

Massless 4D Gravitons from Asymptotically AdS_5 Spacetimes

Francesco Nitti

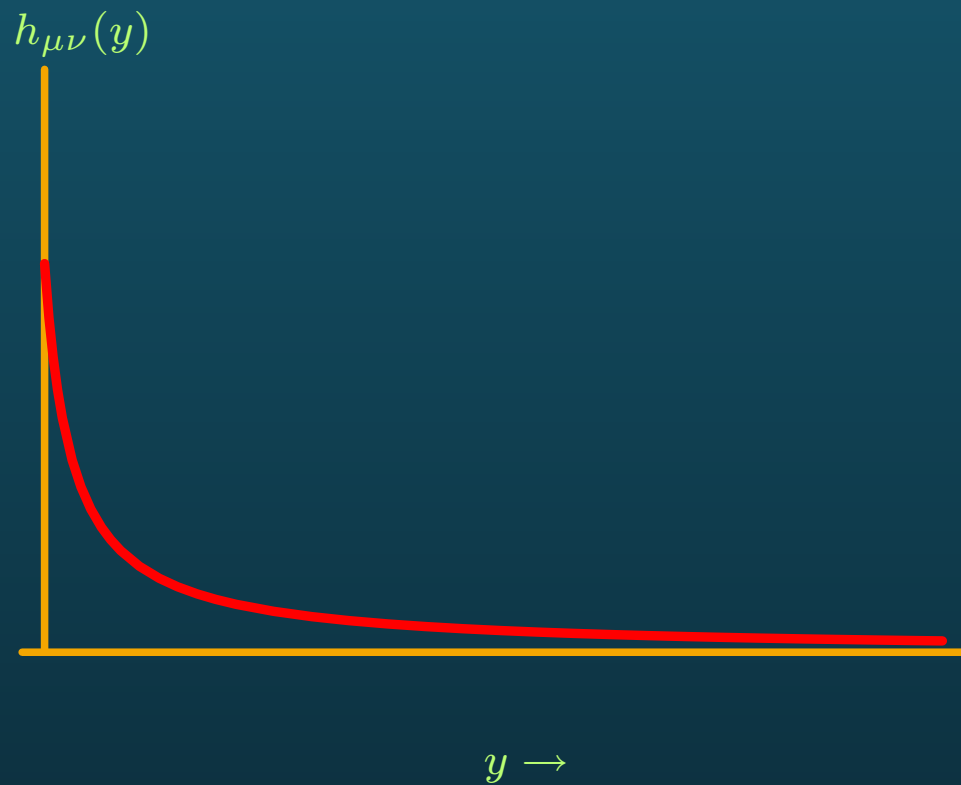
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Based on [hep-th/0611344](#) with E. Kiritsis

