On Type IIA Cosmology From Geometric Fluxes

Marco Zagermann (MPI for Physics, Munich)





Based on: 0812.3551 (Caviezel, Koerber, Körs, Lüst, Wrase, M.Z.) 0806.3458 (Caviezel, Koerber, Körs, Lüst, Tsimpis, M.Z.) (0812.3886 (Flauger, Paban, Robbins, Wrase))

An important problem in string phenomenology:

Moduli stabilization

$$\Rightarrow V(\varphi^{i})$$

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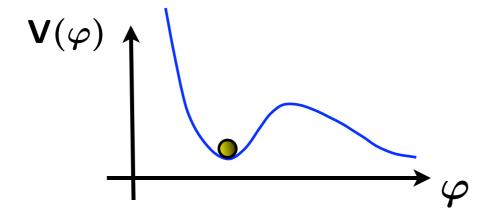
$$\Rightarrow V(\varphi^{i})$$

Particularly interesting:

$$\mathsf{V}(arphi^\mathsf{i}) > \mathsf{0}$$

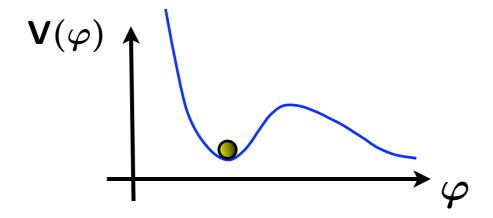
(i) de Sitter vacua

 $(\Lambda > 0 \Rightarrow Today's accelerated expansion)$



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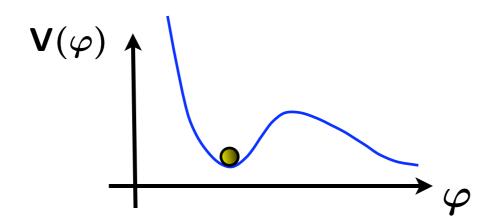


(ii) Slow-roll inflation

$$\mathbf{V}(\varphi,\varphi^{\perp})$$
 (Inflaton) (Stabilized orthogonal fieds)

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(ii) Slow-roll inflation

$$V(\varphi, \varphi^{\perp})$$
 φ $\eta \equiv 0$ (Inflaton) φ (Stabilized orthogonal fieds)

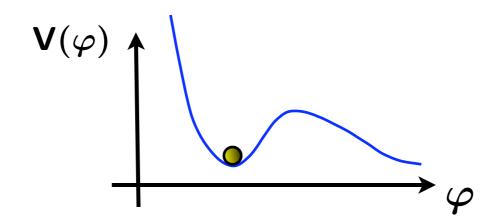
$$\epsilon \equiv rac{1}{2} \, \mathbf{g}^{\mathbf{i} \mathbf{j}} \, rac{(\partial_{arphi^{\mathbf{i}}} \mathsf{V}) \, (\partial_{arphi^{\mathbf{j}}} \mathsf{V})}{\mathsf{V}^2}$$

$$\eta \equiv$$
 Min. eig.val. $\left(rac{oldsymbol{
abla}^i\partial_joldsymbol{V}}{oldsymbol{V}}
ight)$

$$\epsilon, |\eta| \ll 1$$

(i) de Sitter vacua

 $(\Lambda > 0 \Rightarrow Today's accelerated expansion)$



$$\epsilon = 0, \quad \eta > 0$$

(ii) Slow-roll inflation

$$\mathbf{V}(\varphi, \varphi^{\perp})$$
 φ $\eta \equiv \frac{1}{\varphi}$ (Inflaton) (Stabilized orthogonal fieds)

$$\epsilon \equiv rac{1}{2} \, \mathbf{g}^{\mathbf{i} \mathbf{j}} \, rac{(\partial_{arphi^{\mathbf{i}}} \mathsf{V}) \, (\partial_{arphi^{\mathbf{j}}} \mathsf{V})}{\mathsf{V}^2}$$
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A general problem:

