

# *SUSY without the Little Hierarchy*

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based on a collaboration with

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Feb 6th 2009

# MSSM AND THE LITTLE HIERARCHY

- The mass parameters in the Higgs potential are determined by soft SUSY-breaking parameters  $m_s$

$$m_Z^2 = \sum_i \alpha_i m_{s_i}^2$$

- LEP negative search pushes  $m_s$  up
- A very efficient cancellation ?

## MSSM

- $m_h^2|_{\text{tree-level}} < m_Z^2 \Rightarrow$  Large radiative corrections

$$\delta\lambda \simeq (3/4\pi^2)m_t^4/v^2 \log(m_s^2/m_t^2) \Rightarrow m_h^2|_{\text{tree-level}} + 2v^2\delta\lambda$$
$$\delta m_{Hu}^2 = -(3y_t^2/8\pi^2)m_s^2 \ln(\Lambda^2/m_s^2)$$

- $m_h^2 > 115 \text{ GeV} \Rightarrow m_s \gtrsim 1 \text{ TeV} \Rightarrow \Delta = \frac{|\delta m_h^2|}{m_h^2/2} \rightarrow FT < 1\%$

# POSSIBLE SOLUTIONS

$$\delta m_{Hu}^2 = -(3y_t^2/8\pi^2)m_s^2 \log(\Lambda^2/m_s^2)$$

- ❶ Lower the Log
  - 1A Lower  $\Lambda$  (e.g. NMSSM,  $\kappa = 2 \Rightarrow \Lambda = 20$  TeV)
  - 1B Replace  $\Lambda$  with an effective lower scale  $f \sim 1$  TeV (e.g. Little Higgs)
- ❷ Raise up the tree-level Higgs mass,  $\Rightarrow m_s^2 (\downarrow)$ 
  - 2A From F-terms (e.g NMSSM)
  - 2B From D-terms (i.e. extending the gauge structure)
- ❸ 1B+2B= this model: a **Super Little Higgs** with enhanced quartic

# THE MODEL: PRELIMINARY

Why SUSY+LH?

- $SM \longrightarrow \delta m_H^2 \sim \Lambda^2$
- $SUSY \longrightarrow \delta m_H^2 \sim m_s^2 \log \Lambda$
- $LH \longrightarrow \delta m_H^2 \sim f^2 \log \Lambda$

## SUSY+LH

The marriage between SUSY and LH enforce a “double protection” from SUSY and a global symmetry under which the Higgs is a *PseudoGoldstone* boson

$$SUSY + LH \longrightarrow \delta m_H^2 \sim 3m_s^2/8\pi^2 \log f$$

*Bereziani, Chankowky, Falkowski, Pokorski, Wagner '04, '05*

*Csaki, Strumia, Marandella, Shirman '05*

*Birkedal, Chacko, Gaillard '04*

*Roy, Schmaltz '06*

# THE MODEL: $SU(3)_W \times U(1)_X$

The easiest way to embed the LH in a SUSY theory is via the “simplest little Higgs”

$$SU(2)_W \times U(1)_Y \rightarrow SU(3)_W \times U(1)_X$$

The Higgs doublets are promoted to triplets

$$H_{u,d} \rightarrow \mathcal{H}_{u,d} = (H_{u,d}, S_{u,d}) = 3, \bar{3}$$

and cloned in exact copies  $\Phi_{u,d} = 3, \bar{3}$

$$\mathcal{W} = \mathcal{W}_\Phi + \mathcal{W}_\mathcal{H}$$

As far as  $\mathcal{H}$  and  $\Phi$  don't talk each other the Higgs sector has a  $SU(3)_1 \times SU(3)_2 \supset SU(3)_W$

# THE MODEL: DOWN TO THE SM

$$\begin{aligned}\langle \Phi_{u,d} \rangle &= (0, 0, F \sim 10 \text{ TeV}) \Rightarrow SU(3)_1 \times SU(3)_2 \rightarrow SU(2)_1 \times SU(3)_2 \\ \langle |\mathcal{H}_{u,d}| \rangle &= f_{u,d} \sim 1 \text{ TeV} \Rightarrow SU(2)_1 \times SU(3)_2 \rightarrow SU(2)_1 \times SU(2)_2 \\ &16 - 3 - 3 - \textcolor{red}{5} = 5 = \textcolor{red}{4} + \textcolor{red}{1}\end{aligned}$$

