

# Physics of Condensed Matter I

1100-4INZ'PC



Chemical bonding  
and molecules

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# Chemical bonding and molecules

## Born Oppenheimer approximation

Full non-relativistic Hamiltonian of the nuclei and electrons:

$$H(\vec{r}, \vec{R})\Psi(\vec{r}, \vec{R}) = E\Psi(\vec{r}, \vec{R})$$

$$H(\vec{r}, \vec{R}) = \dots$$

$m, \vec{r}_i, i$  – electrons  
 $M_N, \vec{R}_N, Z_N$  - nuclei

Coordinates of electrons subsystem and nuclei subsystem (ions) are mixed, separation of electronic and nuclear variables is impossible.

One should use the Born-Oppenheimer adiabatic approximation

# Chemical bonding and molecules

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$$\begin{aligned} H(\vec{r}, \vec{R}) &= \\ &= -\frac{\hbar^2}{2m} \sum_i \nabla_i^2 - \sum_N \frac{\hbar^2}{2M_N} \nabla_N^2 - \frac{1}{4\pi\epsilon_0} \sum_{N,i} \frac{Z_N e^2}{|\vec{r}_i - \vec{R}_N|} + \\ &+ \frac{1}{4\pi\epsilon_0} \sum_{N < K} \frac{Z_N Z_K e^2}{|\vec{R}_N - \vec{R}_K|} + \frac{1}{4\pi\epsilon_0} \sum_{i < j} \frac{e^2}{|\vec{r}_i - \vec{r}_j|} = \\ &= \hat{T}_e + \hat{T}_N + V(\vec{r}, \vec{R}) + V_e(\vec{r}) + G(\vec{R}) \end{aligned}$$

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## MOLECULAR QUANTUM MECHANICS, FOURTH EDITION

*Peter Atkins  
Ronald Friedman*

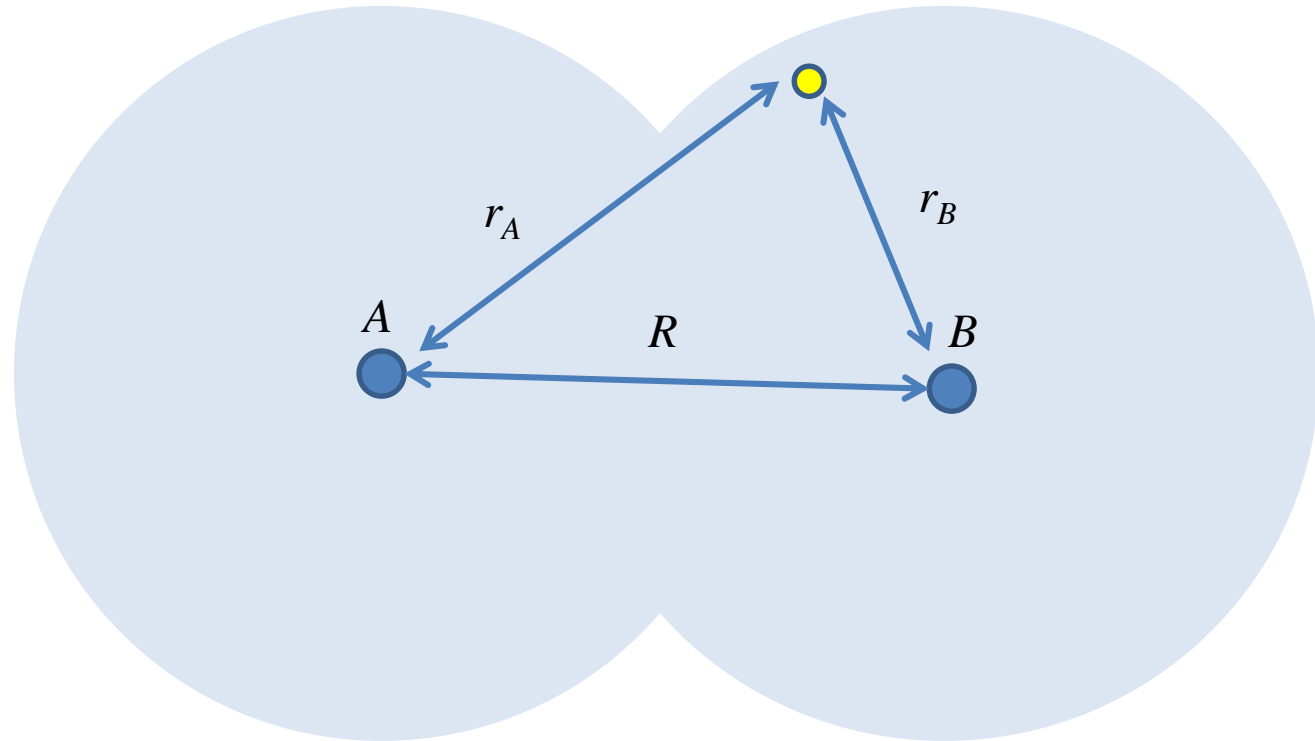
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# Chemical bonding and molecules

