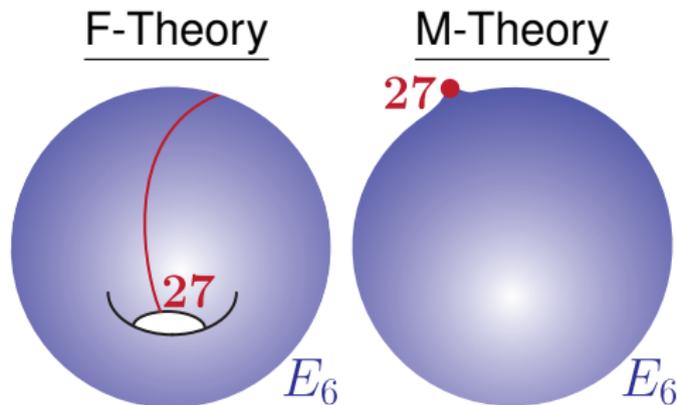
- **Gauge symmetries** arise via co-dim 4 'ADE'-type orbifold singularities
- **Charged, chiral matter** from local **enhancements** of ADE singularities
 - For example, a **27** of E_6 can be supported wherever a co-dim 4 **E_6 -singularity** is enhanced to E_7 .
- M-theory: these are isolated points along the co-dim 4 singular surface (think D6-branes)
- **every** such conical singularity gives rise to ± 1 chiral multiplet, making 'exotic' matter much harder to ignore
 - makes M-theory models relatively more constrained (and hence predictive)
 - softens flavour problems



a conical E_7 -singularity can support a massless **27** of E_6

An Engineer's Toolbox: the Basic Building Blocks

- Gauge symmetries arise via co-dim 4 'ADE'-type orbifold singularities
- Charged, chiral matter from local **enhancements** of ADE singularities
- Superpotential interactions from structures connecting disparate matter-singularities



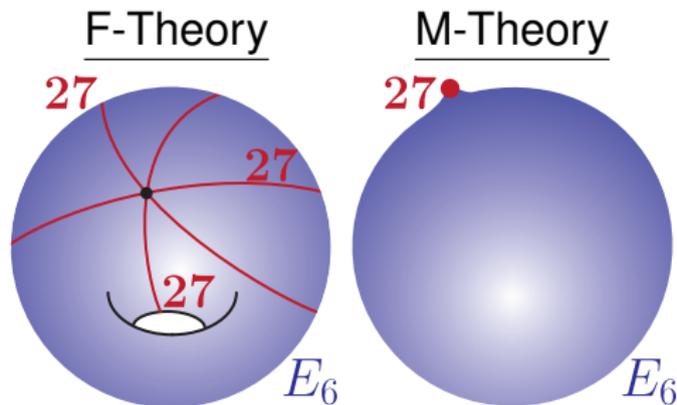
An Engineer's Toolbox: the Basic Building Blocks

- **Gauge symmetries** arise via co-dim 4 'ADE'-type orbifold singularities
- **Charged, chiral matter** from local **enhancements** of ADE singularities
- **Superpotential interactions** from structures connecting disparate matter-singularities

- F-theory: these are triple-intersections of matter-curves

- M-theory: these are supersymmetric three-cycles supporting multiple conical singularities

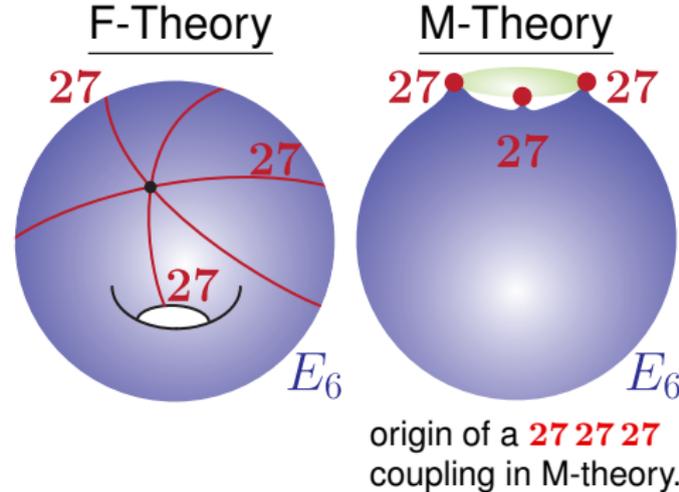
- notice that both structures appear topologically non-generic



origin of a **27 27 27**
coupling in F-theory.

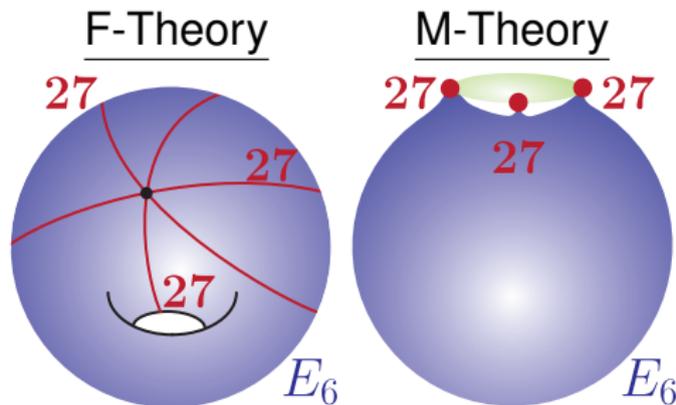
An Engineer's Toolbox: the Basic Building Blocks

- **Gauge symmetries** arise via co-dim 4 'ADE'-type orbifold singularities
- **Charged, chiral matter** from local **enhancements** of ADE singularities
- **Superpotential interactions** from structures connecting disparate matter-singularities
- F-theory: these are triple-intersections of matter-curves
- M-theory: these are super-symmetric three-cycles supporting multiple conical singularities
 - notice that both structures appear topologically non-generic, but they are in fact ubiquitous in ALE-fibrations



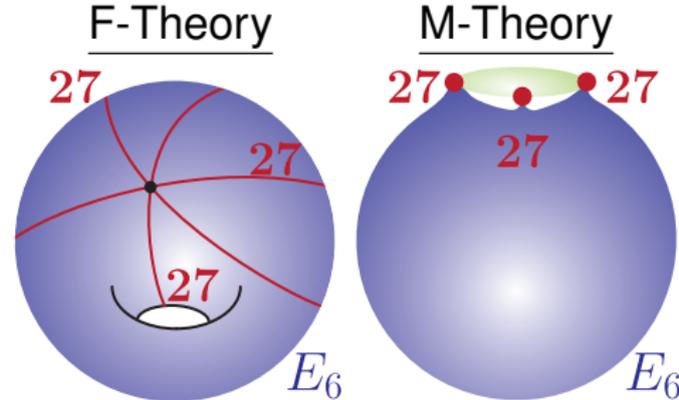
An Engineer's Toolbox: the Basic Building Blocks

- **Gauge symmetries** arise via co-dim 4 'ADE'-type orbifold singularities
- **Charged, chiral matter** from local **enhancements** of ADE singularities
- **Superpotential interactions** from structures connecting disparate matter-singularities
- F-theory: these are triple-intersections of matter-curves
- M-theory: these are supersymmetric three-cycles supporting multiple conical singularities
 - notice that both structures appear topologically non-generic, but they are in fact **ubiquitous** in ALE-fibrations, and for the same reasons



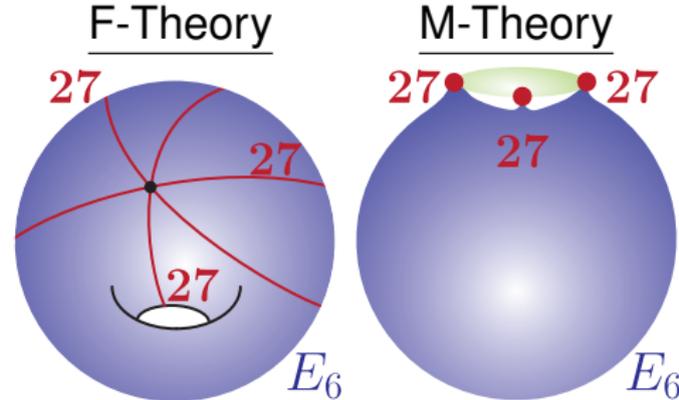
An Engineer's Toolbox: the Basic Building Blocks

- **Gauge symmetries** arise via co-dim 4 'ADE'-type orbifold singularities
- **Charged, chiral matter** from local **enhancements** of ADE singularities
- **Superpotential interactions** from structures connecting disparate matter-singularities
- F-theory: these are triple-intersections of matter-curves
- M-theory: these are supersymmetric three-cycles supporting multiple conical singularities
 - notice that both structures appear topologically non-generic, but they are in fact **ubiquitous** in ALE-fibrations, and for the same reasons



An Engineer's Toolbox: the Basic Building Blocks

- **Gauge symmetries** arise via co-dim 4 'ADE'-type orbifold singularities
- **Charged, chiral matter** from local **enhancements** of ADE singularities
- **Superpotential interactions** from structures connecting disparate matter-singularities
- F-theory: these are triple-intersections of matter-curves
- M-theory: these are supersymmetric three-cycles supporting multiple conical singularities
 - notice that both structures appear topologically non-generic, but they are in fact **ubiquitous** in ALE-fibrations, **and for the same reasons**

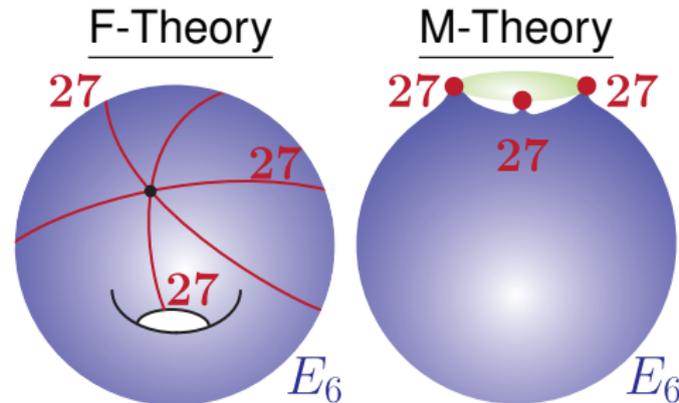


An Engineer's Toolbox: the Basic Building Blocks

- **Gauge symmetries** arise via co-dim 4 'ADE'-type orbifold singularities
- **Charged, chiral matter** from local **enhancements** of ADE singularities
- **Superpotential interactions** from structures connecting disparate matter-singularities

Notice some very generic features of any such superpotential:

- **sparse**
 - ubiquitous D_4 's: think of the extra 'flavour-branes' which must intersect along each matter-curve;
- **hierarchical**
 - coefficients are typically exponentially suppressed

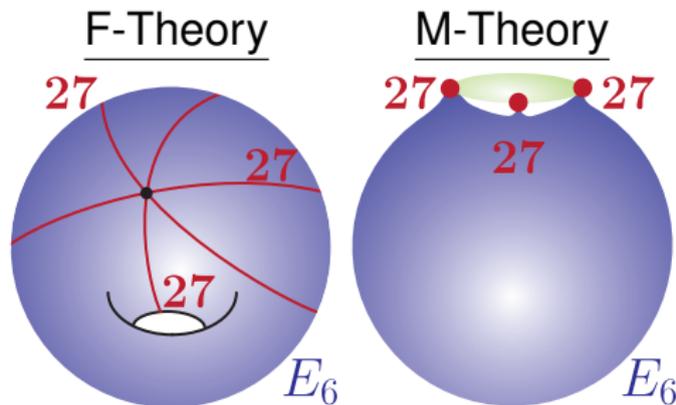


An Engineer's Toolbox: the Basic Building Blocks

- **Gauge symmetries** arise via co-dim 4 'ADE'-type orbifold singularities
- **Charged, chiral matter** from local **enhancements** of ADE singularities
- **Superpotential interactions** from structures connecting disparate matter-singularities

Notice some very generic features of any such superpotential:

- **sparse**
 - ubiquitous U_1 's: think of the extra 'flavour-branes' which must intersect along each matter-curve;
- **hierarchical**
 - coefficients are typically exponentially suppressed

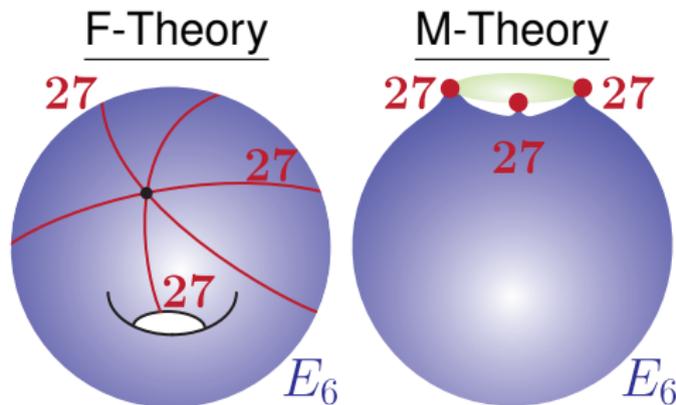


An Engineer's Toolbox: the Basic Building Blocks

- **Gauge symmetries** arise via co-dim 4 'ADE'-type orbifold singularities
- **Charged, chiral matter** from local **enhancements** of ADE singularities
- **Superpotential interactions** from structures connecting disparate matter-singularities

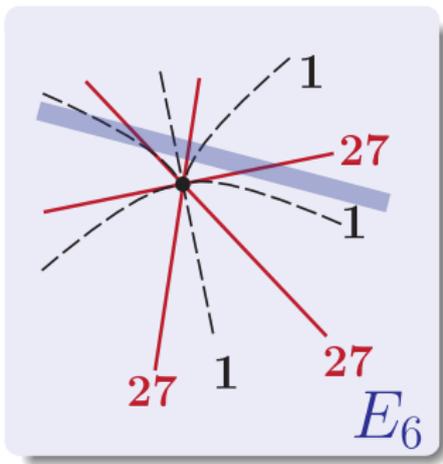
Notice some very generic features of any such superpotential:

- **sparse**
 - ubiquitous U_1 's: think of the extra 'flavour-branes' which must intersect along each matter-curve;
- **hierarchical**
 - coefficients are typically exponentially suppressed



Correspondence of Local Geometries for F-Theory and M-Theory

- For any local model, the most important topological data to have is:
 - a list of all (potentially-massless-)matter-supporting singularities, and
 - the selection rules which determine how matter living along these singularities can interact in the theory
- This data can be encoded in F-theory by a cartoon-collection of mutually-intersecting matter-curves



- These models are constructed explicitly as ALE-fibrations over an appropriate base W , e.g. $\widehat{E}_8(a(W), b(W), 0, 0, 0, 0, 0, 0)$
- Any local geometry constructed in this way for F-theory can be immediately translated into a corresponding geometry for M-theory, with the same matter-singularities and interactions
- It is not hard to classify all possible ‘cartoons’ that can arise from a local ALE-fibration leading to an extremely small landscape of possibilities

Group-Theoretic Classification of Local Fibrations

- Recall that the new light degrees of freedom arising from an ALE-fibration are due to 2-branes wrapping vanishing cycles within the ALE-fibres.
- The correspondence between two-cycles in the ALE-space \widehat{G} and the **root lattice** of the ADE-group G then leads to the following:
 - A generic \widehat{G} -fibred Calabi-Yau four-fold or \widehat{G} -fibred G_2 manifold for which the typical fibre has a singularity of type $H \subset G$ will have one matter-singularity for each vector-like representation in the branching:

Group-Theoretic Classification of Local Fibrations

- Recall that the new light degrees of freedom arising from an ALE-fibration are due to 2-branes wrapping vanishing cycles within the ALE-fibres.
- The correspondence between two-cycles in the ALE-space \widehat{G} and the **root lattice** of the ADE-group G then leads to the following:
 - A generic \widehat{G} -fibred Calabi-Yau four-fold or \widehat{G} -fibred G_2 manifold for which the typical fibre has a singularity of type $H \subset G$ will have one matter-singularity for each vector-like representation in the branching:

Group-Theoretic Classification of Local Fibrations

- Recall that the new light degrees of freedom arising from an ALE-fibration are due to 2-branes wrapping vanishing cycles within the ALE-fibres.
- The correspondence between two-cycles in the ALE-space \widehat{G} and the **root lattice** of the ADE-group G then leads to the following:
 - A generic \widehat{G} -fibred Calabi-Yau four-fold or \widehat{G} -fibred G_2 manifold for which the typical fibre has a singularity of type $H \subset G$ will have one matter-singularity for each vector-like representation in the branching:

Group-Theoretic Classification of Local Fibrations

- Recall that the new light degrees of freedom arising from an ALE-fibration are due to 2-branes wrapping vanishing cycles within the ALE-fibres.
- The correspondence between two-cycles in the ALE-space \widehat{G} and the **root lattice** of the ADE-group G then leads to the following:
 - A generic \widehat{G} -fibred Calabi-Yau four-fold or \widehat{G} -fibred G_2 manifold for which the typical fibre has a singularity of type $H \subset G$ will have one matter-singularity for each vector-like representation in the branching:

$$\text{adj}(G) = \text{adj} \left(H \times \prod_{i=1}^k U_1 \right) \oplus (\mathbf{R}_{\vec{q}} \oplus \overline{\mathbf{R}}_{-\vec{q}}).$$

Group-Theoretic Classification of Local Fibrations

- Recall that the new light degrees of freedom arising from an ALE-fibration are due to 2-branes wrapping vanishing cycles within the ALE-fibres.
- The correspondence between two-cycles in the ALE-space \widehat{G} and the **root lattice** of the ADE-group G then leads to the following:
 - A generic \widehat{G} -fibred Calabi-Yau four-fold or \widehat{G} -fibred G_2 manifold for which the typical fibre has a singularity of type $H \subset G$ will have one matter-singularity for each vector-like representation in the branching:

$$\text{adj}(G) = \text{adj} \left(H \times \prod_{i=1}^k U_1 \right) \oplus (\mathbf{R}_{\vec{q}} \oplus \overline{\mathbf{R}}_{-\vec{q}}).$$

for example, $E_8 \supset E_6 \times U_1^a \times U_1^b$:

$$\begin{aligned} 248 = & \mathbf{78}_{0,0} \oplus \mathbf{1}_{0,0} \oplus \mathbf{1}_{0,0} \oplus \mathbf{27}_{1,0} \oplus \mathbf{27}_{0,1} \oplus \mathbf{27}_{-1,-1} \oplus \mathbf{1}_{1,2} \oplus \mathbf{1}_{2,1} \oplus \mathbf{1}_{1,-1} \\ & \oplus \overline{\mathbf{27}}_{-1,0} \oplus \overline{\mathbf{27}}_{0,-1} \oplus \overline{\mathbf{27}}_{1,1} \oplus \mathbf{1}_{-1,-2} \oplus \mathbf{1}_{-2,-1} \oplus \mathbf{1}_{-1,1} \end{aligned}$$

Group-Theoretic Classification of Local Fibrations

- Recall that the new light degrees of freedom arising from an ALE-fibration are due to 2-branes wrapping vanishing cycles within the ALE-fibres.
- The correspondence between two-cycles in the ALE-space \widehat{G} and the **root lattice** of the ADE-group G then leads to the following:
 - A generic \widehat{G} -fibred Calabi-Yau four-fold or \widehat{G} -fibred G_2 manifold for which the typical fibre has a singularity of type $H \subset G$ will have one matter-singularity for each vector-like representation in the branching:

$$\text{adj}(G) = \text{adj} \left(H \times \prod_{i=1}^k U_1 \right) \oplus (\mathbf{R}_{\vec{q}} \oplus \overline{\mathbf{R}}_{-\vec{q}}).$$

for example, $E_8 \supset E_6 \times U_1^a \times U_1^b$:

$$248 = 78_{0,0} \oplus \mathbf{1}_{0,0} \oplus \mathbf{1}_{0,0} \oplus \mathbf{27}_{1,0} \oplus \mathbf{27}_{0,1} \oplus \mathbf{27}_{-1,-1} \oplus \mathbf{1}_{1,2} \oplus \mathbf{1}_{2,1} \oplus \mathbf{1}_{1,-1} \\ \oplus \overline{\mathbf{27}}_{-1,0} \oplus \overline{\mathbf{27}}_{0,-1} \oplus \overline{\mathbf{27}}_{1,1} \oplus \mathbf{1}_{-1,-2} \oplus \mathbf{1}_{-2,-1} \oplus \mathbf{1}_{-1,1}.$$

Group-Theoretic Classification of Local Fibrations

- Recall that the new light degrees of freedom arising from an ALE-fibration are due to 2-branes wrapping vanishing cycles within the ALE-fibres.
- The correspondence between two-cycles in the ALE-space \widehat{G} and the **root lattice** of the ADE-group G then leads to the following:
 - A generic \widehat{G} -fibred Calabi-Yau four-fold or \widehat{G} -fibred G_2 manifold for which the typical fibre has a singularity of type $H \subset G$ will have one matter-singularity for each vector-like representation in the branching:

$$\text{adj}(G) = \text{adj} \left(H \times \prod_{i=1}^k U_1 \right) \oplus (\mathbf{R}_{\vec{q}} \oplus \overline{\mathbf{R}}_{-\vec{q}}).$$

for example, $E_8 \supset E_6 \times U_1^a \times U_1^b$:

$$248 = 78_{0,0} \oplus \mathbf{1}_{0,0} \oplus \mathbf{1}_{0,0} \oplus \mathbf{27}_{1,0} \oplus \mathbf{27}_{0,1} \oplus \mathbf{27}_{-1,-1} \oplus \mathbf{1}_{1,2} \oplus \mathbf{1}_{2,1} \oplus \mathbf{1}_{1,-1} \\ \oplus \overline{\mathbf{27}}_{-1,0} \oplus \overline{\mathbf{27}}_{0,-1} \oplus \overline{\mathbf{27}}_{1,1} \oplus \mathbf{1}_{-1,-2} \oplus \mathbf{1}_{-2,-1} \oplus \mathbf{1}_{-1,1}.$$

Group-Theoretic Classification of Local Fibrations

- Recall that the new light degrees of freedom arising from an ALE-fibration are due to 2-branes wrapping vanishing cycles within the ALE-fibres.
- The correspondence between two-cycles in the ALE-space \widehat{G} and the **root lattice** of the ADE-group G then leads to the following:
 - A generic \widehat{G} -fibred Calabi-Yau four-fold or \widehat{G} -fibred G_2 manifold for which the typical fibre has a singularity of type $H \subset G$ will have one matter-singularity for each vector-like representation in the branching:

$$\text{adj}(G) = \text{adj} \left(H \times \prod_{i=1}^k U_1 \right) \oplus (\mathbf{R}_{\vec{q}} \oplus \overline{\mathbf{R}}_{-\vec{q}}).$$

for example, $E_8 \supset E_6 \times U_1^a \times U_1^b$:

$$\begin{aligned} 248 = & 78_{0,0} \oplus \mathbf{1}_{0,0} \oplus \mathbf{1}_{0,0} \oplus \mathbf{27}_{1,0} \oplus \mathbf{27}_{0,1} \oplus \mathbf{27}_{-1,-1} \oplus \mathbf{1}_{1,2} \oplus \mathbf{1}_{2,1} \oplus \mathbf{1}_{1,-1} \\ & \oplus \overline{\mathbf{27}}_{-1,0} \oplus \overline{\mathbf{27}}_{0,-1} \oplus \overline{\mathbf{27}}_{1,1} \oplus \mathbf{1}_{-1,-2} \oplus \mathbf{1}_{-2,-1} \oplus \mathbf{1}_{-1,1} \end{aligned}$$

Group-Theoretic Classification of Local Fibrations

- Recall that the new light degrees of freedom arising from an ALE-fibration are due to 2-branes wrapping vanishing cycles within the ALE-fibres.
- The correspondence between two-cycles in the ALE-space \widehat{G} and the **root lattice** of the ADE-group G then leads to the following:
 - A generic \widehat{G} -fibred Calabi-Yau four-fold or \widehat{G} -fibred G_2 manifold for which the typical fibre has a singularity of type $H \subset G$ will have one matter-singularity for each vector-like representation in the branching:

$$\text{adj}(G) = \text{adj} \left(H \times \prod_{i=1}^k U_1 \right) \oplus (\mathbf{R}_{\vec{q}} \oplus \overline{\mathbf{R}}_{-\vec{q}}).$$

for example, $E_8 \supset E_6 \times U_1^a \times U_1^b$:

$$248 = 78_{0,0} \oplus \mathbf{1}_{0,0} \oplus \mathbf{1}_{0,0} \oplus \mathbf{27}_{1,0} \oplus \mathbf{27}_{0,1} \oplus \mathbf{27}_{-1,-1} \oplus \mathbf{1}_{1,2} \oplus \mathbf{1}_{2,1} \oplus \mathbf{1}_{1,-1} \\ \oplus \overline{\mathbf{27}}_{-1,0} \oplus \overline{\mathbf{27}}_{0,-1} \oplus \overline{\mathbf{27}}_{1,1} \oplus \mathbf{1}_{-1,-2} \oplus \mathbf{1}_{-2,-1} \oplus \mathbf{1}_{-1,1}.$$

Group-Theoretic Classification of Local Fibrations

- Recall that the new light degrees of freedom arising from an ALE-fibration are due to 2-branes wrapping vanishing cycles within the ALE-fibres.
- The correspondence between two-cycles in the ALE-space \widehat{G} and the **root lattice** of the ADE-group G then leads to the following:
 - A generic \widehat{G} -fibred Calabi-Yau four-fold or \widehat{G} -fibred G_2 manifold for which the typical fibre has a singularity of type $H \subset G$ will have one matter-singularity for each vector-like representation in the branching:

$$\text{adj}(G) = \text{adj} \left(H \times \prod_{i=1}^k U_1 \right) \oplus (\mathbf{R}_{\vec{q}} \oplus \overline{\mathbf{R}}_{-\vec{q}}).$$

for example, $E_8 \supset E_6 \times U_1^a \times U_1^b$:

$$248 = 78_{0,0} \oplus \mathbf{1}_{0,0} \oplus \mathbf{1}_{0,0} \oplus \mathbf{27}_{1,0} \oplus \mathbf{27}_{0,1} \oplus \mathbf{27}_{-1,-1} \oplus \mathbf{1}_{1,2} \oplus \mathbf{1}_{2,1} \oplus \mathbf{1}_{1,-1} \\ \oplus \overline{\mathbf{27}}_{-1,0} \oplus \overline{\mathbf{27}}_{0,-1} \oplus \overline{\mathbf{27}}_{1,1} \oplus \mathbf{1}_{-1,-2} \oplus \mathbf{1}_{-2,-1} \oplus \mathbf{1}_{-1,1}.$$

Group-Theoretic Classification of Local Fibrations

- Recall that the new light degrees of freedom arising from an ALE-fibration are due to 2-branes wrapping vanishing cycles within the ALE-fibres.
- The correspondence between two-cycles in the ALE-space \widehat{G} and the **root lattice** of the ADE-group G then leads to the following:
 - A generic \widehat{G} -fibred Calabi-Yau four-fold or \widehat{G} -fibred G_2 manifold for which the typical fibre has a singularity of type $H \subset G$ will have one matter-singularity for each vector-like representation in the branching:

$$\text{adj}(G) = \text{adj} \left(H \times \prod_{i=1}^k U_1 \right) \oplus (\mathbf{R}_{\vec{q}} \oplus \overline{\mathbf{R}}_{-\vec{q}}).$$

for example, $E_8 \supset E_6 \times U_1^a \times U_1^b$:

$$248 = 78_{0,0} \oplus 1_{0,0} \oplus 1_{0,0} \oplus 27_{1,0} \oplus 27_{0,1} \oplus 27_{-1,-1} \oplus 1_{1,2} \oplus 1_{2,1} \oplus 1_{1,-1} \\ \oplus \overline{27}_{-1,0} \oplus \overline{27}_{0,-1} \oplus \overline{27}_{1,1} \oplus 1_{-1,-2} \oplus 1_{-2,-1} \oplus 1_{-1,1}.$$

Group-Theoretic Classification of Local Fibrations

- Recall that the new light degrees of freedom arising from an ALE-fibration are due to 2-branes wrapping vanishing cycles within the ALE-fibres.
- The correspondence between two-cycles in the ALE-space \widehat{G} and the **root lattice** of the ADE-group G then leads to the following:
 - A generic $\star \widehat{G}$ -fibred Calabi-Yau four-fold or \widehat{G} -fibred G_2 manifold for which the typical fibre has a singularity of type $H \subset G$ will have one matter-singularity for each vector-like representation in the branching:

$$\text{adj}(G) = \text{adj} \left(H \times \prod_{i=1}^k U_1 \right) \oplus (\mathbf{R}_{\vec{q}} \oplus \overline{\mathbf{R}}_{-\vec{q}}).$$

for example, $E_8 \supset E_6 \times U_1^a \times U_1^b$:

$$248 = 78_{0,0} \oplus \mathbf{1}_{0,0} \oplus \mathbf{1}_{0,0} \oplus \mathbf{27}_{1,0} \oplus \mathbf{27}_{0,1} \oplus \mathbf{27}_{-1,-1} \oplus \mathbf{1}_{1,2} \oplus \mathbf{1}_{2,1} \oplus \mathbf{1}_{1,-1} \\ \oplus \overline{\mathbf{27}}_{-1,0} \oplus \overline{\mathbf{27}}_{0,-1} \oplus \overline{\mathbf{27}}_{1,1} \oplus \mathbf{1}_{-1,-2} \oplus \mathbf{1}_{-2,-1} \oplus \mathbf{1}_{-1,1}.$$

- \star Singularities of representations with charges belonging to any subspace of \vec{q} 's can be excluded together

Group-Theoretic Classification of Local Fibrations

- Recall that the new light degrees of freedom arising from an ALE-fibration are due to 2-branes wrapping vanishing cycles within the ALE-fibres.
- The correspondence between two-cycles in the ALE-space \widehat{G} and the **root lattice** of the ADE-group G then leads to the following:
 - A generic $\star \widehat{G}$ -fibred Calabi-Yau four-fold or \widehat{G} -fibred G_2 manifold for which the typical fibre has a singularity of type $H \subset G$ will have one matter-singularity for each vector-like representation in the branching:

$$\text{adj}(G) = \text{adj} \left(H \times \prod_{i=1}^k U_1 \right) \oplus (\mathbf{R}_{\vec{q}} \oplus \overline{\mathbf{R}}_{-\vec{q}}).$$

for example, $E_8 \supset E_6 \times U_1^a \times U_1^b$:

$$248 = 78_{0,0} \oplus 1_{0,0} \oplus 1_{0,0} \oplus 27_{1,0} \oplus 27_{0,1} \oplus 27_{-1,-1} \oplus 1_{1,2} \oplus 1_{2,1} \oplus 1_{1,-1} \\ \oplus \overline{27}_{-1,0} \oplus \overline{27}_{0,-1} \oplus \overline{27}_{1,1} \oplus 1_{-1,-2} \oplus 1_{-2,-1} \oplus 1_{-1,1}.$$

- \star Singularities of representations with charges belonging to any subspace of \vec{q} 's can be excluded together
- \star It is also possible to take any quotient of the space of \vec{q} 's.

The *Minimality* of E_8 and the Origin of Three Families

- Because the structure of ALE-fibres plays such a critical role in phenomenology, we are naturally led to ask:
 - What is the **lowest-rank** ALE-fibration $\widehat{G}(f_1(W), \dots, f_n(W))$ that is capable of (typically*) giving the interactions of **even a single family**?

The *Minimality* of E_8 and the Origin of Three Families

- Because the structure of ALE-fibres plays such a critical role in phenomenology, we are naturally led to ask:
 - What is the **lowest-rank** ALE-fibration $\widehat{G}(f_1(W), \dots, f_n(W))$ that is capable of (typically*) giving the interactions of **even a single family**?

The Minimality of E_8 and the Origin of Three Families

- Because the structure of ALE-fibres plays such a critical role in phenomenology, we are naturally led to ask:
 - What is the **lowest-rank** ALE-fibration $\widehat{G}(f_1(W), \dots, f_n(W))$ that is capable of (typically*) giving the interactions of **even a single family**?

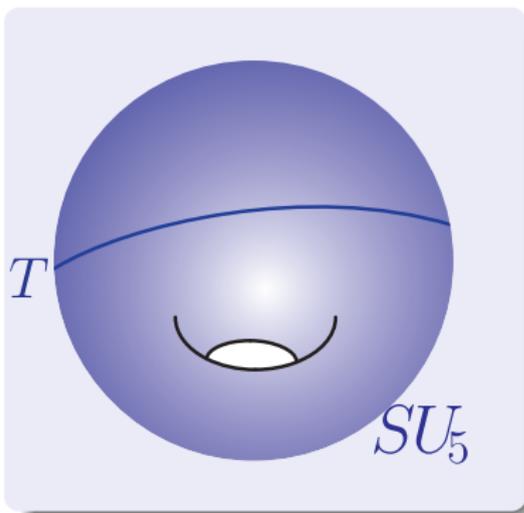


Spectrum of Matter-Curves



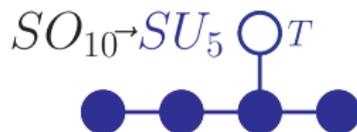
The Minimality of E_8 and the Origin of Three Families

- Because the structure of ALE-fibres plays such a critical role in phenomenology, we are naturally led to ask:
 - What is the **lowest-rank** ALE-fibration $\widehat{G}(f_1(W), \dots, f_n(W))$ that is capable of (typically*) giving the interactions of **even a single family**?



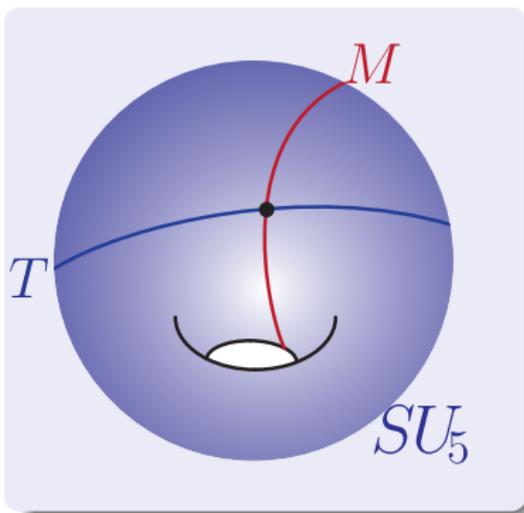
Spectrum of Matter-Curves

$$\frac{10}{T}$$



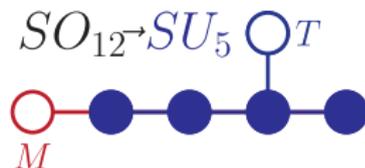
The Minimality of E_8 and the Origin of Three Families

- Because the structure of ALE-fibres plays such a critical role in phenomenology, we are naturally led to ask:
 - What is the **lowest-rank** ALE-fibration $\widehat{G}(f_1(W), \dots, f_n(W))$ that is capable of (typically*) giving the interactions of **even a single family**?