These 4 degrees form a doublet under  $SU(2)_W$  and indeed are the SM Higgs doublet  $\textcolor{red}{H}$

It controls the direction of  $SU(2)_2$  with respect to  $SU(2)_W$

$$\mathcal{H}_u = (H^T, \sqrt{f^2 - |H|^2}) \sin \beta \quad \mathcal{H}_d = (H, \sqrt{f^2 - |H|^2}) \cos \beta$$

If  $v = 0$  there is no EWSB  $\Rightarrow SU(2)_W \times U(1)_Y$  is unbroken

# HIGGS AS A PSEUDO GOLDSTONE BOSON

$H$  is a **pseudoGoldstone**: integrating out the heavy modes  $\Phi$  we generate a potential

$$V_{\mathcal{H}} = V_F^{SU(3)} + V_D \simeq V_F^{SU(3)} + V_D^{SU(2)} + \left( \frac{m^2}{F^2} \right) \sum_{\hat{T}} \hat{D}^2$$

$\hat{D}$  are the D-terms corresponding to the broken generators  
 $m$  is the soft breaking parameter for  $\Phi$ .

## PHYSICAL HIGGS POTENTIAL

- $m^2/F^2 \gg 1 \Rightarrow V_H = m_0^2 |H|^2 + \lambda_0 |H|^4$
- $m^2/F^2 \ll 1 \Rightarrow V_H = \lambda_0 |H|^4$

SUSY + global symmetry  $\Rightarrow$  only a quartic is generated at tree-level

# MATTER CONTENT

The minimal matter content with **anomaly cancellation** and **generation-universal** (*Csaki, Marandella, Strumia and Shirman '05*)

	$SU(3)_c$	$SU(3)_W$	$U(1)_x$	$U(1)_{em}$
$Q$	3	3	0	(+2/3, -1/3, -1/3)
$U$	$\bar{3}$	1	-2/3	-2/3
$2 \times D$	$\bar{3}$	1	+1/3	1/3
$Q'$	$\bar{3}$	3	-1/3	(1/3, -2/3, -2/3)
$\bar{Q}'$	3	$\bar{3}$	+1/3	(-1/3, +2/3, +2/3)
$2 \times L$	1	$\bar{3}$	-1/3	(-1, 0, 0)
$E$	1	$\bar{3}$	+2/3	(0, +1, +1)

- 2 extra heavy **fermionic little-partner of Top** (and Bottom)
- 2 extra heavy leptons and two (almost) **sterile neutrinos**



# RADIATIVE CORRECTIONS

$$\mathcal{W}_{top} = m_{Q'} \bar{Q}' Q' + y_1 \bar{Q}' \mathcal{H}_u U + y_2 Q \mathcal{H}_u Q' + \tilde{y}_1 \bar{Q}' \Phi_u U + \tilde{y}_2 Q \Phi_u Q'.$$

generates 1-loop correction to the Higgs potential once that  $\langle \Phi \rangle = (0, 0, F)$ .

## TOP/STOP LOOPS

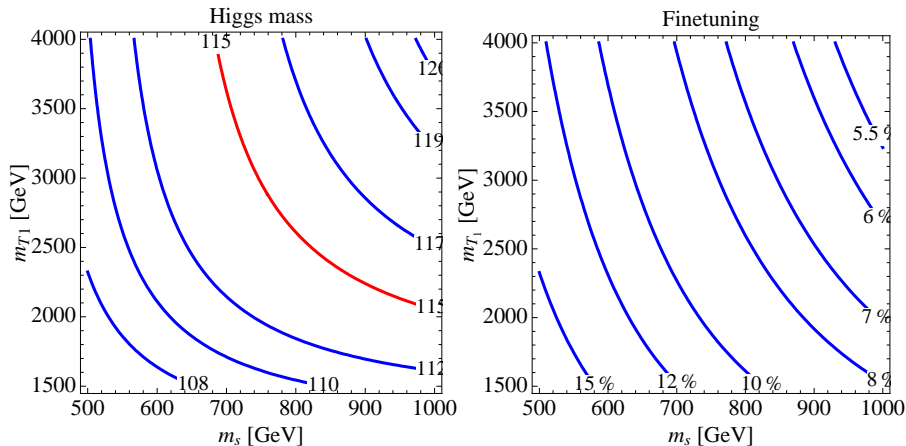
$$V_H = \delta m_H^2 |H|^2 + (\lambda_0 + \delta \lambda) |H|^4$$

