



Constraining Finite-Range Momentum Dependent Effective Interactions

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Skyrme Interaction

Contact interaction

$$v_{\text{Skyrme}}(r) = t_0(1 + x_0 P^\sigma) \delta(r)$$

$$+ t_1(1 + x_1 P^\sigma) (\delta(r) k^2 + k'^2 \delta(r)) / 2$$

$$+ t_2(1 + x_2 P^\sigma) \mathbf{k}'^* \delta(r) \cdot \mathbf{k}$$

$$+ t_3/6(1 + x_3 P_\sigma) \rho^\alpha \delta(r)$$

$$+ iW_0 \hat{\sigma} \cdot [\mathbf{k}'^* \delta(r) \times \mathbf{k}]$$

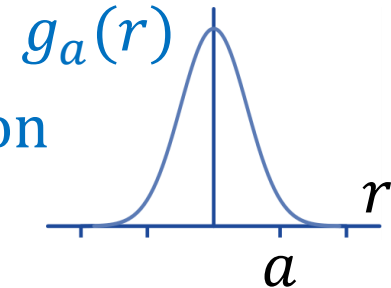
Momentum
Dependent

Density dependent
 $1/6 < \alpha < 2/3$

Spin-orbit

Regularized Interaction

Gaussian interaction



$$v_{\text{Reg}} = t_0(1 + x_0P^\sigma - y_0P^\tau - z_0P^{\sigma\tau})g_a(r)$$

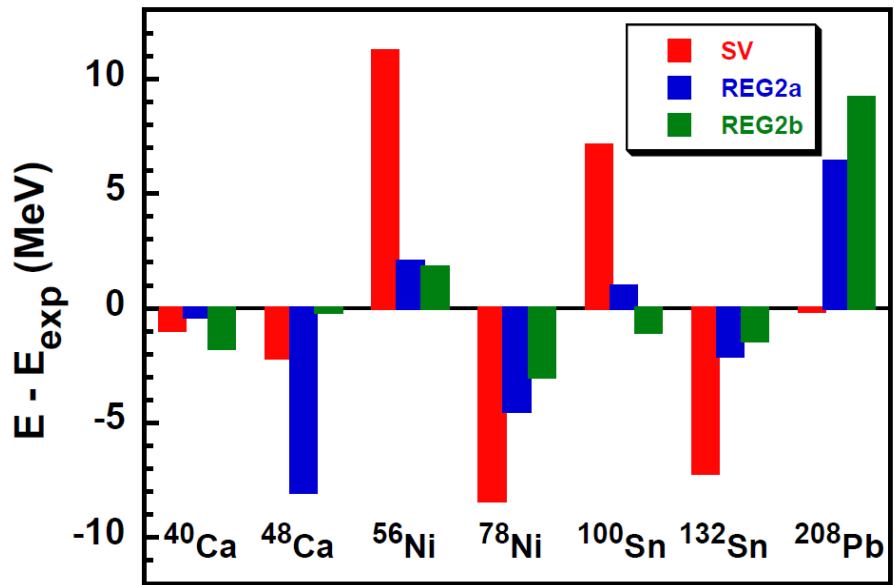
$$+t_1(1 + x_1P^\sigma - y_1P^\tau - z_1P^{\sigma\tau})(g_a(r)k^2 + k'^2g_a(r))/2 + \hat{T}_1$$

$$+t_2(1 + x_2P^\sigma - \underbrace{y_2P^\tau - z_2P^{\sigma\tau}}_{\text{bracket}})k'^*g_a(r) \cdot \mathbf{k} \quad \hat{T}_2$$