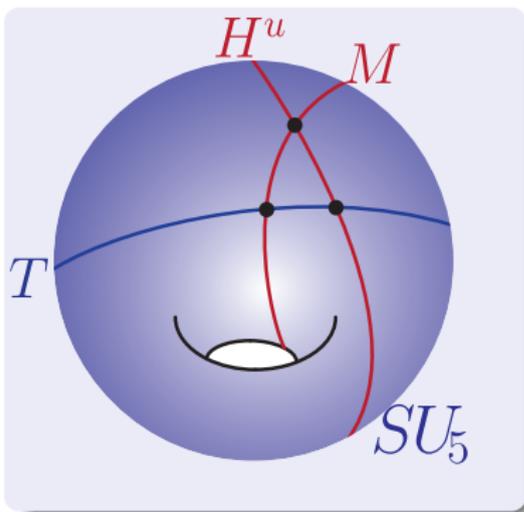
Spectrum of Matter-Curves

$$\frac{10}{T} \quad \frac{\bar{5}}{M}$$



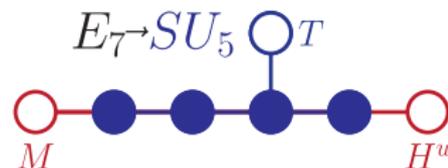
The Minimality of E_8 and the Origin of Three Families

- Because the structure of ALE-fibres plays such a critical role in phenomenology, we are naturally led to ask:
 - What is the **lowest-rank** ALE-fibration $\widehat{G}(f_1(W), \dots, f_n(W))$ that is capable of (typically*) giving the interactions of **even a single family**?



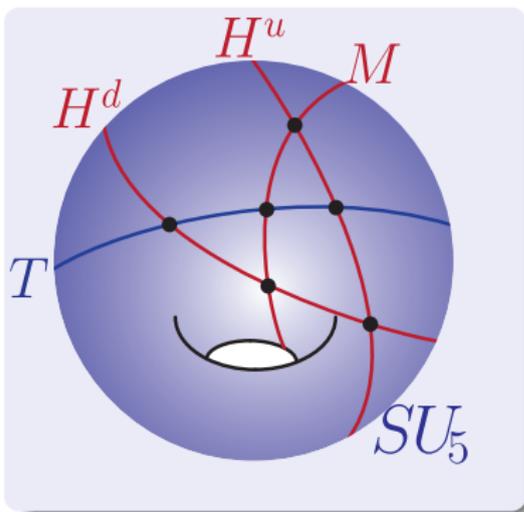
Spectrum of Matter-Curves

$$\frac{10}{T} \quad \frac{\bar{5}}{M} \quad \frac{5}{H^u}$$



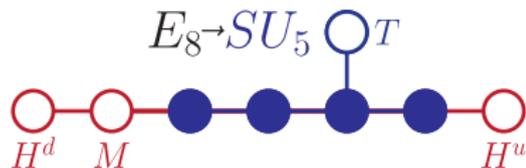
The Minimality of E_8 and the Origin of Three Families

- Because the structure of ALE-fibres plays such a critical role in phenomenology, we are naturally led to ask:
 - What is the **lowest-rank** ALE-fibration $\widehat{G}(f_1(W), \dots, f_n(W))$ that is capable of (typically*) giving the interactions of **even a single family**?



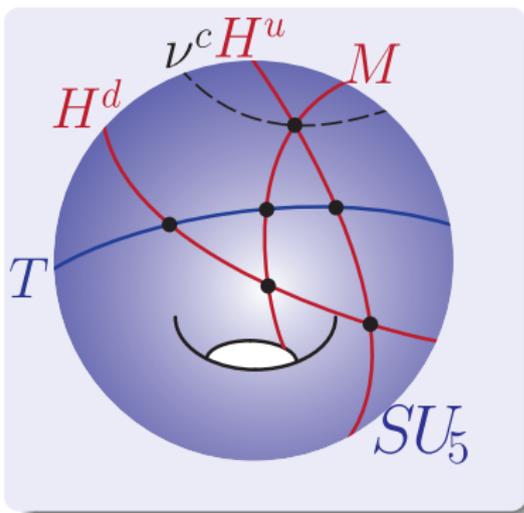
Spectrum of Matter-Curves

$$\frac{10}{T} \quad \frac{\bar{5}}{M} \quad \frac{5}{H^u} \quad \frac{\bar{5}}{H^d}$$



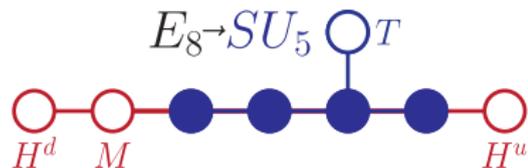
The Minimality of E_8 and the Origin of Three Families

- Because the structure of ALE-fibres plays such a critical role in phenomenology, we are naturally led to ask:
 - What is the **lowest-rank** ALE-fibration $\widehat{G}(f_1(W), \dots, f_n(W))$ that is capable of (typically*) giving the interactions of **even a single family**?



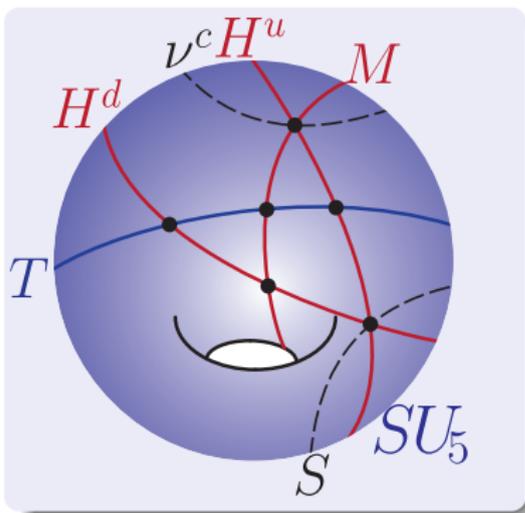
Spectrum of Matter-Curves

$$\frac{10}{T} \quad \frac{\bar{5}}{M} \quad \frac{5}{H^u} \quad \frac{\bar{5}}{H^d} \quad \frac{1}{\nu^c}$$



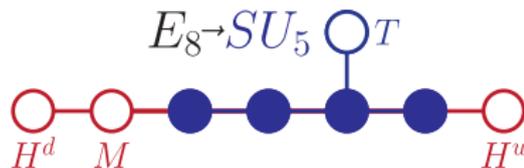
The Minimality of E_8 and the Origin of Three Families

- Because the structure of ALE-fibres plays such a critical role in phenomenology, we are naturally led to ask:
 - What is the **lowest-rank** ALE-fibration $\widehat{G}(f_1(W), \dots, f_n(W))$ that is capable of (typically*) giving the interactions of **even a single family**?



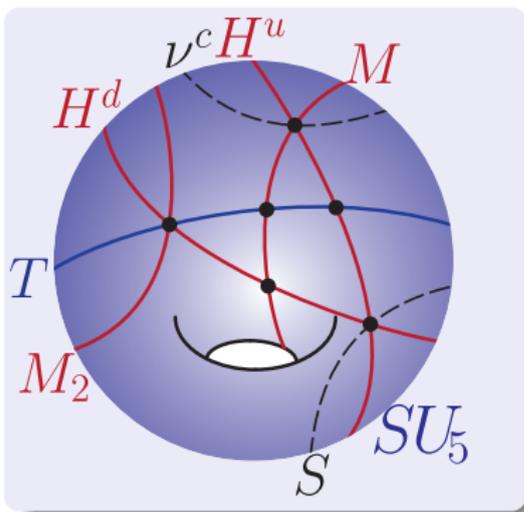
Spectrum of Matter-Curves

$$\frac{10}{T} \quad \frac{\bar{5}}{M} \quad \frac{5}{H^u} \quad \frac{\bar{5}}{H^d} \quad \frac{1}{\nu^c} \quad \frac{1}{S}$$



The Minimality of E_8 and the Origin of Three Families

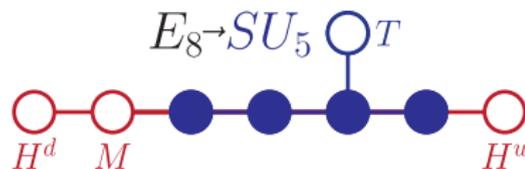
- Because the structure of ALE-fibres plays such a critical role in phenomenology, we are naturally led to ask:
 - What is the **lowest-rank** ALE-fibration $\widehat{G}(f_1(W), \dots, f_n(W))$ that is capable of (typically*) giving the interactions of **even a single family**?



Spectrum of Matter-Curves

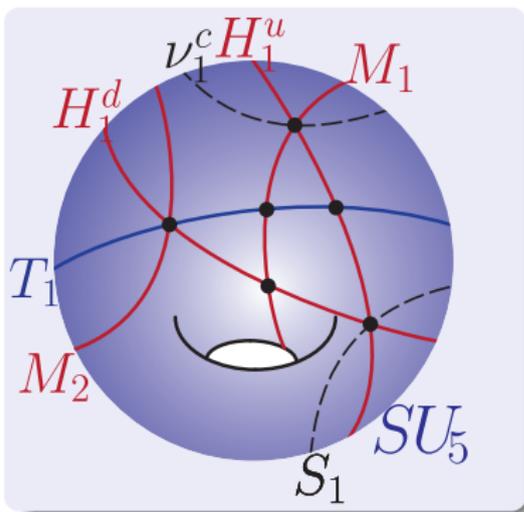
$$\frac{10}{T} \quad \frac{\bar{5}}{M} \quad \frac{5}{H^u} \quad \frac{\bar{5}}{H^d} \quad \frac{1}{\nu^c} \quad \frac{1}{S}$$

M_2



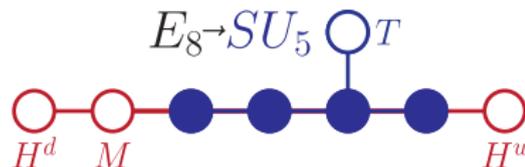
The Minimality of E_8 and the Origin of Three Families

- Because the structure of ALE-fibres plays such a critical role in phenomenology, we are naturally led to ask:
 - What is the **lowest-rank** ALE-fibration $\widehat{G}(f_1(W), \dots, f_n(W))$ that is capable of (typically*) giving the interactions of **even a single family**?



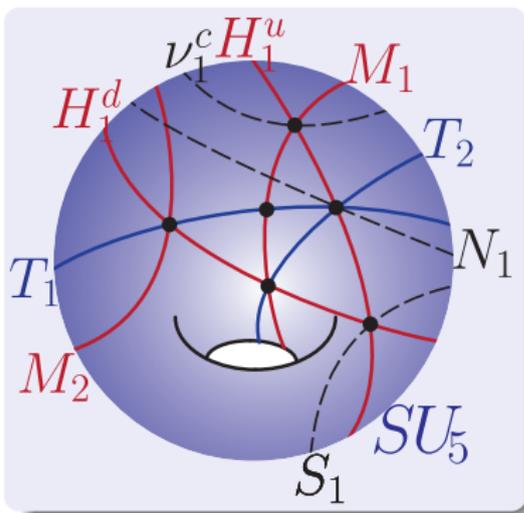
Spectrum of Matter-Curves

$\frac{10}{T_1}$	$\frac{\bar{5}}{M_1}$	$\frac{5}{H_1^u}$	$\frac{\bar{5}}{H_1^d}$	$\frac{1}{\nu_1^c}$	$\frac{1}{S_1}$
	M_2				



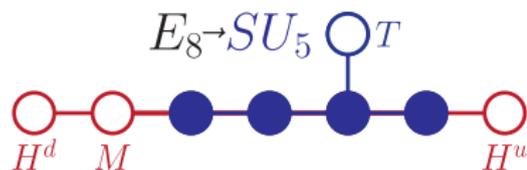
The Minimality of E_8 and the Origin of Three Families

- Because the structure of ALE-fibres plays such a critical role in phenomenology, we are naturally led to ask:
 - What is the **lowest-rank** ALE-fibration $\widehat{G}(f_1(W), \dots, f_n(W))$ that is capable of (typically*) giving the interactions of **even a single family**?



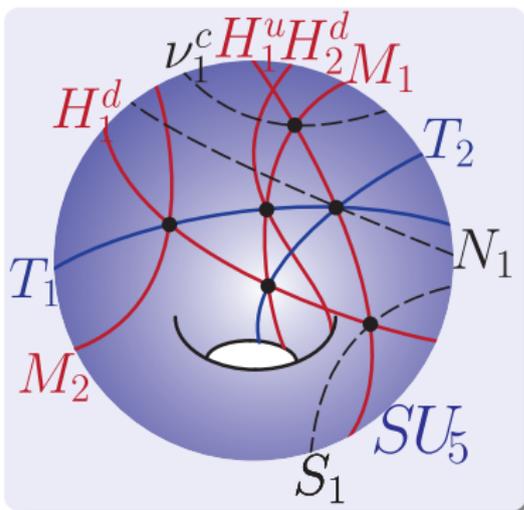
Spectrum of Matter-Curves

$\mathbf{10}$	$\bar{\mathbf{5}}$	$\mathbf{5}$	$\bar{\mathbf{5}}$	$\mathbf{1}$	$\mathbf{1}$	$\mathbf{1}$
T_1	M_1	H_1^u	H_1^d	ν_1^c	S_1	N_1
T_2	M_2					



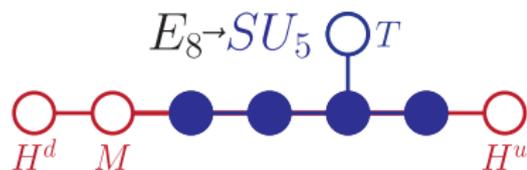
The Minimality of E_8 and the Origin of Three Families

- Because the structure of ALE-fibres plays such a critical role in phenomenology, we are naturally led to ask:
 - What is the **lowest-rank** ALE-fibration $\widehat{G}(f_1(W), \dots, f_n(W))$ that is capable of (typically*) giving the interactions of **even a single family**?



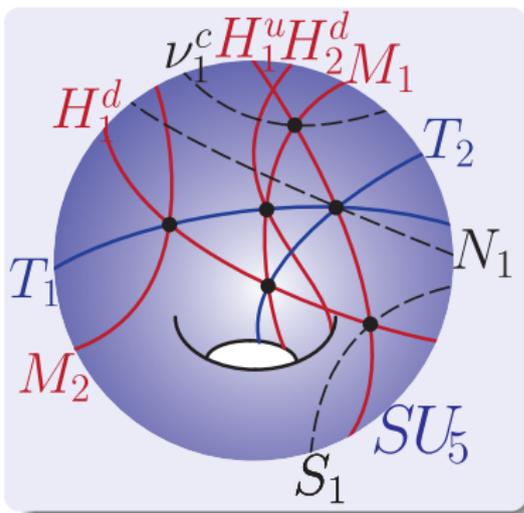
Spectrum of Matter-Curves

$\mathbf{10}$	$\bar{\mathbf{5}}$	$\mathbf{5}$	$\bar{\mathbf{5}}$	$\mathbf{1}$	$\mathbf{1}$	$\mathbf{1}$
T_1	M_1	H_1^u	H_1^d	ν_1^c	S_1	N_1
T_2	M_2		H_2^d			



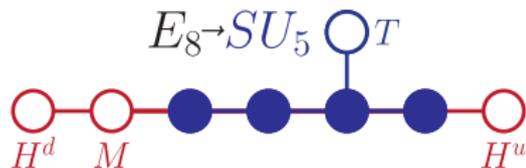
The Minimality of E_8 and the Origin of Three Families

- Because the structure of ALE-fibres plays such a critical role in phenomenology, we are naturally led to ask:
 - What is the **lowest-rank** ALE-fibration $\widehat{G}(f_1(W), \dots, f_n(W))$ that is capable of (typically*) giving the interactions of **even a single family**?



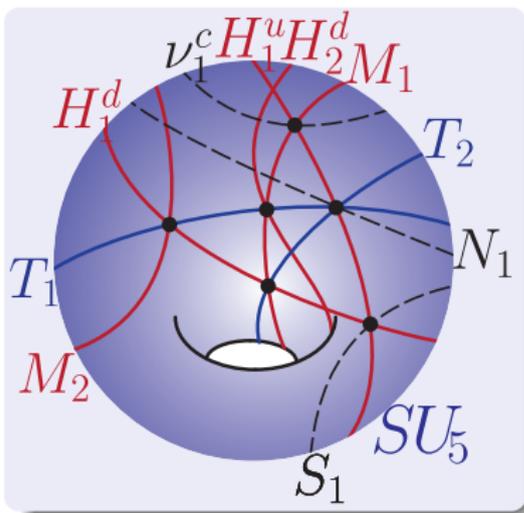
Spectrum of Matter-Curves

<u>10</u>	<u>$\bar{5}$</u>	<u>5</u>	<u>$\bar{5}$</u>	<u>1</u>	<u>1</u>	<u>1</u>
T_1	M_1	H_1^u	H_1^d	ν_1^c	S_1	N_1
T_2	M_2	H_2^u	H_2^d	ν_2^c	S_2	N_2
T_3	M_3					
T_X						



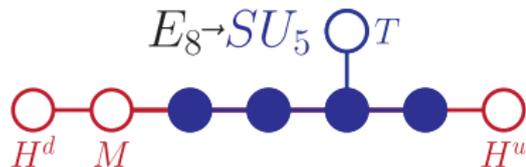
The Minimality of E_8 and the Origin of Three Families

- Because the structure of ALE-fibres plays such a critical role in phenomenology, we are naturally led to ask:
 - What is the **lowest-rank** ALE-fibration $\widehat{G}(f_1(W), \dots, f_n(W))$ that is capable of (typically*) giving the interactions of **even a single family**?



Spectrum of Matter-Curves

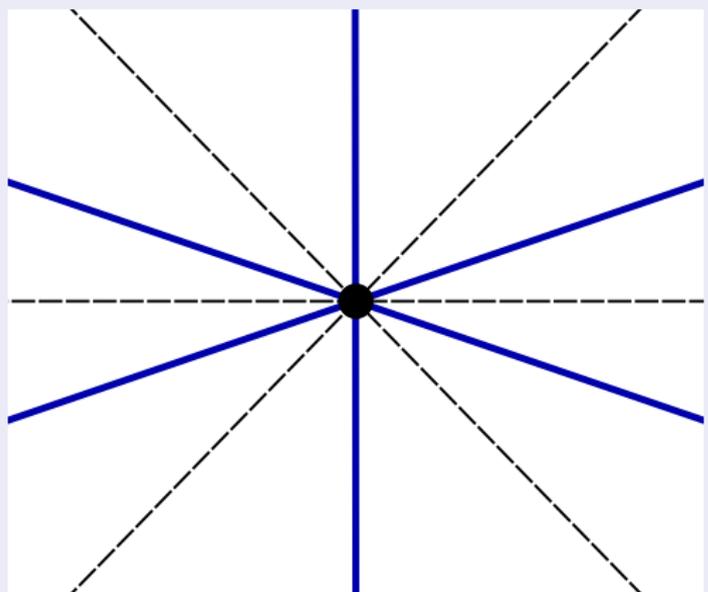
<u>10</u>	<u>$\bar{5}$</u>	<u>5</u>	<u>$\bar{5}$</u>	<u>1</u>	<u>1</u>	<u>1</u>
T_1	M_1	H_1^u	H_1^d	ν_1^c	S_1	N_1
T_2	M_2	H_2^u	H_2^d	ν_2^c	S_2	N_2
T_3	M_3	H_3^u	H_3^d	ν_3^c	S_3	N_3
T_X	M_X			ν_X^c		
T_X^c						



'Unfolding' Three Generations out of E_8

The matter content of a general \widehat{E}_8 -fibred geometry can be derived by thinking in terms of sequential 'unfoldings':

Unfolding $E_6 \rightarrow SO_{10}$

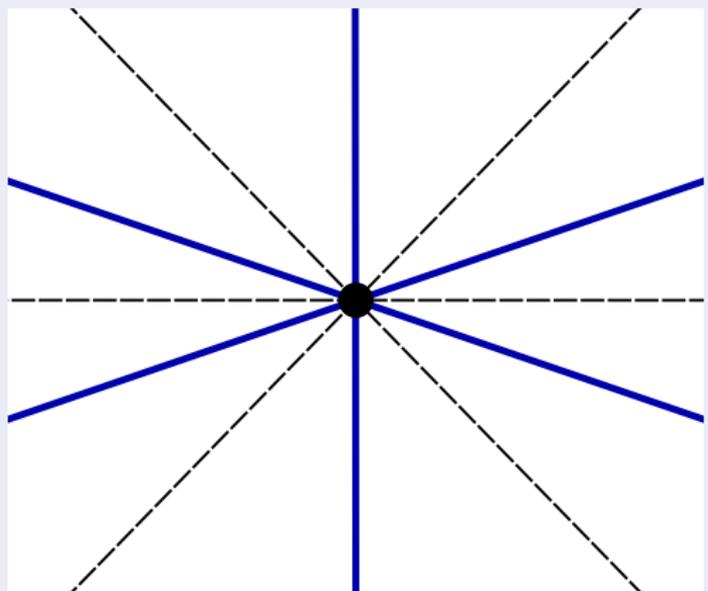


	$E_6 \times$	$U_1^a \times$	U_1^b
T_1	27	1	1
T_2	27	1	-1
T_3	27	-2	0
S_1	1	3	-1
S_3	1	3	1
S_2	1	0	2

'Unfolding' Three Generations out of E_8

The matter content of a general \widehat{E}_8 -fibred geometry can be derived by thinking in terms of sequential 'unfoldings':

Unfolding $E_6 \rightarrow SO_{10}$

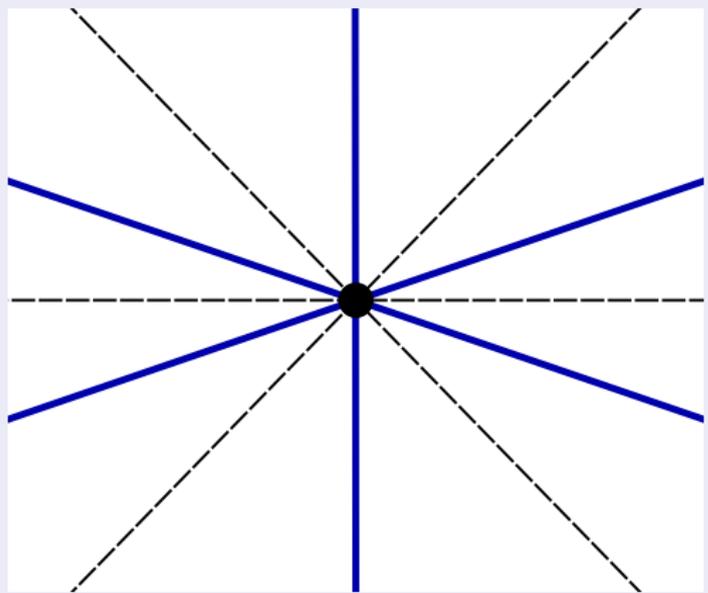


	$E_6 \times$	$U_1^a \times$	U_1^b
T_1	27	1	1
T_2	27	1	-1
T_3	27	-2	0
S_1	1	3	-1
S_3	1	3	1
S_2	1	0	2

'Unfolding' Three Generations out of E_8

The matter content of a general \widehat{E}_8 -fibred geometry can be derived by thinking in terms of sequential 'unfoldings':

Unfolding $E_6 \rightarrow SO_{10}$

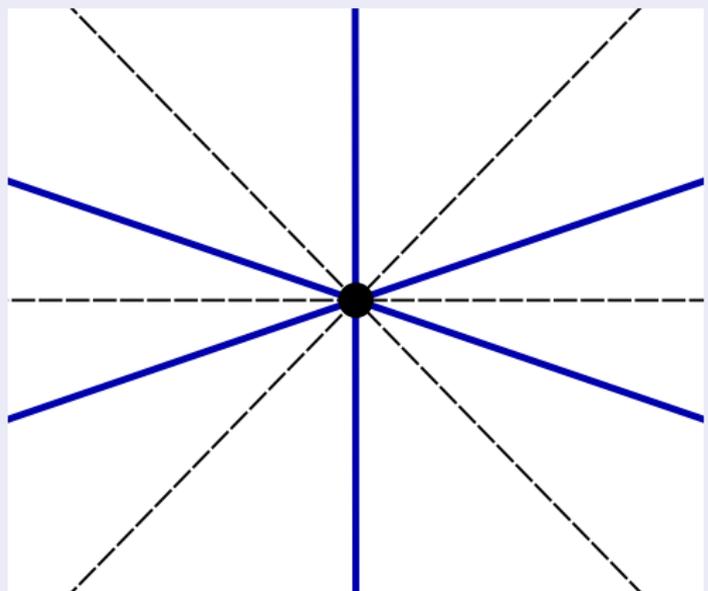


	$E_6 \times$	$U_1^a \times$	U_1^b
T_1	27	1	1
T_2	27	1	-1
T_3	27	-2	0
S_1	1	3	-1
S_3	1	3	1
S_2	1	0	2

'Unfolding' Three Generations out of E_8

The matter content of a general \widehat{E}_8 -fibred geometry can be derived by thinking in terms of sequential 'unfoldings':

Unfolding $E_6 \rightarrow SO_{10}$

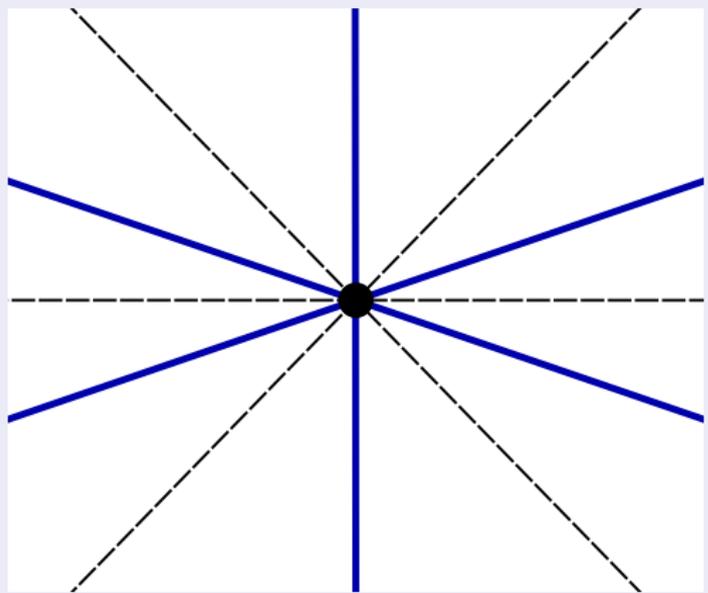


	$E_6 \times$	$U_1^a \times$	U_1^b
T_1	27	1	1
T_2	27	1	-1
T_3	27	-2	0
S_1	1	3	-1
S_3	1	3	1
S_2	1	0	2

'Unfolding' Three Generations out of E_8

The matter content of a general \widehat{E}_8 -fibred geometry can be derived by thinking in terms of sequential 'unfoldings':

Unfolding $E_6 \rightarrow SO_{10}$

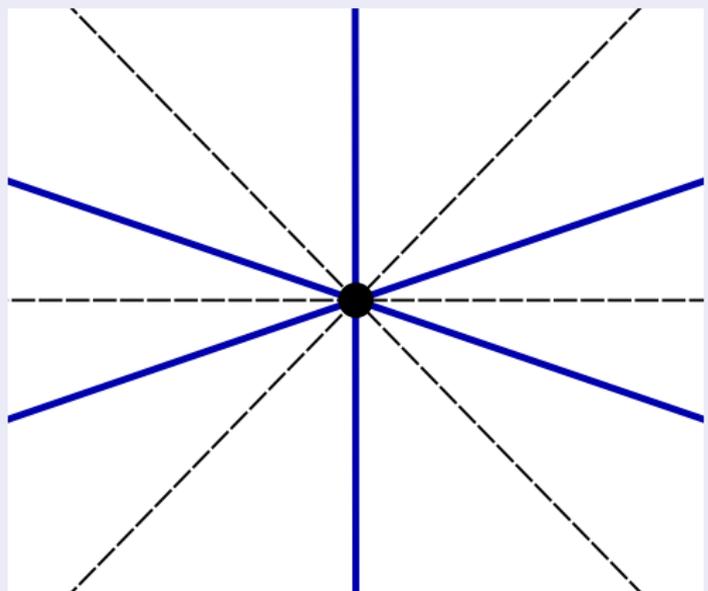


	$E_6 \times$	$U_1^a \times$	U_1^b
T_1	27	1	1
T_2	27	1	-1
T_3	27	-2	0
S_1	1	3	-1
S_3	1	3	1
S_2	1	0	2

'Unfolding' Three Generations out of E_8

The matter content of a general \widehat{E}_8 -fibred geometry can be derived by thinking in terms of sequential 'unfoldings':

Unfolding $E_6 \rightarrow SO_{10}$

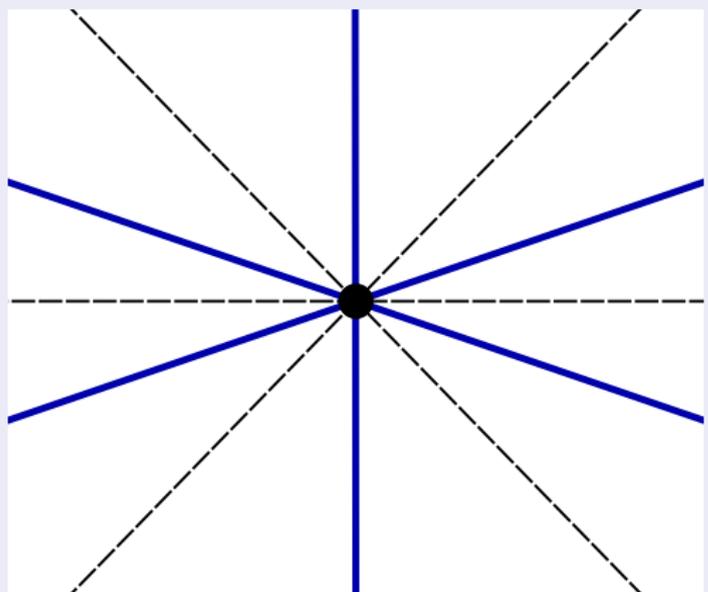


	$E_6 \times$	$U_1^a \times$	U_1^b
T_1	27	1	1
T_2	27	1	-1
T_3	27	-2	0
S_1	1	3	-1
S_3	1	3	1
S_2	1	0	2

'Unfolding' Three Generations out of E_8

The matter content of a general \widehat{E}_8 -fibred geometry can be derived by thinking in terms of sequential 'unfoldings':

Unfolding $E_6 \rightarrow SO_{10}$

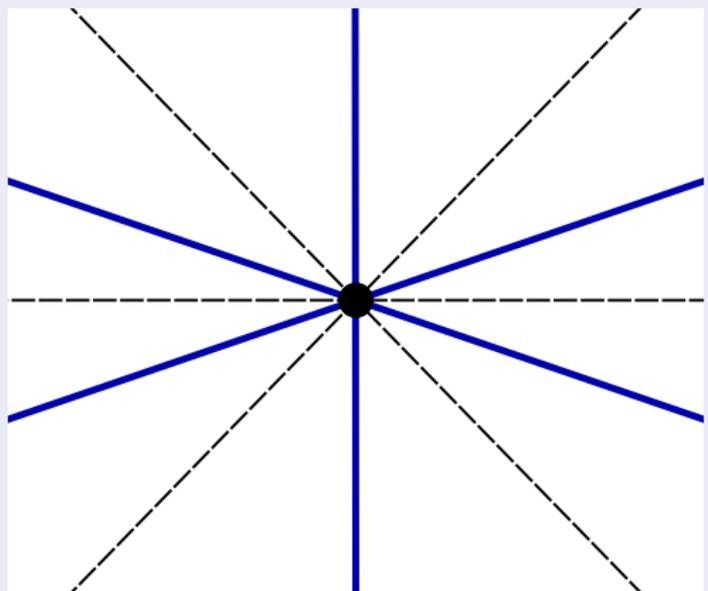


	$E_6 \times$	$U_1^a \times$	U_1^b
T_1	27	1	1
T_2	27	1	-1
T_3	27	-2	0
S_1	1	3	-1
S_3	1	3	1
S_2	1	0	2

'Unfolding' Three Generations out of E_8

The matter content of a general \widehat{E}_8 -fibred geometry can be derived by thinking in terms of sequential 'unfoldings':

Unfolding $E_6 \rightarrow SO_{10}$

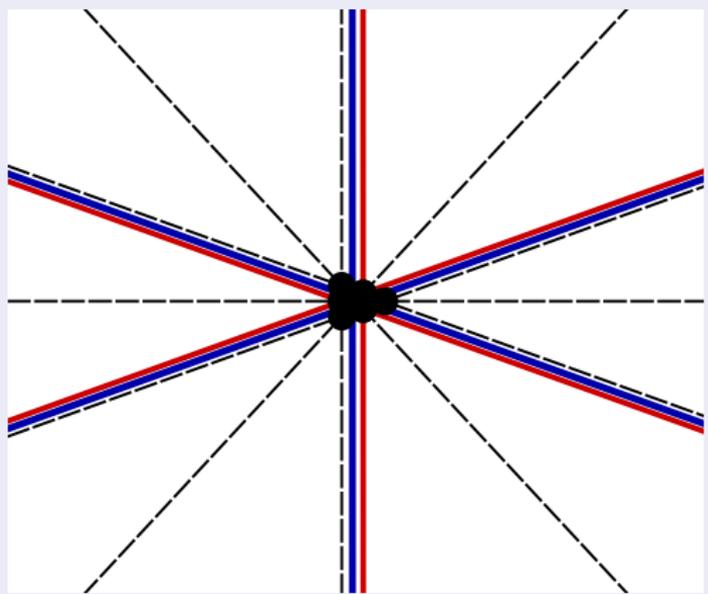


	$E_6 \times$	$U_1^a \times$	U_1^b
T_1	27	1	1
T_2	27	1	-1
T_3	27	-2	0
S_1	1	3	-1
S_3	1	3	1
S_2	1	0	2