$$\delta m_H^2 \simeq - \frac{3y_t^2 \sin^2 \beta}{8\pi^2} \left[ m_s^2 \ln \left( \frac{m_{T_1}^2 + m_s^2}{m_s^2} \right) + m_{T_1}^2 \ln \left( \frac{m_{T_1}^2 + m_s^2}{m_{T_3}^2} \right) \right]$$

$$\delta \lambda \simeq \frac{3y_t^4 \sin^4 \beta}{16\pi^2} \left[ \ln \left( \frac{m_s^2}{m_t^2} \right) + \ln \left( \frac{m_{T_1}^2}{m_s^2 + m_{T_1}^2} \right) + \frac{2m_s^2}{m_{T_1}^2} \ln \left( \frac{m_s^2}{m_s^2 + m_{T_1}^2} \right) \right]$$

# HIGGS MASS AND TUNING

Finetuning around  $5 \div 7\%$  with  $m_h \sim 115 \div 120$  GeV.



*Bellazzini, Pokorski, Rychkov and Varagnolo '08*

# DOPING THE QUARTIC

- “Double protection” works for  $\delta m_H^2$  but the **quartic is still too small**
- How enhance the quartic and retain the “double protection” ?

The idea is to **enlarge the gauge structure** and **break SUSY** in a way that the low-energy D-terms resemble the structure of SUSY D-terms but bigger.

	$SU(3)_W$	$U(1)_x$	$U(1)_z$
$\mathcal{H}_{u,d}$	$3, \bar{3}$	$\pm 1/3$	$\pm q'$
$\Phi_{u,d}$	$3, \bar{3}$	$\pm 1/3$	$\pm q$
$\Psi_{u,d}$	<b>1</b>	<b>0</b>	<b><math>q_\Psi</math></b>

- $\langle \Psi \rangle = \Omega \approx 10 \text{ TeV}$
- $m_\Psi / \Omega \sim 1$  and  $m/F \ll 1$
- $q' = q$

# THE EFFECTIVE D-TERMS

$$V_D \simeq V_D^{SU(2)_W} + \left( \frac{m^2}{F^2} \right) \sum_{\hat{T} \neq T^8} \hat{D}^2 + D_{8,x,y,z}^2$$

$$D_x = D_z, \quad D_y = \frac{1}{\sqrt{3}} D_8 + \frac{2}{3} D_x = |H_u|^2 - |H_d|^2$$

$$D_{8,x,y,z}^2 = a D_y^2 + b D_y D_x + c D_x^2$$

$$V_D = \left( V_D^{SU(2)_W} + a D_y^2 \right) + b D_y D_x + c D_x^2$$

❶  $q = q' \Rightarrow b = c = 0$

❷  $a = \frac{g_y^2}{8} \left[ 1 + a_1 \frac{m_\Psi^2}{\Omega^2} \right] > a^{SM} \Rightarrow m_h^2|_{\text{tree-level}} > m_Z^2 \cos^2 2\beta$

# THE EFFECTIVE D-TERMS

- SUSY limit  $m_\Psi/\Omega \rightarrow 0$

$$\Rightarrow V_D = D_{MSSM}^2 = \left( V_D^{SU(2)_W} + \frac{g_y^2}{8} D_y^2 \right)$$

$$a = a^{sm}(1 + a_1 m_\Psi^2/\Omega^2), \quad b = b_1 m_\Psi^2/\Omega^2, \quad c = c_1 m_\Psi^2/\Omega^2$$