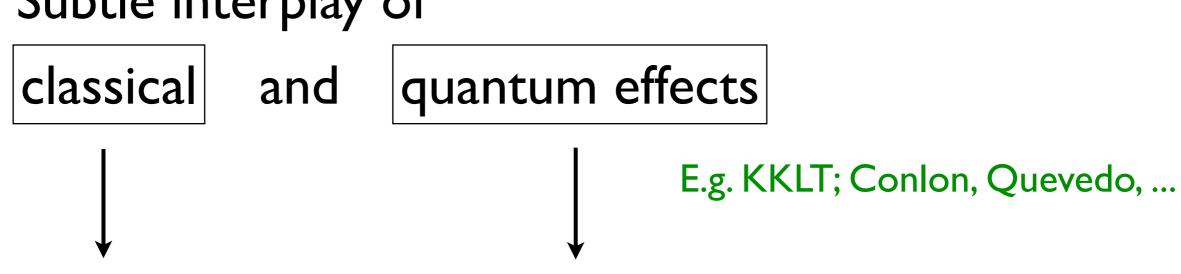
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A general problem:

Typical scalar potentials receive many contributions and corrections

Often:

Subtle interplay of



Easy

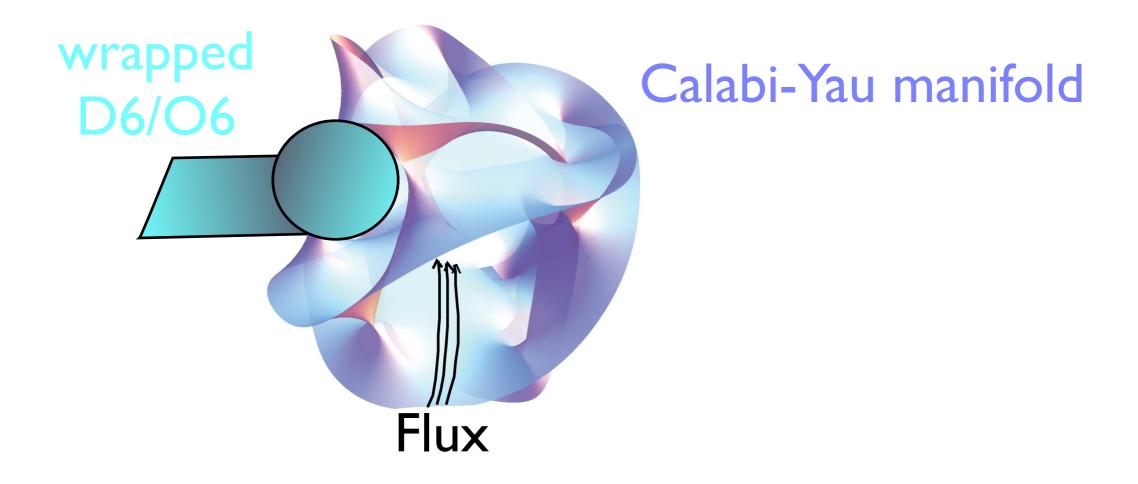
Hard to compute precisely

Cf. McAllister's talk

A nice laboratory:

Type IIA on Calabi-Yau spaces with

- Magnetic fluxes of p-form field strengths
- D6-branes/O6-planes



Observation:

All geometric moduli can be stabilized at tree-level

Grimm, Louis (2004); Kachru, Kashani-Poor (2004)

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In special cases:

- All moduli stabilized
- Parameterically controlled classical regime
- ⇒ Quantum corrections small

Derendinger, Kounnas, Petropoulos, Zwirner (2004, 2005)

Villadoro, Zwirner (2005)

de Wolfe, Giryavets, Kachru, Taylor (2005)