'Unfolding' Three Generations out of E_8

The matter content of a general \widehat{E}_8 -fibred geometry can be derived by thinking in terms of sequential 'unfoldings':

Unfolding $E_6 \rightarrow SO_{10}$

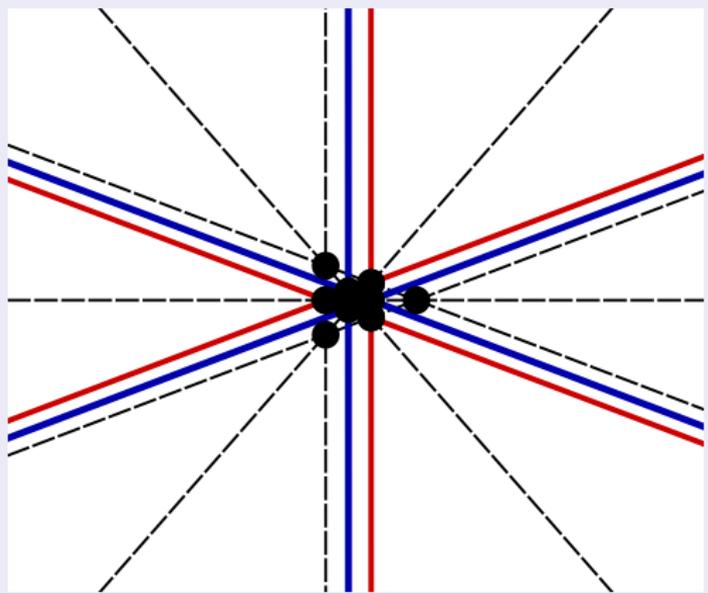


		$SO_{10} \times U_1^a \times U_1^b \times U_1^c$		
T_1	16	1	1	-1
T_2	16	1	-1	-1
T_3	16	-2	0	-1
T_X^c	$\overline{16}$	0	0	-3
H	10	1	1	2
Y_a	10	-1	1	-2
Y_b	10	2	0	-2
X_1	1	-1	-1	4
X_2	1	-1	1	4
N_2^c	1	-2	0	-4
N_3^c	1	3	-1	0
S_1	1	3	1	0
S_2	1	0	2	0

'Unfolding' Three Generations out of E_8

The matter content of a general \widehat{E}_8 -fibred geometry can be derived by thinking in terms of sequential 'unfoldings':

Unfolding $E_6 \rightarrow SO_{10}$

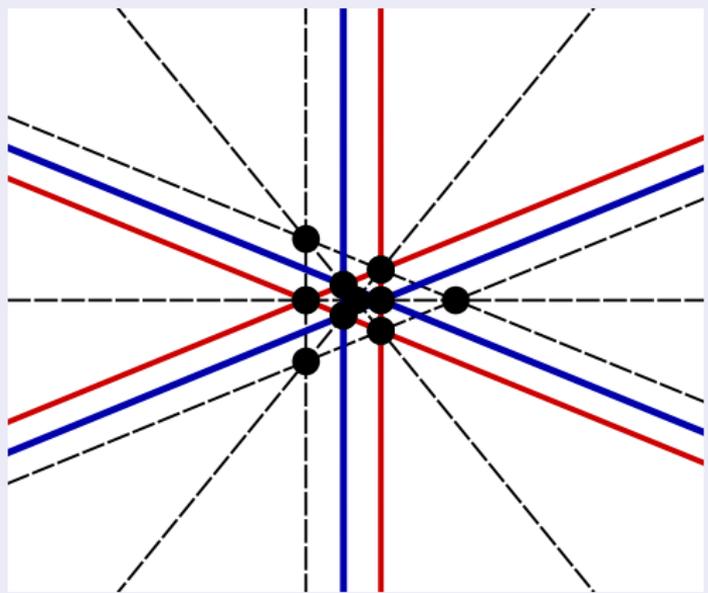


		$SO_{10} \times U_1^a \times U_1^b \times U_1^c$		
T_1	16	1	1	-1
T_2	16	1	-1	-1
T_3	16	-2	0	-1
T_X^c	$\overline{16}$	0	0	-3
H	10	1	1	2
Y_a	10	-1	1	-2
Y_b	10	2	0	-2
X_1	1	-1	-1	4
X_2	1	-1	1	4
N_2^c	1	-2	0	-4
N_3^c	1	3	-1	0
S_1	1	3	1	0
S_2	1	0	2	0

'Unfolding' Three Generations out of E_8

The matter content of a general \widehat{E}_8 -fibred geometry can be derived by thinking in terms of sequential 'unfoldings':

Unfolding $E_6 \rightarrow SO_{10}$

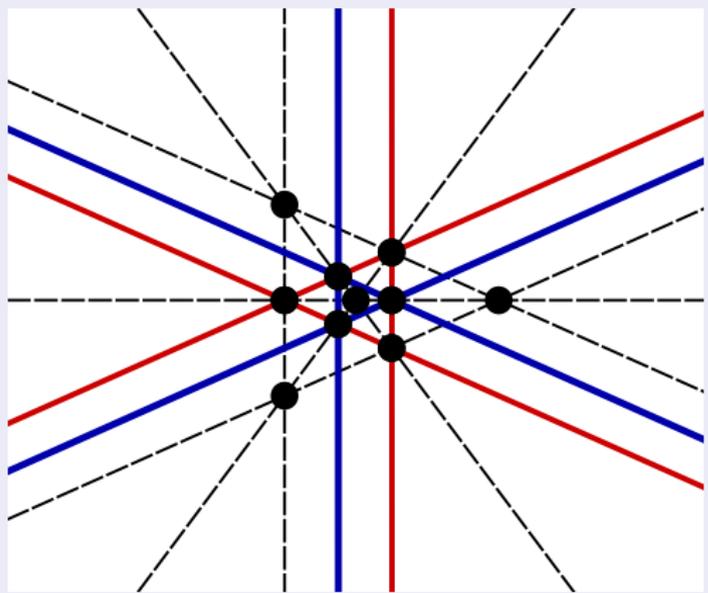


		$SO_{10} \times U_1^a \times U_1^b \times U_1^c$		
T_1	16	1	1	-1
T_2	16	1	-1	-1
T_3	16	-2	0	-1
T_X^c	$\overline{16}$	0	0	-3
H	10	1	1	2
Y_a	10	-1	1	-2
Y_b	10	2	0	-2
X_1	1	-1	-1	4
X_2	1	-1	1	4
N_2^c	1	-2	0	-4
N_3^c	1	3	-1	0
S_1	1	3	1	0
S_2	1	0	2	0

'Unfolding' Three Generations out of E_8

The matter content of a general \widehat{E}_8 -fibred geometry can be derived by thinking in terms of sequential 'unfoldings':

Unfolding $E_6 \rightarrow SO_{10}$

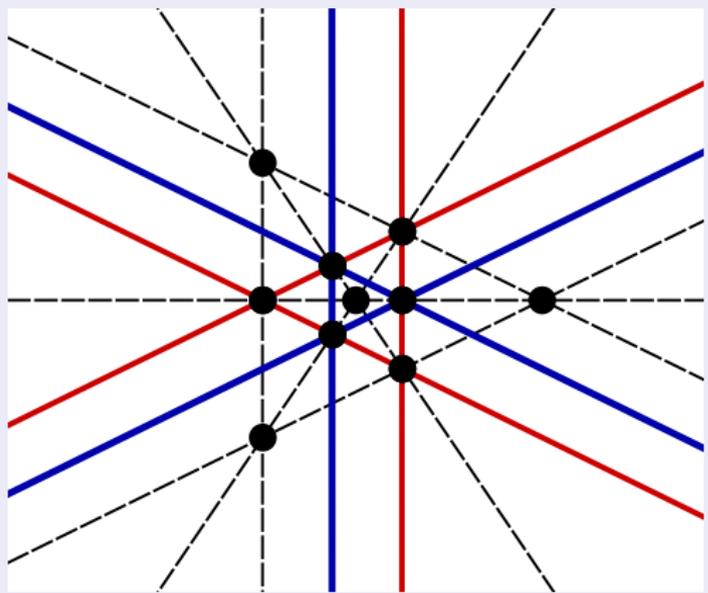


		$SO_{10} \times U_1^a \times U_1^b \times U_1^c$		
T_1	16	1	1	-1
T_2	16	1	-1	-1
T_3	16	-2	0	-1
T_X^c	$\overline{16}$	0	0	-3
H	10	1	1	2
Y_a	10	-1	1	-2
Y_b	10	2	0	-2
X_1	1	-1	-1	4
X_2	1	-1	1	4
N_2^c	1	-2	0	-4
N_3^c	1	3	-1	0
S_1	1	3	1	0
S_2	1	0	2	0

'Unfolding' Three Generations out of E_8

The matter content of a general \widehat{E}_8 -fibred geometry can be derived by thinking in terms of sequential 'unfoldings':

Unfolding $E_6 \rightarrow SO_{10}$

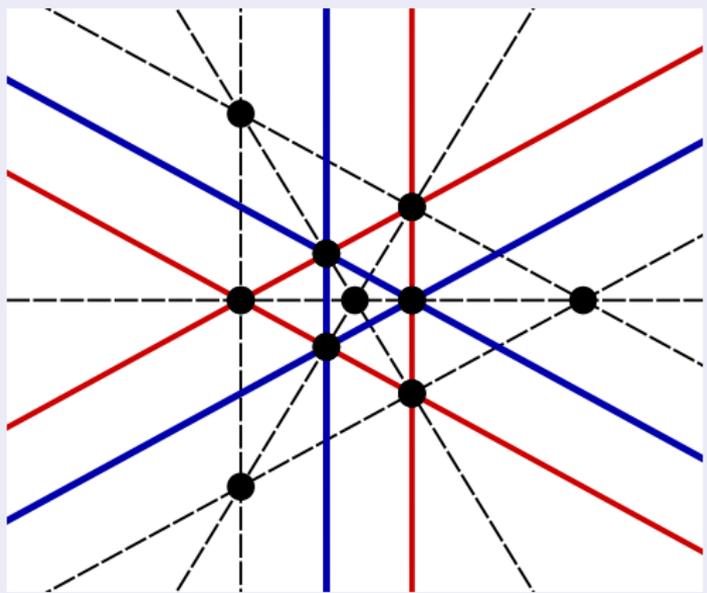


		$SO_{10} \times U_1^a \times U_1^b \times U_1^c$		
T_1	16	1	1	-1
T_2	16	1	-1	-1
T_3	16	-2	0	-1
T_X^c	$\overline{16}$	0	0	-3
H	10	1	1	2
Y_a	10	-1	1	-2
Y_b	10	2	0	-2
X_1	1	-1	-1	4
X_2	1	-1	1	4
N_2^c	1	-2	0	-4
N_3^c	1	3	-1	0
S_1	1	3	1	0
S_2	1	0	2	0

'Unfolding' Three Generations out of E_8

The matter content of a general \widehat{E}_8 -fibred geometry can be derived by thinking in terms of sequential 'unfoldings':

Unfolding $E_6 \rightarrow SO_{10}$

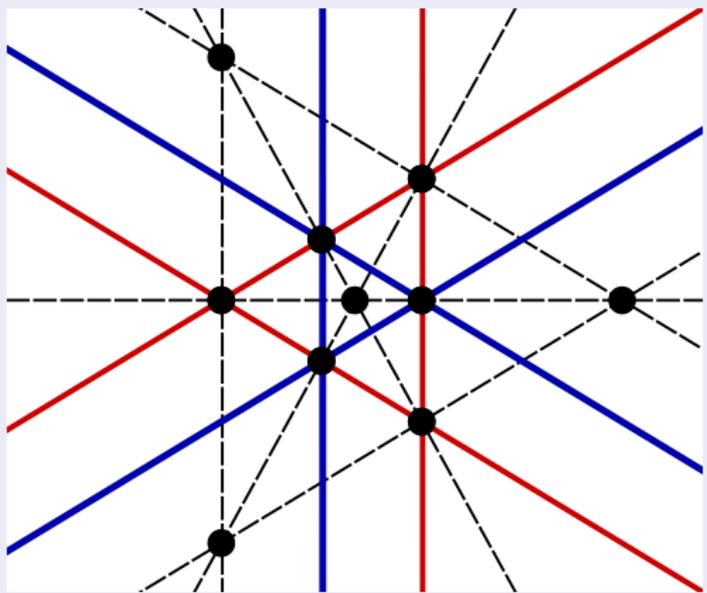


		$SO_{10} \times U_1^a \times U_1^b \times U_1^c$		
T_1	16	1	1	-1
T_2	16	1	-1	-1
T_3	16	-2	0	-1
T_X^c	$\overline{16}$	0	0	-3
H	10	1	1	2
Y_a	10	-1	1	-2
Y_b	10	2	0	-2
X_1	1	-1	-1	4
X_2	1	-1	1	4
N_2^c	1	-2	0	-4
N_3^c	1	3	-1	0
S_1	1	3	1	0
S_2	1	0	2	0

'Unfolding' Three Generations out of E_8

The matter content of a general \widehat{E}_8 -fibred geometry can be derived by thinking in terms of sequential 'unfoldings':

Unfolding $E_6 \rightarrow SO_{10}$

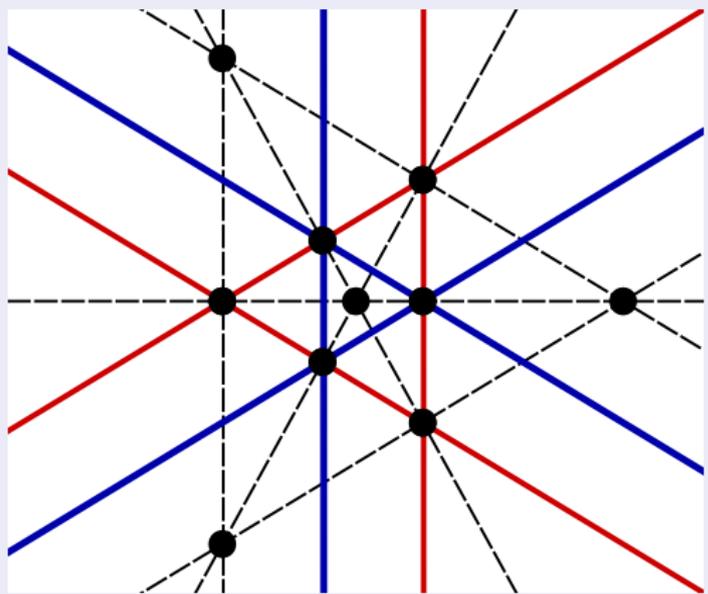


		$SO_{10} \times U_1^a \times U_1^b \times U_1^c$		
T_1	16	1	1	-1
T_2	16	1	-1	-1
T_3	16	-2	0	-1
T_X^c	$\overline{16}$	0	0	-3
H	10	1	1	2
Y_a	10	-1	1	-2
Y_b	10	2	0	-2
X_1	1	-1	-1	4
X_2	1	-1	1	4
N_2^c	1	-2	0	-4
N_3^c	1	3	-1	0
S_1	1	3	1	0
S_2	1	0	2	0

'Unfolding' Three Generations out of E_8

The matter content of a general \widehat{E}_8 -fibred geometry can be derived by thinking in terms of sequential 'unfoldings':

Unfolding $E_6 \rightarrow SO_{10}$

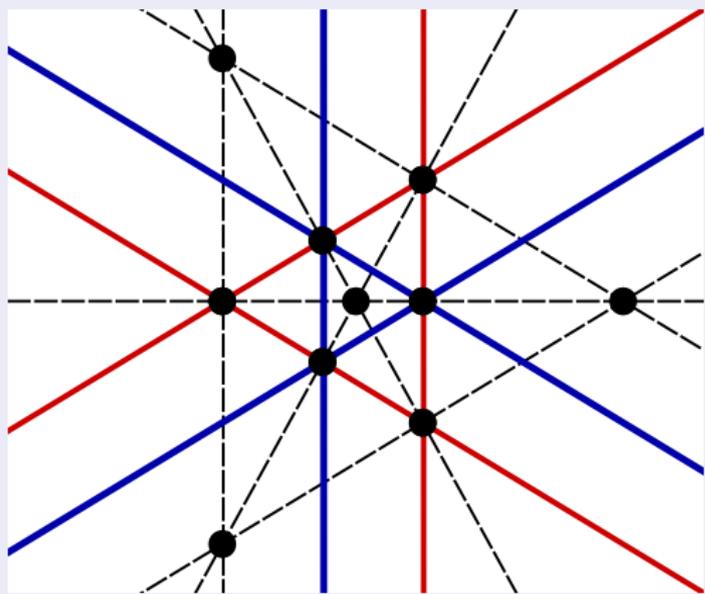


		$SO_{10} \times U_1^a \times U_1^b \times U_1^c$		
T_1	16	1	1	-1
T_2	16	1	-1	-1
T_3	16	-2	0	-1
T_X^c	$\overline{16}$	0	0	-3
H	10	1	1	2
Y_a	10	-1	1	-2
Y_b	10	2	0	-2
X_1	1	-1	-1	4
X_2	1	-1	1	4
N_2^c	1	-2	0	-4
N_3^c	1	3	-1	0
S_1	1	3	1	0
S_2	1	0	2	0

'Unfolding' Three Generations out of E_8

The matter content of a general \widehat{E}_8 -fibred geometry can be derived by thinking in terms of sequential 'unfoldings':

Unfolding $E_6 \rightarrow SO_{10}$

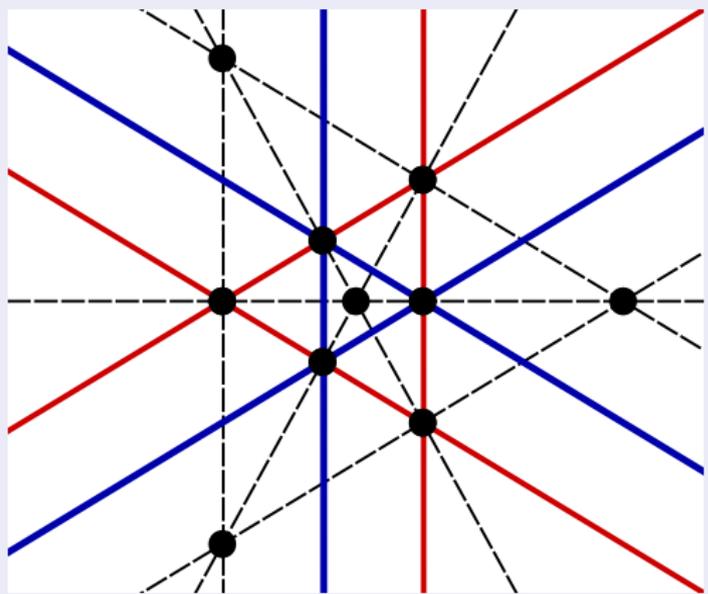


		SO_{10}	U_1^a	U_1^b	U_1^c
T_1	16	1	1	-1	
T_2	16	1	-1	-1	
T_3	16	-2	0	-1	
T_X^c	$\overline{16}$	0	0	-3	
H	10	1	1	2	
Y_a	10	-1	1	-2	
Y_b	10	2	0	-2	
X_1	1	-1	-1	4	
X_2	1	-1	1	4	
N_2^c	1	-2	0	-4	
N_3^c	1	3	-1	0	
S_1	1	3	1	0	
S_2	1	0	2	0	

'Unfolding' Three Generations out of E_8

The matter content of a general \widehat{E}_8 -fibred geometry can be derived by thinking in terms of sequential 'unfoldings':

Unfolding $E_6 \rightarrow SO_{10}$

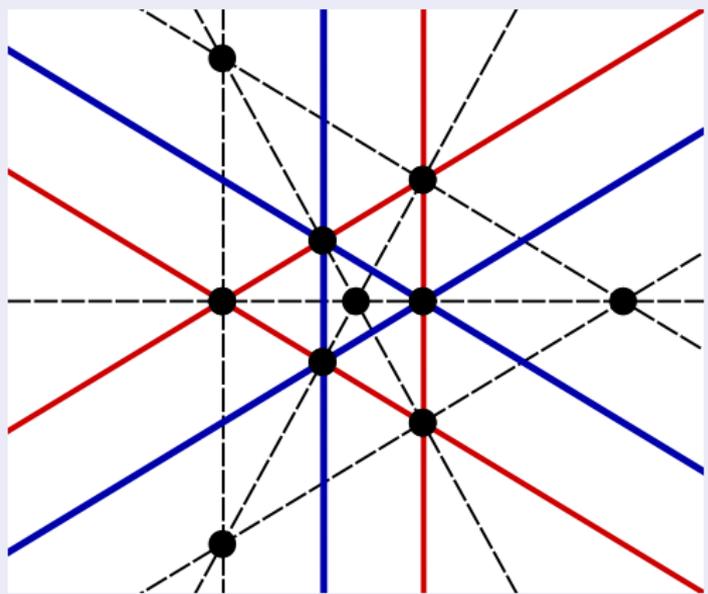


		$SO_{10} \times U_1^a \times U_1^b \times U_1^c$		
T_1	16	1	1	-1
T_2	16	1	-1	-1
T_3	16	-2	0	-1
T_X^c	$\overline{16}$	0	0	-3
H	10	1	1	2
Y_a	10	-1	1	-2
Y_b	10	2	0	-2
X_1	1	-1	-1	4
X_2	1	-1	1	4
N_2^c	1	-2	0	-4
N_3^c	1	3	-1	0
S_1	1	3	1	0
S_2	1	0	2	0

'Unfolding' Three Generations out of E_8

The matter content of a general \widehat{E}_8 -fibred geometry can be derived by thinking in terms of sequential 'unfoldings':

Unfolding $E_6 \rightarrow SO_{10}$

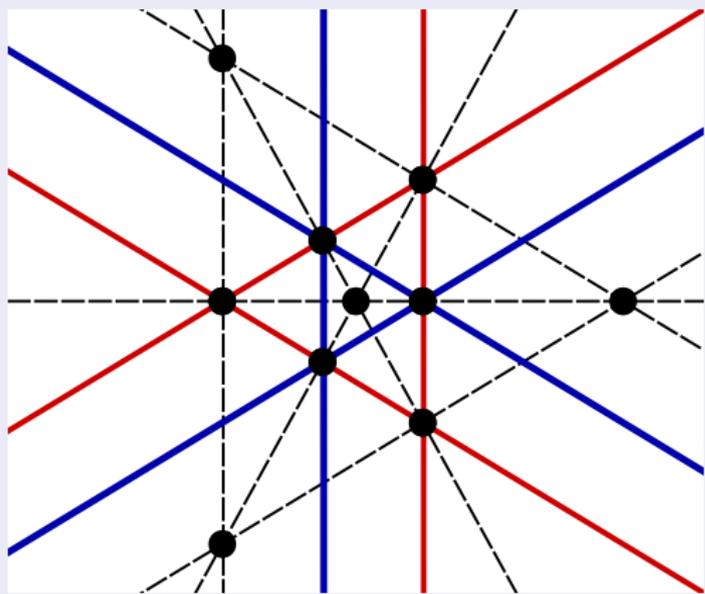


		$SO_{10} \times U_1^a \times U_1^b \times U_1^c$		
T_1	16	1	1	-1
T_2	16	1	-1	-1
T_3	16	-2	0	-1
T_X^c	$\overline{16}$	0	0	-3
H	10	1	1	2
Y_a	10	-1	1	-2
Y_b	10	2	0	-2
X_1	1	-1	-1	4
X_2	1	-1	1	4
N_2^c	1	-2	0	-4
N_3^c	1	3	-1	0
S_1	1	3	1	0
S_2	1	0	2	0

'Unfolding' Three Generations out of E_8

The matter content of a general \widehat{E}_8 -fibred geometry can be derived by thinking in terms of sequential 'unfoldings':

Unfolding $E_6 \rightarrow SO_{10}$

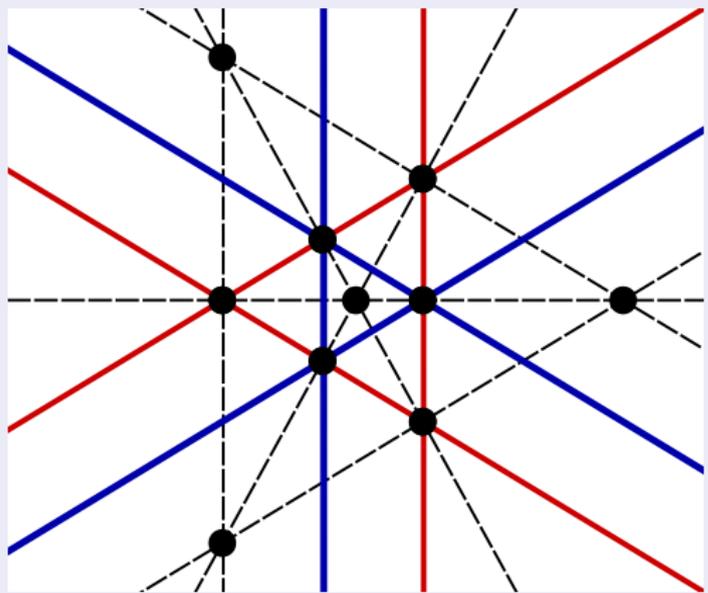


		$SO_{10} \times U_1^a \times U_1^b \times U_1^c$		
T_1	16	1	1	-1
T_2	16	1	-1	-1
T_3	16	-2	0	-1
T_X^c	$\overline{16}$	0	0	-3
H	10	1	1	2
Y_a	10	-1	1	-2
Y_b	10	2	0	-2
X_1	1	-1	-1	4
X_2	1	-1	1	4
N_2^c	1	-2	0	-4
N_3^c	1	3	-1	0
S_1	1	3	1	0
S_2	1	0	2	0

'Unfolding' Three Generations out of E_8

The matter content of a general \widehat{E}_8 -fibred geometry can be derived by thinking in terms of sequential 'unfoldings':

Unfolding $SO_{10} \rightarrow SU_5$

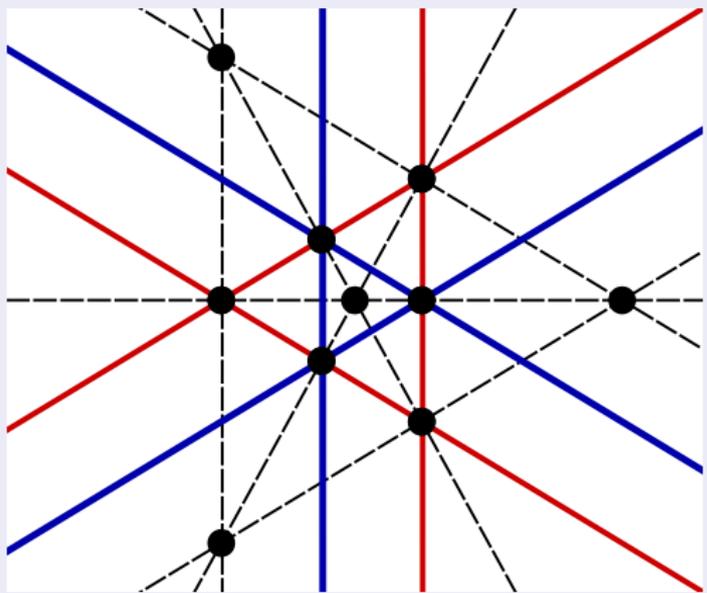


		SO_{10}	U_1^a	U_1^b	U_1^c
T_1	16	1	1	-1	
T_2	16	1	-1	-1	
T_3	16	-2	0	-1	
T_X^c	$\overline{16}$	0	0	-3	
H	10	1	1	2	
Y_a	10	-1	1	-2	
Y_b	10	2	0	-2	
X_1	1	-1	-1	4	
X_2	1	-1	1	4	
N_2^c	1	-2	0	-4	
N_3^c	1	3	-1	0	
S_1	1	3	1	0	
S_2	1	0	2	0	

'Unfolding' Three Generations out of E_8

The matter content of a general \widehat{E}_8 -fibred geometry can be derived by thinking in terms of sequential 'unfoldings':

Unfolding $SO_{10} \rightarrow SU_5$

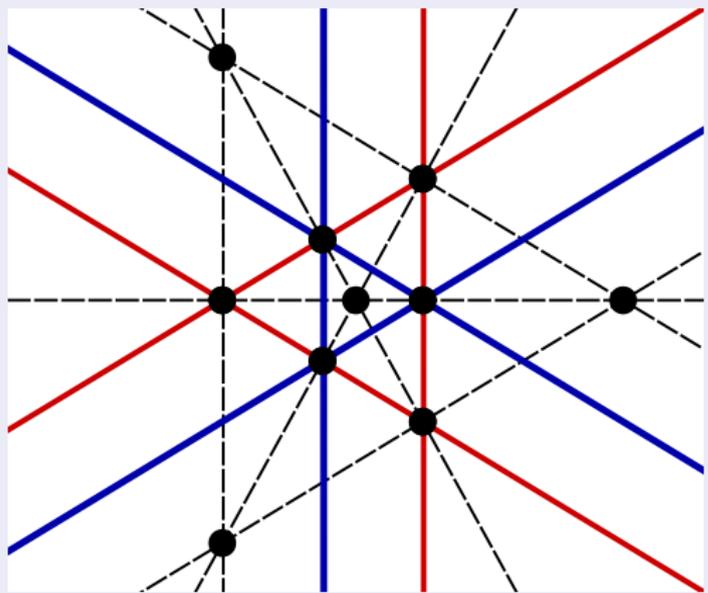


		SO_{10}	U_1^a	U_1^b	U_1^c
T_1	16	1	1	-1	
T_2	16	1	-1	-1	
T_3	16	-2	0	-1	
T_X^c	$\overline{16}$	0	0	-3	
H	10	1	1	2	
Y_a	10	-1	1	-2	
Y_b	10	2	0	-2	
X_1	1	-1	-1	4	
X_2	1	-1	1	4	
N_2^c	1	-2	0	-4	
N_3^c	1	3	-1	0	
S_1	1	3	1	0	
S_2	1	0	2	0	

'Unfolding' Three Generations out of E_8

The matter content of a general \widehat{E}_8 -fibred geometry can be derived by thinking in terms of sequential 'unfoldings':

Unfolding $SO_{10} \rightarrow SU_5$

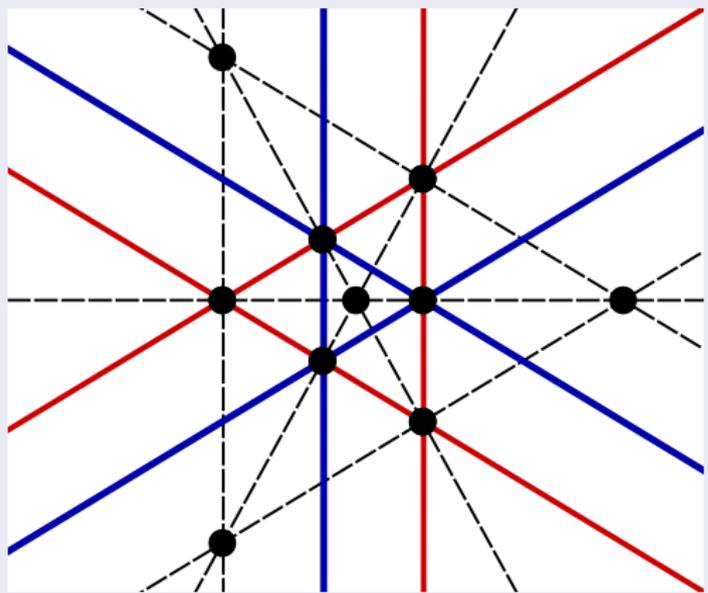


		SO_{10}	U_1^a	U_1^b	U_1^c
T_1	16	1	1	-1	
T_2	16	1	-1	-1	
T_3	16	-2	0	-1	
T_X^c	$\overline{16}$	0	0	-3	
H	10	1	1	2	
Y_a	10	-1	1	-2	
Y_b	10	2	0	-2	
X_1	1	-1	-1	4	
X_2	1	-1	1	4	
N_2^c	1	-2	0	-4	
N_3^c	1	3	-1	0	
S_1	1	3	1	0	
S_2	1	0	2	0	

'Unfolding' Three Generations out of E_8

The matter content of a general \widehat{E}_8 -fibred geometry can be derived by thinking in terms of sequential 'unfoldings':

Unfolding $SO_{10} \rightarrow SU_5$

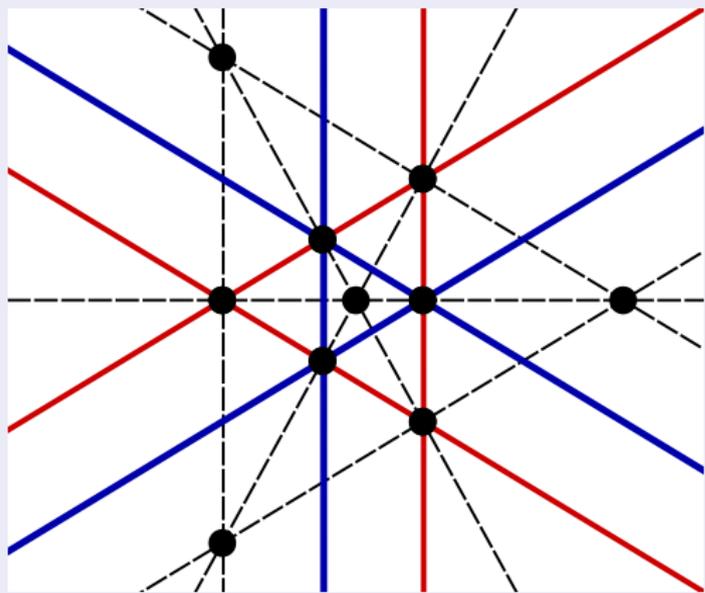


		SO_{10}	U_1^a	U_1^b	U_1^c
T_1	16	1	1	-1	
T_2	16	1	-1	-1	
T_3	16	-2	0	-1	
T_X^c	$\overline{16}$	0	0	-3	
H	10	1	1	2	
Y_a	10	-1	1	-2	
Y_b	10	2	0	-2	
X_1	1	-1	-1	4	
X_2	1	-1	1	4	
N_2^c	1	-2	0	-4	
N_3^c	1	3	-1	0	
S_1	1	3	1	0	
S_2	1	0	2	0	