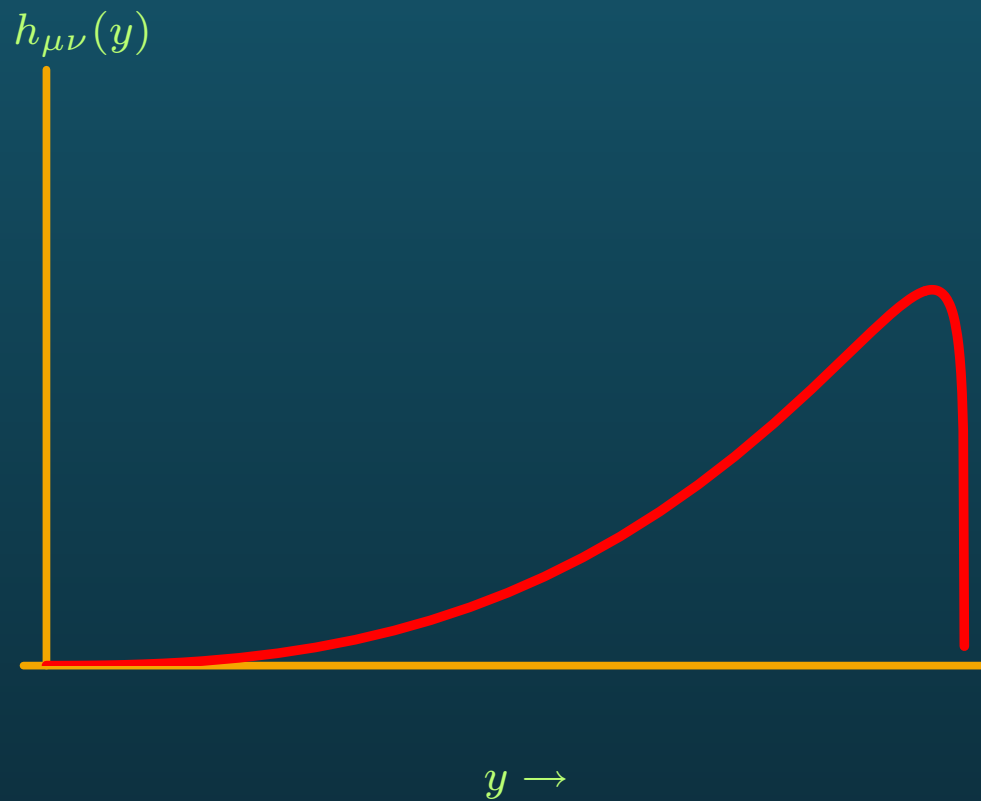
Prologue

RS Model:



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Are there models that exhibit:



Motivation - Holographic Correspondence

Holographic correspondence: Maldacena, '97

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5D theory **with gravity** in Asymptotically AdS_5 spacetime (times some compact factor)

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4D field theory **without gravity**.

- Excitations near the boundary of AdS_5 \Leftrightarrow high energy modes in the FT.
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4D field theory **without gravity**.

- Excitations near the boundary of $AdS_5 \Leftrightarrow$ high energy modes in the FT.
- Radial Evolution away from the boundary \Leftrightarrow RG flow to the IR
- Spectrum of 4D field theory particles = spectrum of **normalizable fluctuations** around the dual 5D geometry.

Witten, '97

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- the graviton wave-function is peaked at the UV cut-off
 \Leftrightarrow graviton is a **fundamental degree of freedom** coupled to the 4D FT at the cut-off

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Can we have 4D graviton localized far from the boundary? (this would be emergent, rather than fundamental, in the dual FT).

Outline

- RS-like Models

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- Einstein-Dilaton Models

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 - Linear Fluctuation Analysis
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- Conclusion and Perspectives

4D Spectrum from 5D

Take asymptotically AdS_5 solution of 5D theory,

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5D fluctuations such that:

- they have a fixed 4D mass: $\square_4 \Phi(x, y) = m^2 \Phi(x, y)$
- are **normalizable** w.r.t. to the radial direction y , i.e. they have a **finite 4D kinetic term**.

correspond to 4D states with mass m^2 in the dual FT.

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We are interested in 4D-massless, y -normalizable fluctuations of the (tensor part of) the 5D metric component, $h_{\mu\nu}(x, y)$

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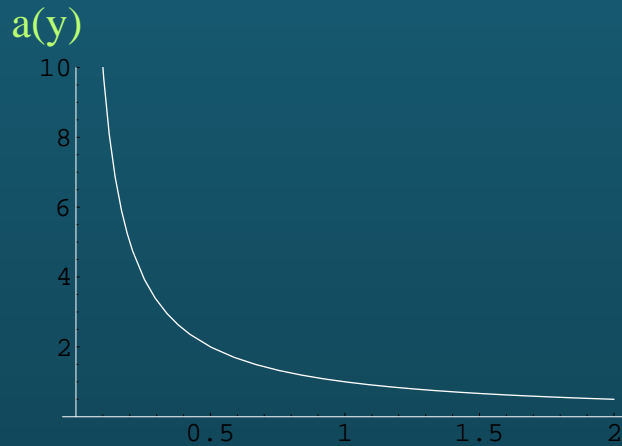
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For tensor spin-2:

$$h_{\mu\nu}(x, y) \sim h_{\mu\nu}^{(0)}(x) + y^4 h_{\mu\nu}^{(4)}(x) + \dots \quad y \rightarrow 0$$

AdS_5 vs. RSII



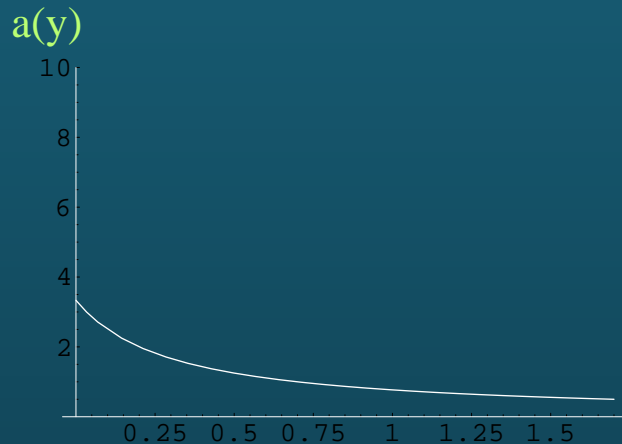
AdS_5 :

$$ds^2 = \frac{1}{(ky)^2} (dy^2 + dx_\mu^2) \quad 0 < y < \infty$$

$$S_{kin}[h^{(0)}] = \int_0^\infty dy \frac{1}{(ky)^3} \int d^4x (\partial h^{(0)})^2 = \infty$$

- In AdS_5 $h_{\mu\nu}^{(0)}$ is **not normalizable** \Rightarrow **not** a state in the 4D FT, rather an external source added to the UV theory.