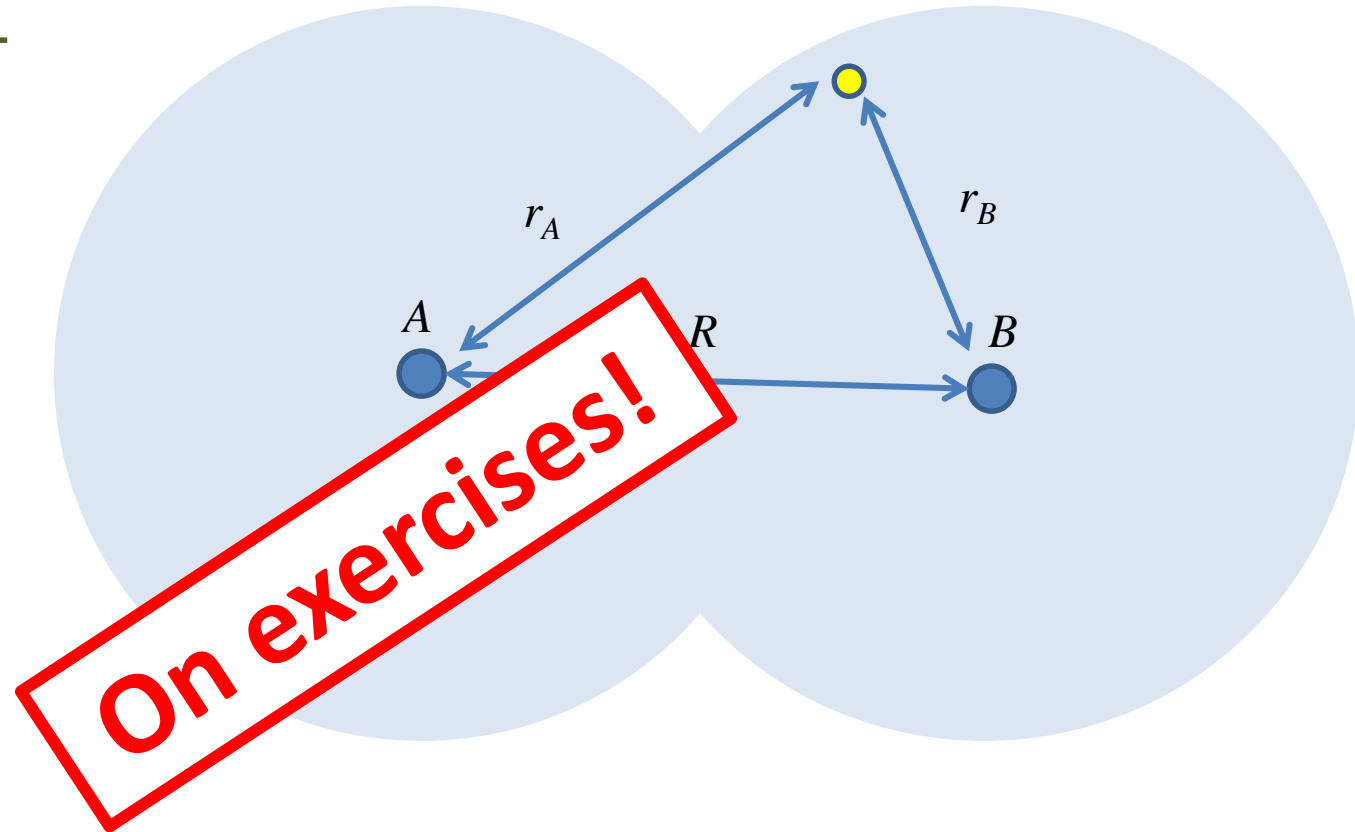
## The molecule $\text{H}_2^+$



$$\hat{H} = -\frac{\hbar^2}{2m_e} \nabla^2 - \frac{\hbar^2}{2M} \nabla_A^2 - \frac{\hbar^2}{2M} \nabla_B^2 - \frac{1}{4\pi\epsilon_0} \frac{e^2}{r_A} - \frac{1}{4\pi\epsilon_0} \frac{e^2}{r_B} + \frac{1}{4\pi\epsilon_0} \frac{e^2}{R}$$

# Chemical bonding and molecules

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# Chemical bonding and molecules

## Born Oppenheimer approximation



Max Born  
(1882-1970)



Jacob R. Oppenheimer  
(1904-1967)

# Chemical bonding and molecules

## Born Oppenheimer approximation

Full non-relativistic Hamiltonian of the nuclei and electrons:

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$m, \vec{r}_i, i$  – electrons  
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Observation: atomic nuclei are tens or even hundreds of thousands heavier than electrons, so the nuclei move more slowly than electrons.