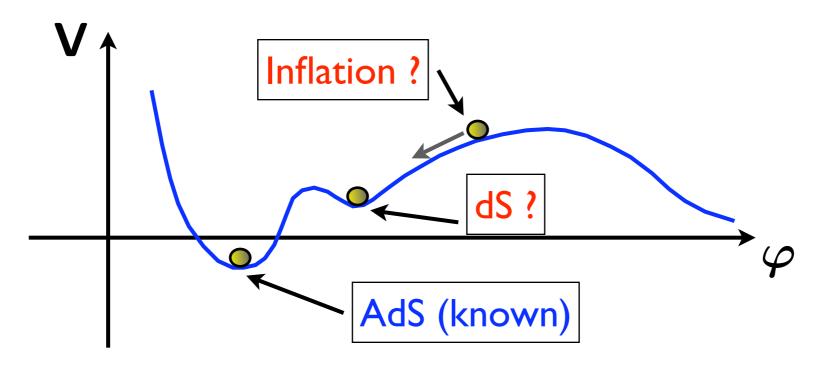
Unfortunately...

All these stabilized vacua have

$$\Lambda$$
<0 (\Rightarrow AdS)

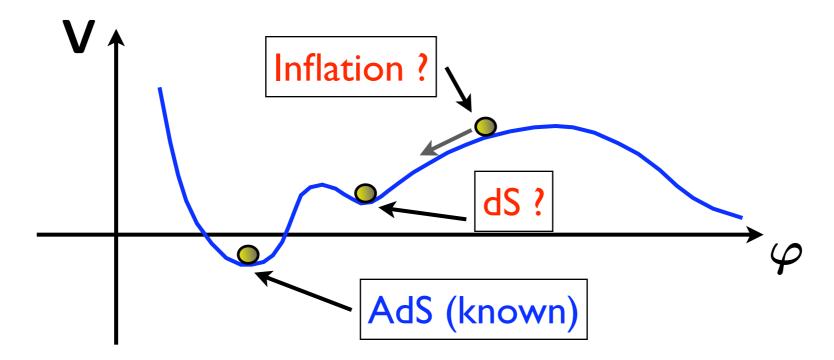
Two possibilities:

(i) Search for dS/inflation away from AdS vacuum

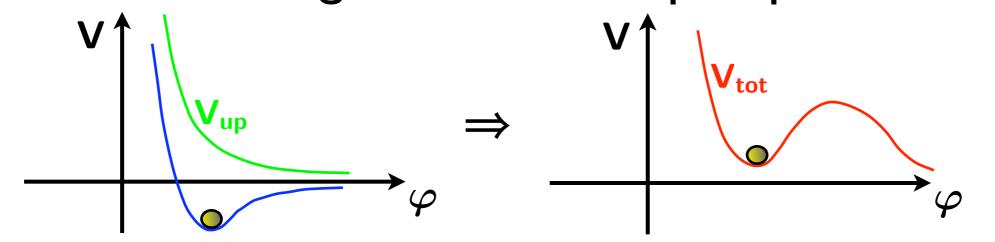


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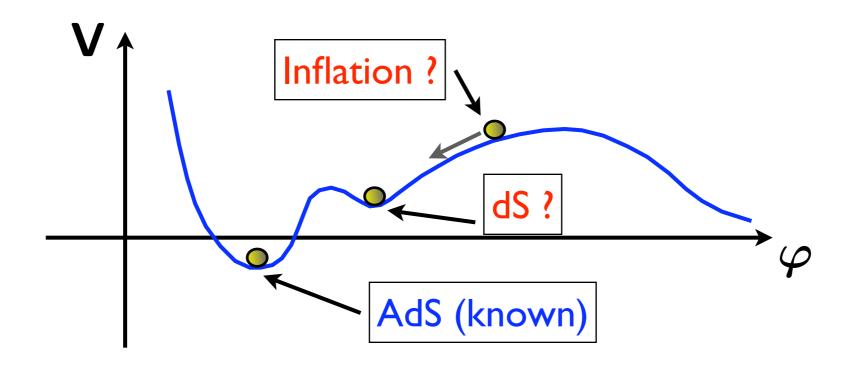
(i) Search for dS/inflation away from AdS vacuum



(ii) Add additional ingredients ⇒ "Uplift potentials"



(i) dS or inflation away from AdS vacuum?