'Unfolding' Three Generations out of E_8

The matter content of a general \widehat{E}_8 -fibred geometry can be derived by thinking in terms of sequential 'unfoldings':

Unfolding $SO_{10} \rightarrow SU_5$

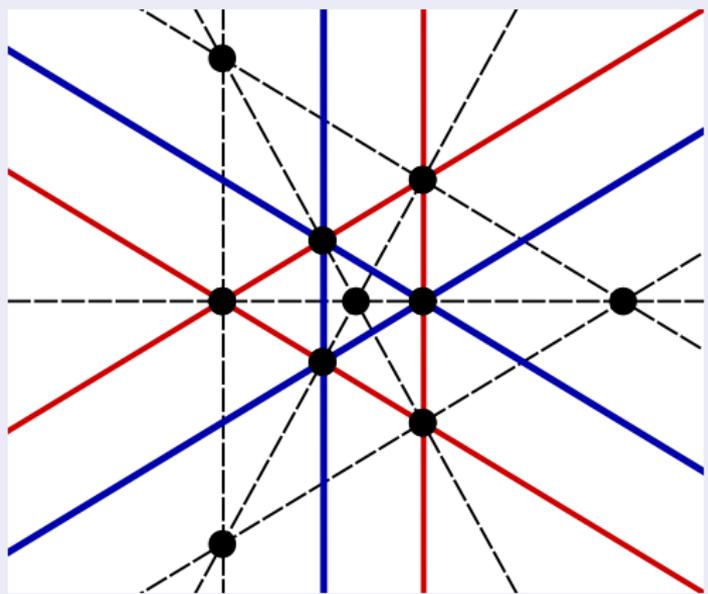


		SO_{10}	U_1^a	U_1^b	U_1^c
T_1	16	1	1	-1	
T_2	16	1	-1	-1	
T_3	16	-2	0	-1	
T_X^c	$\overline{16}$	0	0	-3	
H	10	1	1	2	
Y_a	10	-1	1	-2	
Y_b	10	2	0	-2	
X_1	1	-1	-1	4	
X_2	1	-1	1	4	
N_2^c	1	-2	0	-4	
N_3^c	1	3	-1	0	
S_1	1	3	1	0	
S_2	1	0	2	0	

'Unfolding' Three Generations out of E_8

The matter content of a general \widehat{E}_8 -fibred geometry can be derived by thinking in terms of sequential 'unfoldings':

Unfolding $SO_{10} \rightarrow SU_5$

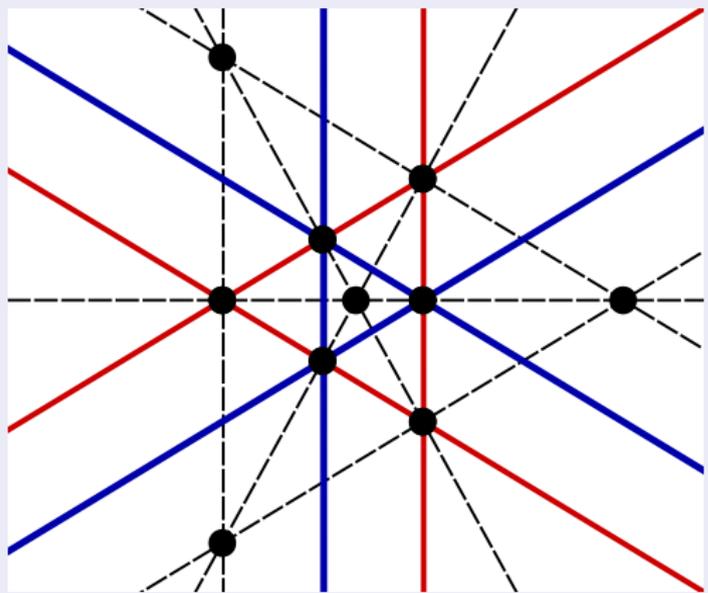


		$SO_{10} \times U_1^a \times U_1^b \times U_1^c$		
T_1	16	1	1	-1
T_2	16	1	-1	-1
T_3	16	-2	0	-1
T_X^c	$\overline{16}$	0	0	-3
H	10	1	1	2
Y_a	10	-1	1	-2
Y_b	10	2	0	-2
X_1	1	-1	-1	4
X_2	1	-1	1	4
N_2^c	1	-2	0	-4
N_3^c	1	3	-1	0
S_1	1	3	1	0
S_2	1	0	2	0

'Unfolding' Three Generations out of E_8

The matter content of a general \widehat{E}_8 -fibred geometry can be derived by thinking in terms of sequential 'unfoldings':

Unfolding $SO_{10} \rightarrow SU_5$

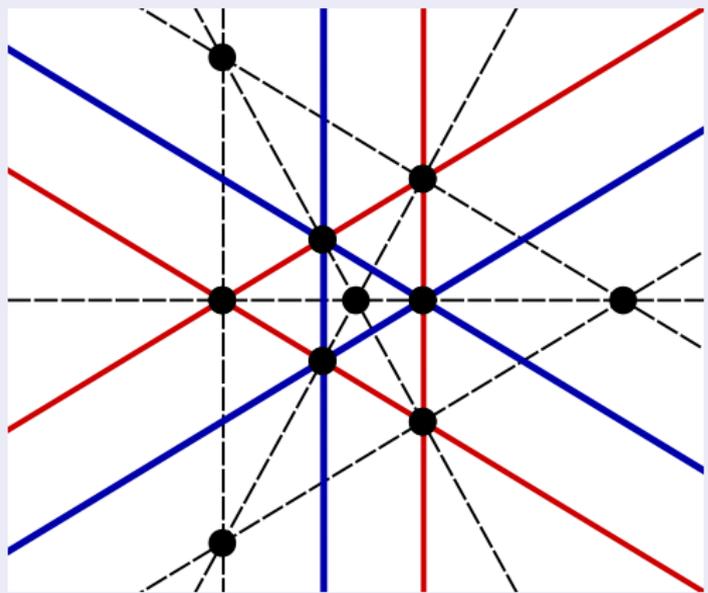


		SO_{10}	U_1^a	U_1^b	U_1^c
T_1	16	1	1	-1	
T_2	16	1	-1	-1	
T_3	16	-2	0	-1	
T_X^c	$\overline{16}$	0	0	-3	
H	10	1	1	2	
Y_a	10	-1	1	-2	
Y_b	10	2	0	-2	
X_1	1	-1	-1	4	
X_2	1	-1	1	4	
N_2^c	1	-2	0	-4	
N_3^c	1	3	-1	0	
S_1	1	3	1	0	
S_2	1	0	2	0	

'Unfolding' Three Generations out of E_8

The matter content of a general \widehat{E}_8 -fibred geometry can be derived by thinking in terms of sequential 'unfoldings':

Unfolding $SO_{10} \rightarrow SU_5$

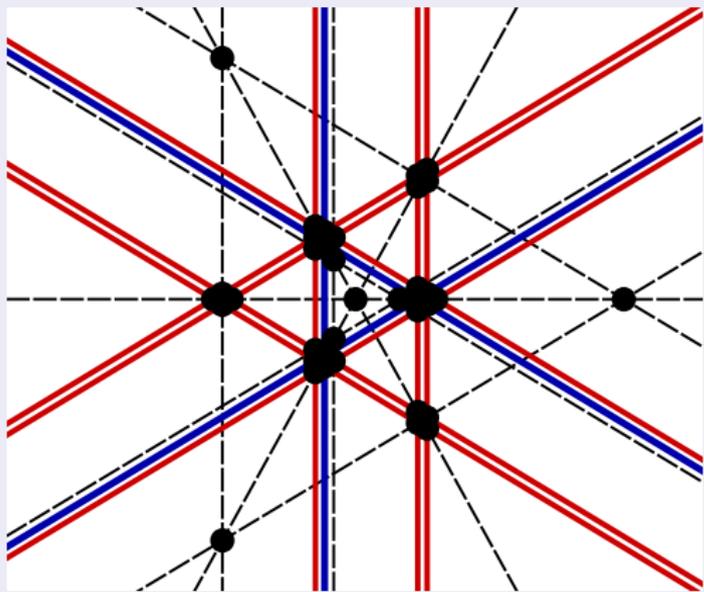


		SO_{10}	U_1^a	U_1^b	U_1^c
T_1	16	1	1	-1	
T_2	16	1	-1	-1	
T_3	16	-2	0	-1	
T_X^c	$\overline{16}$	0	0	-3	
H	10	1	1	2	
Y_a	10	-1	1	-2	
Y_b	10	2	0	-2	
X_1	1	-1	-1	4	
X_2	1	-1	1	4	
N_2^c	1	-2	0	-4	
N_3^c	1	3	-1	0	
S_1	1	3	1	0	
S_2	1	0	2	0	

'Unfolding' Three Generations out of E_8

The matter content of a general \widehat{E}_8 -fibred geometry can be derived by thinking in terms of sequential 'unfoldings':

Unfolding $SO_{10} \rightarrow SU_5$

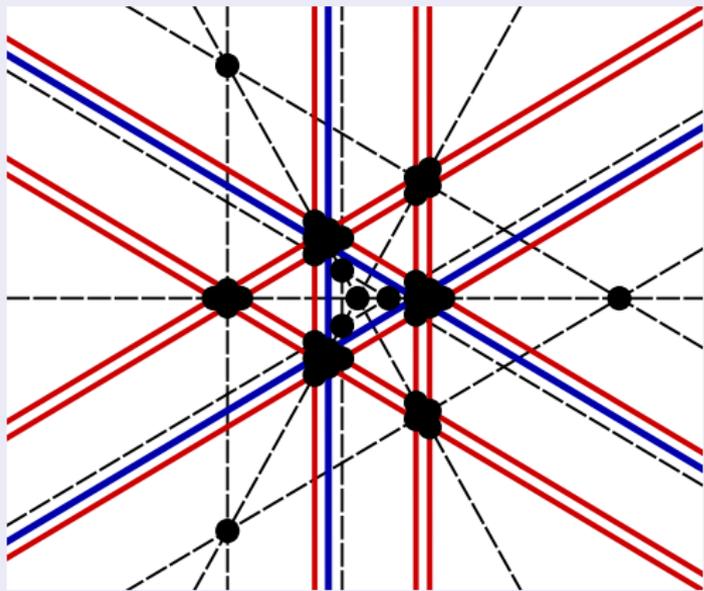


	$SU_5 \times U_1^a \times U_1^b \times U_1^c \times U_1^d$				
T_1	10	1	1	-1	-1
T_2	10	1	-1	-1	-1
T_3	10	-2	0	-1	-1
M_1	$\overline{5}$	1	1	-1	3
M_2	$\overline{5}$	1	-1	-1	3
M_3	$\overline{5}$	-2	0	-1	3
H^u	5	1	1	2	2
H^d	$\overline{5}$	1	1	2	-2
Y_1	5	-1	1	-2	2
Y_1^c	$\overline{5}$	2	0	-2	-2
Y_2	5	2	0	-2	2
Y_2^c	$\overline{5}$	-1	1	-2	-2
X_1	1	-1	-1	4	0
X_2	1	-1	1	4	0
ν_1^b	1	1	1	-1	-5
ν_2^b	1	1	-1	-1	-5
ν_3^b	1	-2	0	-1	-5
N_1^b	1	0	0	-3	5
N_2^b	1	-2	0	-4	0
N_3^b	1	3	-1	0	0
T_X^c	$\overline{10}$	0	0	-3	1
T_X^c	10	0	0	0	4
M_X^c	5	0	0	-3	-3
S_1	1	3	1	0	0
S_2	1	0	2	0	0

'Unfolding' Three Generations out of E_8

The matter content of a general \widehat{E}_8 -fibred geometry can be derived by thinking in terms of sequential 'unfoldings':

Unfolding $SO_{10} \rightarrow SU_5$

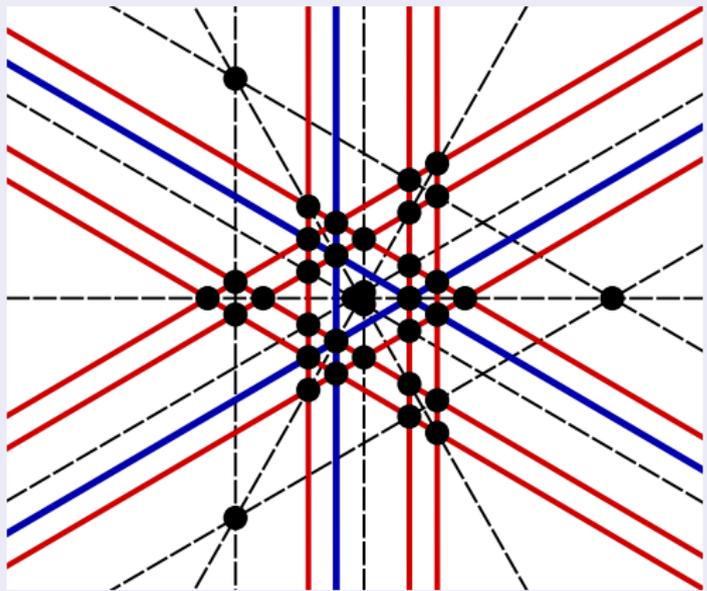


	$SU_5 \times U_1^a \times U_1^b \times U_1^c \times U_1^d$				
T_1	10	1	1	-1	-1
T_2	10	1	-1	-1	-1
T_3	10	-2	0	-1	-1
M_1	$\overline{5}$	1	1	-1	3
M_2	$\overline{5}$	1	-1	-1	3
M_3	$\overline{5}$	-2	0	-1	3
H^u	5	1	1	2	2
H^d	$\overline{5}$	1	1	2	-2
Y_1	5	-1	1	-2	2
Y_1^c	$\overline{5}$	2	0	-2	-2
Y_2	5	2	0	-2	2
Y_2^c	$\overline{5}$	-1	1	-2	-2
X_1	1	-1	-1	4	0
X_2	1	-1	1	4	0
ν_1^b	1	1	1	-1	-5
ν_2^b	1	1	-1	-1	-5
ν_3^b	1	-2	0	-1	-5
N_1^b	1	0	0	-3	5
N_2^b	1	-2	0	-4	0
N_3^b	1	3	-1	0	0
T_X^c	$\overline{10}$	0	0	-3	1
T_X^c	10	0	0	0	4
M_X^c	5	0	0	-3	-3
S_1	1	3	1	0	0
S_2	1	0	2	0	0

'Unfolding' Three Generations out of E_8

The matter content of a general \widehat{E}_8 -fibred geometry can be derived by thinking in terms of sequential 'unfoldings':

Unfolding $SO_{10} \rightarrow SU_5$

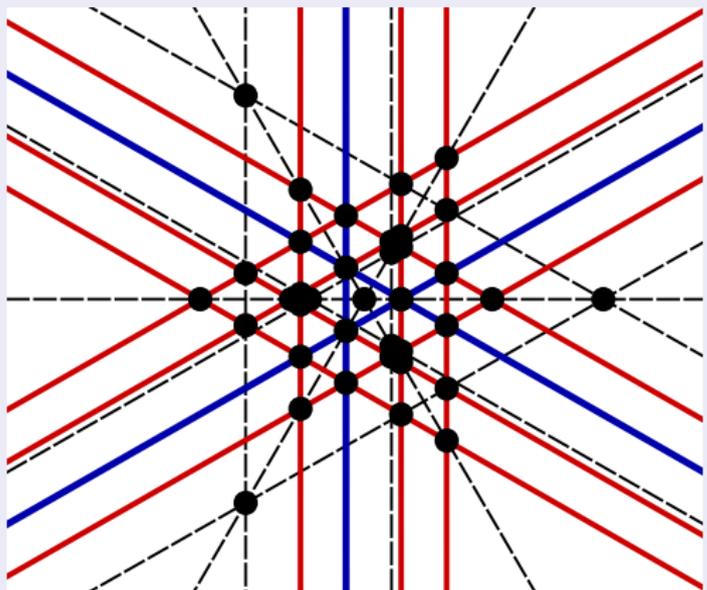


	$SU_5 \times U_1^a \times U_1^b \times U_1^c \times U_1^d$				
T_1	10	1	1	-1	-1
T_2	10	1	-1	-1	-1
T_3	10	-2	0	-1	-1
M_1	$\overline{5}$	1	1	-1	3
M_2	$\overline{5}$	1	-1	-1	3
M_3	$\overline{5}$	-2	0	-1	3
H^u	5	1	1	2	2
H^d	$\overline{5}$	1	1	2	-2
Y_1	5	-1	1	-2	2
Y_1^c	$\overline{5}$	2	0	-2	-2
Y_2	5	2	0	-2	2
Y_2^c	$\overline{5}$	-1	1	-2	-2
X_1	1	-1	-1	4	0
X_2^c	1	-1	1	4	0
ν_1^b	1	1	1	-1	-5
ν_2^c	1	1	-1	-1	-5
ν_3^c	1	-2	0	-1	-5
N_1^b	1	0	0	-3	5
N_2^c	1	-2	0	-4	0
N_3^c	1	3	-1	0	0
T_X^c	$\overline{10}$	0	0	-3	1
T_X^c	10	0	0	0	4
M_X^c	5	0	0	-3	-3
S_1	1	3	1	0	0
S_2	1	0	2	0	0

'Unfolding' Three Generations out of E_8

The matter content of a general \widehat{E}_8 -fibred geometry can be derived by thinking in terms of sequential 'unfoldings':

Unfolding $SO_{10} \rightarrow SU_5$

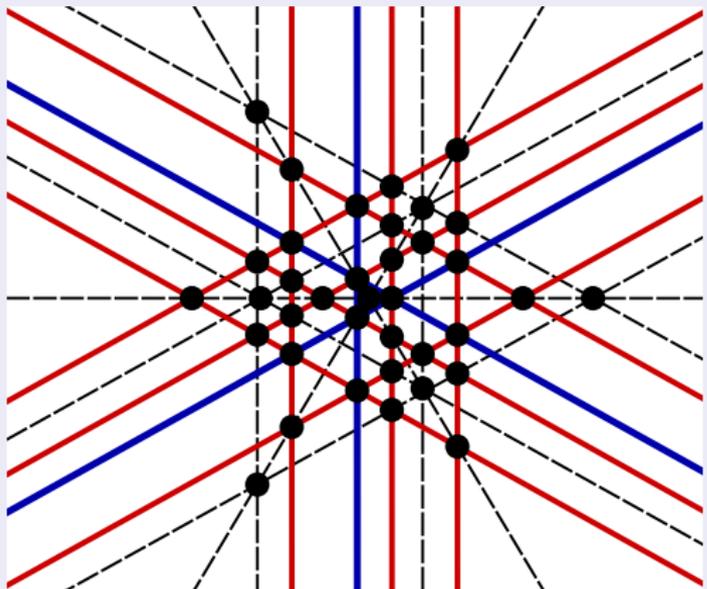


	$SU_5 \times U_1^a \times U_1^b \times U_1^c \times U_1^d$				
T_1	10	1	1	-1	-1
T_2	10	1	-1	-1	-1
T_3	10	-2	0	-1	-1
M_1	$\overline{5}$	1	1	-1	3
M_2	$\overline{5}$	1	-1	-1	3
M_3	$\overline{5}$	-2	0	-1	3
H^u	5	1	1	2	2
H^d	$\overline{5}$	1	1	2	-2
Y_1	5	-1	1	-2	2
Y_1^c	$\overline{5}$	2	0	-2	-2
Y_2	5	2	0	-2	2
Y_2^c	$\overline{5}$	-1	1	-2	-2
X_1	1	-1	-1	4	0
X_2^e	1	-1	1	4	0
ν_1^b	1	1	1	-1	-5
ν_2^e	1	1	-1	-1	-5
ν_3^e	1	-2	0	-1	-5
N_1^b	1	0	0	-3	5
N_2^e	1	-2	0	-4	0
N_3^e	1	3	-1	0	0
T_X^c	$\overline{10}$	0	0	-3	1
T_X^e	10	0	0	0	4
M_X^c	5	0	0	-3	-3
S_1	1	3	1	0	0
S_2	1	0	2	0	0

'Unfolding' Three Generations out of E_8

The matter content of a general \widehat{E}_8 -fibred geometry can be derived by thinking in terms of sequential 'unfoldings':

Unfolding $SO_{10} \rightarrow SU_5$

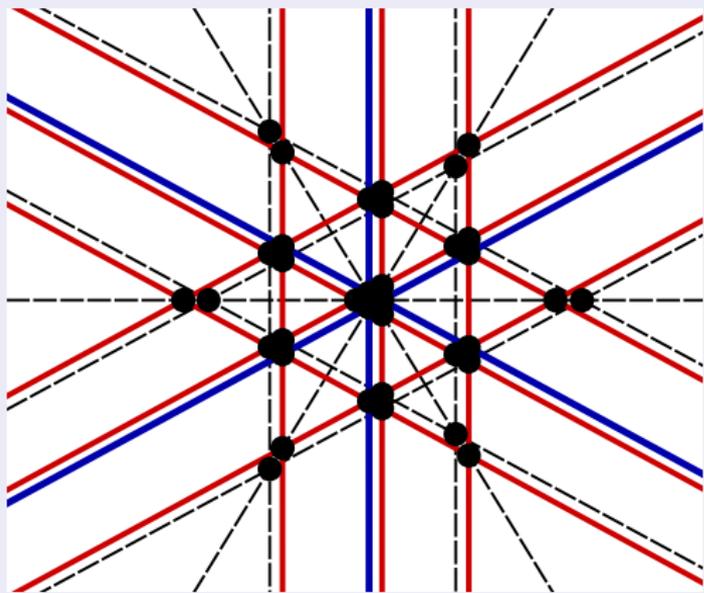


	$SU_5 \times U_1^a \times U_1^b \times U_1^c \times U_1^d$				
T_1	10	1	1	-1	-1
T_2	10	1	-1	-1	-1
T_3	10	-2	0	-1	-1
M_1	$\overline{5}$	1	1	-1	3
M_2	$\overline{5}$	1	-1	-1	3
M_3	$\overline{5}$	-2	0	-1	3
H^u	5	1	1	2	2
H^d	$\overline{5}$	1	1	2	-2
Y_1	5	-1	1	-2	2
Y_1^c	$\overline{5}$	2	0	-2	-2
Y_2	5	2	0	-2	2
Y_2^c	$\overline{5}$	-1	1	-2	-2
X_1	1	-1	-1	4	0
X_2^c	1	-1	1	4	0
ν_1^b	1	1	1	-1	-5
ν_2^c	1	1	-1	-1	-5
ν_3^c	1	-2	0	-1	-5
N_1^b	1	0	0	-3	5
N_2^c	1	-2	0	-4	0
N_3^c	1	3	-1	0	0
T_X^c	$\overline{10}$	0	0	-3	1
T_X^c	10	0	0	0	4
M_X^c	5	0	0	-3	-3
S_1	1	3	1	0	0
S_2	1	0	2	0	0

'Unfolding' Three Generations out of E_8

The matter content of a general \widehat{E}_8 -fibred geometry can be derived by thinking in terms of sequential 'unfoldings':

Unfolding $SO_{10} \rightarrow SU_5$

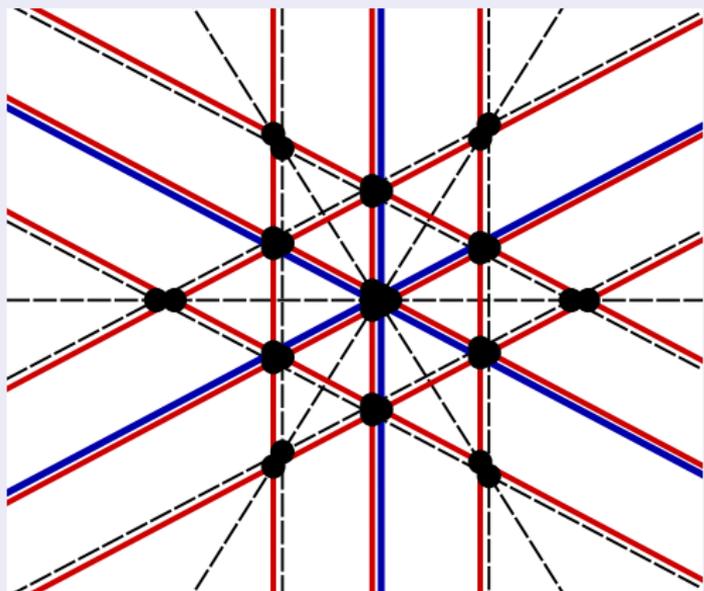


	$SU_5 \times U_1^a \times U_1^b \times U_1^c \times U_1^d$				
T_1	10	1	1	-1	-1
T_2	10	1	-1	-1	-1
T_3	10	-2	0	-1	-1
M_1	$\overline{5}$	1	1	-1	3
M_2	$\overline{5}$	1	-1	-1	3
M_3	$\overline{5}$	-2	0	-1	3
H^u	5	1	1	2	2
H^d	$\overline{5}$	1	1	2	-2
Y_1	5	-1	1	-2	2
Y_1^c	$\overline{5}$	2	0	-2	-2
Y_2	5	2	0	-2	2
Y_2^c	$\overline{5}$	-1	1	-2	-2
X_1	1	-1	-1	4	0
X_2	1	-1	1	4	0
ν_1^b	1	1	1	-1	-5
ν_2^b	1	1	-1	-1	-5
ν_3^b	1	-2	0	-1	-5
N_1^b	1	0	0	-3	5
N_2^b	1	-2	0	-4	0
N_3^b	1	3	-1	0	0
T_X^c	$\overline{10}$	0	0	-3	1
T_X^c	10	0	0	0	4
M_X^c	5	0	0	-3	-3
S_1	1	3	1	0	0
S_2	1	0	2	0	0

'Unfolding' Three Generations out of E_8

The matter content of a general \widehat{E}_8 -fibred geometry can be derived by thinking in terms of sequential 'unfoldings':

Unfolding $SO_{10} \rightarrow SU_5$

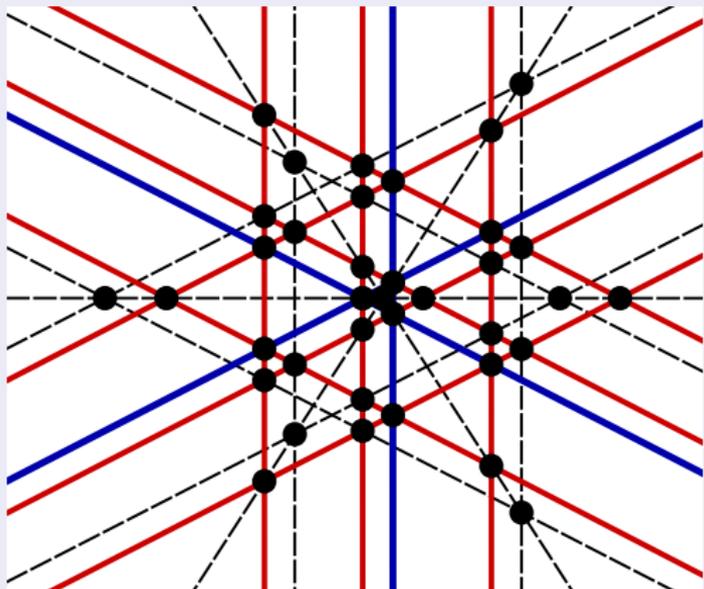


	$SU_5 \times U_1^a \times U_1^b \times U_1^c \times U_1^d$				
T_1	10	1	1	-1	-1
T_2	10	1	-1	-1	-1
T_3	10	-2	0	-1	-1
M_1	$\overline{5}$	1	1	-1	3
M_2	$\overline{5}$	1	-1	-1	3
M_3	$\overline{5}$	-2	0	-1	3
H^u	5	1	1	2	2
H^d	$\overline{5}$	1	1	2	-2
Y_1	5	-1	1	-2	2
Y_1^c	$\overline{5}$	2	0	-2	-2
Y_2	5	2	0	-2	2
Y_2^c	$\overline{5}$	-1	1	-2	-2
X_1	1	-1	-1	4	0
X_2	1	-1	1	4	0
ν_1^l	1	1	1	-1	-5
ν_2^l	1	1	-1	-1	-5
ν_3^l	1	-2	0	-1	-5
N_1^c	1	0	0	-3	5
N_2^c	1	-2	0	-4	0
N_3^c	1	3	-1	0	0
T_X^c	$\overline{10}$	0	0	-3	1
T_X^c	10	0	0	0	4
M_X^c	5	0	0	-3	-3
S_1	1	3	1	0	0
S_2	1	0	2	0	0

'Unfolding' Three Generations out of E_8

The matter content of a general \widehat{E}_8 -fibred geometry can be derived by thinking in terms of sequential 'unfoldings':

Unfolding $SO_{10} \rightarrow SU_5$

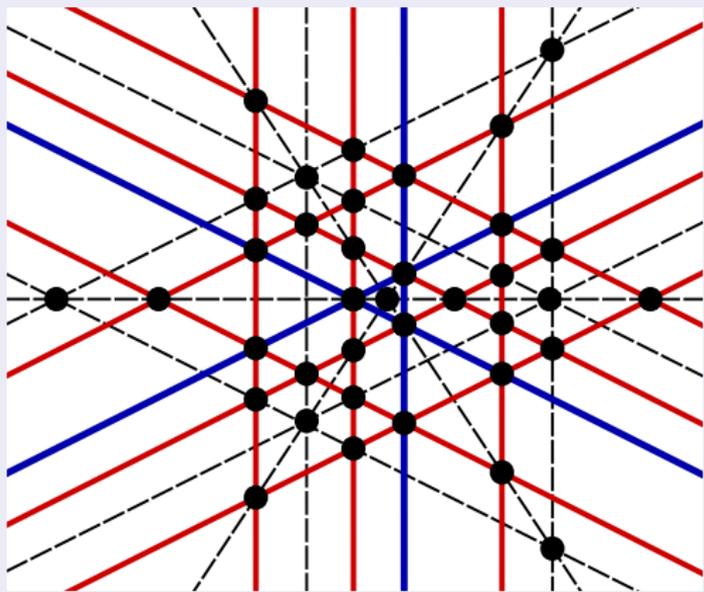


	$SU_5 \times U_1^a \times U_1^b \times U_1^c \times U_1^d$				
T_1	10	1	1	-1	-1
T_2	10	1	-1	-1	-1
T_3	10	-2	0	-1	-1
M_1	$\overline{5}$	1	1	-1	3
M_2	$\overline{5}$	1	-1	-1	3
M_3	$\overline{5}$	-2	0	-1	3
H^u	5	1	1	2	2
H^d	$\overline{5}$	1	1	2	-2
Y_1	5	-1	1	-2	2
Y_1^c	$\overline{5}$	2	0	-2	-2
Y_2	5	2	0	-2	2
Y_2^c	$\overline{5}$	-1	1	-2	-2
X_1	1	-1	-1	4	0
X_2	1	-1	1	4	0
ν_1^b	1	1	1	-1	-5
ν_2^b	1	1	-1	-1	-5
ν_3^b	1	-2	0	-1	-5
N_1^b	1	0	0	-3	5
N_2^b	1	-2	0	-4	0
N_3^b	1	3	-1	0	0
T_X^c	$\overline{10}$	0	0	-3	1
T_X^c	10	0	0	0	4
M_X^c	5	0	0	-3	-3
S_1	1	3	1	0	0
S_2	1	0	2	0	0

'Unfolding' Three Generations out of E_8

The matter content of a general \widehat{E}_8 -fibred geometry can be derived by thinking in terms of sequential 'unfoldings':

Unfolding $SO_{10} \rightarrow SU_5$

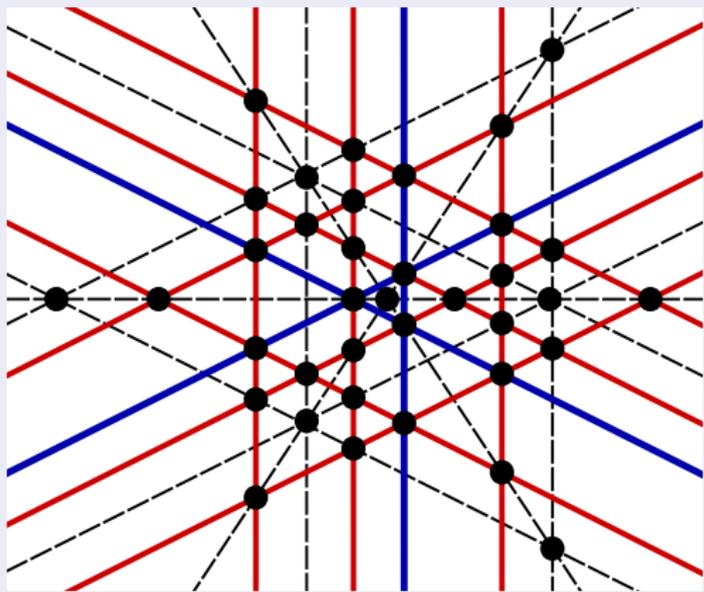


	$SU_5 \times U_1^a \times U_1^b \times U_1^c \times U_1^d$				
T_1	10	1	1	-1	-1
T_2	10	1	-1	-1	-1
T_3	10	-2	0	-1	-1
M_1	$\overline{5}$	1	1	-1	3
M_2	$\overline{5}$	1	-1	-1	3
M_3	$\overline{5}$	-2	0	-1	3
H^u	5	1	1	2	2
H^d	$\overline{5}$	1	1	2	-2
Y_1	5	-1	1	-2	2
Y_1^c	$\overline{5}$	2	0	-2	-2
Y_2	5	2	0	-2	2
Y_2^c	$\overline{5}$	-1	1	-2	-2
X_1	1	-1	-1	4	0
X_2	1	-1	1	4	0
ν_1^b	1	1	1	-1	-5
ν_2^b	1	1	-1	-1	-5
ν_3^b	1	-2	0	-1	-5
N_1^b	1	0	0	-3	5
N_2^b	1	-2	0	-4	0
N_3^b	1	3	-1	0	0
T_X^c	$\overline{10}$	0	0	-3	1
T_X^c	10	0	0	0	4
M_X^c	5	0	0	-3	-3
S_1	1	3	1	0	0
S_2	1	0	2	0	0

