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Inhomogeneities in anisotropic universe

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•MOTIVATION

•VECTOR FIELD – BASIC PICTURE

•EVOLUTION OF AN ANISOTROPIC UNIVERSE

•PERTURBATIONS IN THE CURVATON SCENARIO

• NO SCALAR FIELDS IN MODERN EXPERIMENTS

- AXIS OF EVIL
- •POSSIBLE SOURCE OF INFLATION

• SOURCE OF PRIMORDIAL GRAVITATIONAL WAVES AND VECTOR PERTURBATIONS

LET US CONSIDER
$$g_{\mu\nu} = Diag(1, -a^2, -a^2, -b^2)$$
, THEN

$$G_0^0 = H^2 + 2H\mathcal{H}$$
 $G_3^3 = 2\dot{H} + 3H^2$

$$G_{1}^{1} = G_{2}^{2} = \dot{\mathcal{H}} + \dot{H} + H^{2} + \mathcal{H}^{2} + \mathcal{H}H$$



WE WILL ANALYSE THE EVOLUTION OF $\mathcal{H}-H=\lambda$

LET US INTRODUCE THE VECTOR FIELD IN THE ANIOSOTROPIC UNIVERSE

$$V_{\mu} = (0,0,0,V)$$
 $\delta V_{\mu} = (\delta V_t, \delta \vec{V})$ $\rho_V = \frac{1}{2b^2} (\dot{V}^2 + m^2 V^2)$ $\vec{V} = b \vec{U}$

$$\dot{\vec{V}} + (2H - \mathcal{H})\dot{\vec{V}} + m^2\vec{V} = 0$$

EQUATION OF MOTION

$$\ddot{ec{U}}+(2H+\mathcal{H})\dot{ec{U}}+(m^2+2H\mathcal{H}+\dot{\mathcal{H}})ec{U}=0$$

STRANGE MASS TERM FOR INFLATION!

$$m^2 pprox -2 H \mathcal{H} pprox -2 H^2$$

FOR THE VECTOR FIELD + PERFECT FLUID ONE OBTAINS

$$T_{1}^{1} = T_{2}^{2} = -p_{V} - p_{f} \qquad T_{3}^{3} = +p_{V} - p_{f}$$

THEN $\dot{\lambda} + 3\mathcal{H}\lambda + 2\lambda^2 = 2p_V$, which gives



INFLATION

 $2p_v$

 $\mathcal{H}, p_V = const \Rightarrow \lambda =$

FOR THE QUASI DE SITTER UNIVERSE WE OBTAIN

MORE REALISTIC MODELS GIVES THE SAME RESULT!

NUMBER OF E-FOLDS
$$N_z = \int \mathcal{H} dt$$
 $N_x = \int H dt$

THIS CAN GIVE SOME INFORMATION ABOUT ALLOWED ENERGY DENSITY OF THE VECTOR FIELD

$$\Delta N = \int \lambda \, dt = \frac{2p_V}{3\mathcal{H}} t \Rightarrow \frac{\Delta N}{N_z} = \frac{2p_V}{3\mathcal{H}^2} \propto \frac{p_V}{\rho} \propto U^2$$

DECREASING, WHILE $ec{U}$ IS GLUTTED TO THE WALL OF ITS POTENTIAL $\ mt \ll 1$

$$p_V \approx H\dot{U}U + \frac{1}{2}H^2U^2$$
 $3H\dot{U} \approx -m^2U$

$$\lambda \approx H U_o^2 (1-2.5\,m^2 t^2)$$

DECREASING FOR THE OSCILLATION PHASE $mt \gg 1$

$$p_V \propto \cos(2mt) \Rightarrow < p_V >= 0$$

WHEN \vec{U} OSCILLATES IT EVOLVES LIKE A SCALAR FIELD

$$mU = \sqrt{6}Hcos(mt)$$
 $\dot{U} = -\sqrt{6}Hsin(mt)$

THEN
$$H \approx \frac{2}{3t} \left(1 + \frac{\sin(2mt)}{6mt} \right) \approx \frac{2}{3t}$$

SO
$$\omega \approx 0 \Rightarrow$$
 NO PRESSURE! λ DECREASE LIKE a^{-3}

EXISTENCE OF THE VECTOR FIELD BREAKS ISOTROPY AT THE PERTURBATIONS LEVEL

$$\delta g_{\alpha\beta} = \begin{pmatrix} 2\Phi & -aB_{,i} & -b\chi_{,z} \\ & 2a^2(\Psi\delta^{i}_{j} + E_{,ij}) & -abC_{,zi} \\ & & 2b^2\Sigma \end{pmatrix}$$

GAUGE FIXING: $\Rightarrow E = C = \Psi = 0$

VANISHING OF NON DIAGONAL COMPONENTS

$$\delta T^{i}_{i} = 0 \Rightarrow \Sigma = \Phi + a\dot{B} + a(H + \mathcal{H})B \quad i \neq j$$

THE LIQUID INDEPENDENT EQUATION OF MOTION

 $2\ddot{\Sigma} + 2(H + 2\mathcal{H})\dot{\Sigma} + 2(\mathcal{H} - H)\dot{\Phi} + (\mathcal{H}H + \mathcal{H}^2 - 2H^2 + \dot{\mathcal{H}} - \dot{H})\Phi = 4\delta p_V \simeq 0$

For $\mathcal{H} \cong \mathcal{H}$ we obtain

$$\ddot{\Sigma} + 3 H \dot{\Sigma} = 0 \Rightarrow \Sigma = \Sigma_* = const$$

FROM THE EINSTEINS EQUATION WE HAVE

$$H(\dot{\Phi} + 3Hc_s^2\Phi) = \tau\delta S$$
 $3H^2\Phi = -\frac{1}{2}\delta\rho$

FOR THE VECTOR + RADIATION SCENARIO ONE OBTAINS

$$c_s^2 = \frac{\beta}{a} \Rightarrow \Phi = \Phi_* e^{3\beta/a} \to \Phi_*$$

THEN

$$\Phi_* = \Psi_* \Rightarrow B \propto t^{-4/3}$$

WE RECONSTRUCT THE STANDARD FRW PERTURBATION POWER SPECTRUM!

ONLY Φ IS CHANGED BY THE ENTROPY!

$\Phi_* \to \Phi_f \quad \Rightarrow \quad \Phi_f \neq \Sigma_f$

THE ENTROPY PERTURBATION BREAKS THE INITIAL SYMMETRY!

IT DOES NOT NEED TO BE DANGEROUS. WE SIMPLY NEED SMALL INITIAL ENTROPY PERTURBATION

IT MEANS, THAT INITIALLY WE NEED

$$\frac{\delta U_*}{U_*} \ll \Phi_*$$

•THE SUBDOMINANT VECTOR FIELD SHALL NOT PRODUCE HUGE BACKGROUND ANISOTROPIES

•GOOD CURVATON CANDIDATE

•FUTURE PROBLEMS WITH THE VECTOR PERTURBATIONS POSSIBLE