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No-go theorem:

Classical IIA compactifications with

- $\mathcal{M}^{(6)} = \text{Calabi-Yau} (\rightarrow \text{Ricci-flatness})$
- O6/D6 sources
- p-form fluxes (incl. Romans' mass)

⇒No de Sitter vacua and no slow-roll inflation!

Hertzberg, Kachru, Taylor, Tegmark (2007)

Note: Due to the O6-planes (→ negative tension), this goes beyond no-go theorem by Maldacena-Nuñez (2000)

Cf. also Wesley, Steinhardt (2008)

Sketch of proof:

Consider scaling of potential w.r.t.

$$ho \equiv (extsf{VoI})^{1/3}$$
 $au \equiv \mathrm{e}^{-\phi} \sqrt{ extsf{VoI}} \qquad (\phi = extsf{I0D Dilaton})$

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$$V(
ho, au,\ldots) = V_3 + \sum_{p=0,2,4,6} V_p + V_{06/D6}$$
 $\propto
ho^{-3} au^{-2}$
 $\propto
ho^{3-p} au^{-4}$
 $\propto \pm au^{-3}$

$$\Rightarrow$$
 DV $\equiv (-\rho\partial_{\rho} - 3\tau\partial_{\tau})$ V \geq 9V

$$egin{aligned} \mathsf{DV} &\equiv (-
ho\partial_{
ho} - 3 au\partial_{ au})\mathsf{V} \geq 9\mathsf{V} \ &\epsilon &= \mathsf{V}^{-2}\left[rac{(\mathsf{DV})^2}{39} + (\mathsf{positive})
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$$\Rightarrow$$
 $\epsilon \geq \frac{27}{13}$ whenever > 0 \Rightarrow No inflation No de Sitter

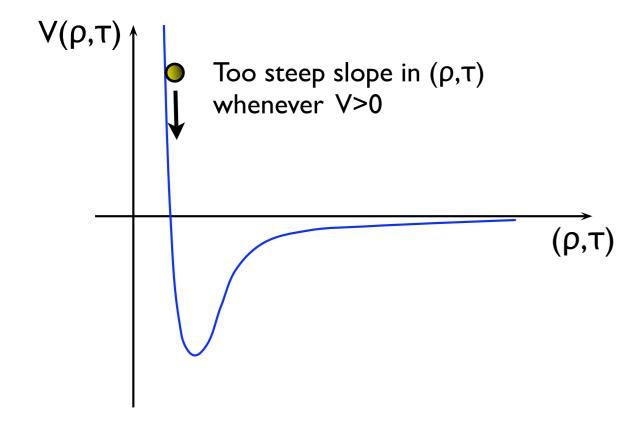


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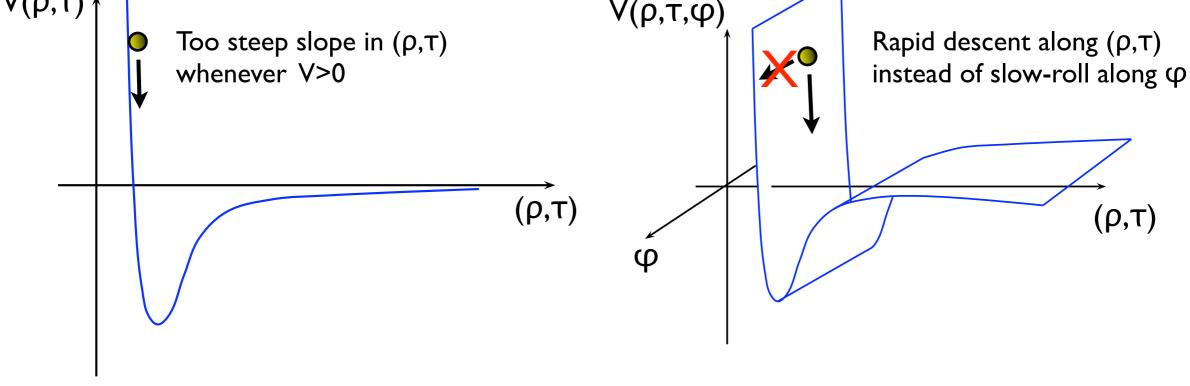