'Unfolding' Three Generations out of E_8

The matter content of a general \widehat{E}_8 -fibred geometry can be derived by thinking in terms of sequential 'unfoldings':

Unfolding $SO_{10} \rightarrow SU_5$

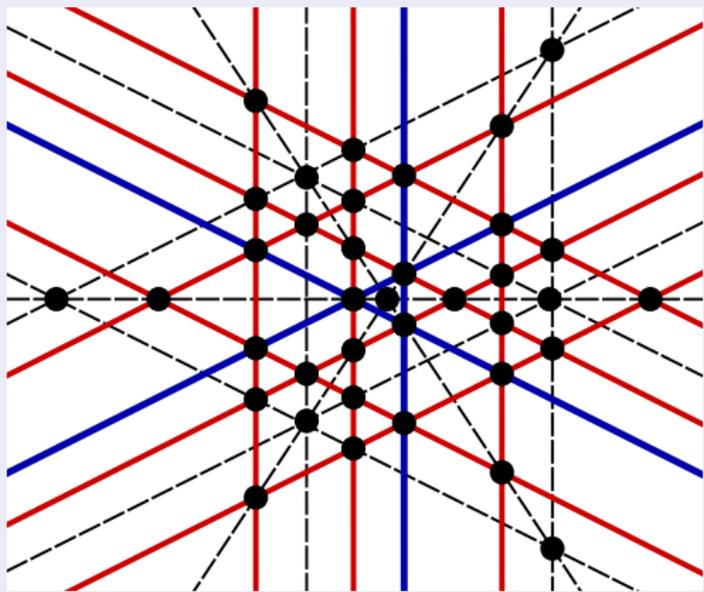


	$SU_5 \times U_1^a \times U_1^b \times U_1^c \times U_1^d$				
T_1	10	1	1	-1	-1
T_2	10	1	-1	-1	-1
T_3	10	-2	0	-1	-1
M_1	$\overline{5}$	1	1	-1	3
M_2	$\overline{5}$	1	-1	-1	3
M_3	$\overline{5}$	-2	0	-1	3
H^u	5	1	1	2	2
H^d	$\overline{5}$	1	1	2	-2
Y_1	5	-1	1	-2	2
Y_1^c	$\overline{5}$	2	0	-2	-2
Y_2	5	2	0	-2	2
Y_2^c	$\overline{5}$	-1	1	-2	-2
X_1	1	-1	-1	4	0
X_2	1	-1	1	4	0
ν_1^b	1	1	1	-1	-5
ν_2^b	1	1	-1	-1	-5
ν_3^b	1	-2	0	-1	-5
N_1^b	1	0	0	-3	5
N_2^b	1	-2	0	-4	0
N_3^b	1	3	-1	0	0
T_X^c	$\overline{10}$	0	0	-3	1
T_X^c	10	0	0	0	4
M_X^c	5	0	0	-3	-3
S_1	1	3	1	0	0
S_2	1	0	2	0	0

'Unfolding' Three Generations out of E_8

The matter content of a general \widehat{E}_8 -fibred geometry can be derived by thinking in terms of sequential 'unfoldings':

Unfolding $SO_{10} \rightarrow SU_5$

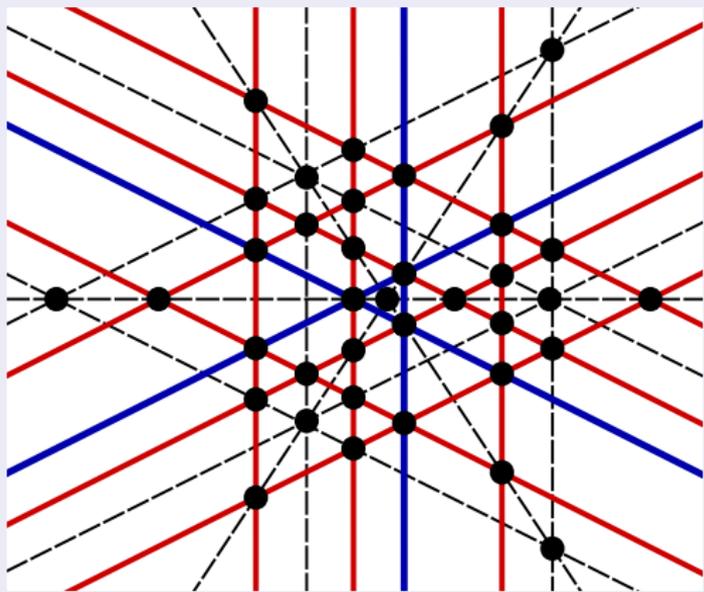


	$SU_5 \times U_1^a \times U_1^b \times U_1^c \times U_1^d$				
T_1	10	1	1	-1	-1
T_2	10	1	-1	-1	-1
T_3	10	-2	0	-1	-1
M_1	$\overline{5}$	1	1	-1	3
M_2	$\overline{5}$	1	-1	-1	3
M_3	$\overline{5}$	-2	0	-1	3
H^u	5	1	1	2	2
H^d	$\overline{5}$	1	1	2	-2
Y_1	5	-1	1	-2	2
Y_1^c	$\overline{5}$	2	0	-2	-2
Y_2	5	2	0	-2	2
Y_2^c	$\overline{5}$	-1	1	-2	-2
X_1	1	-1	-1	4	0
X_2	1	-1	1	4	0
ν_1^b	1	1	1	-1	-5
ν_2^b	1	1	-1	-1	-5
ν_3^b	1	-2	0	-1	-5
N_1^b	1	0	0	-3	5
N_2^b	1	-2	0	-4	0
N_3^b	1	3	-1	0	0
T_X^c	$\overline{10}$	0	0	-3	1
T_X^c	10	0	0	0	4
M_X^c	5	0	0	-3	-3
S_1	1	3	1	0	0
S_2	1	0	2	0	0

'Unfolding' Three Generations out of E_8

The matter content of a general \widehat{E}_8 -fibred geometry can be derived by thinking in terms of sequential 'unfoldings':

Unfolding $SO_{10} \rightarrow SU_5$

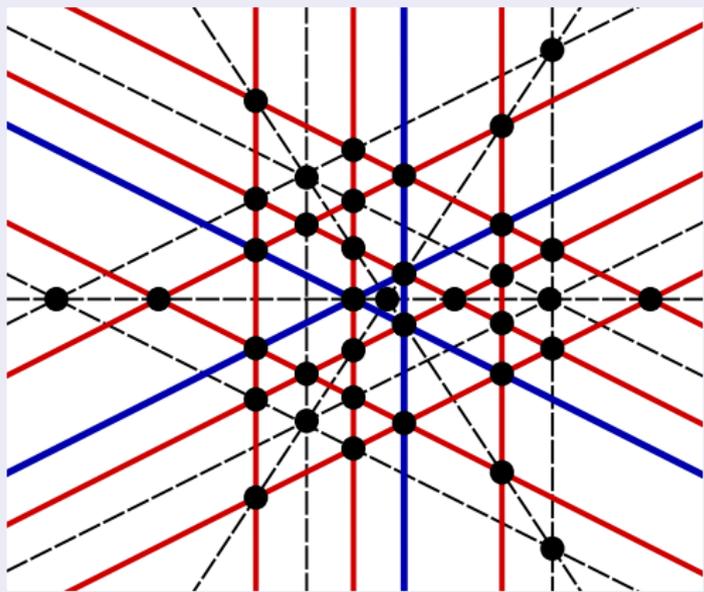


	$SU_5 \times U_1^a \times U_1^b \times U_1^c \times U_1^d$				
T_1	10	1	1	-1	-1
T_2	10	1	-1	-1	-1
T_3	10	-2	0	-1	-1
M_1	$\overline{5}$	1	1	-1	3
M_2	$\overline{5}$	1	-1	-1	3
M_3	$\overline{5}$	-2	0	-1	3
H^u	5	1	1	2	2
H^d	$\overline{5}$	1	1	2	-2
Y_1	5	-1	1	-2	2
Y_1^c	$\overline{5}$	2	0	-2	-2
Y_2	5	2	0	-2	2
Y_2^c	$\overline{5}$	-1	1	-2	-2
X_1	1	-1	-1	4	0
X_2	1	-1	1	4	0
ν_1^b	1	1	1	-1	-5
ν_2^b	1	1	-1	-1	-5
ν_3^b	1	-2	0	-1	-5
N_1^b	1	0	0	-3	5
N_2^b	1	-2	0	-4	0
N_3^b	1	3	-1	0	0
T_X^c	$\overline{10}$	0	0	-3	1
T_X^c	10	0	0	0	4
M_X^c	5	0	0	-3	-3
S_1	1	3	1	0	0
S_2	1	0	2	0	0

'Unfolding' Three Generations out of E_8

The matter content of a general \widehat{E}_8 -fibred geometry can be derived by thinking in terms of sequential 'unfoldings':

Unfolding $SO_{10} \rightarrow SU_5$

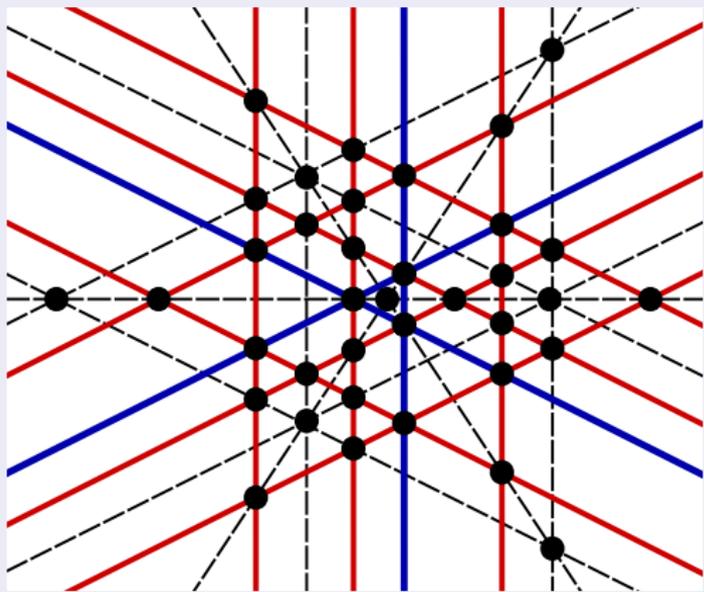


	$SU_5 \times U_1^a \times U_1^b \times U_1^c \times U_1^d$				
T_1	10	1	1	-1	-1
T_2	10	1	-1	-1	-1
T_3	10	-2	0	-1	-1
M_1	$\overline{5}$	1	1	-1	3
M_2	$\overline{5}$	1	-1	-1	3
M_3	$\overline{5}$	-2	0	-1	3
H^u	5	1	1	2	2
H^d	$\overline{5}$	1	1	2	-2
Y_1	5	-1	1	-2	2
Y_1^c	$\overline{5}$	2	0	-2	-2
Y_2	5	2	0	-2	2
Y_2^c	$\overline{5}$	-1	1	-2	-2
X_1	1	-1	-1	4	0
X_2	1	-1	1	4	0
ν_1^b	1	1	1	-1	-5
ν_2^b	1	1	-1	-1	-5
ν_3^b	1	-2	0	-1	-5
N_1^b	1	0	0	-3	5
N_2^b	1	-2	0	-4	0
N_3^b	1	3	-1	0	0
T_X^c	$\overline{10}$	0	0	-3	1
T_X^c	10	0	0	0	4
M_X^c	5	0	0	-3	-3
S_1	1	3	1	0	0
S_2	1	0	2	0	0

Realistic Examples of Phenomenology

This set of matter curves can lead to new, surprisingly realistic models that are nonetheless quite familiar to model builders

Unfolding $SO_{10} \rightarrow SU_5$

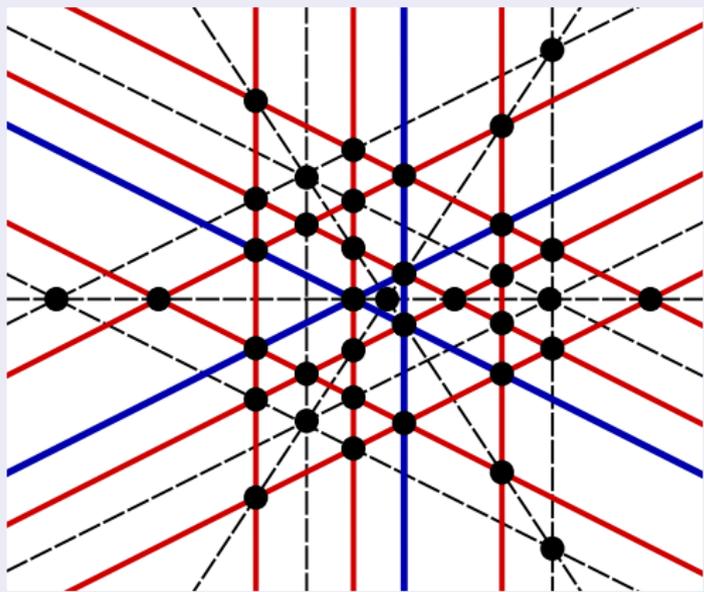


	$SU_5 \times U_1^a \times U_1^b \times U_1^c \times U_1^d$				
T_1	10	1	1	-1	-1
T_2	10	1	-1	-1	-1
T_3	10	-2	0	-1	-1
M_1	$\overline{5}$	1	1	-1	3
M_2	$\overline{5}$	1	-1	-1	3
M_3	$\overline{5}$	-2	0	-1	3
H^u	5	1	1	2	2
H^d	$\overline{5}$	1	1	2	-2
Y_1	5	-1	1	-2	2
Y_1^c	$\overline{5}$	2	0	-2	-2
Y_2	5	2	0	-2	2
Y_2^c	$\overline{5}$	-1	1	-2	-2
X_1	1	-1	-1	4	0
X_2	1	-1	1	4	0
ν_1^b	1	1	1	-1	-5
ν_2^b	1	1	-1	-1	-5
ν_3^b	1	-2	0	-1	-5
N_1^c	1	0	0	-3	5
N_2^c	1	-2	0	-4	0
N_3^c	1	3	-1	0	0
T_X^c	$\overline{10}$	0	0	-3	1
T_X^c	10	0	0	0	4
M_X^c	5	0	0	-3	-3
S_1	1	3	1	0	0
S_2	1	0	2	0	0

Realistic Examples of Phenomenology

This set of matter curves can lead to new, surprisingly realistic models that are nonetheless quite familiar to model builders

Unfolding $SO_{10} \rightarrow SU_5$

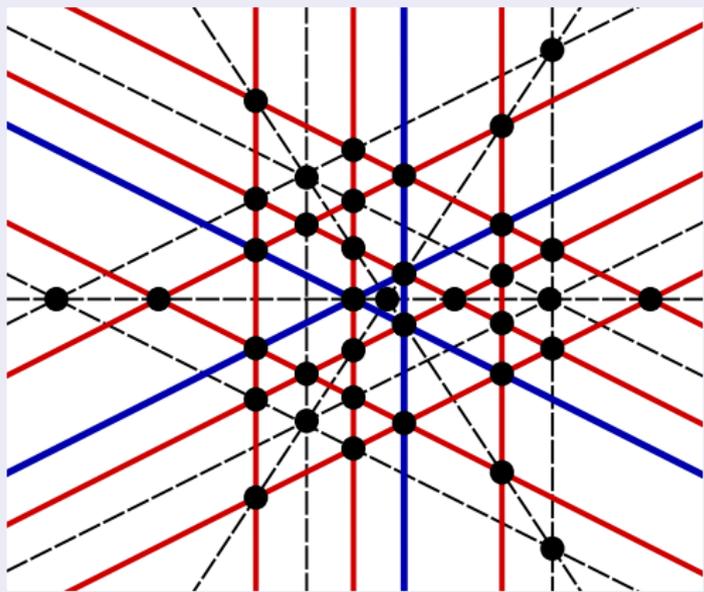


	$SU_5 \times U_1^a \times U_1^b \times U_1^c \times U_1^d$				
T_1	10	1	1	-1	-1
T_2	10	1	-1	-1	-1
T_3	10	-2	0	-1	-1
M_1	$\overline{5}$	1	1	-1	3
M_2	$\overline{5}$	1	-1	-1	3
M_3	$\overline{5}$	-2	0	-1	3
H^u	5	1	1	2	2
H^d	$\overline{5}$	1	1	2	-2
Y_1	5	-1	1	-2	2
Y_1^c	$\overline{5}$	2	0	-2	-2
Y_2	5	2	0	-2	2
Y_2^c	$\overline{5}$	-1	1	-2	-2
X_1	1	-1	-1	4	0
X_2	1	-1	1	4	0
ν_1^b	1	1	1	-1	-5
ν_2^b	1	1	-1	-1	-5
ν_3^b	1	-2	0	-1	-5
N_1^c	1	0	0	-3	5
N_2^c	1	-2	0	-4	0
N_3^c	1	3	-1	0	0
T_X^c	$\overline{10}$	0	0	-3	1
T_X^c	10	0	0	0	4
M_X^c	5	0	0	-3	-3
S_1	1	3	1	0	0
S_2	1	0	2	0	0

Realistic Examples of Phenomenology

This set of matter curves can lead to new, surprisingly realistic models that are nonetheless quite familiar to model builders

Unfolding $SO_{10} \rightarrow SU_5$

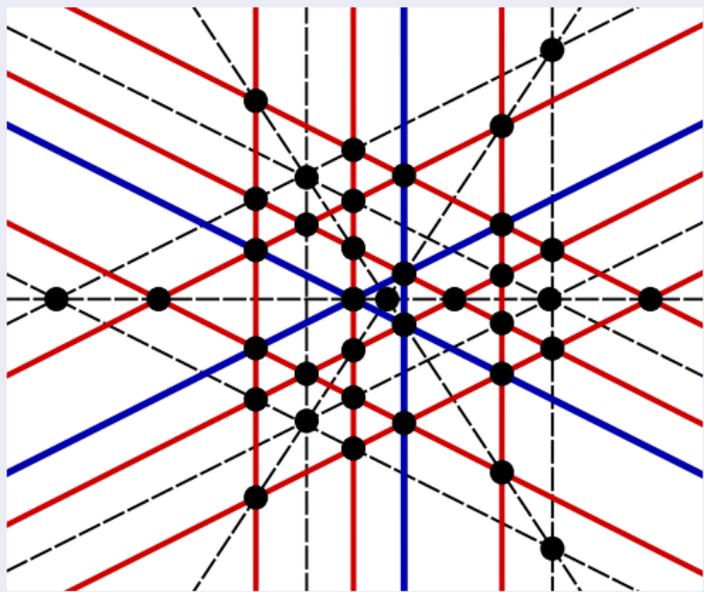


	$SU_5 \times U_1^a \times U_1^b \times U_1^c \times U_1^d$				
T_1	10	1	1	-1	-1
T_2	10	1	-1	-1	-1
T_3	10	-2	0	-1	-1
M_1	$\overline{5}$	1	1	-1	3
M_2	$\overline{5}$	1	-1	-1	3
M_3	$\overline{5}$	-2	0	-1	3
H^u	5	1	1	2	2
H^d	$\overline{5}$	1	1	2	-2
Y_1	5	-1	1	-2	2
Y_1^c	$\overline{5}$	2	0	-2	-2
Y_2	5	2	0	-2	2
Y_2^c	$\overline{5}$	-1	1	-2	-2
X_1	1	-1	-1	4	0
X_2	1	-1	1	4	0
ν_1^b	1	1	1	-1	-5
ν_2^b	1	1	-1	-1	-5
ν_3^b	1	-2	0	-1	-5
N_1^c	1	0	0	-3	5
N_2^c	1	-2	0	-4	0
N_3^c	1	3	-1	0	0
T_X^c	$\overline{10}$	0	0	-3	1
T_X^c	10	0	0	0	4
M_X^c	5	0	0	-3	-3
S_1	1	3	1	0	0
S_2	1	0	2	0	0

Realistic Examples of Phenomenology

This set of matter curves can lead to new, surprisingly realistic models that are nonetheless quite familiar to model builders

Unfolding $SO_{10} \rightarrow SU_5$

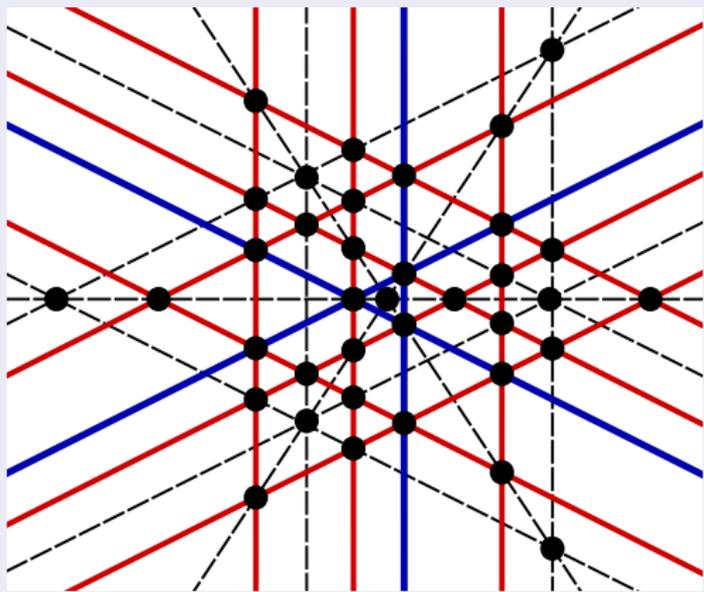


	$SU_5 \times U_1^a \times U_1^b \times U_1^c \times U_1^d$				
T_1	10	1	1	-1	-1
T_2	10	1	-1	-1	-1
T_3	10	-2	0	-1	-1
M_1	$\overline{5}$	1	1	-1	3
M_2	$\overline{5}$	1	-1	-1	3
M_3	$\overline{5}$	-2	0	-1	3
H^u	5	1	1	2	2
H^d	$\overline{5}$	1	1	2	-2
Y_1	5	-1	1	-2	2
Y_1^c	$\overline{5}$	2	0	-2	-2
Y_2	5	2	0	-2	2
Y_2^c	$\overline{5}$	-1	1	-2	-2
X_1	1	-1	-1	4	0
X_2	1	-1	1	4	0
ν_1^b	1	1	1	-1	-5
ν_2^b	1	1	-1	-1	-5
ν_3^b	1	-2	0	-1	-5
N_1^c	1	0	0	-3	5
N_2^c	1	-2	0	-4	0
N_3^c	1	3	-1	0	0
T_X^c	$\overline{10}$	0	0	-3	1
T_X^c	10	0	0	0	4
M_X^c	5	0	0	-3	-3
S_1	1	3	1	0	0
S_2	1	0	2	0	0

Realistic Examples of Phenomenology

This set of matter curves can lead to new, surprisingly realistic models that are nonetheless quite familiar to model builders

Unfolding $SO_{10} \rightarrow SU_5$



	$SU_5 \times U_1^a \times U_1^b \times U_1^c \times U_1^d$				
T_1	10	1	1	-1	-1
T_2	10	1	-1	-1	-1
T_3	10	-2	0	-1	-1
M_1	$\overline{5}$	1	1	-1	3
M_2	$\overline{5}$	1	-1	-1	3
M_3	$\overline{5}$	-2	0	-1	3
H^u	5	1	1	2	2
H^d	$\overline{5}$	1	1	2	-2
Y_1	5	-1	1	-2	2
Y_1^c	$\overline{5}$	2	0	-2	-2
Y_2	5	2	0	-2	2
Y_2^c	$\overline{5}$	-1	1	-2	-2
X_1	1	-1	-1	4	0
X_2	1	-1	1	4	0
ν_1^b	1	1	1	-1	-5
ν_2^b	1	1	-1	-1	-5
ν_3^b	1	-2	0	-1	-5
N_1^c	1	0	0	-3	5
N_2^c	1	-2	0	-4	0
N_3^c	1	3	-1	0	0
T_X^c	$\overline{10}$	0	0	-3	1
T_X^c	10	0	0	0	4
M_X^c	5	0	0	-3	-3
S_1	1	3	1	0	0
S_2	1	0	2	0	0

Realistic Examples of Phenomenology

This set of matter curves can lead to new, surprisingly realistic models that are nonetheless quite familiar to model builders

- The freedom to ‘conjugate’ fields is quite unusual in traditional unified model building, but leads to powerful new phenomenological mechanisms.
- For any non-trivial choice of fluxes, there are *always* anomalous U_1 -symmetries, which become Higgsed by the Green-Schwarz mechanism, which also generates mass-terms for some fields: ‘vacuum realignment.’
- By choosing fluxes appropriately, one can find models that are surprisingly realistic in both M-Theory and F-Theory.

	$SU_5 \times U_1^a \times U_1^b \times U_1^c \times U_1^d$				
T_1	10	1	1	-1	-1
T_2	10	1	-1	-1	-1
T_3	10	-2	0	-1	-1
M_1	$\overline{5}$	1	1	-1	3
M_2	$\overline{5}$	1	-1	-1	3
M_3	$\overline{5}$	-2	0	-1	3
H^u	5	1	1	2	2
H^d	5	1	1	2	-2
Y_1	5	-1	1	-2	2
Y_1^c	$\overline{5}$	2	0	-2	-2
Y_2	5	2	0	-2	2
Y_2^c	$\overline{5}$	-1	1	-2	-2
X_1	1	-1	-1	4	0
X_2	1	-1	1	4	0
ν_1^t	1	1	1	-1	-5
ν_2^t	1	1	-1	-1	-5
ν_3^t	1	-2	0	-1	-5
N_1^t	1	0	0	-3	5
N_2^t	1	-2	0	-4	0
N_3^t	1	3	-1	0	0
T_X^c	$\overline{10}$	0	0	-3	1
T_X	10	0	0	0	4
M_X	5	0	0	-3	-3
S_1	1	3	1	0	0
S_2	1	0	2	0	0

Realistic Examples of Phenomenology

This set of matter curves can lead to new, surprisingly realistic models that are nonetheless quite familiar to model builders

- The freedom to ‘conjugate’ fields is quite unusual in traditional unified model building, but leads to powerful new phenomenological mechanisms.
- For any non-trivial choice of fluxes, there are **always** anomalous U_1 -symmetries, which become Higgsed by the Green-Schwarz mechanism, which also generates mass-terms for some fields: ‘**vacuum realignment.**’
- By choosing fluxes appropriately, one can find models that are surprisingly realistic in both M-Theory and F-Theory.

	$SU_5 \times U_1^a \times U_1^b \times U_1^c \times U_1^d$				
T_1	10	1	1	-1	-1
T_2	10	1	-1	-1	-1
T_3	10	-2	0	-1	-1
M_1	$\bar{5}$	1	1	-1	3
M_2	$\bar{5}$	1	-1	-1	3
M_3	$\bar{5}$	-2	0	-1	3
H^u	5	1	1	2	2
H^d	$\bar{5}$	1	1	2	-2
Y_1	5	-1	1	-2	2
Y_1^c	$\bar{5}$	2	0	-2	-2
Y_2	5	2	0	-2	2
Y_2^c	$\bar{5}$	-1	1	-2	-2
X_1	1	-1	-1	4	0
X_2	1	-1	1	4	0
ν_1^c	1	1	1	-1	-5
ν_2^c	1	1	-1	-1	-5
ν_3^c	1	-2	0	-1	-5
N_1^c	1	0	0	-3	5
N_2^c	1	-2	0	-4	0
N_3^c	1	3	-1	0	0
T_X^c	$\bar{10}$	0	0	-3	1
T_X	10	0	0	0	4
M_X	5	0	0	-3	-3
S_1	1	3	1	0	0
S_2	1	0	2	0	0

Realistic Examples of Phenomenology

This set of matter curves can lead to new, surprisingly realistic models that are nonetheless quite familiar to model builders

- The freedom to ‘conjugate’ fields is quite unusual in traditional unified model building, but leads to powerful new phenomenological mechanisms.
- For any non-trivial choice of fluxes, there are **always** anomalous U_1 -symmetries, which become Higgsed by the Green-Schwarz mechanism, which also generates mass-terms for some fields: ‘**vacuum realignment.**’
- By choosing fluxes appropriately, one can find models that are surprisingly realistic in both M-Theory and F-Theory.

	$SU_5 \times U_1^a \times U_1^b \times U_1^c \times U_1^d$				
T_1	10	1	1	-1	-1
T_2	10	1	-1	-1	-1
T_3	10	-2	0	-1	-1
M_1	$\bar{5}$	1	1	-1	3
M_2	$\bar{5}$	1	-1	-1	3
M_3	$\bar{5}$	-2	0	-1	3
H^u	5	1	1	2	2
H^d	$\bar{5}$	1	1	2	-2
Y_1	5	-1	1	-2	2
Y_1^c	$\bar{5}$	2	0	-2	-2
Y_2	5	2	0	-2	2
Y_2^c	$\bar{5}$	-1	1	-2	-2
X_1	1	-1	-1	4	0
X_2	1	-1	1	4	0
ν_1^t	1	1	1	-1	-5
ν_2^t	1	1	-1	-1	-5
ν_3^t	1	-2	0	-1	-5
N_1^t	1	0	0	-3	5
N_2^t	1	-2	0	-4	0
N_3^t	1	3	-1	0	0
T_X^c	10	0	0	-3	1
T_X^c	10	0	0	0	4
M_X^c	5	0	0	-3	-3
S_1	1	3	1	0	0
S_2	1	0	2	0	0

Re-Folding the Geometry to Enforce Couplings

- The choice of fluxes listed in the table on the right would generate a superpotential of the form:

	$SU_5 \times U_1^a \times U_1^b \times U_1^c \times U_1^d$				
T_1	$\mathbf{10}$	1	1	-1	-1
T_2	$\mathbf{10}$	1	-1	-1	-1
T_3	$\mathbf{10}$	-2	0	-1	-1
M_1	$\overline{\mathbf{5}}$	1	1	-1	3
M_2	$\overline{\mathbf{5}}$	1	-1	-1	3
M_3	$\overline{\mathbf{5}}$	-2	0	-1	3
H^u	$\mathbf{5}$	1	1	2	2
H^d	$\mathbf{5}$	1	1	2	-2
Y_1	$\mathbf{5}$	-1	1	-2	2
Y_1^c	$\overline{\mathbf{5}}$	2	0	-2	-2
Y_2	$\mathbf{5}$	2	0	-2	2
Y_2^c	$\overline{\mathbf{5}}$	-1	1	-2	-2
X_1	$\mathbf{1}$	-1	-1	4	0
X_2	$\mathbf{1}$	-1	1	4	0
ν_1^t	$\mathbf{1}$	1	1	-1	-5
ν_2^t	$\mathbf{1}$	1	-1	-1	-5
ν_3^t	$\mathbf{1}$	-2	0	-1	-5
N_1^t	$\mathbf{1}$	0	0	-3	5
N_2^t	$\mathbf{1}$	-2	0	-4	0
N_3^t	$\mathbf{1}$	3	-1	0	0
(T_X^c)	$\overline{\mathbf{10}}$	0	0	-3	1
(T_X)	$\mathbf{10}$	0	0	0	4
(M_X^c)	$\mathbf{5}$	0	0	-3	-3
(S_1)	$\mathbf{1}$	3	1	0	0
(S_2)	$\mathbf{1}$	0	2	0	0

Re-Folding the Geometry to Enforce Couplings

- The choice of fluxes listed in the table on the right would generate a superpotential of the form:

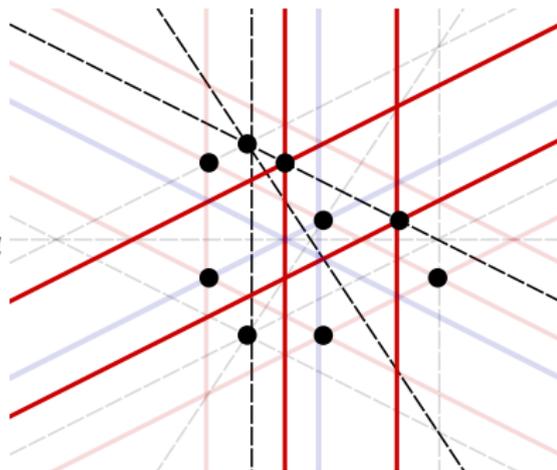
$$\begin{aligned}
 W = & T_2 T_3 H^u + T_2 M_3 H^d + T_3 M_2 H^d \\
 & + H^u M_2 \nu_3^c + H^u M_3 \nu_2^c + X_1 Y_1 Y_1^c \\
 & + X_1 Y_2 Y_2^c + X_1 N_1^c \nu_1^c + X_2 N_1^c \nu_2^c \\
 & + X_2 N_2^c N_3^c.
 \end{aligned}$$

	$SU_5 \times U_1^a \times U_1^b \times U_1^c \times U_1^d$				
T_1	$\mathbf{10}$	1	1	-1	-1
T_2	$\mathbf{10}$	1	-1	-1	-1
T_3	$\mathbf{10}$	-2	0	-1	-1
M_1	$\overline{\mathbf{5}}$	1	1	-1	3
M_2	$\overline{\mathbf{5}}$	1	-1	-1	3
M_3	$\overline{\mathbf{5}}$	-2	0	-1	3
H^u	$\mathbf{5}$	1	1	2	2
H^d	$\overline{\mathbf{5}}$	1	1	2	-2
Y_1	$\mathbf{5}$	-1	1	-2	2
Y_1^c	$\overline{\mathbf{5}}$	2	0	-2	-2
Y_2	$\mathbf{5}$	2	0	-2	2
Y_2^c	$\overline{\mathbf{5}}$	-1	1	-2	-2
X_1	$\mathbf{1}$	-1	-1	4	0
X_2	$\mathbf{1}$	-1	1	4	0
ν_1^c	$\mathbf{1}$	1	1	-1	-5
ν_2^c	$\mathbf{1}$	1	-1	-1	-5
ν_3^c	$\mathbf{1}$	-2	0	-1	-5
N_1^c	$\mathbf{1}$	0	0	-3	5
N_2^c	$\mathbf{1}$	-2	0	-4	0
N_3^c	$\mathbf{1}$	3	-1	0	0
(T_X^c)	$\overline{\mathbf{10}}$	0	0	-3	1
(T_X)	$\mathbf{10}$	0	0	0	4
(M_X^c)	$\mathbf{5}$	0	0	-3	-3
(S_1)	$\mathbf{1}$	3	1	0	0
(S_2)	$\mathbf{1}$	0	2	0	0

Re-Folding the Geometry to Enforce Couplings

- The choice of fluxes listed in the table on the right would generate a superpotential of the form:

$$\begin{aligned}
 W = & T_2 T_3 H^u + T_2 M_3 H^d + T_3 M_2 H^d \\
 & + H^u M_2 \nu_3^c + H^u M_3 \nu_2^c + X_1 Y_1 Y_1^c \\
 & + X_1 Y_2 Y_2^c + X_1 N_1^c \nu_1^c + X_2 N_1^c \nu_2^c \\
 & + X_2 N_2^c N_3^c.
 \end{aligned}$$

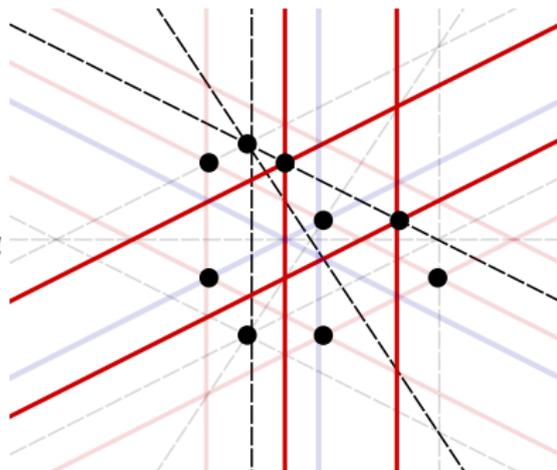


Re-Folding the Geometry to Enforce Couplings

- The choice of fluxes listed in the table on the right would generate a superpotential of the form:

$$\begin{aligned}
 W = & T_2 T_3 H^u + T_2 M_3 H^d + T_3 M_2 H^d \\
 & + H^u M_2 \nu_3^c + H^u M_3 \nu_2^c + X_1 Y_1 Y_1^c \\
 & + X_1 Y_2 Y_2^c + X_1 N_1^c \nu_1^c + X_2 N_1^c \nu_2^c \\
 & + X_2 N_2^c N_3^c.
 \end{aligned}$$

Notice that T_1 and M_1 do not appear in the superpotential at all!