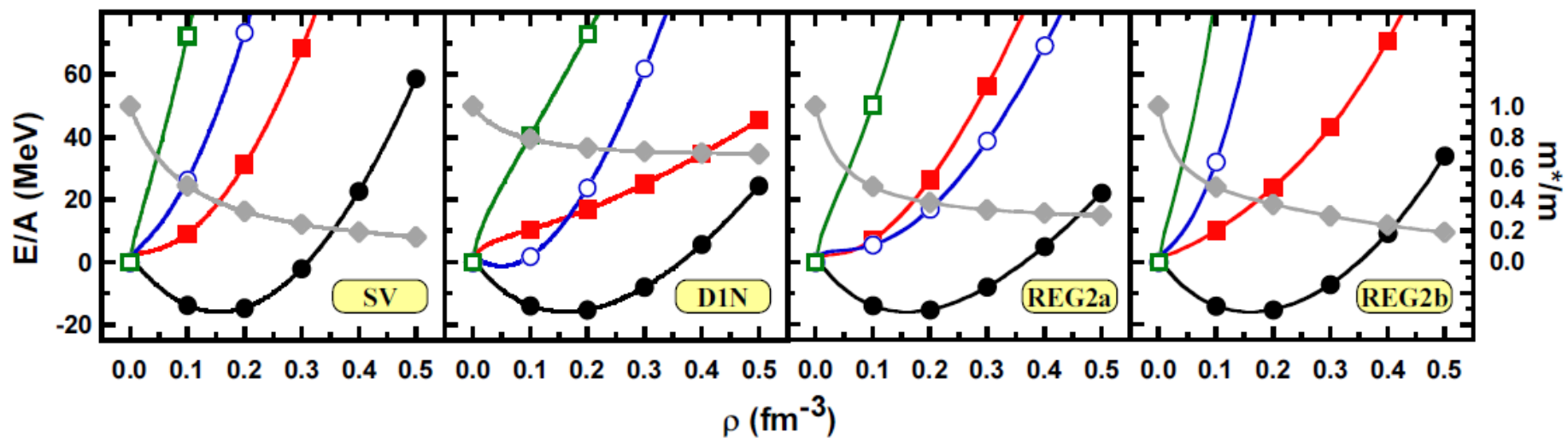
- $P^\tau \neq \pm 1$
- coefficients y_i, z_i not reabsorbed
- more parameters at a given order

	Ref.	REG2a
K_∞ (MeV)	230	230.00
L (MeV)	75	100.2
E/A (sat.) (MeV)	-16	-16
ρ_{sat} (fm)	0.16	0.160
J (MeV)	32	32

