

# Non-perturbative effects and wall crossing in 4d $N=1,2$ string vacua

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Based on I. Garcia-Etxebarria, A.U., arXiv:0711.1430  
I. Garcia-Etxebarria, F. Marchesano, A.U. arXiv:0805.0713  
A.U. arXiv:0808.2918  
A. Collinucci, P. Soler, A.U., arXiv:0904.1133

## Motivation

D-brane instantons as important ingredients in String Pheno  
Moduli stabilization, violation of global  $U(1)$ 's,...

Need good knowledge of their properties, e.g.

- Understand D-brane instanton effects, globally in moduli space

  - Solve certain puzzles involving wall crossing

  - Provide powerful criteria for the generation of superpotentials

- Understand instanton sum, beyond individual contributions

  - Multi-instantons (aka poly-inst)

  - Multiply wrapped instantons

- Understand connection with topological string

  - Wall crossing behaviour manifest

  - New computational tools in 4d  $N=2, 1$

## D-brane instantons and wall crossing

- Non-perturbative 4d effective action from D-brane instantons
  - BPS instantons  $\Rightarrow$  4d F-terms (“holomorphic”) e.g. superpotential
  - Non-BPS instantons  $\Rightarrow$  4d D-terms (“non-holomorphic”)
- List of BPS D-brane instantons can jump discontinuously by moving one real parameter closed string modulus: **Walls of BPS stability**

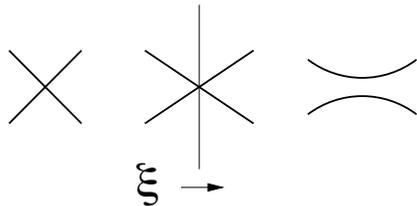
Closed string modulus couples as FI-term

Fayet model on worldvolume of D-brane instanton

Marginal stability:

BPS  $\Rightarrow$  non-BPS

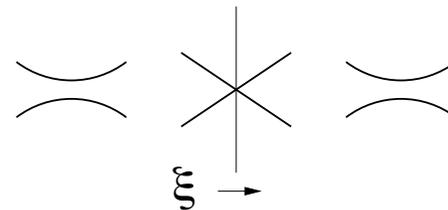
$$V_D = (|\phi|^2 - \xi)^2$$



Threshold stability:

BPS  $\Rightarrow$  split BPS

$$V_D = (|\phi_1|^2 - |\phi_2|^2 - \xi)^2$$



- Fate of non-perturbative terms upon wall crossing?



Clue to the resolution:

Number of (exact) fermion zero modes for D-brane instanton

- Determines the kind of 4d superspace interaction
- Is topological (continuous upon change of parameters)
- Must include Goldstinos, at least 2 for BPS, 4 for non-BPS

### Application to non-perturbative superpotentials

- Generated by BPS instantons with exactly 2 fermion zero modes
- Instantons contributing to superpotential can never become non-BPS

Not enough fermion zero modes to account for 4 required goldstinos

⇒ Safe against marginal wall crossing

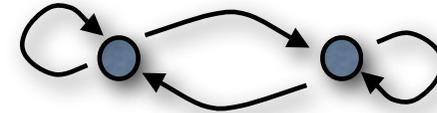
⇒ Powerful criterion: Any instanton which can become non-BPS cannot contribute to superpotential

- What about threshold stability?

# Threshold stability and non-pert. superpotentials

📌 BPS instanton splits into two BPS instantons

2-instanton process: 0-dim quiver theory



Ex: Translational Goldstones  $x_1, x_2$ ; “Goldstinos”  $\theta_1, \tilde{\theta}_1, \theta_2$ ;  
bi-fundamental hyperm.  $\Phi_{12}, \Phi_{21}$  ie  $\varphi_{12}, \varphi_{21}, \chi_{12}, \chi_{21}$

Contribution to superpotential localize onto  $x_1=x_2$ , couplings are

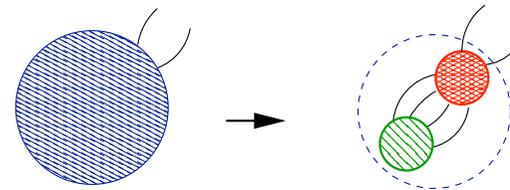
$$(\chi_{12}(\theta_1 - \theta_2))\varphi_{12}^* - (\chi_{21}(\theta_1 - \theta_2))\varphi_{21}^* + (\bar{\chi}_{12}\tilde{\theta})\varphi_{12} - (\bar{\chi}_{21}\tilde{\theta})\varphi_{21} \\ \chi_{12}\varphi_{21}\chi_{12}\varphi_{21} + 2\chi_{12}\chi_{21}\varphi_{12}\varphi_{21} + \varphi_{12}\chi_{21}\varphi_{12}\chi_{21} + \text{h.c.}$$

