

Multi-Reference Calculations for Odd-Mass Nuclei

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- 1 Outline of the method
- 2 Application to ^{25}Mg
- 3 Conclusion and outlook

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We define an EDF (\equiv effective Hamiltonian).



We create a set of one-quasiparticle states:

$$\{|\Phi_a\rangle, a = \dots\}$$



We project each of them on the good quantum numbers:

$$\{|JMNZP\epsilon, a\rangle, a, J, P, \epsilon = \dots\}.$$



We diagonalize the (effective) Hamiltonian between the projected states: $\{|JMNZP\xi\rangle, J, P, \xi = \dots\}$.



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$$\mathcal{E}^{nuc}[\rho, \kappa, \kappa^*]^{ab} = \frac{\langle \Phi_a | \hat{H} | \Phi_b \rangle}{\langle \Phi_a | \Phi_b \rangle}$$

$$\rho^{ab} = \frac{\langle \Phi_a | \hat{a}^\dagger \hat{a} | \Phi_b \rangle}{\langle \Phi_a | \Phi_b \rangle}, \quad \kappa^{ab} = \frac{\langle \Phi_a | \hat{a} \hat{a} | \Phi_b \rangle}{\langle \Phi_a | \Phi_b \rangle}, \quad \kappa_t^{ba^*} = \frac{\langle \Phi_a | \hat{a}^\dagger \hat{a}^\dagger | \Phi_b \rangle}{\langle \Phi_a | \Phi_b \rangle}.$$

- $|\Phi_a\rangle, |\Phi_b\rangle$: different quasiparticle states.
 $\langle \Phi_a | \Phi_b \rangle \neq 0$ (condition to use the EWT of Balian-Brézin)
- \mathcal{E}^{nuc} **directly and uniquely** determined by \hat{H} .
 \Rightarrow respect the **Pauli principle**.

$$\hat{H} = \hat{K}^{(1)} + \hat{V}_{Coul}^{(2)} + \hat{V}_{Sky}^{(2-4)}$$

- $\hat{K}^{(1)}$: kinetic energy (+ CoM corr.).
- $\hat{V}_{Coul}^{(2)}$: Coulomb interaction.
- $\hat{V}_{Sky}^{(2-4)}$: Skyrme pseudo-potential. **Phenomenological.**

$$\hat{V}_{Sky}^{(2-4)} = \hat{V}_{Sky}^{(2)} + \hat{V}_{Sky}^{(3)} + \hat{V}_{Sky}^{(4)}$$

- $\hat{V}_{Sky}^{(2)} = t_0 (1 + x_0 \hat{\Gamma}_{12}^\sigma) \hat{\delta}_{r_1 r_2} + \frac{t_1}{2} (1 + x_1 \hat{\Gamma}_{12}^\sigma) \left(\hat{k}_{12}'^2 \hat{\delta}_{r_1 r_2} + \hat{\delta}_{r_1 r_2} \hat{k}_{12}^2 \right) + t_2 (1 + x_2 \hat{\Gamma}_{12}^\sigma) \hat{k}_{12}' \hat{\delta}_{r_1 r_2} \cdot \hat{k}_{12} + iW_0 (\hat{\sigma}_1 + \hat{\sigma}_2) \hat{k}_{12}' \hat{\delta}_{r_1 r_2} \times \hat{k}_{12}$
- $\hat{V}_{Sky}^{(3)} = u_0 (\hat{\delta}_{r_1 r_3} \hat{\delta}_{r_2 r_3} + \hat{\delta}_{r_3 r_2} \hat{\delta}_{r_1 r_3} + \hat{\delta}_{r_2 r_1} \hat{\delta}_{r_3 r_1})$
- $\hat{V}_{Sky}^{(4)} = v_0 (\hat{\delta}_{r_1 r_3} \hat{\delta}_{r_2 r_3} \hat{\delta}_{r_3 r_4} + \dots)$
- **9** parameters.