AdS_5 vs. RSII



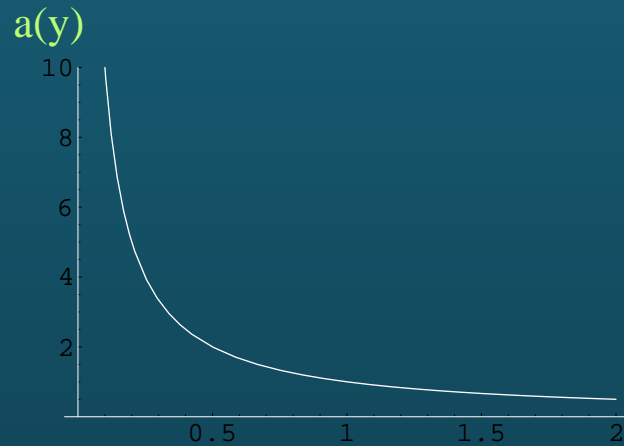
Slice of AdS_5 :

$$ds^2 = \frac{1}{(1 + ky)^2} (dy^2 + dx_\mu^2) \quad 0 < y < \infty$$

$$S_{kin}[h^{(0)}] = \int_0^\infty dy \frac{1}{(1 + ky)^3} \int d^4x \left(\partial h^{(0)} \right)^2 < \infty$$

- In **RSII** it becomes **normalizable** \Rightarrow the source gets a kinetic term and becomes dynamical. it is promoted to a fundamental d.o.f of the UV theory.

AdS_5 vs. RSII



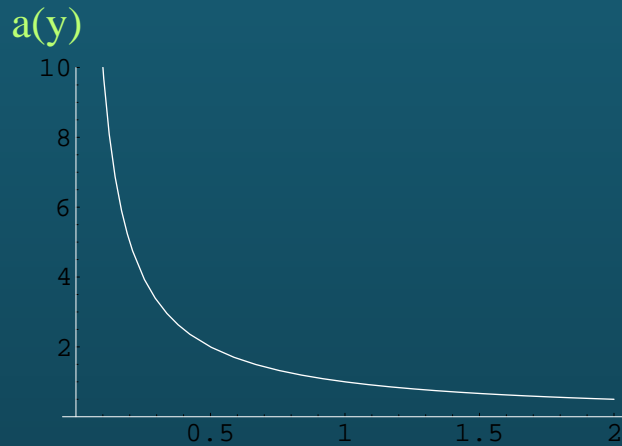
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- A normalizable $h_{\mu\nu}^{(4)}(x)$ would correspond to a low-energy excitation (“glueball”), since it has support in the IR (large y) region.

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- A normalizable $h_{\mu\nu}^{(4)}(x)$ would correspond to a low-energy excitation (“glueball”), since it has support in the IR (large y) region. Can we realize a set-up in which the only normalizable massless tensor mode is $h_{\mu\nu}^{(4)}(x)$?

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Gherghetta, Peloso, Poppitz, '05

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- theory does not “run”: it is conformal all the way to the boundary
- does 5D massive gravity have a holographic interpretation anyway?

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Take solution **asymptotically** AdS_5 in the UV ($y \rightarrow 0$):

$$a(y) \sim \frac{1}{ky}; \quad \Phi_0(y) \sim \text{const}; \quad V(\Phi_0(y)) \sim 2\Lambda$$

Spin-2 Fluctuations

$$ds^2 = a^2(y) [dy^2 + (\eta_{\mu\nu} + h_{\mu\nu}(x, y)) dx^\mu dx^\nu] \quad h_{\mu}^{\mu} = \partial^\mu h_{\mu\nu} = 0$$

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Look for **normalizable** solution with given 4D mass m :

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Get equation for the profile $h(y)$:

$$h''(y) + 3\frac{a'}{a}h'(y) + m^2 h(y) = 0$$

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What kinds of $B(y)$ provide a zero-energy state?

Background

$B(y)$ can be chosen (almost) arbitrarily:

$$ds^2 = dr^2 + e^{-4B(r)/3} \eta_{\mu\nu} dx^\mu dx^\nu, \quad (B = -3/2 \log a)$$

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$$\Rightarrow \quad \psi^{UV}(y) \sim y^{-3/2}, \quad \psi^{IR}(y) \sim y^{5/2}.$$

ψ^{IR} is normalizable, $\psi^{UV}(y)$ is not.