All fermions couple except for the overall Goldstinos  $\theta_1+\theta_2$

Pull down interactions in  $\exp(-S_{\text{inst}})$  and soak up zero modes

We recover 
$$S_{4d} \simeq \int d^4 d^2\theta e^{-(T_1+T_2)} = \int d^4 d^2\theta e^{-T}$$

📌 Decay products combine into  
2-instanton process reconstructing  
amplitude before decay



## Marginal stability and 4d N=1 higher F-terms

### 📍 Global picture for instantons generating higher F-terms

- Non-perturbative higher F-terms are continuous across general lines of **marginal** stability (BPS  $\Rightarrow$  non-BPS)
- Consistent w/ standard wisdom of BPS=F-term, non-BPS=D-term

Inst. amplitude as 4d operator in non-trivial Beasley-Witten cohomology:

- Locally in moduli space, can be written as a D-term
- Obstruction (localized on BPS locus) to write as global D-term

Holomorphy of higher F-term throughout moduli space:

- Full instanton amplitude and expression at BPS locus differ by a D-term

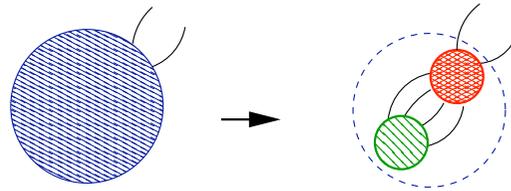
$$S_{4d} \simeq \int d^4x d^2\theta e^{-T} \bar{D}\bar{\Sigma} \bar{D}\bar{\Sigma} + (\text{D-term}) \simeq \int d^4x d^4\theta e^{-T} f(\Sigma, \bar{\Sigma})$$

$\longleftarrow$  core F-term (BPS locus) ↑ Away from BPS locus

### 📍 Subsequent beautiful analysis in 4d N=2 by Gaiotto, Moore, Neitzke

## The topological string connection

📌 4d non-perturbative terms are insensitive to the stability of D-brane instantons (decay products of unstable instanton reconstruct it)



📌 Reminiscent of D-branes in topological strings, defined by holomorphic conditions, without imposing stability condition (impose worldvolume F-term, but don't impose worldvolume D-term)

Non-perturbative F-terms computable in category of holomorphic branes

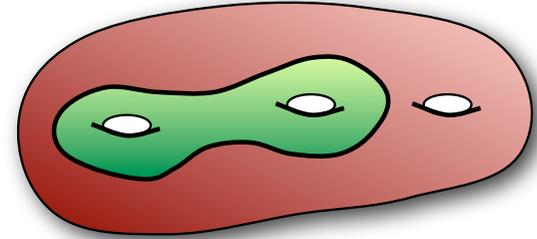
📌 Plan:

- Topological strings and non-perturbative effects in 4d,  $N=2$
- Lessons for 4d,  $N=1$

# GV interpretation of topological strings [Gopakumar, Vafa '98]

- Topological string partition function

$$Z_{\text{top}} = \sum_g \lambda^{2g-2} F_g(t)$$



- Computes F-terms for v.m. in 4d N=2 IIA on CY [Antoniadis, Gava, Narain, Taylor '93]

$$S_{4d} = \int d^4x d^4\theta \sum_g F_g(t) \mathcal{W}^{2g} = \int d^4x F_g(t) R_+^2 F_+^{2g-2} + \dots$$

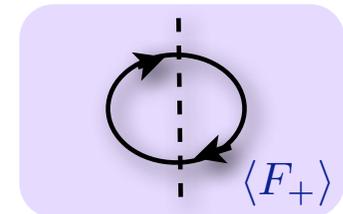
- Turn on graviphoton background  $\langle F_+ \rangle = \lambda \rightarrow S_{4d} = \int d^4x Z_{\text{top}}(\lambda) R_+^2$

$\Rightarrow$  Schwinger diagram for electric charges: D2/D0-branes

Perturbative  $F_g$  from integrating out genus  $g$  D2/D0 particles

$\Rightarrow$  Non-pert. creation of 4d D2/D0-particles by Schwinger effect

- Interpretation in M-theory on CY  $\times S^1$  as wrapped genus  $g$  M2 particles with momentum on  $S^1$



$$Z_{\text{top,n.p.}} = \sum_{r, m, (n_{\mathbf{k}}^r)} n_{\mathbf{k}}^r \int_{\epsilon}^{\infty} \frac{ds}{s} \left(2 \sin \frac{s}{2}\right)^{2r-2} \exp \left[ -2\pi \frac{s}{\lambda} (k_i t_i + i m) \right]$$