- SLyMR0 parametrization.
Sadoudi *et al.* Physica Scripta T154 014013 (2013).

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We calculate observables.

- Number parity: $\hat{\Pi}_A|\Phi_a\rangle = \pi_a|\Phi_a\rangle$ ($\hat{\Pi}_A = e^{-i\pi\hat{A}}$)

Subgroup of D_{2h}^{TD} :

- Signature: $\hat{R}_x|\Phi_a\rangle = \eta_a|\Phi_a\rangle$ ($\hat{R}_x = e^{-i\pi\hat{J}_x}$)

- Parity: $\hat{P}|\Phi_a\rangle = p_a|\Phi_a\rangle$

- y -Time Simplex: $\hat{S}_y^T|\Phi_a\rangle = |\Phi_a\rangle$ ($\hat{S}_y^T = \hat{R}_y\hat{P}\hat{T}$)

- odd A : $\pi_a = -1$, $\eta_a = \pm i$, $p_a = \pm 1$

$$\text{Minimization: } \delta \mathcal{E}^{nuc}[\rho, \kappa, \kappa^*]^{aa} = 0$$

Constraints using Lagrange parameters:

- Neutron number: $\langle \Phi_a | \hat{N} | \Phi_a \rangle = N$
- Proton number: $\langle \Phi_a | \hat{Z} | \Phi_a \rangle = Z$
- Quadrupole deformation: $\langle \Phi_a | \hat{Q} | \Phi_a \rangle = Q$

- Self-consistent problem: solved by an iterative procedure.
- Solved for different values of Q and/or $\hat{\beta}_a^\dagger$

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Conservation of the neutron and proton numbers:

- $U(1)_N \times U(1)_Z$
- Broken by: pairing correlations.

⇒ Projection on neutron and proton numbers.

Conservation of total angular momentum:

- $SU(2)_A$
- Broken by: quadrupole deformation.

⇒ Projection on total angular momentum.

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We calculate observables.

- $|\Lambda M \xi\rangle = \sum_{i=1}^{\Omega_I} \sum_{\epsilon=1}^{\Omega_i^\Lambda} F_\xi^\Lambda(i, \epsilon) |\Lambda M \epsilon, i\rangle$

$$\Lambda \equiv (J, N, Z, P)$$

Ω_I : set of states $|\Phi_i\rangle$

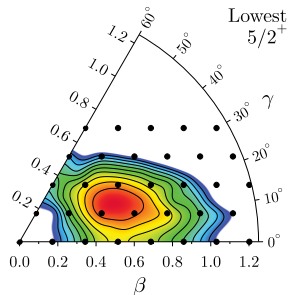
Ω_i^Λ : set of projected states given (Λ, i) .

- i : deformation and blocked quasiparticle.

$$\frac{\delta}{\delta F_\xi^{\Lambda*}(i, \epsilon)} \left(\frac{\langle \Lambda M \xi | \hat{H} | \Lambda M \xi \rangle}{\langle \Lambda M \xi | \Lambda M \xi \rangle} \right) = 0 \implies F_\xi^\Lambda(i, \epsilon) \text{ et } E_\xi^\Lambda(\Omega_I)$$