(Notice: **both** are normalizable in RS, where $y > 1/\Lambda$)

IR Asymptotics

We have one candidate Zero-Mode, normalizable around $y = 0$:

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Two distinct cases:

- y range extends to $+\infty \Rightarrow \Psi_{IR} \rightarrow \infty$ as $y \rightarrow \infty$, **not normalizable**
- spacetime ends at $y = y_0$ (singularity)

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example with a singularity:

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- close to y_0 :

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Singular Case

We have one candidate Zero-Mode:

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example with a singularity:

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- These arise only if the **5th dimension terminates**, and only if suitable b.c. are imposed
- We found cases with **no other scalar or vector** massless degrees of freedom. This is an advantage over previous attempts.
- Our analysis indicates how one can relax the requirement of an exactly massless, strictly 4D state, to try to overcome the problems with the singularity and/or the boundary conditions in the IR.

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Concrete Example:

$$B(y) = \frac{3}{2} \log ky - \alpha \log(1 - y/y_0), \quad 0 < \alpha < 3/2$$

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\Rightarrow this is the boundary conditions we need to impose on the fluctuations to keep zero-mode is in the spectrum

Logarithmic $B(y)$ - 4D Planck Scale

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$$S = \frac{1}{2k_5^2} \int dy \frac{a^3(y)}{a^2(y_b)} (\partial_\rho h_{\mu\nu}(y))^2 + \int_{y=y_b} h_{\mu\nu}(y_b) T^{\mu\nu}$$

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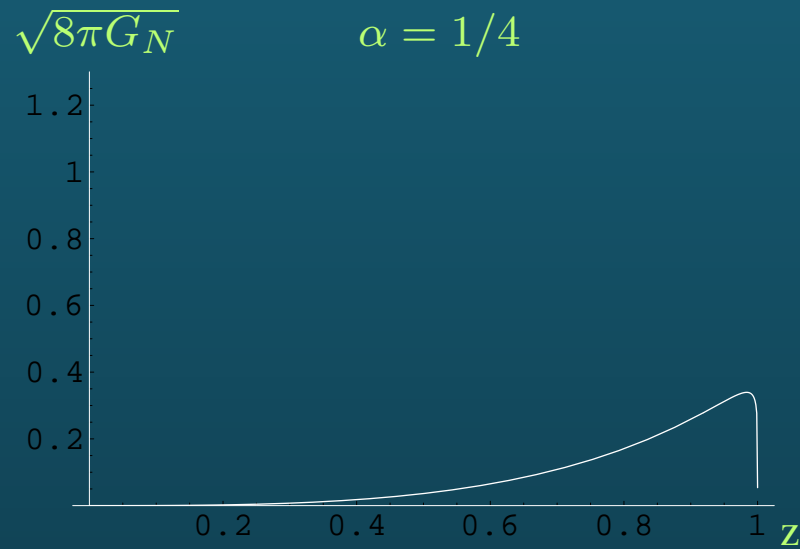
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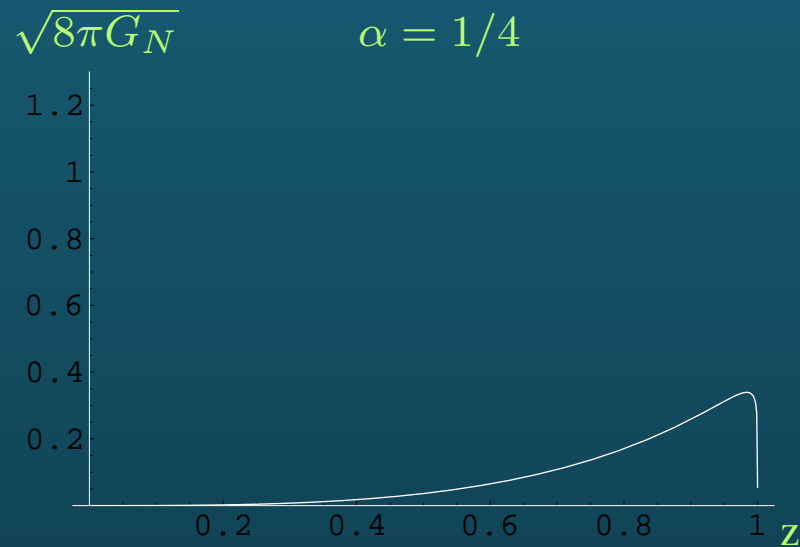
$$z_b \equiv y_b/y_0$$