Re-Folding the Geometry to Enforce Couplings

- The choice of fluxes listed in the table on the right would generate a superpotential of the form:

$$\begin{aligned}
 W = & T_2 T_3 H^u + T_2 M_3 H^d + T_3 M_2 H^d \\
 & + H^u M_2 \nu_3^c + H^u M_3 \nu_2^c + X_1 Y_1 Y_1^c \\
 & + X_1 Y_2 Y_2^c + X_1 N_1^c \nu_1^c + X_2 N_1^c \nu_2^c \\
 & + X_2 N_2^c N_3^c.
 \end{aligned}$$

Notice that T_1 and M_1 do not appear in the superpotential at all!

	$SU_5 \times U_1^a$	$\times U_1^b$	$\times U_1^c$	$\times U_1^d$	
T_1	$\mathbf{10}$	1	1	-1	-1
T_2	$\mathbf{10}$	1	-1	-1	-1
T_3	$\mathbf{10}$	-2	0	-1	-1
M_1	$\overline{\mathbf{5}}$	1	1	-1	3
M_2	$\overline{\mathbf{5}}$	1	-1	-1	3
M_3	$\overline{\mathbf{5}}$	-2	0	-1	3
H^u	$\mathbf{5}$	1	1	2	2
H^d	$\overline{\mathbf{5}}$	1	1	2	-2
Y_1	$\mathbf{5}$	-1	1	-2	2
Y_1^c	$\overline{\mathbf{5}}$	2	0	-2	-2
Y_2	$\mathbf{5}$	2	0	-2	2
Y_2^c	$\overline{\mathbf{5}}$	-1	1	-2	-2
X_1	$\mathbf{1}$	-1	-1	4	0
X_2	$\mathbf{1}$	-1	1	4	0
ν_1^c	$\mathbf{1}$	1	1	-1	-5
ν_2^c	$\mathbf{1}$	1	-1	-1	-5
ν_3^c	$\mathbf{1}$	-2	0	-1	-5
N_1^c	$\mathbf{1}$	0	0	-3	5
N_2^c	$\mathbf{1}$	-2	0	-4	0
N_3^c	$\mathbf{1}$	3	-1	0	0
(T_X)	$\overline{\mathbf{10}}$	0	0	-3	1
(T_X^c)	$\mathbf{10}$	0	0	0	4
(M_X^c)	$\mathbf{5}$	0	0	-3	-3
(S_1)	$\mathbf{1}$	3	1	0	0
(S_2)	$\mathbf{1}$	0	2	0	0

Re-Folding the Geometry to Enforce Couplings

- The choice of fluxes listed in the table on the right would generate a superpotential of the form:

$$\begin{aligned}
 W = & T_2 T_3 H^u + T_2 M_3 H^d + T_3 M_2 H^d \\
 & + H^u M_2 \nu_3^c + H^u M_3 \nu_2^c + X_1 Y_1 Y_1^c \\
 & + X_1 Y_2 Y_2^c + X_1 N_1^c \nu_1^c + X_2 N_1^c \nu_2^c \\
 & + X_2 N_2^c N_3^c.
 \end{aligned}$$

Consider the operator $T_1 T_1 H^u$, it has charges $(3, 3, 0, 0)$ under $U_1^a \times U_1^b \times U_1^c \times U_1^d$.

	$SU_5 \times U_1^a$	$\times U_1^b$	$\times U_1^c$	$\times U_1^d$	
T_1	$\mathbf{10}$	1	1	-1	-1
T_2	$\mathbf{10}$	1	-1	-1	-1
T_3	$\mathbf{10}$	-2	0	-1	-1
M_1	$\bar{\mathbf{5}}$	1	1	-1	3
M_2	$\bar{\mathbf{5}}$	1	-1	-1	3
M_3	$\bar{\mathbf{5}}$	-2	0	-1	3
H^u	$\bar{\mathbf{5}}$	1	1	2	2
H^d	$\bar{\mathbf{5}}$	1	1	2	-2
Y_1	$\bar{\mathbf{5}}$	-1	1	-2	2
Y_1^c	$\bar{\mathbf{5}}$	2	0	-2	-2
Y_2	$\bar{\mathbf{5}}$	2	0	-2	2
Y_2^c	$\bar{\mathbf{5}}$	-1	1	-2	-2
X_1	$\mathbf{1}$	-1	-1	4	0
X_2	$\mathbf{1}$	-1	1	4	0
ν_1^c	$\mathbf{1}$	1	1	-1	-5
ν_2^c	$\mathbf{1}$	1	-1	-1	-5
ν_3^c	$\mathbf{1}$	-2	0	-1	-5
N_1^c	$\mathbf{1}$	0	0	-3	5
N_2^c	$\mathbf{1}$	-2	0	-4	0
N_3^c	$\mathbf{1}$	3	-1	0	0
(T_X)	$\mathbf{10}$	0	0	-3	1
(T_X)	$\mathbf{10}$	0	0	0	4
(M_X)	$\bar{\mathbf{5}}$	0	0	-3	-3
(S_1)	$\mathbf{1}$	3	1	0	0
(S_2)	$\mathbf{1}$	0	2	0	0

Re-Folding the Geometry to Enforce Couplings

- The choice of fluxes listed in the table on the right would generate a superpotential of the form:

$$\begin{aligned}
 W = & T_2 T_3 H^u + T_2 M_3 H^d + T_3 M_2 H^d \\
 & + H^u M_2 \nu_3^c + H^u M_3 \nu_2^c + X_1 Y_1 Y_1^c \\
 & + X_1 Y_2 Y_2^c + X_1 N_1^c \nu_1^c + X_2 N_1^c \nu_2^c \\
 & + X_2 N_2^c N_3^c.
 \end{aligned}$$

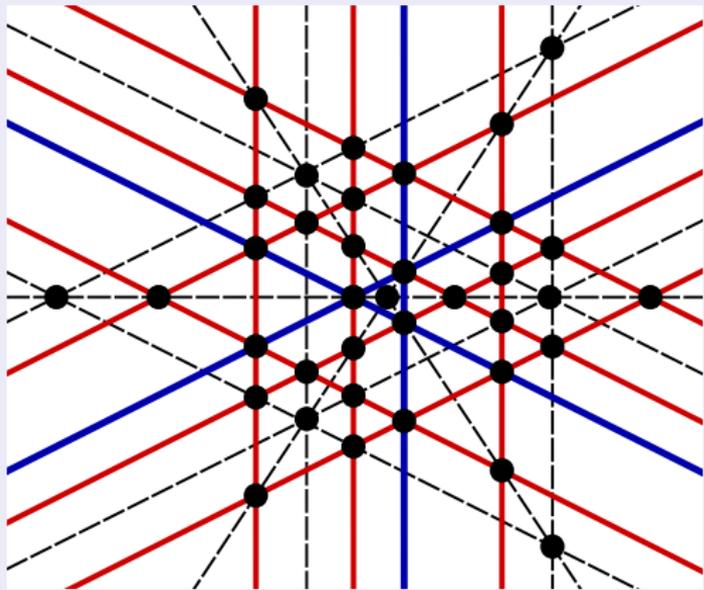
Consider the operator $T_1 T_1 H^u$, it has charges $(3, 3, 0, 0)$ under $U_1^a \times U_1^b \times U_1^c \times U_1^d$.

	$SU_5 \times U_1^a$	$\times U_1^b$	$\times U_1^c$	$\times U_1^d$	
T_1	$\mathbf{10}$	1	1	-1	-1
T_2	$\mathbf{10}$	1	-1	-1	-1
T_3	$\mathbf{10}$	-2	0	-1	-1
M_1	$\bar{\mathbf{5}}$	1	1	-1	3
M_2	$\bar{\mathbf{5}}$	1	-1	-1	3
M_3	$\bar{\mathbf{5}}$	-2	0	-1	3
H^u	$\bar{\mathbf{5}}$	1	1	2	2
H^d	$\bar{\mathbf{5}}$	1	1	2	-2
Y_1	$\bar{\mathbf{5}}$	-1	1	-2	2
Y_1^c	$\bar{\mathbf{5}}$	2	0	-2	-2
Y_2	$\bar{\mathbf{5}}$	2	0	-2	2
Y_2^c	$\bar{\mathbf{5}}$	-1	1	-2	-2
X_1	$\mathbf{1}$	-1	-1	4	0
X_2	$\mathbf{1}$	-1	1	4	0
ν_1^c	$\mathbf{1}$	1	1	-1	-5
ν_2^c	$\mathbf{1}$	1	-1	-1	-5
ν_3^c	$\mathbf{1}$	-2	0	-1	-5
N_1^c	$\mathbf{1}$	0	0	-3	5
N_2^c	$\mathbf{1}$	-2	0	-4	0
N_3^c	$\mathbf{1}$	3	-1	0	0
(T_X)	$\mathbf{10}$	0	0	-3	1
(T_X)	$\mathbf{10}$	0	0	0	4
(M_X)	$\bar{\mathbf{5}}$	0	0	-3	-3
(S_1)	$\mathbf{1}$	3	1	0	0
(S_2)	$\mathbf{1}$	0	2	0	0

A Local Diamond Ring in F-Theory

To enforce the gauge-invariance of $T_1 T_1 H^u$, we should impose the **geometric** condition $b = -a$: $E_8(c+d, 2a+c+d, -3c+d, -4d, 0, 0, 0, 0)$

Constraining the Moduli Space

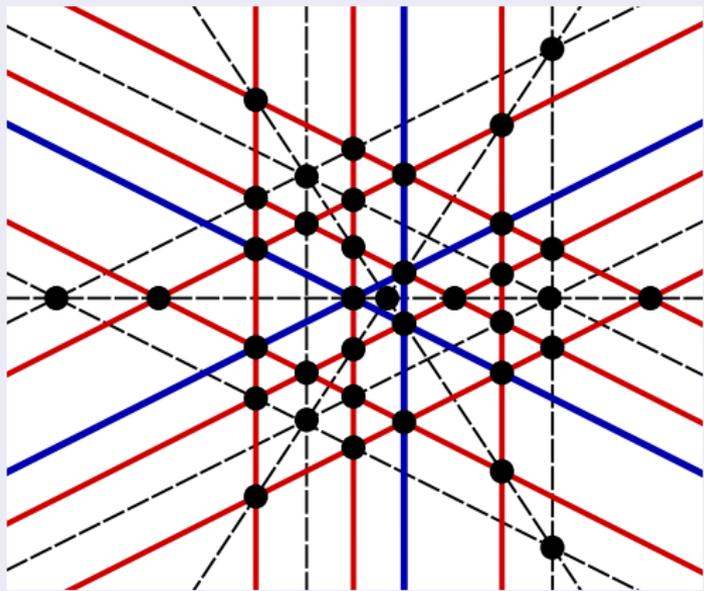


	$SU_5 \times U_1^a \times U_1^b \times U_1^c \times U_1^d$				
T_1	10	1	1	-1	-1
T_2	10	1	-1	-1	-1
T_3	10	-2	0	-1	-1
M_1	$\overline{5}$	1	1	-1	3
M_2	$\overline{5}$	1	-1	-1	3
M_3	$\overline{5}$	-2	0	-1	3
H^u	5	1	1	2	2
H^d	$\overline{5}$	1	1	2	-2
Y_1	5	-1	1	-2	2
Y_1^c	$\overline{5}$	2	0	-2	-2
Y_2	5	2	0	-2	2
Y_2^c	$\overline{5}$	-1	1	-2	-2
X_1	1	-1	-1	4	0
X_2	1	-1	1	4	0
ν_1^b	1	1	1	-1	-5
ν_2^b	1	1	-1	-1	-5
ν_3^b	1	-2	0	-1	-5
N_1^b	1	0	0	-3	5
N_2^b	1	-2	0	-4	0
N_3^b	1	3	-1	0	0
T_X^c	$\overline{10}$	0	0	-3	1
T_X^c	10	0	0	0	4
M_X^c	5	0	0	-3	-3
S_1	1	3	1	0	0
S_2	1	0	2	0	0

A Local Diamond Ring in F-Theory

To enforce the gauge-invariance of $T_1 T_1 H^u$, we should impose the **geometric** condition $b = -a$: $E_8(c+d, 2a+c+d, -3c+d, -4d, 0, 0, 0, 0)$

Constraining the Moduli Space

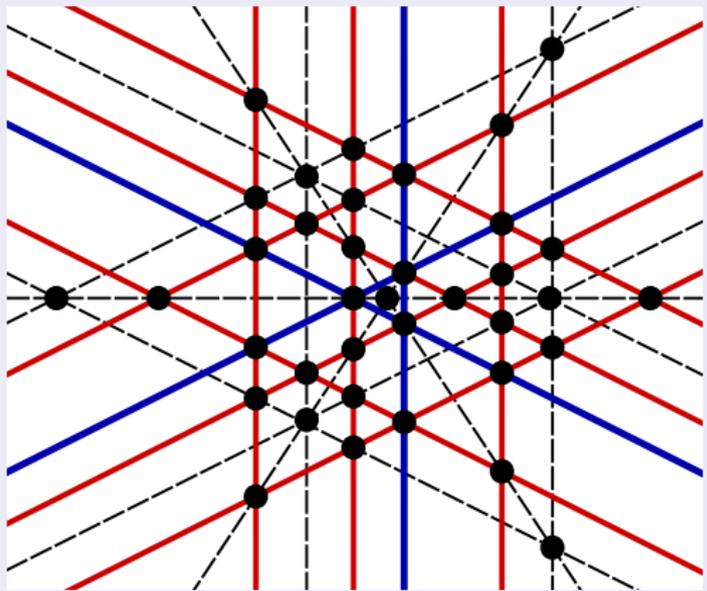


	$SU_5 \times U_1^a \times U_1^b \times U_1^c \times U_1^d$				
T_1	10	1	1	-1	-1
T_2	10	1	-1	-1	-1
T_3	10	-2	0	-1	-1
M_1	$\overline{5}$	1	1	-1	3
M_2	$\overline{5}$	1	-1	-1	3
M_3	$\overline{5}$	-2	0	-1	3
H^u	5	1	1	2	2
H^d	$\overline{5}$	1	1	2	-2
Y_1	5	-1	1	-2	2
Y_1^c	$\overline{5}$	2	0	-2	-2
Y_2	5	2	0	-2	2
Y_2^c	$\overline{5}$	-1	1	-2	-2
X_1	1	-1	-1	4	0
X_2	1	-1	1	4	0
ν_1^b	1	1	1	-1	-5
ν_2^b	1	1	-1	-1	-5
ν_3^b	1	-2	0	-1	-5
N_1^b	1	0	0	-3	5
N_2^b	1	-2	0	-4	0
N_3^b	1	3	-1	0	0
T_X^c	$\overline{10}$	0	0	-3	1
T_X^c	10	0	0	0	4
M_X^c	5	0	0	-3	-3
S_1	1	3	1	0	0
S_2	1	0	2	0	0

A Local Diamond Ring in F-Theory

To enforce the gauge-invariance of $T_1 T_1 H^u$, we should impose the **geometric** condition $b = -a$: $E_8(c+d, 2a+c+d, -3c+d, -4d, 0, 0, 0, 0)$

Constraining the Moduli Space

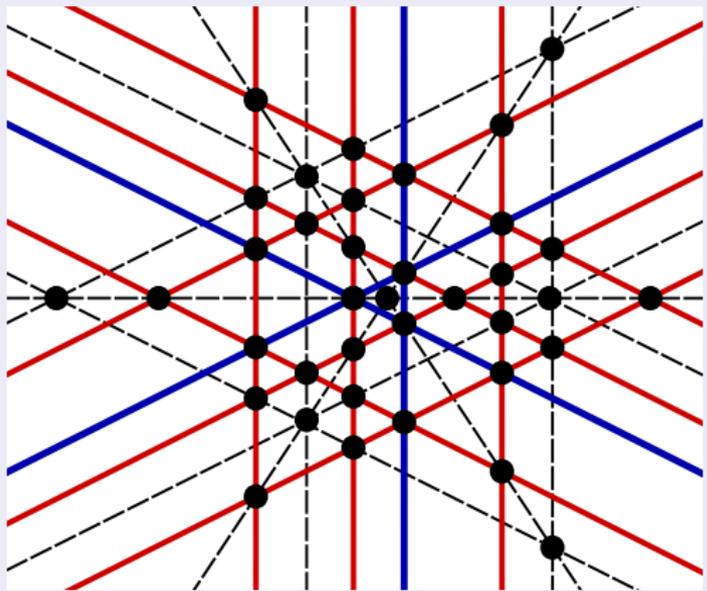


	$SU_5 \times U_1^a \times U_1^b \times U_1^c \times U_1^d$				
T_1	10	1	1	-1	-1
T_2	10	1	-1	-1	-1
T_3	10	-2	0	-1	-1
M_1	$\overline{5}$	1	1	-1	3
M_2	$\overline{5}$	1	-1	-1	3
M_3	$\overline{5}$	-2	0	-1	3
H^u	5	1	1	2	2
H^d	$\overline{5}$	1	1	2	-2
Y_1	5	-1	1	-2	2
Y_1^c	$\overline{5}$	2	0	-2	-2
Y_2	5	2	0	-2	2
Y_2^c	$\overline{5}$	-1	1	-2	-2
X_1	1	-1	-1	4	0
X_2	1	-1	1	4	0
ν_1^b	1	1	1	-1	-5
ν_2^b	1	1	-1	-1	-5
ν_3^b	1	-2	0	-1	-5
N_1^b	1	0	0	-3	5
N_2^b	1	-2	0	-4	0
N_3^b	1	3	-1	0	0
T_X^c	$\overline{10}$	0	0	-3	1
T_X^c	10	0	0	0	4
M_X^c	5	0	0	-3	-3
S_1	1	3	1	0	0
S_2	1	0	2	0	0

A Local Diamond Ring in F-Theory

To enforce the gauge-invariance of $T_1 T_1 H^u$, we should impose the **geometric** condition $b = -a$: $E_8(c+d, 2a+c+d, -3c+d, -4d, 0, 0, 0, 0)$

Constraining the Moduli Space

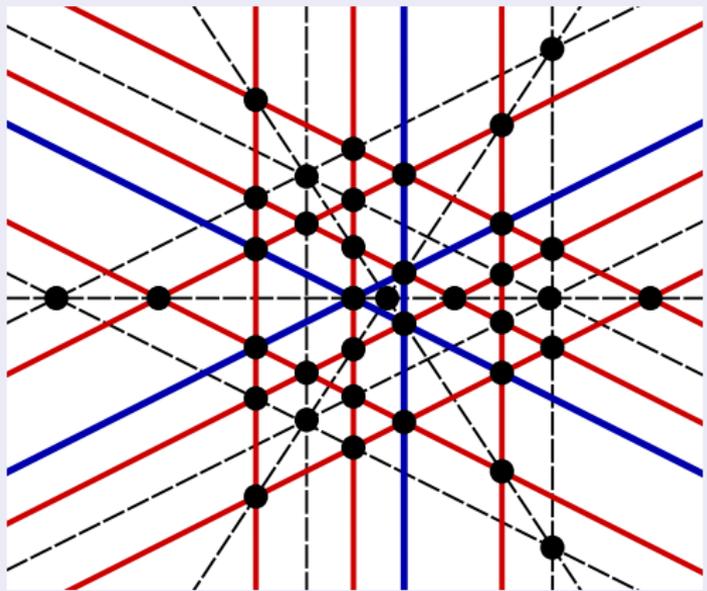


	$SU_5 \times U_1^a \times U_1^b \times U_1^c \times U_1^d$				
T_1	10	1	1	-1	-1
T_2	10	1	-1	-1	-1
T_3	10	-2	0	-1	-1
M_1	$\overline{5}$	1	1	-1	3
M_2	$\overline{5}$	1	-1	-1	3
M_3	$\overline{5}$	-2	0	-1	3
H^u	5	1	1	2	2
H^d	$\overline{5}$	1	1	2	-2
Y_1	5	-1	1	-2	2
Y_1^c	$\overline{5}$	2	0	-2	-2
Y_2	5	2	0	-2	2
Y_2^c	$\overline{5}$	-1	1	-2	-2
X_1	1	-1	-1	4	0
X_2	1	-1	1	4	0
ν_1^t	1	1	1	-1	-5
ν_2^t	1	1	-1	-1	-5
ν_3^t	1	-2	0	-1	-5
N_1^t	1	0	0	-3	5
N_2^t	1	-2	0	-4	0
N_3^t	1	3	-1	0	0
T_X^c	$\overline{10}$	0	0	-3	1
T_X^u	10	0	0	0	4
M_X^c	5	0	0	-3	-3
S_1	1	3	1	0	0
S_2	1	0	2	0	0

A Local Diamond Ring in F-Theory

To enforce the gauge-invariance of $T_1 T_1 H^u$, we should impose the **geometric** condition $b = -a$: $E_8(c+d, 2a+c+d, -3c+d, -4d, 0, 0, 0, 0)$

Constraining the Moduli Space

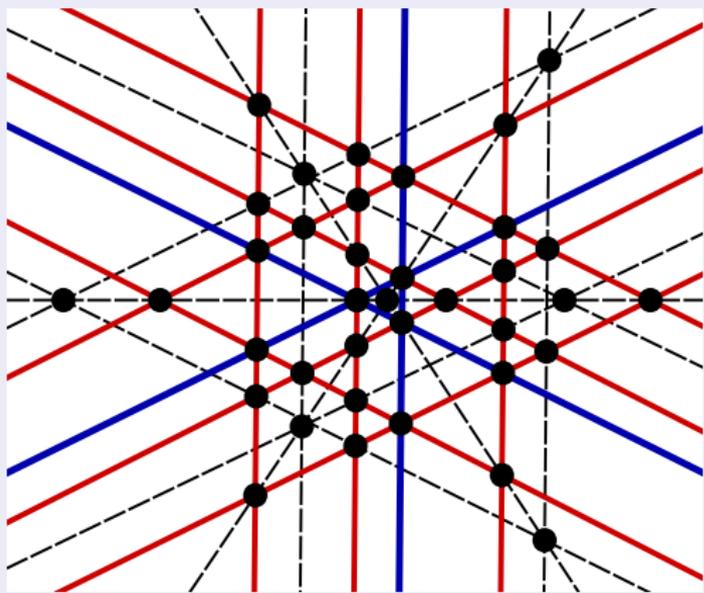


	$SU_5 \times U_1^a \times U_1^b \times U_1^c \times U_1^d$				
T_1	10	1	1	-1	-1
T_2	10	1	-1	-1	-1
T_3	10	-2	0	-1	-1
M_1	$\overline{5}$	1	1	-1	3
M_2	$\overline{5}$	1	-1	-1	3
M_3	$\overline{5}$	-2	0	-1	3
H^u	5	1	1	2	2
H^d	$\overline{5}$	1	1	2	-2
Y_1	5	-1	1	-2	2
Y_1^c	$\overline{5}$	2	0	-2	-2
Y_2	5	2	0	-2	2
Y_2^c	$\overline{5}$	-1	1	-2	-2
X_1	1	-1	-1	4	0
X_2	1	-1	1	4	0
ν_1^b	1	1	1	-1	-5
ν_2^b	1	1	-1	-1	-5
ν_3^b	1	-2	0	-1	-5
N_1^b	1	0	0	-3	5
N_2^b	1	-2	0	-4	0
N_3^b	1	3	-1	0	0
T_X^c	$\overline{10}$	0	0	-3	1
T_X^c	10	0	0	0	4
M_X^c	5	0	0	-3	-3
S_1	1	3	1	0	0
S_2	1	0	2	0	0

A Local Diamond Ring in F-Theory

To enforce the gauge-invariance of $T_1 T_1 H^u$, we should impose the **geometric** condition $b = -a$: $E_8(c+d, 2a+c+d, -3c+d, -4d, 0, 0, 0, 0)$

Constraining the Moduli Space

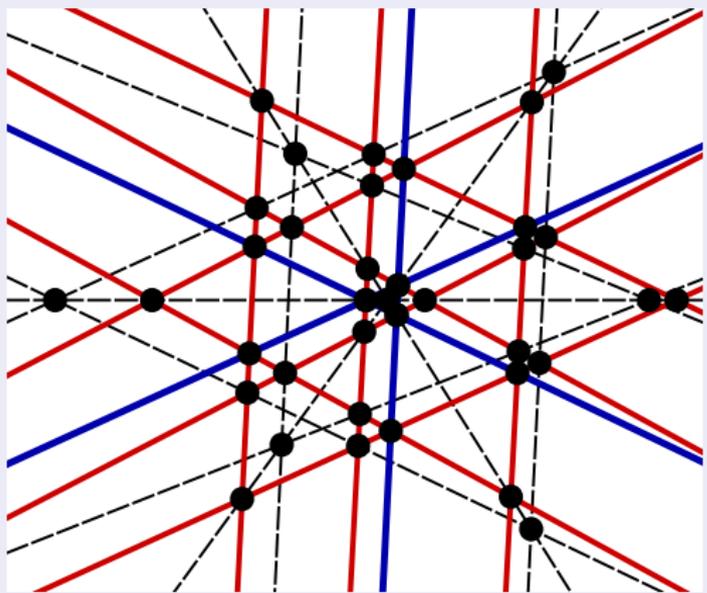


	$SU_5 \times U_1^a \times U_1^b \times U_1^c \times U_1^d$				
T_1	10	1	1	-1	-1
T_2	10	1	-1	-1	-1
T_3	10	-2	0	-1	-1
M_1	$\overline{5}$	1	1	-1	3
M_2	$\overline{5}$	1	-1	-1	3
M_3	$\overline{5}$	-2	0	-1	3
H^u	5	1	1	2	2
H^d	$\overline{5}$	1	1	2	-2
Y_1	5	-1	1	-2	2
Y_1^c	$\overline{5}$	2	0	-2	-2
Y_2	5	2	0	-2	2
Y_2^c	$\overline{5}$	-1	1	-2	-2
X_1	1	-1	-1	4	0
X_2	1	-1	1	4	0
ν_1^b	1	1	1	-1	-5
ν_2^b	1	1	-1	-1	-5
ν_3^b	1	-2	0	-1	-5
N_1^b	1	0	0	-3	5
N_2^b	1	-2	0	-4	0
N_3^b	1	3	-1	0	0
T_X^c	$\overline{10}$	0	0	-3	1
T_X^c	10	0	0	0	4
M_X^c	5	0	0	-3	-3
S_1	1	3	1	0	0
S_2	1	0	2	0	0

A Local Diamond Ring in F-Theory

To enforce the gauge-invariance of $T_1 T_1 H^u$, we should impose the **geometric** condition $b = -a$: $E_8(c+d, 2a+c+d, -3c+d, -4d, 0, 0, 0, 0)$

Constraining the Moduli Space

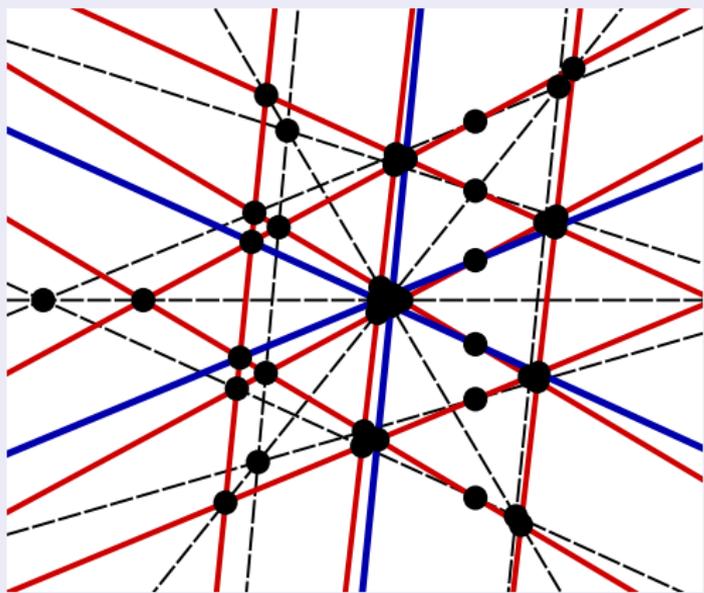


	$SU_5 \times U_1^a \times U_1^b \times U_1^c \times U_1^d$				
T_1	10	1	1	-1	-1
T_2	10	1	-1	-1	-1
T_3	10	-2	0	-1	-1
M_1	$\overline{5}$	1	1	-1	3
M_2	$\overline{5}$	1	-1	-1	3
M_3	$\overline{5}$	-2	0	-1	3
H^u	5	1	1	2	2
H^d	$\overline{5}$	1	1	2	-2
Y_1	5	-1	1	-2	2
Y_1^c	$\overline{5}$	2	0	-2	-2
Y_2	5	2	0	-2	2
Y_2^c	$\overline{5}$	-1	1	-2	-2
X_1	1	-1	-1	4	0
X_2	1	-1	1	4	0
ν_1^b	1	1	1	-1	-5
ν_2^b	1	1	-1	-1	-5
ν_3^b	1	-2	0	-1	-5
N_1^b	1	0	0	-3	5
N_2^b	1	-2	0	-4	0
N_3^b	1	3	-1	0	0
T_X^c	$\overline{10}$	0	0	-3	1
T_X^c	10	0	0	0	4
M_X^c	5	0	0	-3	-3
S_1	1	3	1	0	0
S_2	1	0	2	0	0

A Local Diamond Ring in F-Theory

To enforce the gauge-invariance of $T_1 T_1 H^u$, we should impose the **geometric** condition $b = -a$: $E_8(c+d, 2a+c+d, -3c+d, -4d, 0, 0, 0, 0)$

Constraining the Moduli Space

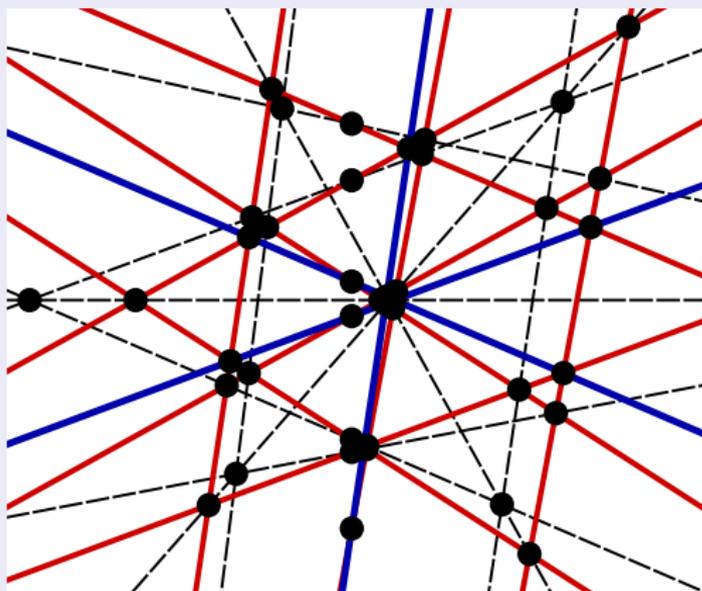


	$SU_5 \times U_1^a \times U_1^b \times U_1^c \times U_1^d$				
T_1	10	1	1	-1	-1
T_2	10	1	-1	-1	-1
T_3	10	-2	0	-1	-1
M_1	$\overline{5}$	1	1	-1	3
M_2	$\overline{5}$	1	-1	-1	3
M_3	$\overline{5}$	-2	0	-1	3
H^u	5	1	1	2	2
H^d	$\overline{5}$	1	1	2	-2
Y_1	5	-1	1	-2	2
Y_1^c	$\overline{5}$	2	0	-2	-2
Y_2	5	2	0	-2	2
Y_2^c	$\overline{5}$	-1	1	-2	-2
X_1	1	-1	-1	4	0
X_2	1	-1	1	4	0
ν_1^c	1	1	1	-1	-5
ν_2^c	1	1	-1	-1	-5
ν_3^c	1	-2	0	-1	-5
N_1^c	1	0	0	-3	5
N_2^c	1	-2	0	-4	0
N_3^c	1	3	-1	0	0
T_X^c	$\overline{10}$	0	0	-3	1
T_X^c	10	0	0	0	4
M_X^c	5	0	0	-3	-3
S_1	1	3	1	0	0
S_2	1	0	2	0	0