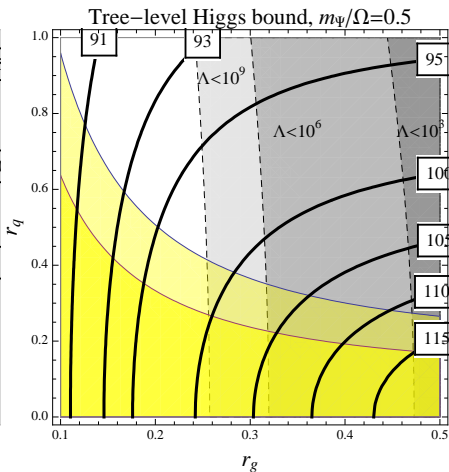
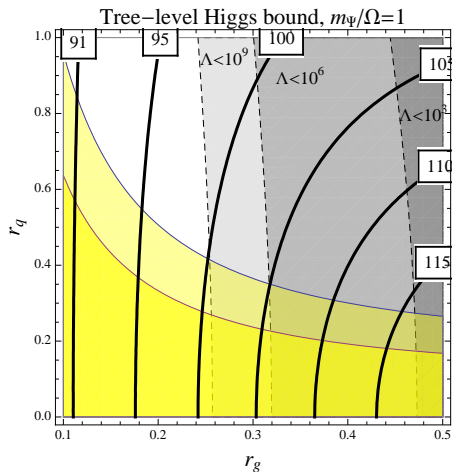
- Hard breaking  $m_\Psi/\Omega \rightarrow \infty$

$$\Rightarrow Y, \tilde{Y} = T^8/\sqrt{3} + Z/3q \quad \text{conserved} \Rightarrow D_{8,x,y,z}^2 = D_y^2 + D_{\tilde{y}}^2$$

$$YH_{u,d} \propto \tilde{Y}H_{u,d} \Rightarrow D_{\tilde{y}}^2 \propto D_y^2 \Rightarrow b = c = 0$$

$$V_D = V_D^{SU(2)_W} + \left( 1 + a_1 \frac{m_\Psi^2}{\Omega^2} \right) D_y^2$$

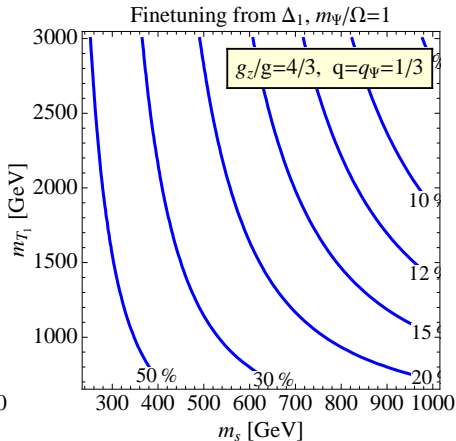
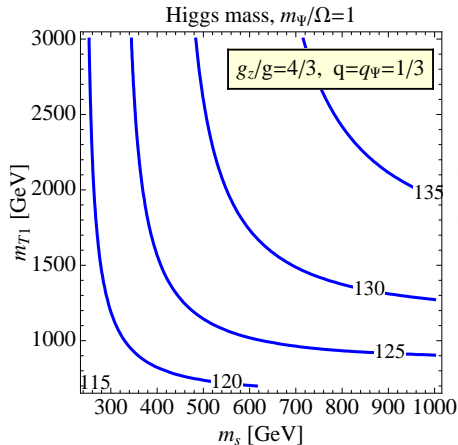
# THE TREE-LEVEL HIGG MASS