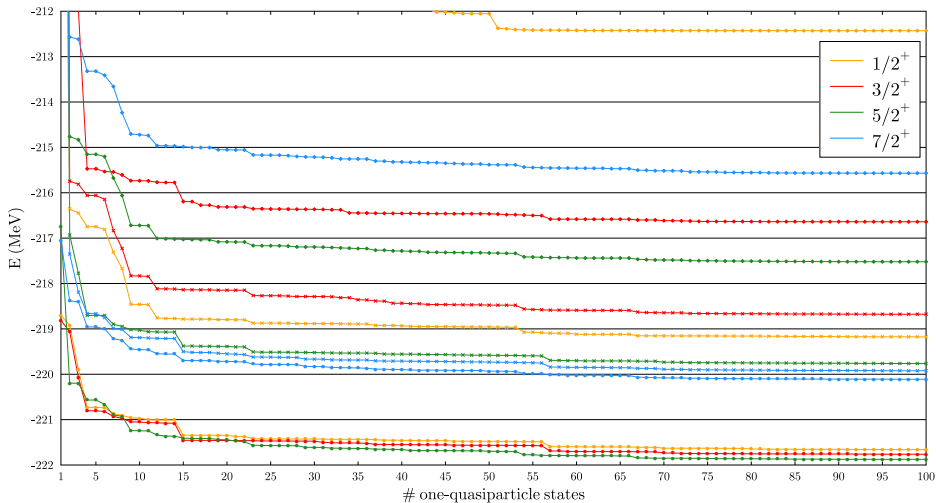
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- Proof of principle.
- Light nucleus with a simple structure.
- Phys. Rev. Lett. **113** 162501 (2014)

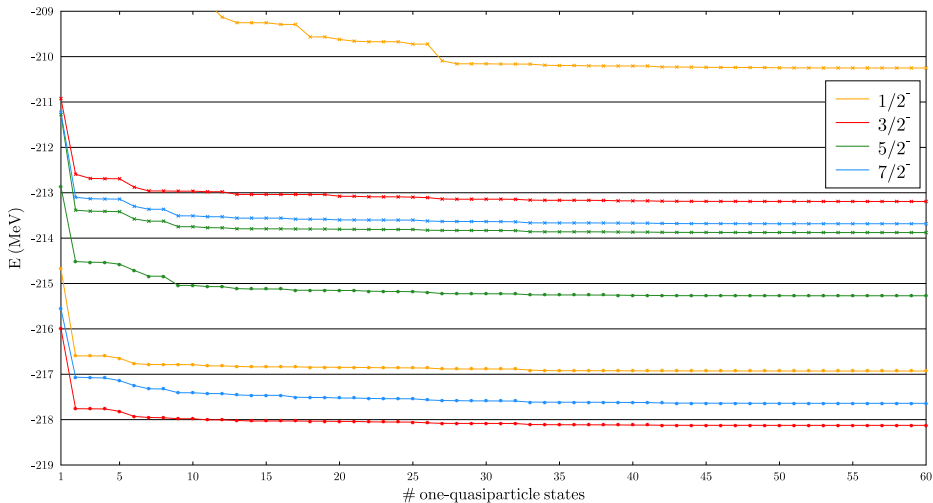


- Discretization mesh (q_1, q_2): 40 fm^2
- Several 1qp states at each deformation.
- Total number of one-quasiparticle states used:
 - positive parity: 100 states.
 - negative parity: 60 states.