Close Shell masses

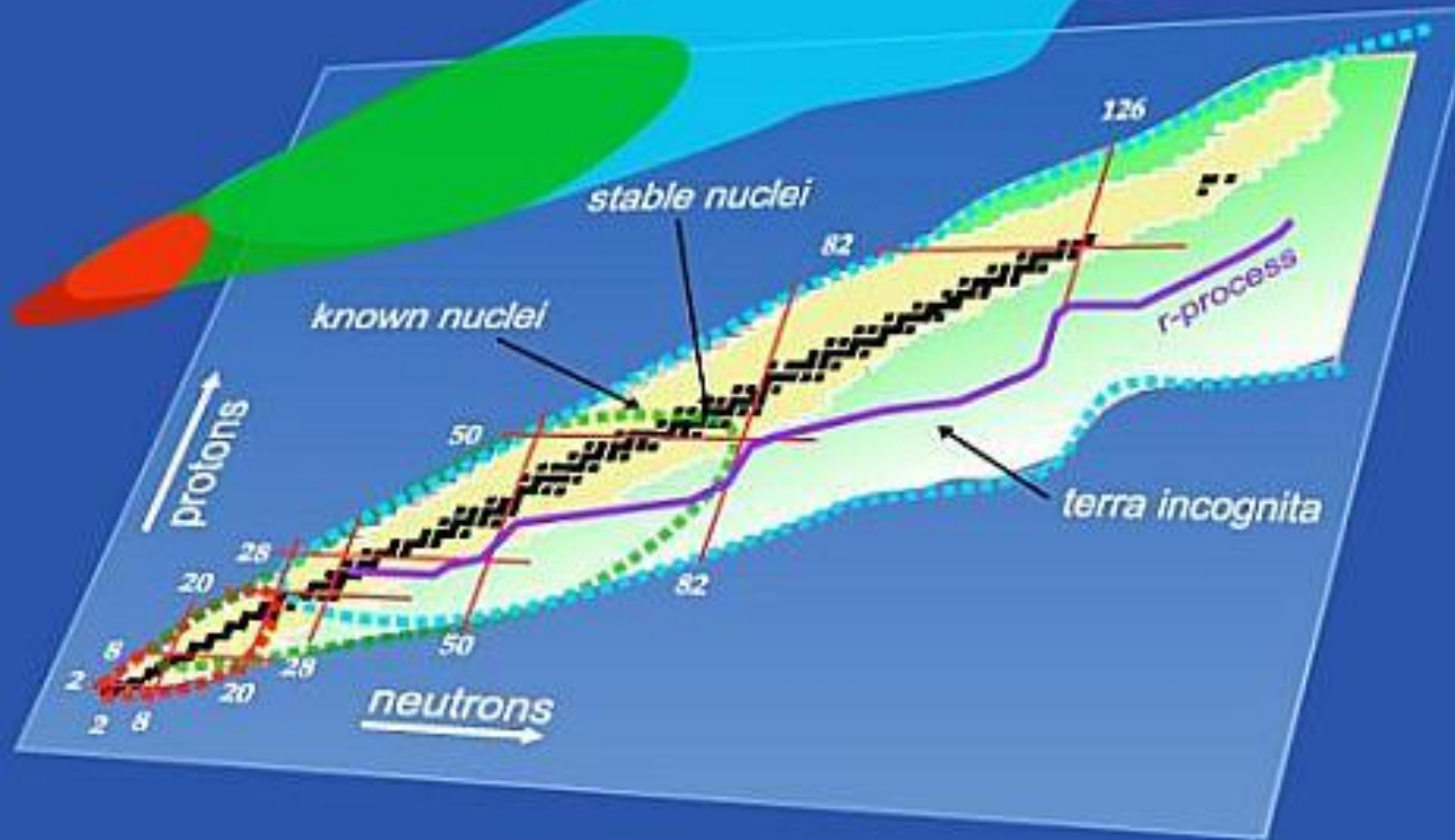


Infinite Matter Properties



Nuclear Landscape

- Ab initio
- Configuration Interaction
- Density Functional Theory



from UNEDF website

Pairing Matrix Elements to constrain properties of open shell finite nuclei

Gomez, Prieto, Navarro, NPA 549 (1992) 125

$$V_{nlj}^{ST} = \langle (nlj, nlj)_{J=0} | v_{Reg}^{ST} | (nlj, nlj)_{J=0} \rangle$$

one can gauge the matrix element on the pairing energy of a close shell + 2v nucleus. In $V_{f_{7/2}^2}^{ST}$ case, ^{42}Ca

Pairing Matrix Elements to constrain properties of open shell finite nuclei

Gomez, Prieto, Navarro, NPA 549 (1992) 125

$$V_{nlj}^{ST} = \langle (nlj, nlj)_{J=0} | v_{Reg}^{ST} | (nlj, nlj)_{J=0} \rangle$$

$$V_{f_{7/2}}^{ST}$$

Wavefunction

$$\psi_{nljq}^{S=0}(\vec{r}_1\sigma_1, \vec{r}_2\sigma_2) = \sum_{m_{l_1}} \frac{\sqrt{2j+1}}{2(2l+1)} (-)^{l-m_{l_1}} \delta_{\sigma_1-\sigma_2} Y_{m_{l_1}}^l(\hat{r}_1) Y_{-m_{l_1}}^l(\hat{r}_2) R_{nljq}(r_1) R_{nljq}(r_2)$$

Multipole expansion

$$g_a(\vec{r}_1 - \vec{r}_2) = \frac{e^{-|\vec{r}_1 - \vec{r}_2|^2/a^2}}{(a\sqrt{\pi})^3} = \frac{4\pi}{(a\sqrt{\pi})^3} e^{-(r_1^2+r_2^2)/a^2} \sum_{LM} i_L \left(2 \frac{r_1 r_2}{a^2}\right) Y_M^{L*}(\hat{r}_1) Y_M^L(\hat{r}_2)$$

The M.E. is an integral over a **total of 6 Spherical Harmonics** with gradients!