A Local Diamond Ring in F-Theory

To enforce the gauge-invariance of $T_1 T_1 H^u$, we should impose the **geometric** condition $b = -a$: $E_8(c+d, 2a+c+d, -3c+d, -4d, 0, 0, 0, 0)$

Constraining the Moduli Space

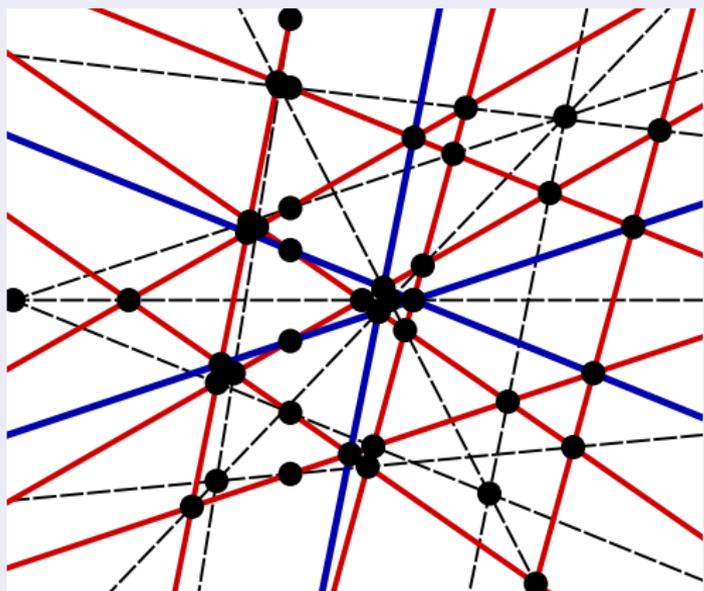


	$SU_5 \times U_1^a \times U_1^b \times U_1^c \times U_1^d$				
T_1	10	1	1	-1	-1
T_2	10	1	-1	-1	-1
T_3	10	-2	0	-1	-1
M_1	$\overline{5}$	1	1	-1	3
M_2	$\overline{5}$	1	-1	-1	3
M_3	$\overline{5}$	-2	0	-1	3
H^u	5	1	1	2	2
H^d	$\overline{5}$	1	1	2	-2
Y_1	5	-1	1	-2	2
Y_1^c	$\overline{5}$	2	0	-2	-2
Y_2	5	2	0	-2	2
Y_2^c	$\overline{5}$	-1	1	-2	-2
X_1	1	-1	-1	4	0
X_2^c	1	-1	1	4	0
ν_1^b	1	1	1	-1	-5
ν_2^c	1	1	-1	-1	-5
ν_3^c	1	-2	0	-1	-5
N_1^b	1	0	0	-3	5
N_2^c	1	-2	0	-4	0
N_3^c	1	3	-1	0	0
T_X^c	$\overline{10}$	0	0	-3	1
T_X^c	10	0	0	0	4
M_X^c	5	0	0	-3	-3
S_1	1	3	1	0	0
S_2	1	0	2	0	0

A Local Diamond Ring in F-Theory

To enforce the gauge-invariance of $T_1 T_1 H^u$, we should impose the **geometric** condition $b = -a$: $E_8(c+d, 2a+c+d, -3c+d, -4d, 0, 0, 0, 0)$

Constraining the Moduli Space

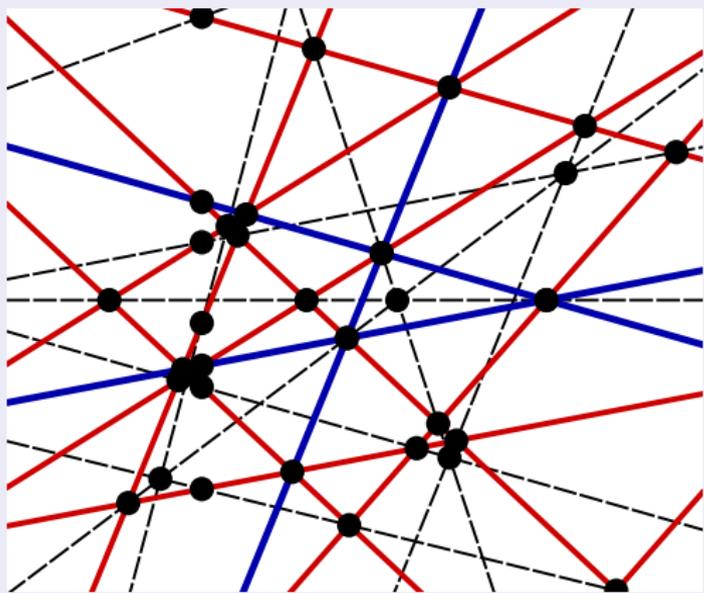


	$SU_5 \times U_1^a \times U_1^b \times U_1^c \times U_1^d$				
T_1	10	1	1	-1	-1
T_2	10	1	-1	-1	-1
T_3	10	-2	0	-1	-1
M_1	$\overline{5}$	1	1	-1	3
M_2	$\overline{5}$	1	-1	-1	3
M_3	$\overline{5}$	-2	0	-1	3
H^u	5	1	1	2	2
H^d	$\overline{5}$	1	1	2	-2
Y_1	5	-1	1	-2	2
Y_1^c	$\overline{5}$	2	0	-2	-2
Y_2	5	2	0	-2	2
Y_2^c	$\overline{5}$	-1	1	-2	-2
X_1	1	-1	-1	4	0
X_2	1	-1	1	4	0
ν_1^c	1	1	1	-1	-5
ν_2^c	1	1	-1	-1	-5
ν_3^c	1	-2	0	-1	-5
N_1^c	1	0	0	-3	5
N_2^c	1	-2	0	-4	0
N_3^c	1	3	-1	0	0
T_X^c	$\overline{10}$	0	0	-3	1
T_X	10	0	0	0	4
M_X^c	5	0	0	-3	-3
S_1	1	3	1	0	0
S_2	1	0	2	0	0

A Local Diamond Ring in F-Theory

To enforce the gauge-invariance of $T_1 T_1 H^u$, we should impose the **geometric** condition $b = -a$: $E_8(c+d, 2a+c+d, -3c+d, -4d, 0, 0, 0, 0)$

Constraining the Moduli Space

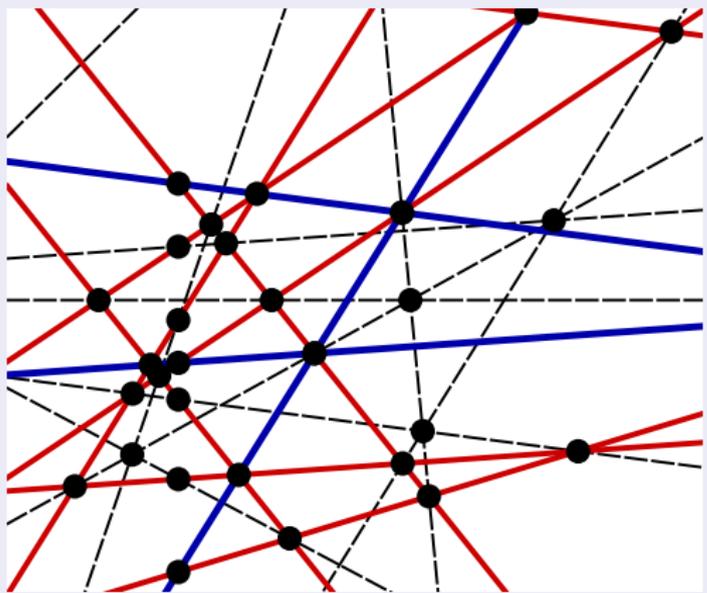


	$SU_5 \times U_1^a \times U_1^b \times U_1^c \times U_1^d$				
T_1	10	1	1	-1	-1
T_2	10	1	-1	-1	-1
T_3	10	-2	0	-1	-1
M_1	$\overline{5}$	1	1	-1	3
M_2	$\overline{5}$	1	-1	-1	3
M_3	$\overline{5}$	-2	0	-1	3
H^u	5	1	1	2	2
H^d	$\overline{5}$	1	1	2	-2
Y_1	5	-1	1	-2	2
Y_1^c	$\overline{5}$	2	0	-2	-2
Y_2	5	2	0	-2	2
Y_2^c	$\overline{5}$	-1	1	-2	-2
X_1	1	-1	-1	4	0
X_2	1	-1	1	4	0
ν_1^t	1	1	1	-1	-5
ν_2^t	1	1	-1	-1	-5
ν_3^t	1	-2	0	-1	-5
N_1^t	1	0	0	-3	5
N_2^t	1	-2	0	-4	0
N_3^t	1	3	-1	0	0
T_X	$\overline{10}$	0	0	-3	1
T_X^c	10	0	0	0	4
M_X^c	5	0	0	-3	-3
S_1	1	3	1	0	0
S_2	1	0	2	0	0

A Local Diamond Ring in F-Theory

To enforce the gauge-invariance of $T_1 T_1 H^u$, we should impose the **geometric** condition $b = -a$: $E_8(c+d, 2a+c+d, -3c+d, -4d, 0, 0, 0, 0)$

Constraining the Moduli Space

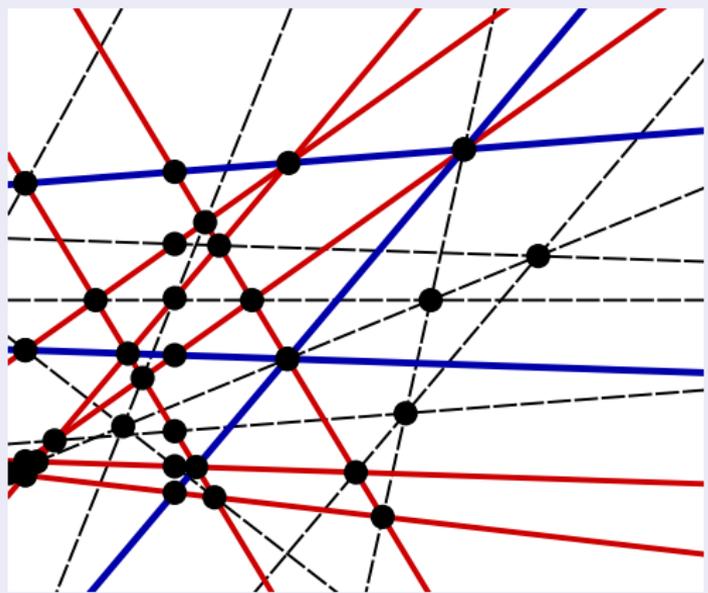


	$SU_5 \times U_1^a \times U_1^b \times U_1^c \times U_1^d$				
T_1	10	1	1	-1	-1
T_2	10	1	-1	-1	-1
T_3	10	-2	0	-1	-1
M_1	$\overline{5}$	1	1	-1	3
M_2	$\overline{5}$	1	-1	-1	3
M_3	$\overline{5}$	-2	0	-1	3
H^u	5	1	1	2	2
H^d	$\overline{5}$	1	1	2	-2
Y_1	5	-1	1	-2	2
Y_1^c	$\overline{5}$	2	0	-2	-2
Y_2	5	2	0	-2	2
Y_2^c	$\overline{5}$	-1	1	-2	-2
X_1	1	-1	-1	4	0
X_2	1	-1	1	4	0
ν_1^t	1	1	1	-1	-5
ν_2^t	1	1	-1	-1	-5
ν_3^t	1	-2	0	-1	-5
N_1^t	1	0	0	-3	5
N_2^t	1	-2	0	-4	0
N_3^t	1	3	-1	0	0
T_X^c	$\overline{10}$	0	0	-3	1
T_X^c	10	0	0	0	4
M_X^c	5	0	0	-3	-3
S_1	1	3	1	0	0
S_2	1	0	2	0	0

A Local Diamond Ring in F-Theory

To enforce the gauge-invariance of $T_1 T_1 H^u$, we should impose the **geometric** condition $b = -a$: $E_8(c+d, 2a+c+d, -3c+d, -4d, 0, 0, 0, 0)$

Constraining the Moduli Space

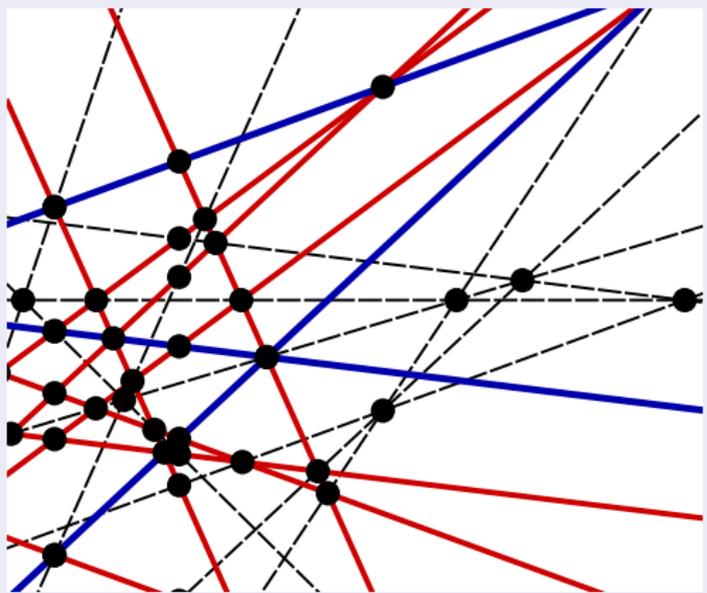


	$SU_5 \times U_1^a \times U_1^b \times U_1^c \times U_1^d$				
T_1	10	1	1	-1	-1
T_2	10	1	-1	-1	-1
T_3	10	-2	0	-1	-1
M_1	$\overline{5}$	1	1	-1	3
M_2	$\overline{5}$	1	-1	-1	3
M_3	$\overline{5}$	-2	0	-1	3
H^u	5	1	1	2	2
H^d	$\overline{5}$	1	1	2	-2
Y_1	5	-1	1	-2	2
Y_1^c	$\overline{5}$	2	0	-2	-2
Y_2	5	2	0	-2	2
Y_2^c	$\overline{5}$	-1	1	-2	-2
X_1	1	-1	-1	4	0
X_2	1	-1	1	4	0
ν_1^b	1	1	1	-1	-5
ν_2^b	1	1	-1	-1	-5
ν_3^b	1	-2	0	-1	-5
N_1^b	1	0	0	-3	5
N_2^b	1	-2	0	-4	0
N_3^b	1	3	-1	0	0
T_X^c	$\overline{10}$	0	0	-3	1
T_X^c	10	0	0	0	4
M_X^c	5	0	0	-3	-3
S_1	1	3	1	0	0
S_2	1	0	2	0	0

A Local Diamond Ring in F-Theory

To enforce the gauge-invariance of $T_1 T_1 H^u$, we should impose the **geometric** condition $b = -a$: $E_8(c+d, 2a+c+d, -3c+d, -4d, 0, 0, 0, 0)$

Constraining the Moduli Space

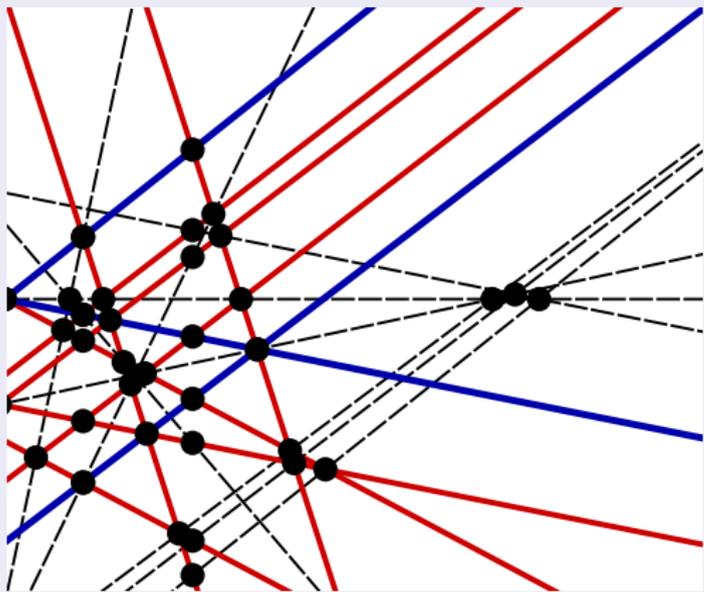


	$SU_5 \times U_1^a \times U_1^b \times U_1^c \times U_1^d$				
T_1	10	1	1	-1	-1
T_2	10	1	-1	-1	-1
T_3	10	-2	0	-1	-1
M_1	$\overline{5}$	1	1	-1	3
M_2	$\overline{5}$	1	-1	-1	3
M_3	$\overline{5}$	-2	0	-1	3
H^u	5	1	1	2	2
H^d	$\overline{5}$	1	1	2	-2
Y_1	5	-1	1	-2	2
Y_1^c	$\overline{5}$	2	0	-2	-2
Y_2	5	2	0	-2	2
Y_2^c	$\overline{5}$	-1	1	-2	-2
X_1	1	-1	-1	4	0
X_2^c	1	-1	1	4	0
ν_1^b	1	1	1	-1	-5
ν_2^b	1	1	-1	-1	-5
ν_3^b	1	-2	0	-1	-5
N_2^b	1	0	0	-3	5
N_2^c	1	-2	0	-4	0
N_3^c	1	3	-1	0	0
T_X^c	$\overline{10}$	0	0	-3	1
T_X^c	10	0	0	0	4
M_X^c	5	0	0	-3	-3
S_1	1	3	1	0	0
S_2	1	0	2	0	0

A Local Diamond Ring in F-Theory

To enforce the gauge-invariance of $T_1 T_1 H^u$, we should impose the **geometric** condition $b = -a$: $E_8(c+d, 2a+c+d, -3c+d, -4d, 0, 0, 0, 0)$

Constraining the Moduli Space

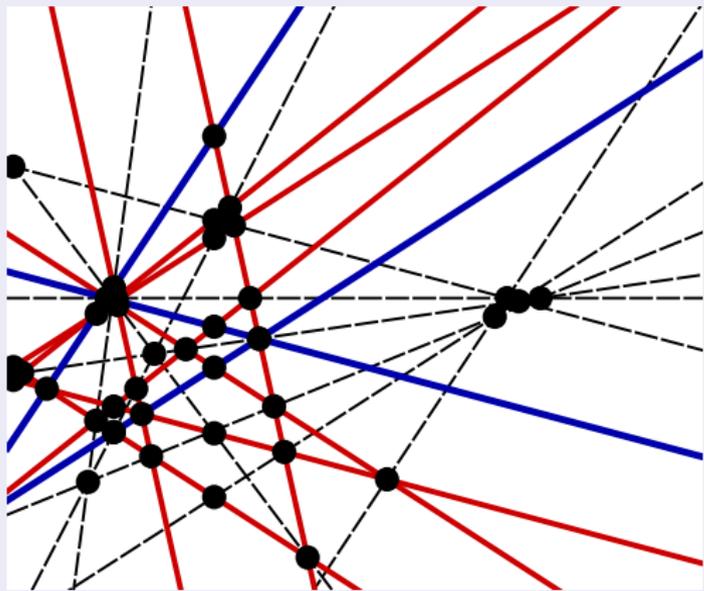


	$SU_5 \times U_1^a \times U_1^b \times U_1^c \times U_1^d$				
T_1	10	1	1	-1	-1
T_2	10	1	-1	-1	-1
T_3	10	-2	0	-1	-1
M_1	$\overline{5}$	1	1	-1	3
M_2	$\overline{5}$	1	-1	-1	3
M_3	$\overline{5}$	-2	0	-1	3
H^u	5	1	1	2	2
H^d	$\overline{5}$	1	1	2	-2
Y_1	5	-1	1	-2	2
Y_1^c	$\overline{5}$	2	0	-2	-2
Y_2	5	2	0	-2	2
Y_2^c	$\overline{5}$	-1	1	-2	-2
X_1	1	-1	-1	4	0
X_2	1	-1	1	4	0
ν_1^t	1	1	1	-1	-5
ν_2^t	1	1	-1	-1	-5
ν_3^t	1	-2	0	-1	-5
N_1^t	1	0	0	-3	5
N_2^t	1	-2	0	-4	0
N_3^t	1	3	-1	0	0
T_X^c	$\overline{10}$	0	0	-3	1
T_X	10	0	0	0	4
M_X^c	5	0	0	-3	-3
S_1	1	3	1	0	0
S_2	1	0	2	0	0

A Local Diamond Ring in F-Theory

To enforce the gauge-invariance of $T_1 T_1 H^u$, we should impose the **geometric** condition $b = -a$: $E_8(c+d, 2a+c+d, -3c+d, -4d, 0, 0, 0, 0)$

Constraining the Moduli Space

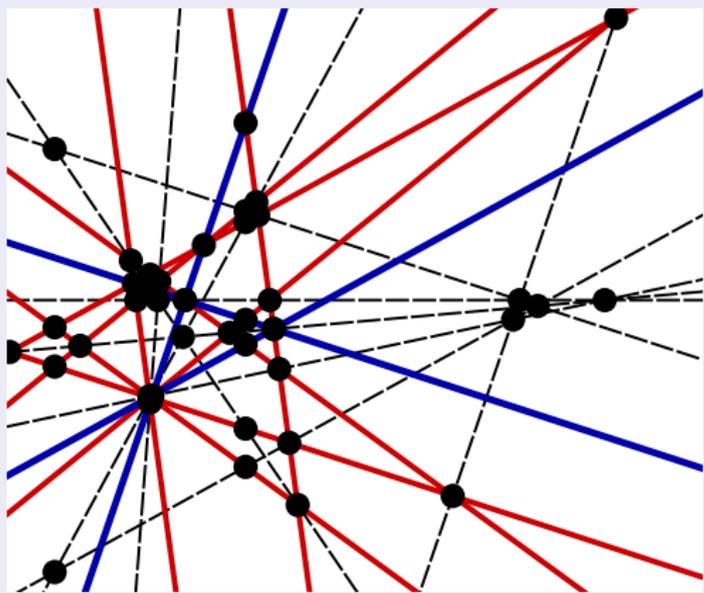


	$SU_5 \times U_1^a \times U_1^b \times U_1^c \times U_1^d$				
T_1	10	1	1	-1	-1
T_2	10	1	-1	-1	-1
T_3	10	-2	0	-1	-1
M_1	$\overline{5}$	1	1	-1	3
M_2	$\overline{5}$	1	-1	-1	3
M_3	$\overline{5}$	-2	0	-1	3
H^u	5	1	1	2	2
H^d	$\overline{5}$	1	1	2	-2
Y_1	5	-1	1	-2	2
Y_1^c	$\overline{5}$	2	0	-2	-2
Y_2	5	2	0	-2	2
Y_2^c	$\overline{5}$	-1	1	-2	-2
X_1	1	-1	-1	4	0
X_2	1	-1	1	4	0
ν_1^t	1	1	1	-1	-5
ν_2^t	1	1	-1	-1	-5
ν_3^t	1	-2	0	-1	-5
N_1^t	1	0	0	-3	5
N_2^t	1	-2	0	-4	0
N_3^t	1	3	-1	0	0
T_X^c	$\overline{10}$	0	0	-3	1
T_X^u	10	0	0	0	4
M_X^c	5	0	0	-3	-3
S_1	1	3	1	0	0
S_2	1	0	2	0	0

A Local Diamond Ring in F-Theory

To enforce the gauge-invariance of $T_1 T_1 H^u$, we should impose the **geometric** condition $b = -a$: $E_8(c+d, 2a+c+d, -3c+d, -4d, 0, 0, 0, 0)$

Constraining the Moduli Space

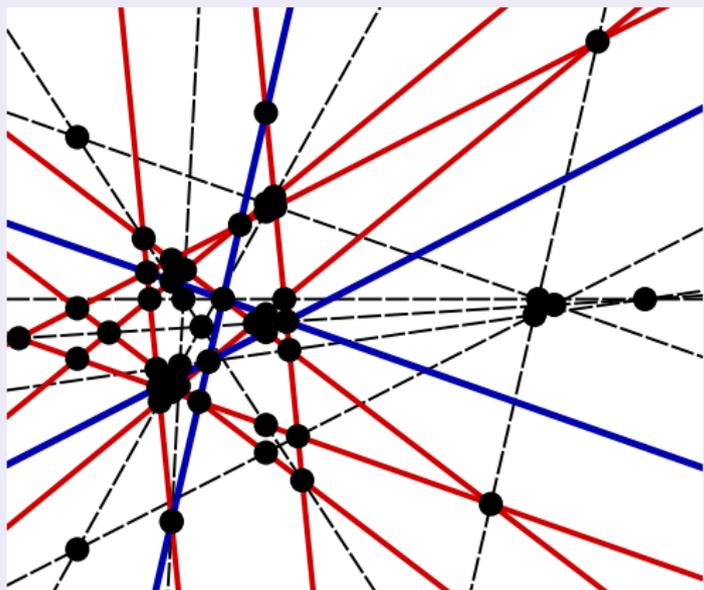


	$SU_5 \times U_1^a \times U_1^b \times U_1^c \times U_1^d$				
T_1	10	1	1	-1	-1
T_2	10	1	-1	-1	-1
T_3	10	-2	0	-1	-1
M_1	$\overline{5}$	1	1	-1	3
M_2	$\overline{5}$	1	-1	-1	3
M_3	$\overline{5}$	-2	0	-1	3
H^u	5	1	1	2	2
H^d	$\overline{5}$	1	1	2	-2
Y_1	5	-1	1	-2	2
Y_1^c	$\overline{5}$	2	0	-2	-2
Y_2	5	2	0	-2	2
Y_2^c	$\overline{5}$	-1	1	-2	-2
X_1	1	-1	-1	4	0
X_2	1	-1	1	4	0
ν_1^b	1	1	1	-1	-5
ν_2^b	1	1	-1	-1	-5
ν_3^b	1	-2	0	-1	-5
N_1^b	1	0	0	-3	5
N_2^b	1	-2	0	-4	0
N_3^b	1	3	-1	0	0
T_X^c	$\overline{10}$	0	0	-3	1
T_X^c	10	0	0	0	4
M_X^c	5	0	0	-3	-3
S_1	1	3	1	0	0
S_2	1	0	2	0	0

A Local Diamond Ring in F-Theory

To enforce the gauge-invariance of $T_1 T_1 H^u$, we should impose the **geometric** condition $b = -a$: $E_8(c+d, 2a+c+d, -3c+d, -4d, 0, 0, 0, 0)$

Constraining the Moduli Space

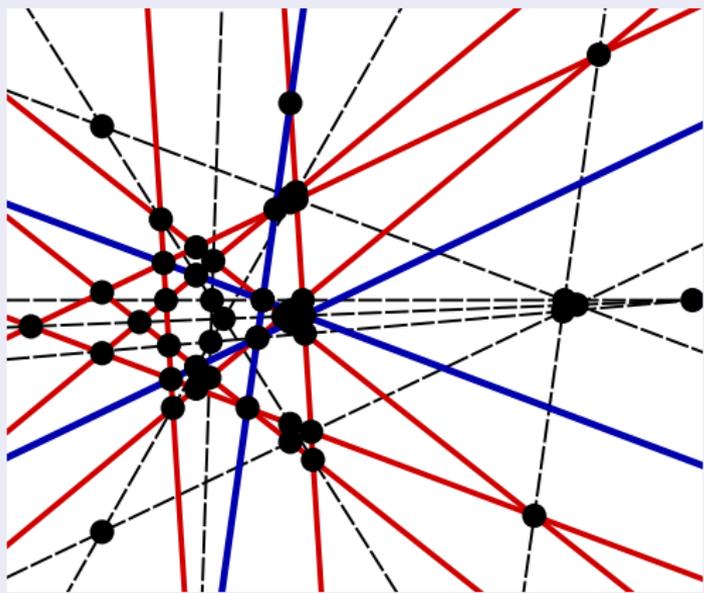


	$SU_5 \times U_1^a \times U_1^b \times U_1^c \times U_1^d$				
T_1	10	1	1	-1	-1
T_2	10	1	-1	-1	-1
T_3	10	-2	0	-1	-1
M_1	$\overline{5}$	1	1	-1	3
M_2	$\overline{5}$	1	-1	-1	3
M_3	$\overline{5}$	-2	0	-1	3
H^u	5	1	1	2	2
H^d	$\overline{5}$	1	1	2	-2
Y_1	5	-1	1	-2	2
Y_1^c	$\overline{5}$	2	0	-2	-2
Y_2	5	2	0	-2	2
Y_2^c	$\overline{5}$	-1	1	-2	-2
X_1	1	-1	-1	4	0
X_2	1	-1	1	4	0
ν_1^b	1	1	1	-1	-5
ν_2^b	1	1	-1	-1	-5
ν_3^b	1	-2	0	-1	-5
N_1^b	1	0	0	-3	5
N_2^b	1	-2	0	-4	0
N_3^b	1	3	-1	0	0
T_X^c	$\overline{10}$	0	0	-3	1
T_X^c	10	0	0	0	4
M_X^c	5	0	0	-3	-3
S_1	1	3	1	0	0
S_2	1	0	2	0	0

A Local Diamond Ring in F-Theory

To enforce the gauge-invariance of $T_1 T_1 H^u$, we should impose the **geometric** condition $b = -a$: $E_8(c+d, 2a+c+d, -3c+d, -4d, 0, 0, 0, 0)$

Constraining the Moduli Space

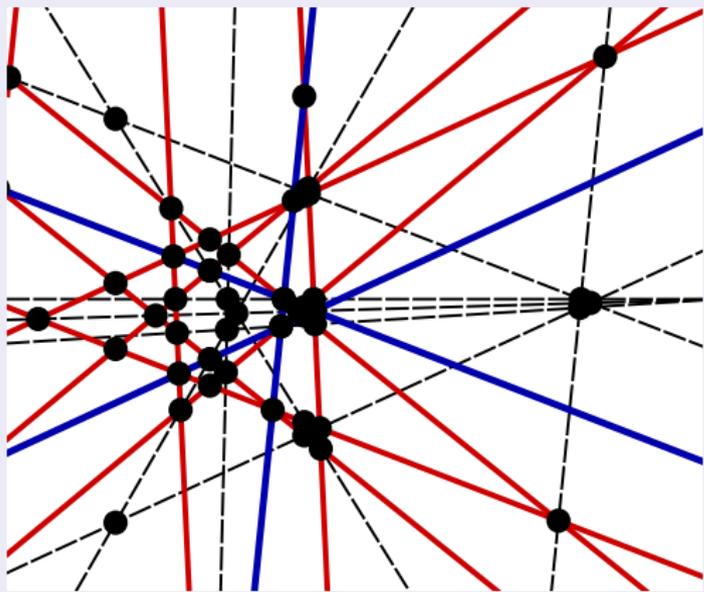


	$SU_5 \times U_1^a \times U_1^b \times U_1^c \times U_1^d$				
T_1	10	1	1	-1	-1
T_2	10	1	-1	-1	-1
T_3	10	-2	0	-1	-1
M_1	$\overline{5}$	1	1	-1	3
M_2	$\overline{5}$	1	-1	-1	3
M_3	$\overline{5}$	-2	0	-1	3
H^u	5	1	1	2	2
H^d	$\overline{5}$	1	1	2	-2
Y_1	5	-1	1	-2	2
Y_1^c	$\overline{5}$	2	0	-2	-2
Y_2	5	2	0	-2	2
Y_2^c	$\overline{5}$	-1	1	-2	-2
X_1	1	-1	-1	4	0
X_2	1	-1	1	4	0
ν_1^b	1	1	1	-1	-5
ν_2^b	1	1	-1	-1	-5
ν_3^b	1	-2	0	-1	-5
N_1^b	1	0	0	-3	5
N_2^b	1	-2	0	-4	0
N_3^b	1	3	-1	0	0
T_X^c	$\overline{10}$	0	0	-3	1
T_X^c	10	0	0	0	4
M_X^c	5	0	0	-3	-3
S_1	1	3	1	0	0
S_2	1	0	2	0	0

A Local Diamond Ring in F-Theory

To enforce the gauge-invariance of $T_1 T_1 H^u$, we should impose the **geometric** condition $b = -a$: $E_8(c+d, 2a+c+d, -3c+d, -4d, 0, 0, 0, 0)$

Constraining the Moduli Space

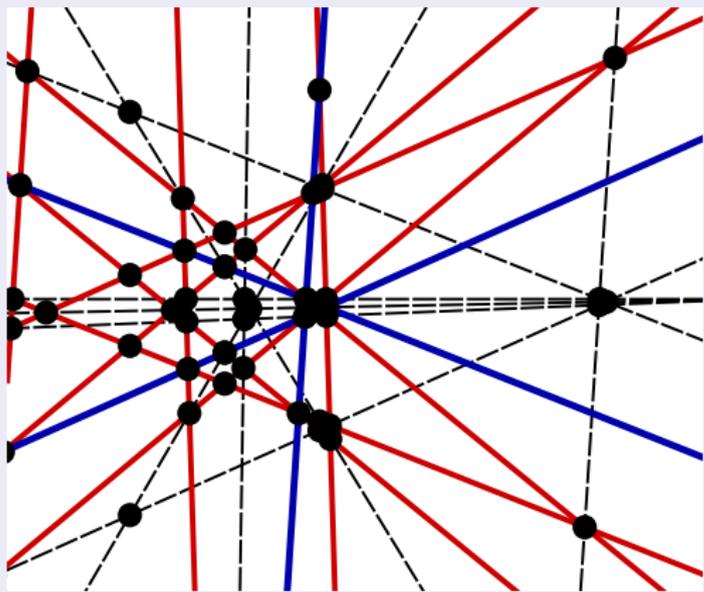


	$SU_5 \times U_1^a \times U_1^b \times U_1^c \times U_1^d$				
T_1	10	1	1	-1	-1
T_2	10	1	-1	-1	-1
T_3	10	-2	0	-1	-1
M_1	$\overline{5}$	1	1	-1	3
M_2	$\overline{5}$	1	-1	-1	3
M_3	$\overline{5}$	-2	0	-1	3
H^u	5	1	1	2	2
H^d	$\overline{5}$	1	1	2	-2
Y_1	5	-1	1	-2	2
Y_1^c	$\overline{5}$	2	0	-2	-2
Y_2	5	2	0	-2	2
Y_2^c	$\overline{5}$	-1	1	-2	-2
X_1	1	-1	-1	4	0
X_2	1	-1	1	4	0
ν_1^b	1	1	1	-1	-5
ν_2^b	1	1	-1	-1	-5
ν_3^b	1	-2	0	-1	-5
N_1^b	1	0	0	-3	5
N_2^b	1	-2	0	-4	0
N_3^b	1	3	-1	0	0
T_X^c	$\overline{10}$	0	0	-3	1
T_X^c	10	0	0	0	4
M_X^c	5	0	0	-3	-3
S_1	1	3	1	0	0
S_2	1	0	2	0	0