# Chemical bonding and molecules

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## Born Oppenheimer approximation

Atomic nuclei are tens or even hundreds of thousands heavier than electrons, so the nuclei move more slowly than electrons.

**Infinitely heavy nuclei** - omit the kinetic energy of the nuclei

$$H_{el}(\vec{r}, \vec{R}) \Psi_{el}^k(\vec{r}, \vec{R}) = E_{el}^k \Psi_{el}^k(\vec{r}, \vec{R})$$

$$[\hat{T}_e + V(\vec{r}, \vec{R}) + V_e(\vec{r})]$$

$\vec{R}$  – is treated as fixed parameter

$k$  – a set of quantum numbers of a multielectron quantum state

$E_{el}^k(\vec{R})$  – electron energies of different states  $k$  as a function of the positions of the nuclei

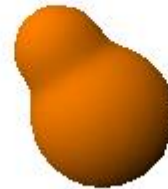
$$T_N = 0$$



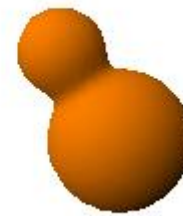
1A



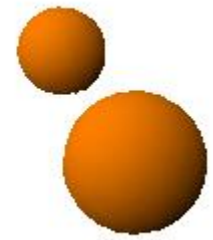
2A



3A



4A



5A

# Chemical bonding and molecules

## Born Oppenheimer approximation

First we solve simplified Hamiltonian for the fixed configuration of atoms (when nuclei do not move). This is so-called **electronic Hamiltonian or clamped nucleus Hamiltonian** (*hamiltonian elektronowy*)

$$H_{el}(\vec{r}, \vec{R})\Psi_{el}^k(\vec{r}, \vec{R}) = E_{el}^k \Psi_{el}^k(\vec{r}, \vec{R})$$

for each instantaneous (*chwilowa*) position of the ion  $\vec{R}$  electrons are in quantum states  $\Psi_{el}^k(\vec{r}, \vec{R})$  corresponding to the global potential of the configuration of ions

$$H_{el}(\vec{r}, \vec{R})\Psi_{el}^k(\vec{r}, \vec{R}) = [\hat{T}_e + V(\vec{r}, \vec{R}) + V_e(\vec{r})]\Psi_{el}^k(\vec{r}, \vec{R}) = E_{el}^k(\vec{R}) \Psi_{el}^k(\vec{r}, \vec{R})$$

Multi-electron wave functions  $\Psi_{el}^k(\vec{r}, \vec{R})$  depend on the positions of all the electrons, and are parameterized by instantaneous positions of all nuclei (ions)  $\vec{R}$ . Index  $k$  represents a set of quantum numbers of a multielectron quantum state. Energies  $E_{el}^k(\vec{R})$  depend on the parameters  $\vec{R}$ .

These functions will be the base of the final solution – they contain electron-electron kinetic and potential energy  $T_e + V(\vec{r}, \vec{R}) + V_e(\vec{r})$

# Chemical bonding and molecules

## Born Oppenheimer approximation

$$H_{el}(\vec{r}, \vec{R})\Psi_{el}^k(\vec{r}, \vec{R}) = [\hat{T}_e + V(\vec{r}, \vec{R}) + V_e(\vec{r})]\Psi_{el}^k(\vec{r}, \vec{R}) = E_{el}^k(\vec{R})\Psi_{el}^k(\vec{r}, \vec{R})$$

Expanding the wavefunction  $\Psi(\vec{r}, \vec{R})$  in basis  $\Psi_{el}^k(\vec{r}, \vec{R})$ :  $\Psi(\vec{r}, \vec{R}) = \sum_k \chi^k(\vec{R})\Psi_{el}^k(\vec{r}, \vec{R})$

$$[\hat{T}_N + \hat{T}_e + V(\vec{r}, \vec{R}) + V_e(\vec{r}) + G(\vec{R})]\Psi(\vec{r}, \vec{R}) = E\Psi(\vec{r}, \vec{R})$$