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$$\Rightarrow \boxed{\epsilon \geq \frac{27}{13}} \text{ whenever } \boxed{V > 0} \Rightarrow \boxed{\text{No inflation No de Sitter}}$$



Possible caveats:

Quantum corrections

E.g. Saueressig, Theis, Vandoren (2005); Palti, Tasinato, Ward (2008)

- Additional classical ingredients
 - "Geometric fluxes" ("Torsion") (= geometric twisting away from CY $\Rightarrow R_{mn} \neq 0$)
 - O4-planes
 - D8-branes
 - NS5-branes
 - KK5-monopoles
 - "Nongeometric fluxes"

Silverstein (2007):

For classical de Sitter vacua add, e.g.:

- Geometric fluxes
 (Particular twisted torus)
- KK5-monopoles
- Fractional Chern-Simons invariants

Some issues to keep in mind:

- High SUSY breaking scale
- No large mass gap to KK modes
- Backreaction under control?

What is the minimal controllable setup?

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Best understood extra ingredient:

Geometric fluxes

⇒ Are they sufficient?

Three recent works:

(i) Haque, Shiu, Underwood, Van Riet (2008)

$$\mathcal{M}^{(6)} = (\text{Nil}_3 \times \text{Nil}_3')/\mathcal{O}$$

(ii) Caviezel, Koerber, Körs, Lüst, Wrase, M.Z. (2008)

 $\mathcal{M}^{(6)}$ = Cosets with SU(3)-structure

(iii) Flauger, Paban, Robbins, Wrase (2008)

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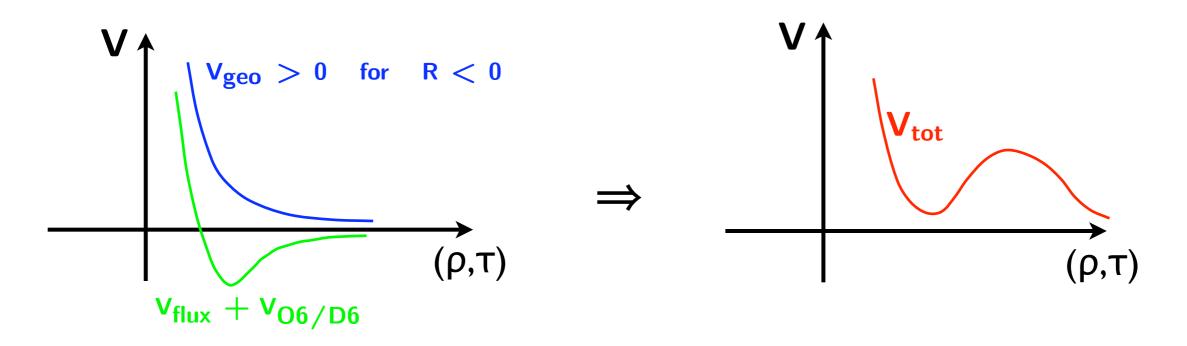
Geometric fluxes

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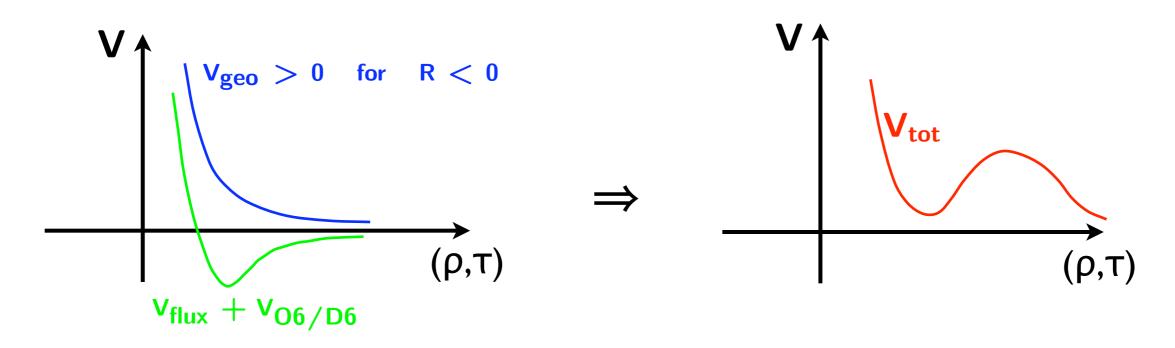


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But: Are all orthogonal field directions also ok?