Constraints:  $Z'$  and  $\Lambda_{Landau} = 10^3, 10^6, 10^9$  TeV

$$r_g = qg_z/g, \quad r_q = q_\psi/3q$$

# HIGGS MASS AND TUNING I

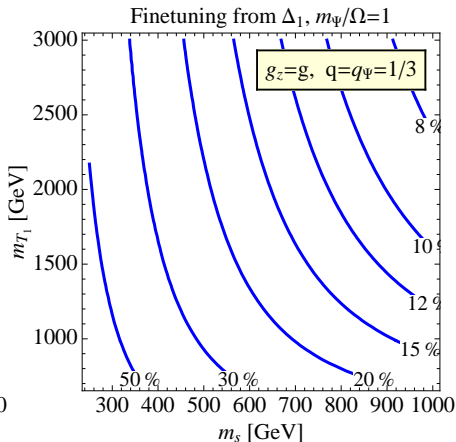
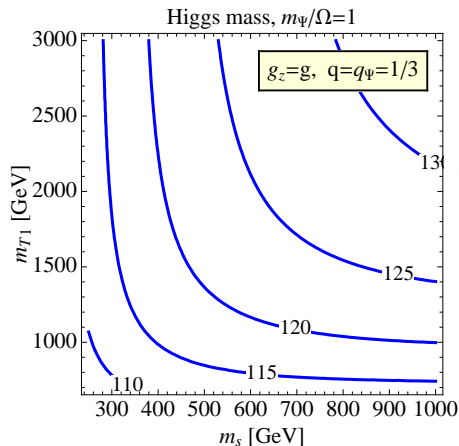


$$\Lambda_{Landau} = 10^3 \text{ TeV} \quad m_{stop} \approx 400 \text{ GeV}$$

$$m_{T1} \gtrsim 700 \div 1000 \text{ GeV}$$

Bellazzini, Csaki, Delgado, Weiler '09

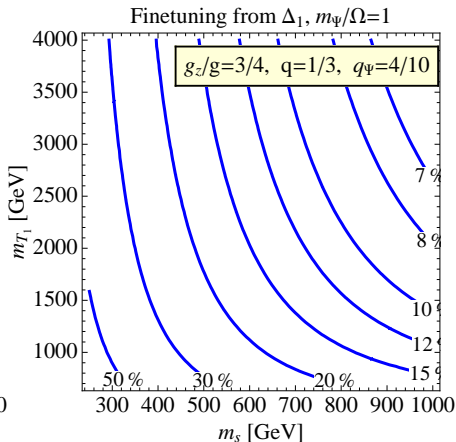
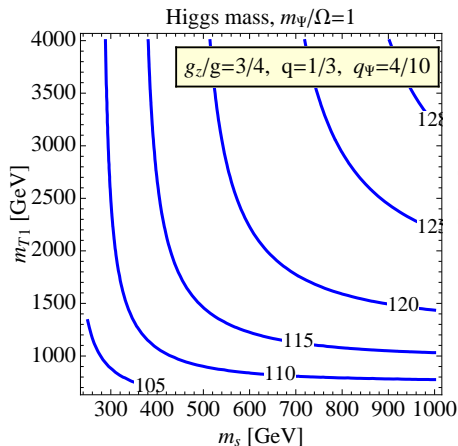
# HIGGS MASS AND TUNING II



$$\Lambda_{Landau} = 10^6 \text{ TeV}$$

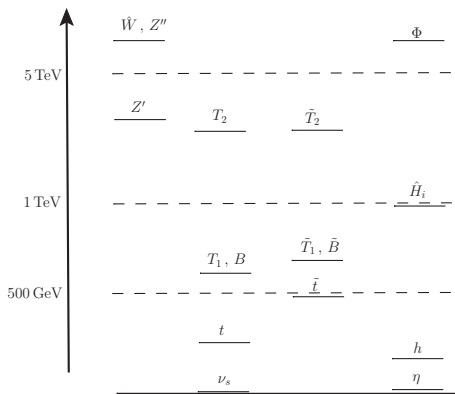


# HIGGS MASS AND TUNING III



$$\Lambda_{Landau} = 10^9 \text{ TeV}$$

# CONCLUSIONS



- SUSY+LH  $\Rightarrow$  “double protection”
- non decoupling D-terms  $\Rightarrow$  quartic enhancement
- no tuning

- fermionic “little” partner  $T_1$  of top at 700 GeV
- $Z'$  at 3.5 TeV or more
- sterile neutrinos at  $\sim$  KeV.