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$$[\hat{T}_N + H_{el}(\vec{r}, \vec{R}) + \cancel{G(\vec{R})}] \sum_k \chi^k(\vec{R})|\Psi_{el}^k(\vec{r}, \vec{R})\rangle = E(\vec{R}) \sum_k \chi^k(\vec{R})|\Psi_{el}^k(\vec{r}, \vec{R})\rangle \quad / * \langle \Psi_{el}^n(\vec{r}, \vec{R}) |$$

$$\langle \Psi_{el}^n(\vec{r}, \vec{R}) | [\hat{T}_N + H_{el}(\vec{r}, \vec{R})] | \sum_k \chi^k(\vec{R}) |\Psi_{el}^k(\vec{r}, \vec{R})\rangle \rangle = \langle \Psi_{el}^n(\vec{r}, \vec{R}) | E(\vec{R}) | \sum_k \chi^k(\vec{R}) |\Psi_{el}^k(\vec{r}, \vec{R})\rangle \rangle$$

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$$[\hat{T}_N + E_{el}(\vec{r}, \vec{R})]\chi^n(\vec{R}) + \sum_k \chi^k(\vec{R}) \langle \Psi_{el}^n(\vec{r}, \vec{R}) | \hat{T}_N | \Psi_{el}^k(\vec{r}, \vec{R}) \rangle = E(\vec{R})\chi^n(\vec{R})$$

Electronic states are mixed!

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$$[\hat{T}_N + E_{el}(\vec{r}, \vec{R})]\chi^n(\vec{R}) + \sum_k \chi^k(\vec{R}) \langle \Psi_{el}^n(\vec{r}, \vec{R}) | \hat{T}_N | \Psi_{el}^k(\vec{r}, \vec{R}) \rangle = E(\vec{R})\chi^n(\vec{R})$$

$$[\hat{T}_N + E_{el}(\vec{r}, \vec{R})]\chi^n(\vec{R}) + \sum_{k \neq n} \chi^k(\vec{R}) \langle \Psi_{el}^n(\vec{r}, \vec{R}) | \hat{T}_N | \Psi_{el}^k(\vec{r}, \vec{R}) \rangle + \chi^n(\vec{R}) \langle \Psi_{el}^n(\vec{r}, \vec{R}) | \hat{T}_N | \Psi_{el}^n(\vec{r}, \vec{R}) \rangle = E(\vec{R})\chi^n(\vec{R})$$

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### Approximation 1: Adiabatic

The movement of electrons is so fast that during a small change of the nuclei position electrons immediately adapt to the new conditions. This means that the electrons do not change their state  $\Psi_{el}^k$  upon movement of the nuclei. Mathematically:  $\hat{T}_N$  operator acting on  $\Psi_{el}^k$  does not change it into different function  $\Psi_{el}^k, n \neq k$



# Chemical bonding and molecules

## Born Oppenheimer approximation

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$$\langle \Psi_{el}^n(\vec{r}, \vec{R}) | [\hat{T}_N + H_{el}(\vec{r}, \vec{R})] | \sum_k \chi^k(\vec{R})\Psi_{el}^k(\vec{r}, \vec{R}) \rangle = \langle \Psi_{el}^n(\vec{r}, \vec{R}) | E(\vec{R}) | \sum_k \chi^k(\vec{R})\Psi_{el}^k(\vec{r}, \vec{R}) \rangle$$

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### Approximation 1: Adiabatic

Thus:

$$\Psi(\vec{r}, \vec{R}) = \sum_k \chi^k(\vec{R})\Psi_{el}^k(\vec{r}, \vec{R}) \approx \chi^n(\vec{R})\Psi_{el}^n(\vec{r}, \vec{R})$$

# Chemical bonding and molecules

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Expanding the wavefunction  $\Psi(\vec{r}, \vec{R})$  in basis  $\Psi_{el}^k(\vec{r}, \vec{R})$ :  $\Psi(\vec{r}, \vec{R}) = \sum_k \chi^k(\vec{R})\Psi_{el}^k(\vec{r}, \vec{R})$

$$\langle \Psi_{el}^n(\vec{r}, \vec{R}) | [\hat{T}_N + H_{el}(\vec{r}, \vec{R})] | \sum_k \chi^k(\vec{R}) \Psi_{el}^k(\vec{r}, \vec{R}) \rangle = \langle \Psi_{el}^n(\vec{r}, \vec{R}) | E(\vec{R}) | \sum_k \chi^k(\vec{R}) \Psi_{el}^k(\vec{r}, \vec{R}) \rangle$$