Convergence analysis



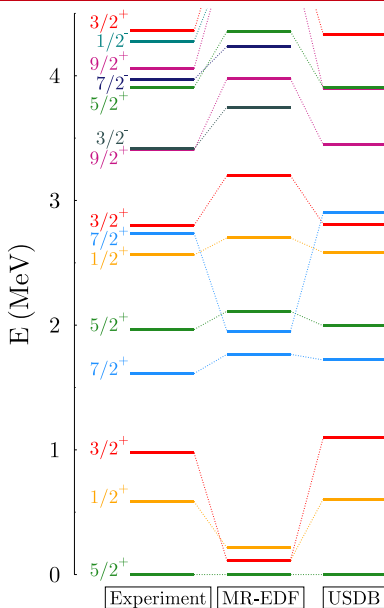
Convergence analysis II



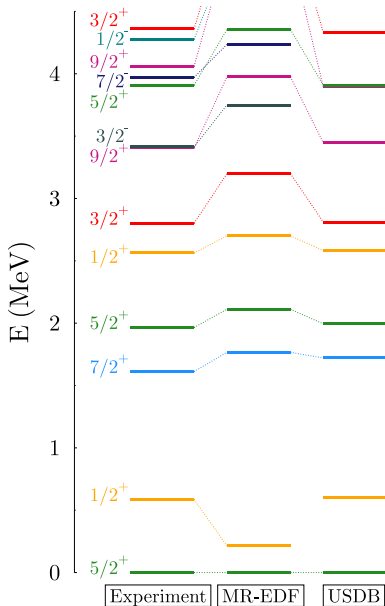
	J^π	Binding energy (MeV)	Q_s (e fm ²)	μ (μ_N)
Experiment	$\frac{5}{2}^+$	-205.587	20.1(3)	-0.85545(8)
MR-EDF	$\frac{5}{2}^+$	-221.875	23.25	-1.054

- No effective charge or effective g -factor!

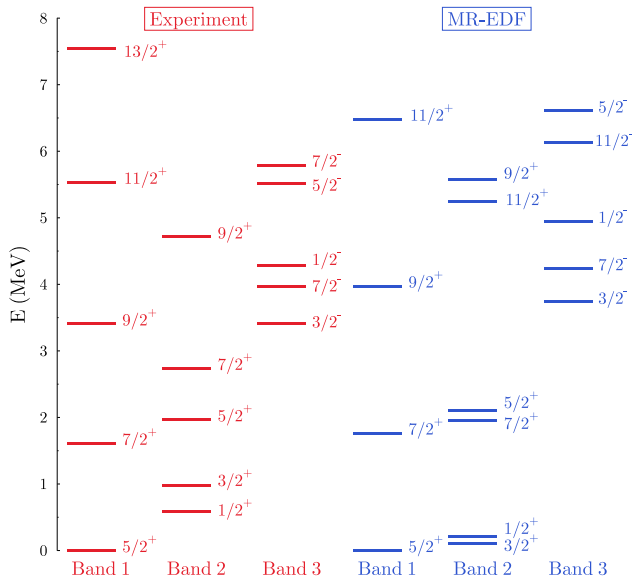
- Experiment: Nuclear Data Sheets **110** 1691 (2009)

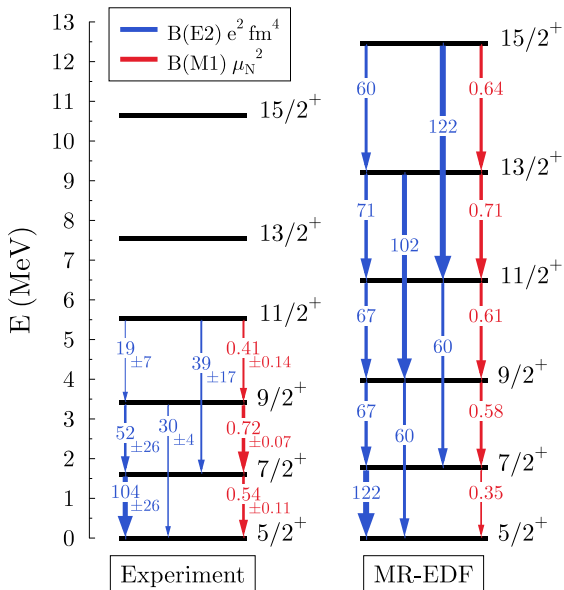


- Shell Model: USDB with NushellX.
- Experiment: Nuclear Data Sheets **110** 1691 (2009)
- Not as good as Shell Model but less parameters!
And not fitted specifically to the sd shell!
- Negative parity states!



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- Overall reasonable description ...
- ... especially considering the limited quality of SLyMR0.
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Goals:

- ✓ Treatment of even-even and odd-even (and even odd-odd) nuclei on the same footing.
- ✓ MR-EDF calculations with a Hamiltonian-based functional.
- ✓ Spectroscopy of odd-mass nuclei.

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 - **underway**: Skyrme with gradient three-body terms.

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- More calculations and heavier nuclei.
 - speedup of the programs.