Local NLO contribution

Difference between Momentum dependent operators

$$\widehat{T}_1 - \widehat{T}_2 = (k'_{12}{}^2 + k_{12}^2)/2 - \vec{k}'_{12} \cdot \vec{k}_{12} = \left(\vec{k}'_{12} - \vec{k}_{12} \right)^2 / 2$$

Local part of the interaction

$$\left[\widehat{T}_1 - \widehat{T}_2, \delta(r'_1 - r_1) \delta(r'_2 - r_2) \right] = 0$$

Commutates with locality deltas

$$\left(\widehat{T}_1 - \widehat{T}_2 \right) \delta(r'_1 - r_1) \delta(r'_2 - r_2) g_a(r_1 - r_2) \rightarrow - \frac{\left(\vec{\nabla}_1 - \vec{\nabla}_2 \right)^2}{2} g_a(r_1 - r_2)$$

thus is the laplacian of the gaussian

That is conveniently related to the

Derivative respect to the range

$$\frac{\left(\vec{\nabla}_1 - \vec{\nabla}_2 \right)^2}{2} g_a(r_1 - r_2) = \frac{1}{a} \frac{\partial}{\partial a} g_a(r_1 - r_2)$$

Local N^nLO contribution to matrix elements for regularized interaction

$$\left(\frac{(\vec{\nabla}_1 - \vec{\nabla}_2)^2}{2} \right)^n g_a(r_1 - r_2) = \left(\frac{1}{a} \frac{\partial}{\partial a} \right)^n g_a(r_1 - r_2)$$

The matrix element for a local regularized interaction ($t_2 = -t_1$),

$$V_{N^nLO}[t^{(n)}] = \left(-\frac{1}{a} \frac{\partial}{\partial a} \right)^n V_{LO}[t^{(n)}]$$

	Ref.	REG2a	New
K_∞ (MeV)	230	230.00	230.51
L (MeV)	75	100.20	84.87
E/A (sat.) (MeV)	-16	-16.00	-16.17
ρ_{sat} (fm)	0.16	0.160	0.160
J (MeV)	32	32.00	33.31
$\langle \psi_{f_{7/2}^2} V_{Loc}^{S=0} \psi_{f_{7/2}^2} \rangle$ (MeV)	-3	-0.54	-1.76
$\langle \psi_{f_{7/2}^2} V_{Loc}^{S=1} \psi_{f_{7/2}^2} \rangle$ (MeV)	$\gtrsim 0$	0.12	0.13

Pairing Energy ^{44}Ca
-9.153 MeV.

One Body
Density Funct.

$$E_{p-p}^{ST}[\rho] = \frac{1}{2} \sum_{\substack{nlj \\ n'l'j'}} (-)^{l'+l} \sqrt{(2j+1)(2j'+1)} (u \cdot v)_{nlj} (u \cdot v)_{n'l'j'}$$