$$[\hat{T}_N + H_{el}(\vec{r}, \vec{R})]\chi^n(\vec{R}) + \sum_k \chi^k(\vec{R}) \langle \Psi_{el}^n(\vec{r}, \vec{R}) | \hat{T}_N | \Psi_{el}^k(\vec{r}, \vec{R}) \rangle = E(\vec{R})\chi^n(\vec{R})$$

$$[\hat{T}_N + E_{el}(\vec{r}, \vec{R})]\chi^n(\vec{R}) + \sum_{k \neq n} \chi^k(\vec{R}) \langle \Psi_{el}^n(\vec{r}, \vec{R}) | \hat{T}_N | \Psi_{el}^k(\vec{r}, \vec{R}) \rangle + \chi^n(\vec{R}) \langle \Psi_{el}^n(\vec{r}, \vec{R}) | \hat{T}_N | \Psi_{el}^n(\vec{r}, \vec{R}) \rangle = E(\vec{R})\chi^n(\vec{R})$$

Approximation 2:

changing position of the nuclei weakly affects the state of electrons.

Mathematically:  $\hat{T}_N$  operator acting on  $\Psi_{el}^k$  gives  $\hat{T}_N |\Psi_{el}^n(\vec{r}, \vec{R})\rangle \approx 0$

# Chemical bonding and molecules


## Born Oppenheimer approximation

$$H_{el}(\vec{r}, \vec{R})\Psi_{el}^k(\vec{r}, \vec{R}) = [\hat{T}_e + V(\vec{r}, \vec{R}) + V_e(\vec{r})]\Psi_{el}^k(\vec{r}, \vec{R}) = E_{el}^k(\vec{R}) \Psi_{el}^k(\vec{r}, \vec{R})$$

Expanding the wavefunction  $\Psi(\vec{r}, \vec{R})$  in basis  $\Psi_{el}^k(\vec{r}, \vec{R})$ : 
$$\Psi(\vec{r}, \vec{R}) = \sum_k \chi^k(\vec{R})\Psi_{el}^k(\vec{r}, \vec{R})$$

$$\Psi(\vec{r}, \vec{R}) \approx \chi^n(\vec{R})\Psi_{el}^n(\vec{r}, \vec{R})$$

$$[\hat{T}_N + E_{el}(\vec{r}, \vec{R})]\chi^n(\vec{R}) + \sum_{k \neq n} \chi^k(\vec{R})\langle \Psi_{el}^n(\vec{r}, \vec{R}) | \hat{T}_N | \Psi_{el}^k(\vec{r}, \vec{R}) \rangle + \chi^n(\vec{R})\langle \Psi_{el}^n(\vec{r}, \vec{R}) | \hat{T}_N | \Psi_{el}^n(\vec{r}, \vec{R}) \rangle = E(\vec{R})\chi^n(\vec{R})$$


$$E(\vec{R}) \rightarrow E^n$$

$$[\hat{T}_N + E_{el}(\vec{r}, \vec{R}) + G(\vec{R})]\chi^n(\vec{R}) = E^n\chi^n(\vec{R})$$

effective potential

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$$\Psi(\vec{r}, \vec{R}) = \sum_k \chi^k(\vec{R})\Psi_{el}^k(\vec{r}, \vec{R})$$

$$\Psi(\vec{r}, \vec{R}) \approx \chi^n(\vec{R})\Psi_{el}^n(\vec{r}, \vec{R})$$

$$[\hat{T}_N + E_{el}(\vec{r}, \vec{R})]\chi^n(\vec{R}) + \sum_{k \neq n} \chi^k(\vec{R})\langle \Psi_{el}^n(\vec{r}, \vec{R}) | \hat{T}_N | \Psi_{el}^k(\vec{r}, \vec{R}) \rangle + \chi^n(\vec{R})\langle \Psi_{el}^n(\vec{r}, \vec{R}) | \hat{T}_N | \Psi_{el}^n(\vec{r}, \vec{R}) \rangle = E(\vec{R})\chi^n(\vec{R})$$

Does it remind you of something?

$$[\hat{T}_N + E_{el}(\vec{r}, \vec{R}) + G(\vec{R})]\chi^n(\vec{R}) = E^n\chi^n(\vec{R})$$

effective potential

# Chemical bonding and molecules

## Born Oppenheimer approximation

$$H_{el}(\vec{r}, \vec{R})\Psi_{el}^k(\vec{r}, \vec{R}) = [\hat{T}_e + V(\vec{r}, \vec{R}) + V_e(\vec{r})]\Psi_{el}^k(\vec{r}, \vec{R}) = E_{el}^k(\vec{R}) \Psi_{el}^k(\vec{r}, \vec{R})$$