A Local Diamond Ring in F-Theory

To enforce the gauge-invariance of $T_1 T_1 H^u$, we should impose the **geometric** condition $b = -a$: $E_8(c+d, 2a+c+d, -3c+d, -4d, 0, 0, 0, 0)$

Constraining the Moduli Space

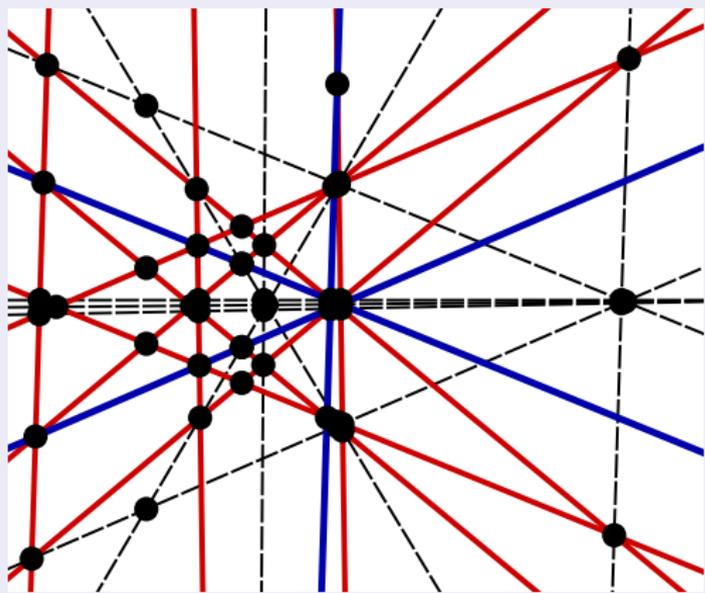


	$SU_5 \times U_1^a \times U_1^b \times U_1^c \times U_1^d$				
T_1	10	1	1	-1	-1
T_2	10	1	-1	-1	-1
T_3	10	-2	0	-1	-1
M_1	$\overline{5}$	1	1	-1	3
M_2	$\overline{5}$	1	-1	-1	3
M_3	$\overline{5}$	-2	0	-1	3
H^u	5	1	1	2	2
H^d	$\overline{5}$	1	1	2	-2
Y_1	5	-1	1	-2	2
Y_1^c	$\overline{5}$	2	0	-2	-2
Y_2	5	2	0	-2	2
Y_2^c	$\overline{5}$	-1	1	-2	-2
X_1	1	-1	-1	4	0
X_2	1	-1	1	4	0
ν_1^t	1	1	1	-1	-5
ν_2^t	1	1	-1	-1	-5
ν_3^t	1	-2	0	-1	-5
N_1^t	1	0	0	-3	5
N_2^t	1	-2	0	-4	0
N_3^t	1	3	-1	0	0
T_X^c	$\overline{10}$	0	0	-3	1
T_X^u	10	0	0	0	4
M_X^c	5	0	0	-3	-3
S_1	1	3	1	0	0
S_2	1	0	2	0	0

A Local Diamond Ring in F-Theory

To enforce the gauge-invariance of $T_1 T_1 H^u$, we should impose the **geometric** condition $b = -a$: $E_8(c+d, 2a+c+d, -3c+d, -4d, 0, 0, 0, 0)$

Constraining the Moduli Space

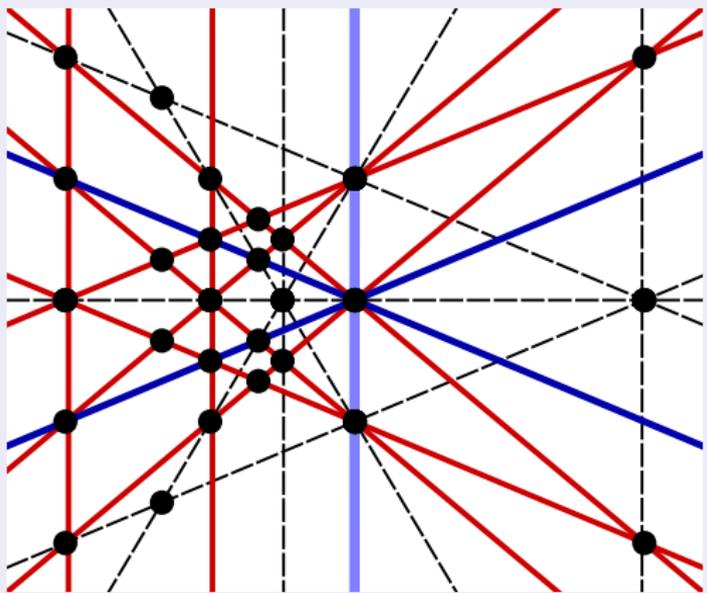


	$SU_5 \times U_1^a \times U_1^b \times U_1^c \times U_1^d$				
T_1	10	1	1	-1	-1
T_2	10	1	-1	-1	-1
T_3	10	-2	0	-1	-1
M_1	$\overline{5}$	1	1	-1	3
M_2	$\overline{5}$	1	-1	-1	3
M_3	$\overline{5}$	-2	0	-1	3
H^u	$\overline{5}$	1	1	2	2
H^d	$\overline{5}$	1	1	2	-2
Y_1	$\overline{5}$	-1	1	-2	2
Y_1^c	$\overline{5}$	2	0	-2	-2
Y_2	$\overline{5}$	2	0	-2	2
Y_2^c	$\overline{5}$	-1	1	-2	-2
X_1	1	-1	-1	4	0
X_2	1	-1	1	4	0
ν_1^b	1	1	1	-1	-5
ν_2^b	1	1	-1	-1	-5
ν_3^b	1	-2	0	-1	-5
N_2^b	1	0	0	-3	5
N_2^c	1	-2	0	-4	0
N_3^c	1	3	-1	0	0
T_X^c	$\overline{10}$	0	0	-3	1
T_X^b	10	0	0	0	4
M_X^c	$\overline{5}$	0	0	-3	-3
S_1	1	3	1	0	0
S_2	1	0	2	0	0

A Local Diamond Ring in F-Theory

To enforce the gauge-invariance of $T_1 T_1 H^u$, we should impose the **geometric** condition $b = -a$: $E_8(c+d, 2a+c+d, -3c+d, -4d, 0, 0, 0, 0)$

Constraining the Moduli Space

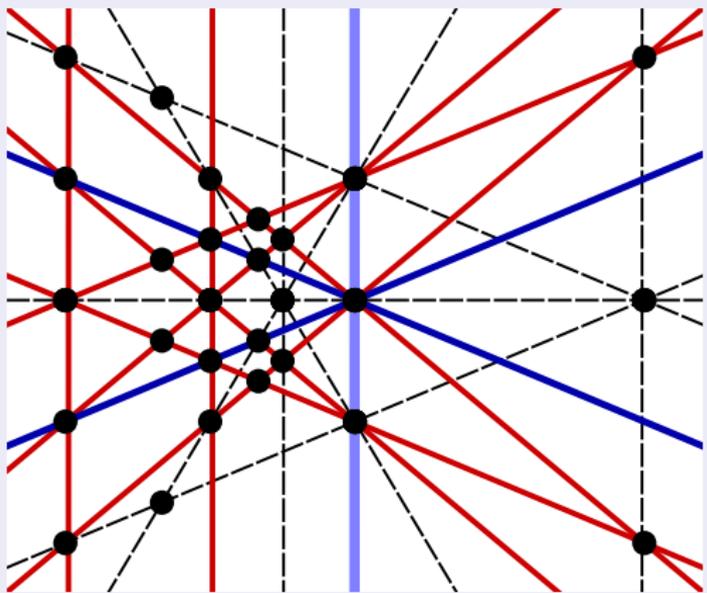


	$SU_5 \times U_1^a \times U_1^c \times U_1^d$			
T_1	10	0	-1	-1
T_2	10	2	-1	-1
T_3	10	-2	-1	-1
M_1	$\overline{5}$	0	-1	3
M_2	$\overline{5}$	2	-1	3
M_3	$\overline{5}$	-2	-1	3
H^u	5	0	2	2
H^d	$\overline{5}$	0	2	-2
Y_1	5	-2	-2	2
Y_1^c	$\overline{5}$	2	-2	-2
Y_2	5	2	-2	2
Y_2^c	$\overline{5}$	-2	-2	-2
X_1	1	0	4	0
X_2	1	-2	4	0
ν_1^t	1	0	-1	-5
ν_2^t	1	2	-1	-5
ν_3^t	1	-2	-1	-5
N_1^t	1	0	-3	5
N_2^t	1	-2	-4	0
N_3^t	1	4	0	0
T_X^c	$\overline{10}$	0	-3	1
T_X	10	0	0	4
M_X^c	5	0	-3	-3
S_1	1	2	0	0

A Local Diamond Ring in F-Theory

To enforce the gauge-invariance of $T_1 T_1 H^u$, we should impose the **geometric** condition $b = -a$: $E_8(c+d, 2a+c+d, -3c+d, -4d, 0, 0, 0, 0)$

Constraining the Moduli Space

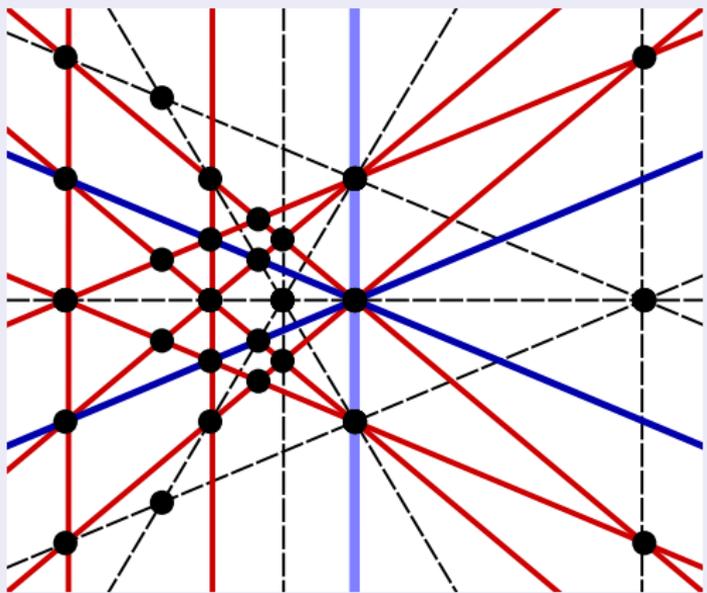


	$SU_5 \times U_1^a \times U_1^c \times U_1^d$			
T_1	10	0	-1	-1
T_2	10	2	-1	-1
T_3	10	-2	-1	-1
M_1	$\overline{5}$	0	-1	3
M_2	$\overline{5}$	2	-1	3
M_3	$\overline{5}$	-2	-1	3
H^u	5	0	2	2
H^d	$\overline{5}$	0	2	-2
Y_1	5	-2	-2	2
Y_1^c	$\overline{5}$	2	-2	-2
Y_2	5	2	-2	2
Y_2^c	$\overline{5}$	-2	-2	-2
X_1	1	0	4	0
X_2	1	-2	4	0
ν_1^t	1	0	-1	-5
ν_2^t	1	2	-1	-5
ν_3^t	1	-2	-1	-5
N_1^t	1	0	-3	5
N_2^t	1	-2	-4	0
N_3^t	1	4	0	0
T_X^c	$\overline{10}$	0	-3	1
T_X	10	0	0	4
M_X^c	5	0	-3	-3
S_1	1	2	0	0

A Local Diamond Ring in F-Theory

To enforce the gauge-invariance of $T_1 T_1 H^u$, we should impose the **geometric** condition $b = -a$: $E_8(c+d, 2a+c+d, -3c+d, -4d, 0, 0, 0, 0)$

Constraining the Moduli Space

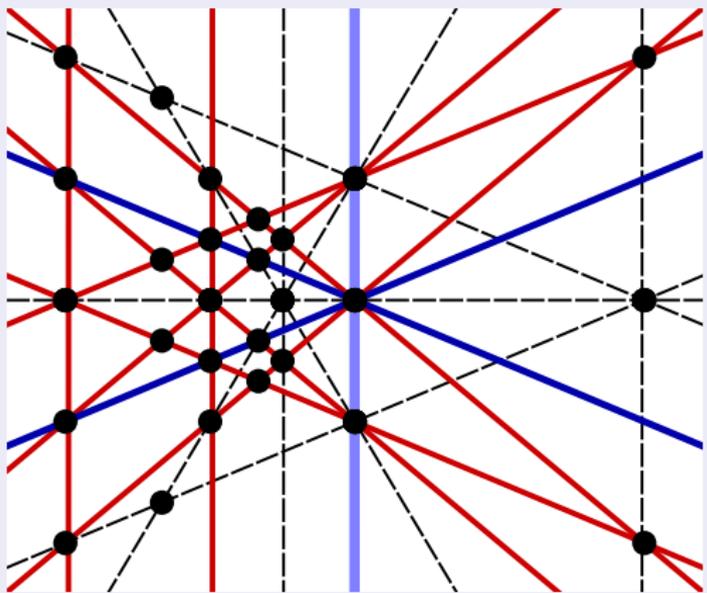


	$SU_5 \times U_1^a \times U_1^c \times U_1^d$			
T_1	10	0	-1	-1
T_2	10	2	-1	-1
T_3	10	-2	-1	-1
M_1	$\overline{5}$	0	-1	3
M_2	$\overline{5}$	2	-1	3
M_3	$\overline{5}$	-2	-1	3
H^u	5	0	2	2
H^d	$\overline{5}$	0	2	-2
Y_1	5	-2	-2	2
Y_1^c	$\overline{5}$	2	-2	-2
Y_2	5	2	-2	2
Y_2^c	$\overline{5}$	-2	-2	-2
X_1	1	0	4	0
X_2	1	-2	4	0
ν_1^t	1	0	-1	-5
ν_2^t	1	2	-1	-5
ν_3^t	1	-2	-1	-5
N_1^t	1	0	-3	5
N_2^t	1	-2	-4	0
N_3^t	1	4	0	0
T_X^c	$\overline{10}$	0	-3	1
T_X	10	0	0	4
M_X^c	5	0	-3	-3
S_1	1	2	0	0

A Local Diamond Ring in F-Theory

To enforce the gauge-invariance of $T_1 T_1 H^u$, we should impose the **geometric** condition $b = -a$: $E_8(c+d, 2a+c+d, -3c+d, -4d, 0, 0, 0, 0)$

Constraining the Moduli Space

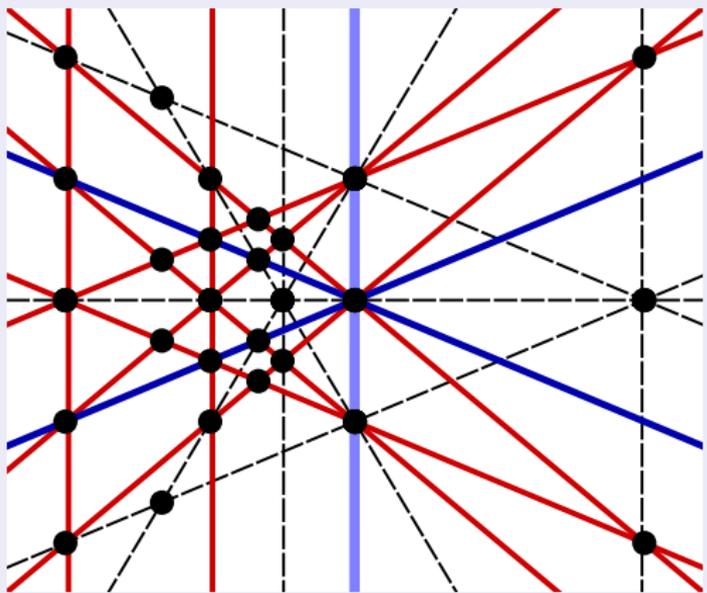


	$SU_5 \times U_1^a \times U_1^c \times U_1^d$			
T_1	10	0	-1	-1
T_2	10	2	-1	-1
T_3	10	-2	-1	-1
M_1	$\overline{5}$	0	-1	3
M_2	$\overline{5}$	2	-1	3
M_3	$\overline{5}$	-2	-1	3
H^u	5	0	2	2
H^d	$\overline{5}$	0	2	-2
Y_1	5	-2	-2	2
Y_1^c	$\overline{5}$	2	-2	-2
Y_2	5	2	-2	2
Y_2^c	$\overline{5}$	-2	-2	-2
X_1	1	0	4	0
X_2	1	-2	4	0
ν_1^t	1	0	-1	-5
ν_2^t	1	2	-1	-5
ν_3^t	1	-2	-1	-5
N_1^t	1	0	-3	5
N_2^t	1	-2	-4	0
N_3^t	1	4	0	0
T_X^c	$\overline{10}$	0	-3	1
T_X	10	0	0	4
M_X^c	5	0	-3	-3
S_1	1	2	0	0

A Local Diamond Ring in F-Theory

To enforce the gauge-invariance of $T_1 T_1 H^u$, we should impose the **geometric** condition $b = -a$: $E_8(c+d, 2a+c+d, -3c+d, -4d, 0, 0, 0, 0)$

Constraining the Moduli Space

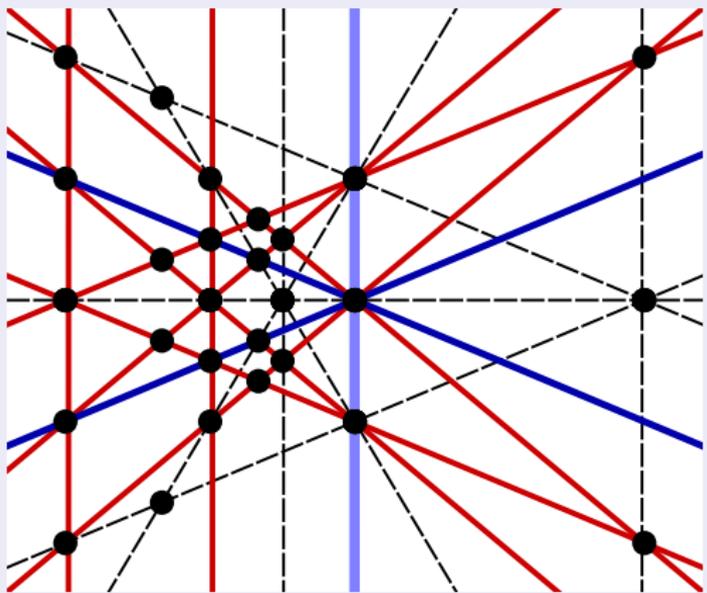


	$SU_5 \times U_1^a \times U_1^c \times U_1^d$			
T_1	10	0	-1	-1
T_2	10	2	-1	-1
T_3	10	-2	-1	-1
M_1	$\overline{5}$	0	-1	3
M_2	$\overline{5}$	2	-1	3
M_3	$\overline{5}$	-2	-1	3
H^u	5	0	2	2
H^d	$\overline{5}$	0	2	-2
Y_1	5	-2	-2	2
Y_1^c	$\overline{5}$	2	-2	-2
Y_2	5	2	-2	2
Y_2^c	$\overline{5}$	-2	-2	-2
X_1	1	0	4	0
X_2	1	-2	4	0
ν_1^t	1	0	-1	-5
ν_2^t	1	2	-1	-5
ν_3^t	1	-2	-1	-5
N_1^t	1	0	-3	5
N_2^t	1	-2	-4	0
N_3^t	1	4	0	0
T_X^c	$\overline{10}$	0	-3	1
T_X	10	0	0	4
M_X^c	5	0	-3	-3
S_1	1	2	0	0

A Local Diamond Ring in F-Theory

To enforce the gauge-invariance of $T_1 T_1 H^u$, we should impose the **geometric** condition $b = -a$: $E_8(c+d, 2a+c+d, -3c+d, -4d, 0, 0, 0, 0)$

Constraining the Moduli Space



	$SU_5 \times U_1^a \times U_1^c \times U_1^d$			
T_1	10	0	-1	-1
T_2	10	2	-1	-1
T_3	10	-2	-1	-1
M_1	$\overline{5}$	0	-1	3
M_2	$\overline{5}$	2	-1	3
M_3	$\overline{5}$	-2	-1	3
H^u	5	0	2	2
H^d	$\overline{5}$	0	2	-2
Y_1	5	-2	-2	2
Y_1^c	$\overline{5}$	2	-2	-2
Y_2	5	2	-2	2
Y_2^c	$\overline{5}$	-2	-2	-2
X_1	1	0	4	0
X_2	1	-2	4	0
ν_1^t	1	0	-1	-5
ν_2^t	1	2	-1	-5
ν_3^t	1	-2	-1	-5
N_1^t	1	0	-3	5
N_2^t	1	-2	-4	0
N_3^t	1	4	0	0
T_X^c	$\overline{10}$	0	-3	1
T_X	10	0	0	4
M_X^c	5	0	-3	-3
S_1	1	2	0	0

A Local Diamond Ring in F-Theory

To enforce the gauge-invariance of $T_1 T_1 H^u$, we should impose the **geometric** condition $b = -a$: $E_8(c+d, 2a+c+d, -3c+d, -4d, 0, 0, 0, 0)$

Once this condition is imposed, we obtain the following superpotential,

$$\begin{aligned}
 W = & T_1 T_1 H^u + T_2 T_3 H^u + T_1 M_1 H^d \\
 & + T_2 M_3 H^d + T_3 M_2 H^d + H^u M_1 \nu_1^c \\
 & + H^u M_2 \nu_3^c + H^u M_3 \nu_2^c + X_1 Y_1 Y_1^c \\
 & + X_1 Y_2 Y_2^c + X_1 N_1^c \nu_1^c + X_2 N_1^c \nu_2^c \\
 & + X_2 N_2^c N_3^c.
 \end{aligned}$$

	$SU_5 \times U_1^a \times U_1^c \times U_1^d$			
T_1	10	0	-1	-1
T_2	10	2	-1	-1
T_3	10	-2	-1	-1
M_1	$\overline{5}$	0	-1	3
M_2	$\overline{5}$	2	-1	3
M_3	$\overline{5}$	-2	-1	3
H^u	5	0	2	2
H^d	$\overline{5}$	0	2	-2
Y_1	5	-2	-2	2
Y_1^c	$\overline{5}$	2	-2	-2
Y_2	5	2	-2	2
Y_2^c	$\overline{5}$	-2	-2	-2
X_1	1	0	4	0
X_2	1	-2	4	0
ν_1^c	1	0	-1	-5
ν_2^c	1	2	-1	-5
ν_3^c	1	-2	-1	-5
N_1^c	1	0	-3	5
N_2^c	1	-2	-4	0
N_3^c	1	4	0	0
T_X^c	$\overline{10}$	0	-3	1
T_X	10	0	0	4
M_X^c	5	0	-3	-3
S_1	1	2	0	0

A Local Diamond Ring in F-Theory

To enforce the gauge-invariance of $T_1 T_1 H^u$, we should impose the **geometric** condition $b = -a$: $E_8(c+d, 2a+c+d, -3c+d, -4d, 0, 0, 0, 0)$

Once this condition is imposed, we obtain the following superpotential,

$$\begin{aligned}
 W = & T_1 T_1 H^u + T_2 T_3 H^u + T_1 M_1 H^d \\
 & + T_2 M_3 H^d + T_3 M_2 H^d + H^u M_1 \nu_1^c \\
 & + H^u M_2 \nu_3^c + H^u M_3 \nu_2^c + X_1 Y_1 Y_1^c \\
 & + X_1 Y_2 Y_2^c + X_1 N_1^c \nu_1^c + X_2 N_1^c \nu_2^c \\
 & + X_2 N_2^c N_3^c.
 \end{aligned}$$

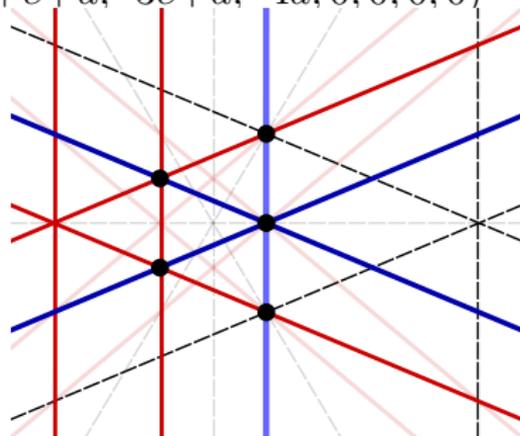
	$SU_5 \times U_1^a \times U_1^c \times U_1^d$			
T_1	10	0	-1	-1
T_2	10	2	-1	-1
T_3	10	-2	-1	-1
M_1	$\overline{5}$	0	-1	3
M_2	$\overline{5}$	2	-1	3
M_3	$\overline{5}$	-2	-1	3
H^u	5	0	2	2
H^d	$\overline{5}$	0	2	-2
Y_1	5	-2	-2	2
Y_1^c	$\overline{5}$	2	-2	-2
Y_2	5	2	-2	2
Y_2^c	$\overline{5}$	-2	-2	-2
X_1	1	0	4	0
X_2	1	-2	4	0
ν_1^c	1	0	-1	-5
ν_2^c	1	2	-1	-5
ν_3^c	1	-2	-1	-5
N_1^c	1	0	-3	5
N_2^c	1	-2	-4	0
N_3^c	1	4	0	0
T_X^c	$\overline{10}$	0	-3	1
T_X	10	0	0	4
M_X^c	5	0	-3	-3
S_1	1	2	0	0

A Local Diamond Ring in F-Theory

To enforce the gauge-invariance of $T_1 T_1 H^u$, we should impose the **geometric** condition $b = -a$: $E_8(c+d, 2a+c+d, -3c+d, -4d, 0, 0, 0, 0)$

Once this condition is imposed, we obtain the following superpotential,

$$\begin{aligned}
 W = & T_1 T_1 H^u + T_2 T_3 H^u + T_1 M_1 H^d \\
 & + T_2 M_3 H^d + T_3 M_2 H^d + H^u M_1 \nu_1^c \\
 & + H^u M_2 \nu_3^c + H^u M_3 \nu_2^c + X_1 Y_1 Y_1^c \\
 & + X_1 Y_2 Y_2^c + X_1 N_1^c \nu_1^c + X_2 N_1^c \nu_2^c \\
 & + X_2 N_2^c N_3^c.
 \end{aligned}$$

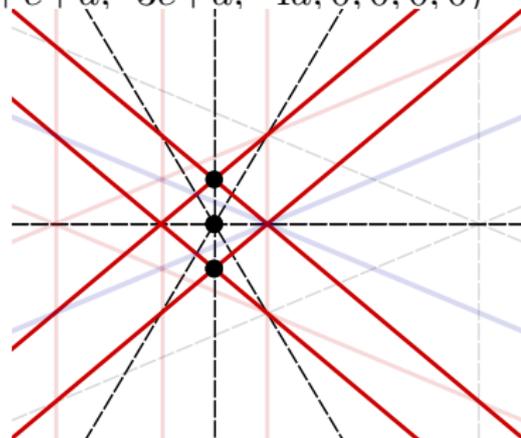


A Local Diamond Ring in F-Theory

To enforce the gauge-invariance of $T_1 T_1 H^u$, we should impose the **geometric** condition $b = -a$: $E_8(c+d, 2a+c+d, -3c+d, -4d, 0, 0, 0, 0)$

Once this condition is imposed, we obtain the following superpotential,

$$\begin{aligned}
 W = & T_1 T_1 H^u + T_2 T_3 H^u + T_1 M_1 H^d \\
 & + T_2 M_3 H^d + T_3 M_2 H^d + H^u M_1 \nu_1^c \\
 & + H^u M_2 \nu_3^c + H^u M_3 \nu_2^c + X_1 Y_1 Y_1^c \\
 & + X_1 Y_2 Y_2^c + X_1 N_1^c \nu_1^c + X_2 N_1^c \nu_2^c \\
 & + X_2 N_2^c N_3^c.
 \end{aligned}$$



A Local Diamond Ring in F-Theory

To enforce the gauge-invariance of $T_1 T_1 H^u$, we should impose the **geometric** condition $b = -a$: $E_8(c+d, 2a+c+d, -3c+d, -4d, 0, 0, 0, 0)$

Once this condition is imposed, we obtain the following superpotential,

$$\begin{aligned}
 W = & T_1 T_1 H^u + T_2 T_3 H^u + T_1 M_1 H^d \\
 & + T_2 M_3 H^d + T_3 M_2 H^d + H^u M_1 \nu_1^c \\
 & + H^u M_2 \nu_3^c + H^u M_3 \nu_2^c + X_1 Y_1 Y_1^c \\
 & + X_1 Y_2 Y_2^c + X_1 N_1^c \nu_1^c + X_2 N_1^c \nu_2^c \\
 & + X_2 N_2^c N_3^c.
 \end{aligned}$$

Notice that the operator $X_1^\dagger H^u H^d$ is gauge invariant.

	$SU_5 \times U_1^a \times U_1^c \times U_1^d$			
T_1	10	0	-1	-1
T_2	10	2	-1	-1
T_3	10	-2	-1	-1
M_1	$\overline{5}$	0	-1	3
M_2	$\overline{5}$	2	-1	3
M_3	$\overline{5}$	-2	-1	3
H^u	5	0	2	2
H^d	$\overline{5}$	0	2	-2
Y_1	5	-2	-2	2
Y_1^c	$\overline{5}$	2	-2	-2
Y_2	5	2	-2	2
Y_2^c	$\overline{5}$	-2	-2	-2
X_1	1	0	4	0
X_2	1	-2	4	0
ν_1^c	1	0	-1	-5
ν_2^c	1	2	-1	-5
ν_3^c	1	-2	-1	-5
N_1^c	1	0	-3	5
N_2^c	1	-2	-4	0
N_3^c	1	4	0	0
T_X^c	$\overline{10}$	0	-3	1
T_X	10	0	0	4
M_X^c	5	0	-3	-3
S_1	1	2	0	0

Conclusions and Future Directions

- We have seen that **explicit**, **purely local** phenomenological models **with three generations** exist in both F-theory and M-theory, but only just-so.
- These models are ‘in principle’ explicit enough to be exhaustively studied, and even make **falsifiable** predictions for low-energy physics.
 - To what extent is this claim true in both F-theory and M-theory?
 - How many assumptions are required in each case?
- What about moduli stabilization? Are there compact manifolds with these local patches? How is the locally continuous landscape of fibrations quantized by compactification?
- ...

Conclusions and Future Directions

- We have seen that **explicit**, **purely local** phenomenological models **with three generations** exist in both F-theory and M-theory, but only just-so.
- These models are ‘in principle’ explicit enough to be exhaustively studied, and even make **falsifiable** predictions for low-energy physics.
 - To what extent is this claim true in both F-theory and M-theory?
 - How many assumptions are required in each case?
- What about moduli stabilization? Are there compact manifolds with these local patches? How is the locally continuous landscape of fibrations quantized by compactification?
- ...

Conclusions and Future Directions

- We have seen that **explicit**, **purely local** phenomenological models **with three generations** exist in both F-theory and M-theory, but only just-so.
- These models are ‘in principle’ explicit enough to be exhaustively studied, and even make **falsifiable** predictions for low-energy physics.
 - To what extent is this claim true in both F-theory and M-theory?
 - How many assumptions are required in each case?
- What about moduli stabilization? Are there compact manifolds with these local patches? How is the locally continuous landscape of fibrations quantized by compactification?
- ...

Conclusions and Future Directions

- We have seen that **explicit**, **purely local** phenomenological models **with three generations** exist in both F-theory and M-theory, but only just-so.
- These models are ‘in principle’ explicit enough to be exhaustively studied, and even make **falsifiable** predictions for low-energy physics.
 - To what extent is this claim true in both F-theory and M-theory?
 - How many assumptions are required in each case?
- What about moduli stabilization? Are there compact manifolds with these local patches? How is the locally continuous landscape of fibrations quantized by compactification?
- ...

Conclusions and Future Directions

- We have seen that **explicit**, **purely local** phenomenological models **with three generations** exist in both F-theory and M-theory, but only just-so.
- These models are ‘in principle’ explicit enough to be exhaustively studied, and even make **falsifiable** predictions for low-energy physics.
 - To what extent is this claim true in both F-theory and M-theory?
 - How many assumptions are required in each case?
- What about moduli stabilization? Are there compact manifolds with these local patches? How is the locally continuous landscape of fibrations quantized by compactification?
- ...

Conclusions and Future Directions

- We have seen that **explicit**, **purely local** phenomenological models **with three generations** exist in both F-theory and M-theory, but only just-so.
- These models are ‘in principle’ explicit enough to be exhaustively studied, and even make **falsifiable** predictions for low-energy physics.
 - To what extent is this claim true in both F-theory and M-theory?
 - How many assumptions are required in each case?
- What about moduli stabilization? Are there compact manifolds with these local patches? How is the locally continuous landscape of fibrations quantized by compactification?
- ...

Conclusions and Future Directions

- We have seen that **explicit**, **purely local** phenomenological models **with three generations** exist in both F-theory and M-theory, but only just-so.
- These models are ‘in principle’ explicit enough to be exhaustively studied, and even make **falsifiable** predictions for low-energy physics.
 - To what extent is this claim true in both F-theory and M-theory?
 - How many assumptions are required in each case?
- What about moduli stabilization? Are there compact manifolds with these local patches? How is the locally continuous landscape of fibrations quantized by compactification?
- ...

Conclusions and Future Directions

- We have seen that **explicit**, **purely local** phenomenological models **with three generations** exist in both F-theory and M-theory, but only just-so.
- These models are ‘in principle’ explicit enough to be exhaustively studied, and even make **falsifiable** predictions for low-energy physics.
 - To what extent is this claim true in both F-theory and M-theory?
 - How many assumptions are required in each case?
- What about moduli stabilization? Are there compact manifolds with these local patches? How is the locally continuous landscape of fibrations quantized by compactification?
- ...