Logarithmic $B(y)$ - 4D Planck Scale II



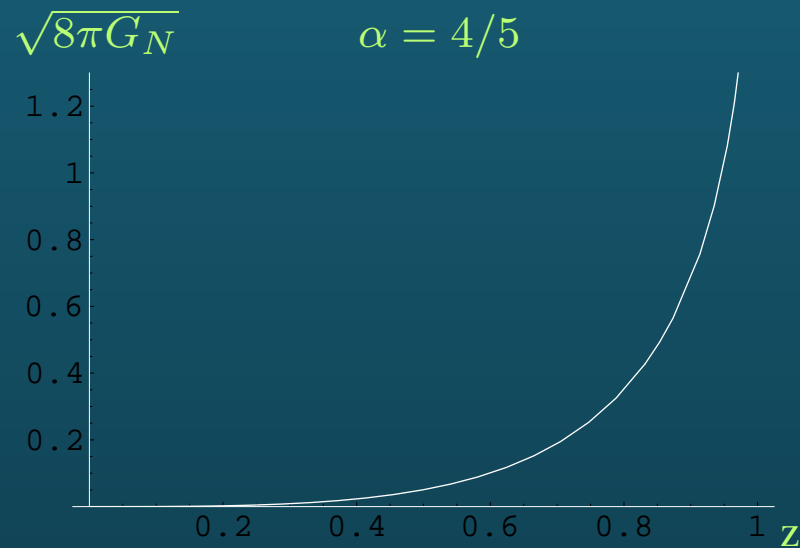
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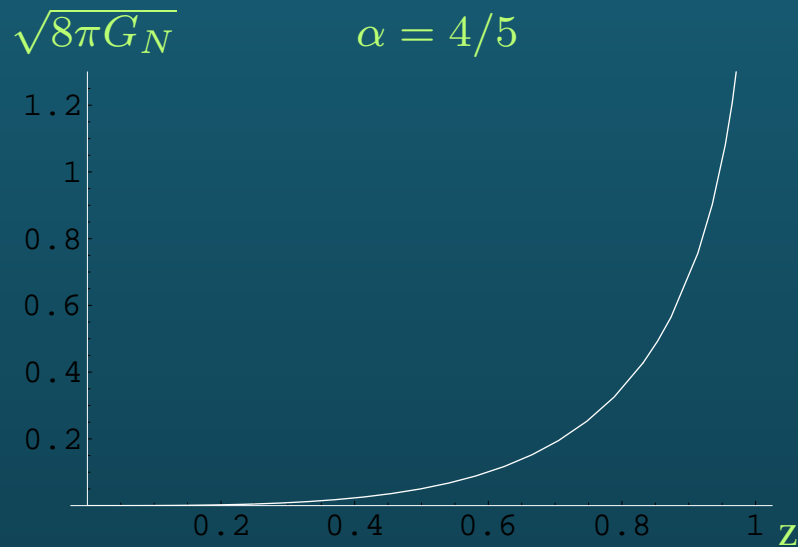
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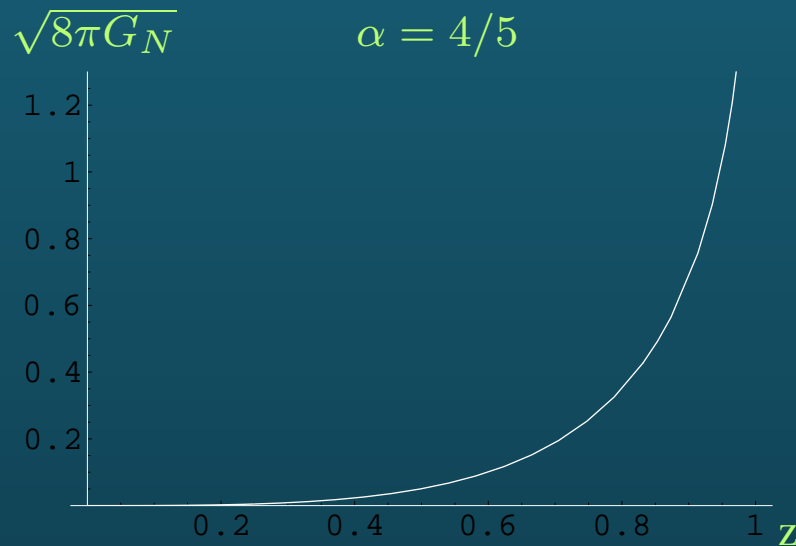
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KK scale masses: $m_{kk}^2 \sim 1/y_0^2$