Expanding the wavefunction  $\Psi(\vec{r}, \vec{R})$  in basis  $\Psi_{el}^k(\vec{r}, \vec{R})$ : 
$$\Psi(\vec{r}, \vec{R}) = \sum_k \chi^k(\vec{R})\Psi_{el}^k(\vec{r}, \vec{R})$$

$$\Psi(\vec{r}, \vec{R}) \approx \chi^n(\vec{R})\Psi_{el}^n(\vec{r}, \vec{R})$$

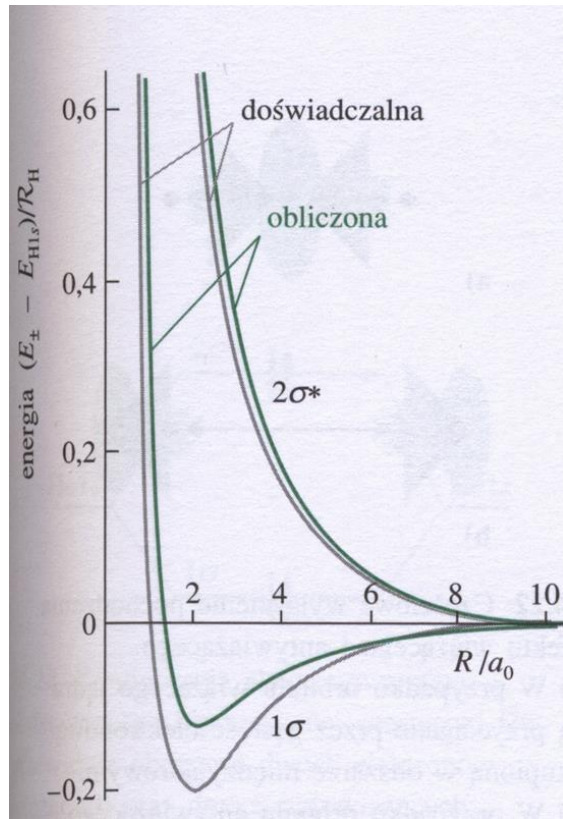
$$[\hat{T}_N + E_{el}(\vec{r}, \vec{R})]\chi^n(\vec{R}) + \sum_{k \neq n} \chi^k(\vec{R})\langle \Psi_{el}^n(\vec{r}, \vec{R}) | \hat{T}_N | \Psi_{el}^k(\vec{r}, \vec{R}) \rangle + \chi^n(\vec{R})\langle \Psi_{el}^n(\vec{r}, \vec{R}) | \hat{T}_N | \Psi_{el}^n(\vec{r}, \vec{R}) \rangle = E(\vec{R})\chi^n(\vec{R})$$

Schrodinger equation of the motion of nuclei with repulsive potential  $G(\vec{R})$ :

$$[\hat{T}_N + E_{el}(\vec{r}, \vec{R}) + G(\vec{R})]\chi^n(\vec{R}) = E^n\chi^n(\vec{R})$$

effective potential

# Chemical bonding and molecules



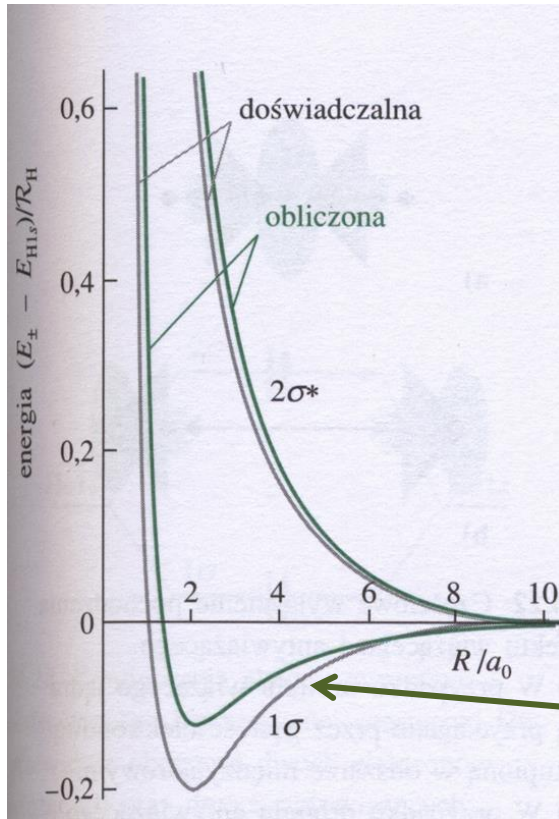
$\chi^n(\vec{R})$  is the wave function describing the motion of nuclei (ions) in their mutual interaction potential  $G(\vec{R})$  adiabatic electron contribution to the energy of the motion of nuclei (ions)  $E_{el}^k(\vec{R})$