⇒ Need a setup in which

$$\mathsf{V} = \mathsf{V}(
ho, au, \hspace{-0.1cm} \overline{\hspace{-0.1cm} \hspace{-0.1cm} \hspace{-0.1cm} \hspace{-0.1cm} } \hspace{-0.1cm})$$

is well understood

A well-controlled setup:

Caviezel, Koerber, Körs, Lüst, Tsimpis, M.Z. (2008)

$$\mathcal{M}^{(6)}$$
 = Coset space G/H with (G-invariant) SU(3)-structure (+ orientifolding)

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SU(3)-structure:

- $\Rightarrow \mathcal{M}^{(6)}$ admits globally well-defined spinor η
- \Rightarrow 4D, N=I supergravity action
- \Rightarrow For $\nabla \eta \neq 0$: No CY $\Rightarrow R_{mn} \neq 0 \Rightarrow V_{geo} \neq 0$

Cf. Marchesano's talk

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Early work: Gurrieri, Louis, Micu, Waldram; Dall'Agata, Prezas; Lüst, Tsimpis; Behrndt, Cvetič; Gauntlett, Martelli, Waldram; Graña, Minasian, Petrini, Tomasiello;...

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Problem:

$$\mathsf{J}_{\mathsf{mn}} \equiv \mathsf{i} \eta_+^\dagger \gamma_{\mathsf{mn}} \eta_+ \qquad \Omega_{\mathsf{mnp}} \equiv \eta_-^\dagger \gamma_{\mathsf{mnp}} \eta_+$$

$$\nabla \eta \neq 0 \Rightarrow dJ, d\Omega \neq 0 \Rightarrow Expansion basis, moduli?$$

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Coset space structure:

- ⇒ Natural expansion basis: G-invariant forms
- ⇒ Explicit 4D action (consistent truncation)

Cassani, Kashani-Poor (2009)

Restriction to semisimple and Abelian group factors

$$\frac{G_2}{SU(3)}, \quad \frac{Sp(2)}{S(U(2)\times U(1))}, \quad \frac{SU(3)}{U(1)\times U(1)}, \quad \frac{SU(3)\times U(1)}{SU(2)}, \quad \frac{SU(2)^2}{U(1)}\times U(1)$$

$$SU(2)\times U(1)^3, \quad SU(2)\times SU(2)$$

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$$\mathsf{SU}(2) \times \mathsf{U}(1)^3, \quad \mathsf{SU}(2) \times \mathsf{SU}(2)$$

$$\mathsf{R} < 0 \text{ possible} \Rightarrow \mathsf{Evade old no-go!}$$

⇒ de Sitter or inflation? Or are there new no-go's?

Cf. Flauger, Paban, Robbins, Wrase (2008)

$$V_{geo}=0,$$

$$V_{geo} = 0, \qquad V = V(\tau, \rho, \ldots)$$

$$\Rightarrow \epsilon \geq 27/13$$

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,

Refined no-go:
$$V_{geo} \neq 0$$
, $V = V(\tau, \sigma, \ldots)$

Violates old no-go

Different Kähler modulus

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If:

(i)
$$\kappa_{ijk} = \kappa_{0ab} \Rightarrow$$

$$\sigma \equiv \sqrt{rac{
ho^3}{\mathsf{k}^0}}$$

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$$\kappa_{ijk} = \kappa_{0ab} \Rightarrow \sigma \equiv \sqrt{\frac{\rho^3}{k^0}}$$

$$\sigma \equiv \sqrt{rac{
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(ii)
$$-\sigma\partial_{\sigma}\left|2\tau^{2}\rho^{3}\mathsf{V}_{\mathrm{geo}}\right|\geq0$$

$$\Rightarrow$$

$$\Rightarrow$$
 $\epsilon \geq 2$ for

Our cosets:

- (i) $\kappa_{ijk} = \kappa_{0ab}$ is always satisfied
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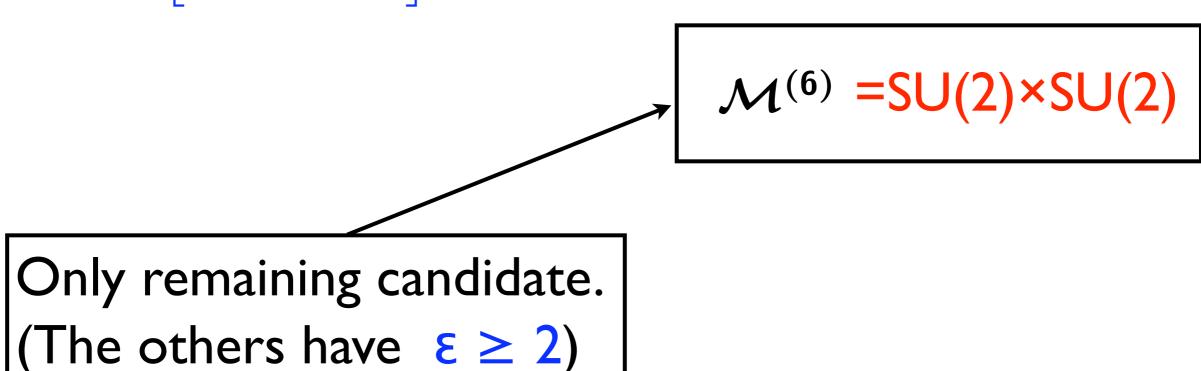
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 is always satisfied except for

$$\mathcal{M}^{(6)} = SU(2) \times SU(2)$$

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 is always satisfied except for



$SU(2)\times SU(2)$

Numerically: $\varepsilon \approx 0$ with V > 0

But:

$$\eta \leq -2.4$$

(Large tachyonic direction)

Interestingly:

- Tachyon is combination of all moduli
- Not the tachyon of

Covi, Gomez-Reino, Gross, Louis, Palma, Scrucca (2008)

Also:

- NS5, D4, D8 can not be added in these models
- No F-term uplifting à la "O'KKLT" possible
 Cf. Kallosh, Linde (2006), Kallosh, Soroush (2006)
- KK5-Monopole ? ⇒ Drastic modification of
 geometry
 See also Villadoro, Zwirner (2007)

⇒ So far nothing really worked...

Conclusions

Type IIA on CY + p-form fluxes + D6/O6:

- Tree-level moduli stabilization in AdS
- No-go against dS and inflation (HKKT)

```
\rightarrow V > 0 is too steep in (\rho,T)
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Type IIA on CY + p-form fluxes + D6/O6:

- Tree-level moduli stabilization in AdS
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$$\rightarrow$$
 V > 0 is too steep in (ρ,T)

- ⇒ Quantum effects or/and additional classical ingredients
 - Best understood: Geometric fluxes (deviation from CY)

Studied cosets with SU(3)-structure

⇒ Refined no-go in
$$(\sigma,\tau)$$
: $\epsilon \geq 2$ except for $SU(2)\times SU(2)$

SU(2)×SU(2):
$$\varepsilon \approx 0$$
, but $\eta \leq -2.4$

Consistent with other works:

Haque, Shiu, Underwood, Van Riet (2008) Flauger, Paban, Robbins, Wrase (2008)

 \Rightarrow Geometric fluxes may help in (ρ,T) -plane, but certainly do not automatically take care of all moduli \Rightarrow Many dangerous directions

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⇒ de Sitter and inflation in IIA?

- More general manifolds? cf. Dall'Agata, Villadoro, Zwirner (2009)
- More general classical ingredients (e.g. Silverstein)?
- Quantum effects (e.g. Saueressig, Theis, Vandoren; Palti, Tasinato, Ward)?

Cf. Talks by Villadoro, de Carlos?