## D-brane instantons from M-theory on T2

- GV give non perturbative effects for v.m. from 4d D2/D0-particles
- Relate to D-brane instantons by  $S^1$  compactification and T-duality

[Ooguri, Vafa '96]

### Non-pert. effects on hyperm. moduli space of dual 4d IIB on CY

- IIA on  $CY \times S^1$  computed as M-theory on  $CY \times T^2$   
(genus  $g$  M2 particles with momentum on  $T^2$ ), then shrink  $T^2$  for IIB

Manifest  $SL(2, \mathbb{Z})$  on type IIB side

Includes D1/D(-1)-brane instantons (general  $(p, q)$  string instantons)

- Not much known on higher F-terms for hm's, focus on genus 0

### Corrections to hm metric

$$Z_{M2} = \frac{1}{4\pi} \sum_{\mathbf{k}} n_{k_a}^{(0)} \sum_{(m,n) \neq (0,0)} \frac{\tau_2^{3/2}}{|m + \tau n|^{3/2}} (1 + 2\pi |m + \tau n| k_a t^a) e^{-S_{k_1, k_2}}$$

$$\downarrow$$

$$Z_{D-inst} = \frac{1}{4\pi} \sum_{k_a} n_{k_a} \sum_{m \neq 0, n \in \mathbb{Z}} \frac{|z + q|^{1/2}}{|p|^{3/2}} \left[ 1 + \sum_{k=1}^{\infty} \frac{\Gamma(3/2 + k)}{k! \Gamma(3/2 - k)} (4\pi |p\tau_2| |z + q|)^{-k} \right] e^{-S_{(p,q)}}$$

Reproduces [Robles-Llana, Rocek, Saueressig, Theis, Vandoren, '06]

## Continuity under wall crossing

📌 GV includes only instantons with zero mutual “intersection number” (e.g. no D1, D(-1) but no D3, D5) i.e. no magnetic charges in Schwinger

⇒ Only threshold stability walls

⇒ Enough to discuss continuity of superpotential (in 4d N=1 setup)

📌 Continuity across threshold walls from continuity of GV invariants

📌 More in detail, as a BPS D-brane particle splits into two BPS ones

2-particle system: 1-dim quiver quantum mechanics [Denef '02]

2-particle system has one bound state at threshold

Bound state in Schwinger loop ensures continuity



📌 Interesting interpretation from the particle-instanton dictionary

Particle bound state in Schwinger loop maps to 2-instanton process

# Implications and tools for $N=1$

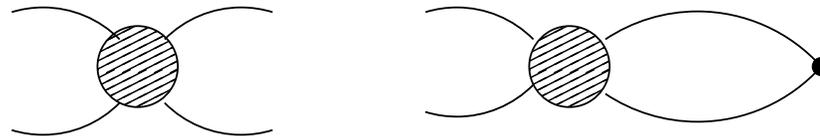
Need to reduce to  $N=1$  to really discuss superpotentials

📌 Turn on  $N=2 \rightarrow N=1$  fluxes

- Superpotential from flux lifting of fermion zero modes

[Bergshoeff, Kallosh, Kashani-Poor, Sorkin, Tomasiello; ...;  
Billo, Ferro, Frau, Fucito, Lerda, Morales]

- In effective theory, just dress non-pert.  $N=2$  hm metric with flux superpotential to obtain non.pert superpotential of  $N=1$  flux model



[A.U.]

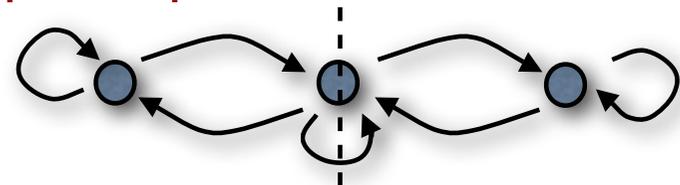
📌 Introduce orientifold planes

- Topological string may produce useful tools [work in progress]

- Partial lesson to study threshold walls: particle-instanton dictionary

2 -particle system: orientifolded 1-dim quiver quantum mechanics

Can check existence of bound state



## Conclusions

We have studied certain formal properties of D-brane instantons

- Showed holomorphy across stability walls
- Criteria for the generation of superpotentials

Instantons which can become non-BPS cannot generate superp.

(can iff misalignment breaks spacetime susy)

- Described relation to topological string in 4d  $N=2$

Resum of D-brane instanton corrections to hm metric

Wall crossing behaviour manifest

Particle-instanton T-duality: New computational tools in 4d  $N=2, 1$

[in progress]

Expect further applications to  $N=1$  model building