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effective potential

# Chemical bonding and molecules



$\chi^n(\vec{R})$  is the wave function describing the motion of nuclei (ions) in their mutual interaction potential  $G(\vec{R})$  adiabatic electron contribution to the energy of the motion of nuclei (ions)  $E_{el}^k(\vec{R})$

Born-Oppenheimer approximation **is not fulfilled** when the potential energy surfaces of two electronic states are too close.

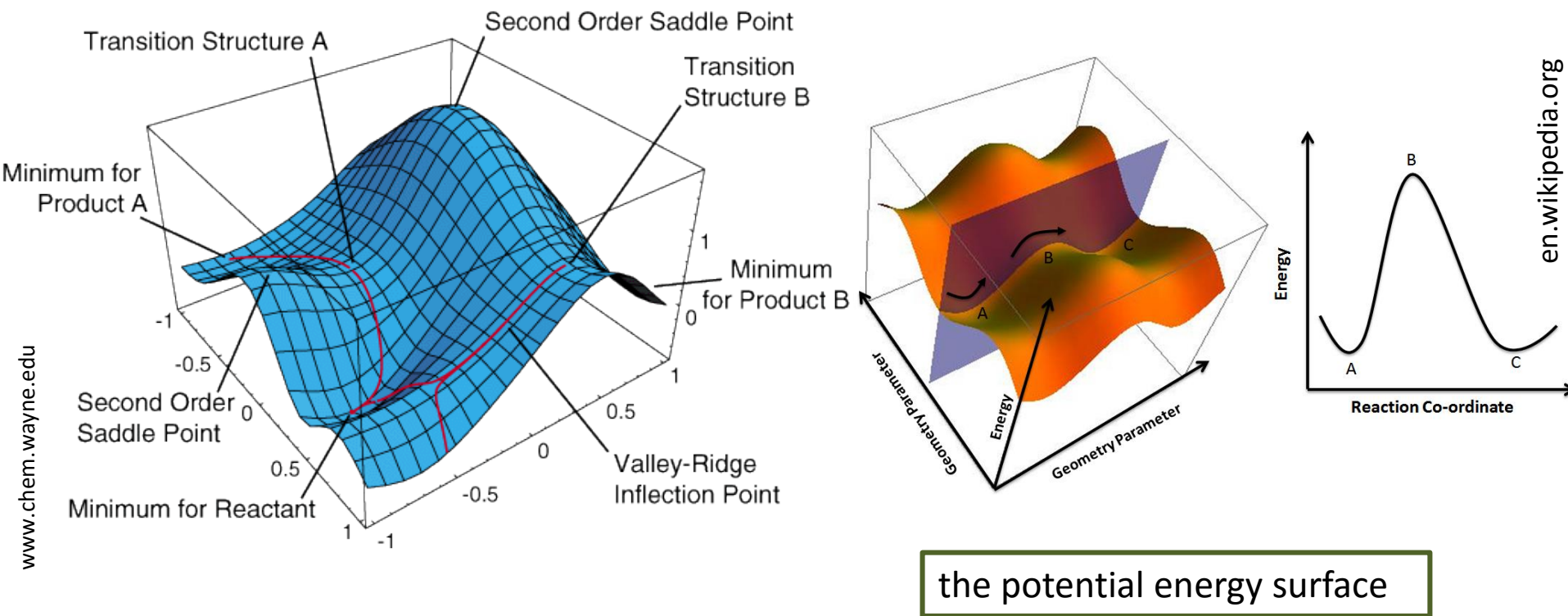
the potential energy surface

Schrodinger equation of the motion of nuclei with repulsive potential  $G(\vec{R})$ :

$$[\hat{T}_N + E_{el}(\vec{r}, \vec{R}) + G(\vec{R})]\chi^n(\vec{R}) = E^n \chi^n(\vec{R})$$

effective potential

# Chemical bonding and molecules



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effective potential



# Chemical bonding and molecules

## Approximations

$$[\hat{T}_N + E_{el}(\vec{r}, \vec{R}) + G(\vec{R})]\chi^n(\vec{R}) = E^n \chi^n(\vec{R})$$

The kinetic energy separates on vibration (oscillation) and rotation energy – we assume "small" oscillations and slow speed of rotation.

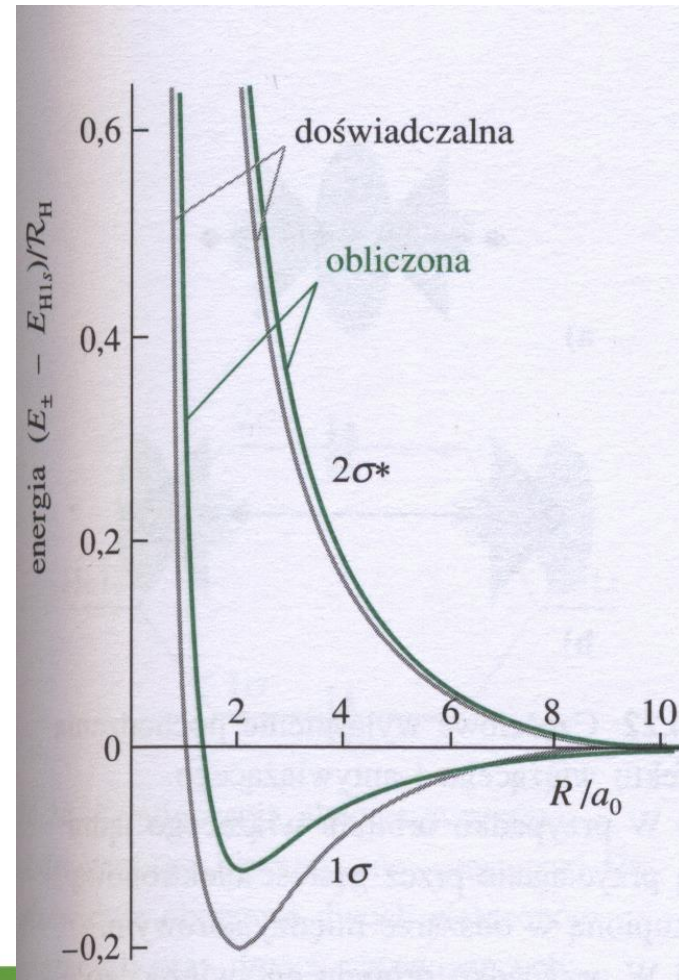
$$[\hat{T}_{osc} + \hat{T}_{rot} + E_{el}(\vec{r}, \vec{R}) + G(\vec{R})]\chi^n(\vec{R}) = E^n \chi^n(\vec{R})$$

Operators act on different coordinates: we can separate the variables:

$$\begin{aligned}\chi^n(\vec{R}) &= \chi_{osc}^n(R) \chi_{rot}^n(\theta, \varphi) \\ E^n &= E_{osc}^n + E_{rot}^n\end{aligned}$$

Altogether:

$$\begin{aligned}\Psi(\vec{r}, \vec{R}) &= \chi^n(\vec{R}) \Psi_{el}^n(\vec{r}, \vec{R}) = \chi_{osc}^n(R) \chi_{rot}^n(\theta, \varphi) \Psi_{el}^n(\vec{r}, \vec{R}) \\ E^n &= E_{osc}^n + E_{rot}^n + E_{el}\end{aligned}$$



# Chemical bonding and molecules

## Approximations

$$[\hat{T}_N + E_{el}(\vec{r}, \vec{R}) + G(\vec{R})]\chi^n(\vec{R}) = E^n \chi^n(\vec{R})$$

The kinetic energy separates on vibration (oscillation) and rotation energy – we assume "small" oscillations and slow speed of rotation.

$$[\hat{T}_{osc} + \hat{T}_{rot} + E_{el}(\vec{r}, \vec{R}) + G(\vec{R})]\chi^n(\vec{R}) = E^n \chi^n(\vec{R})$$

Operators act on different coordinates: we can separate the variables:

$$\chi^n(\vec{R}) = \chi_{osc}^n(R) \chi_{rot}^n(\theta, \varphi)$$
$$E^n = E_{osc}^n + E_{rot}^n$$

- One by one we will discuss:
- electronic structure
  - rotations
  - oscillations

Altogether:

$$\Psi(\vec{r}, \vec{R}) = \chi^n(\vec{R}) \Psi_{el}^n(\vec{r}, \vec{R}) = \chi_{osc}^n(R) \chi_{rot}^n(\theta, \varphi) \Psi_{el}^n(\vec{r}, \vec{R})$$
$$E^n = E_{osc}^n + E_{rot}^n + E_{el}$$