$$\langle (n'l'j', n'l'j')_{00} | v_{Reg}^{ST} | (nlj, nlj)_{00} \rangle$$

Two Body
Matrix Element

In the case of pairing for close shell + 2v

$$n'l'j' = nlj$$

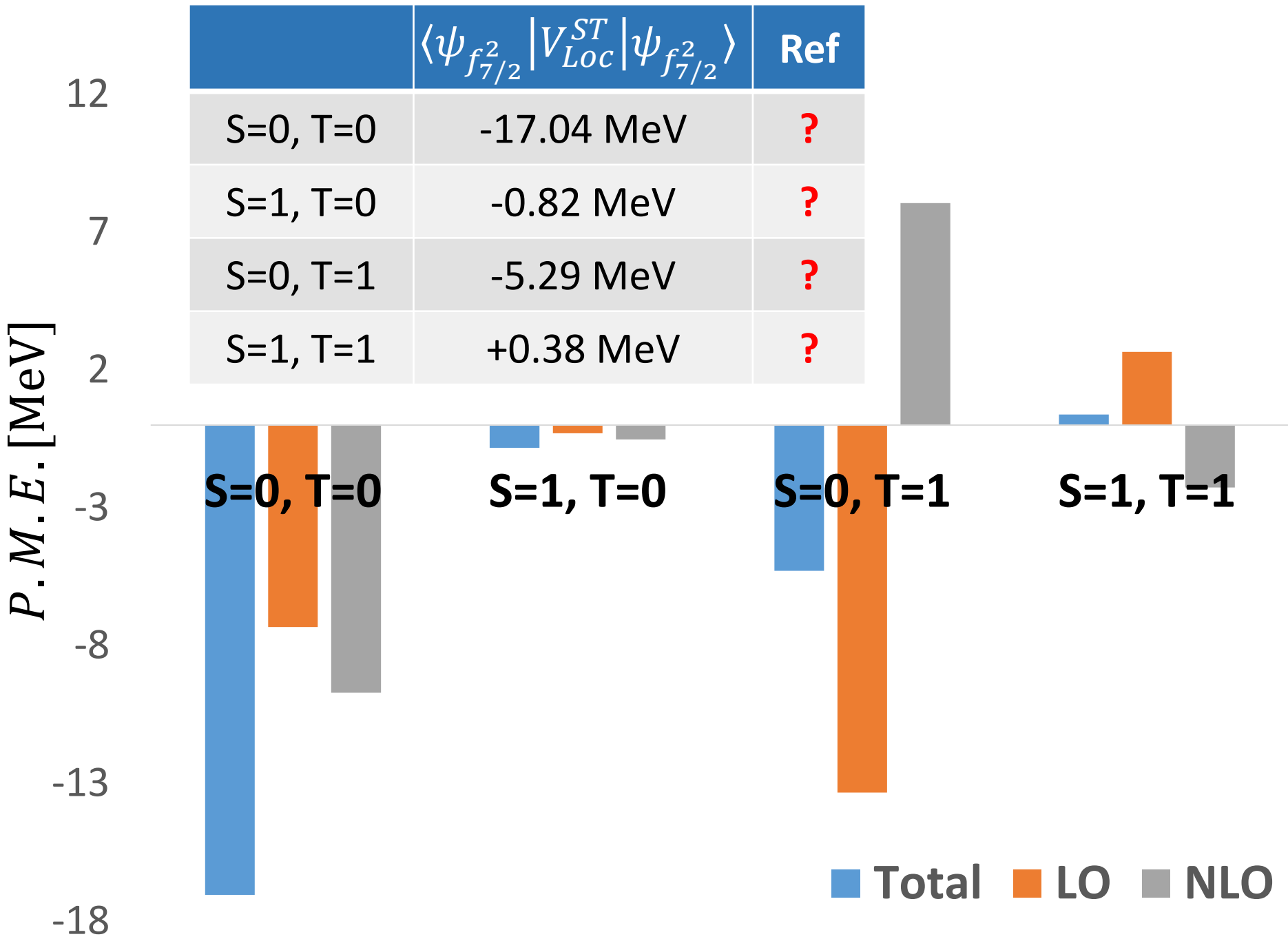
One Body

Density Funct.

$$E_{p-p}^{ST}[\rho] = \frac{1}{2} (-)^{l'+l} \sqrt{(2j+1)(2j'+1)} (u \cdot v)_{nlj}^2$$

$$\langle (nlj, nlj)_{00} | v_{Reg}^{ST} | (nlj, nlj)_{00} \rangle$$

Two Body
Matrix Element



Landau Parameters are the coefficient of the p-h interaction expanded in the legendre polynomial basis for the different spin-isospin channels, calculated at the Fermi surface.

$$\sum_l f_l^{(\alpha)} P_l(\vec{k} \cdot \vec{k}') = v_{p-h}(\vec{k}, \vec{k}') \quad \text{Fourier Transform}$$

channel

Particle-hole interaction
 $= v(1 - P^x P^\sigma P^\tau)$

$$\sum_l f_l^{(\alpha)} P_l(\vec{k} \cdot \vec{k}') = v_{p-h}(\vec{k}, \vec{k}')$$

$$v_{\text{Reg}}^{LO}(r_{12}, r'_{12}) = t_0(1 + x_0 P^\sigma - y_0 P^\tau - z_0 P^{\sigma\tau}) g_a(r_1 - r_2) \\ (\delta(r_1 - r'_1) \delta(r_2 - r'_2) - \delta(r_1 - r'_2) \delta(r_2 - r'_1) P^\sigma P^\tau)$$

Direct term

Exchange term

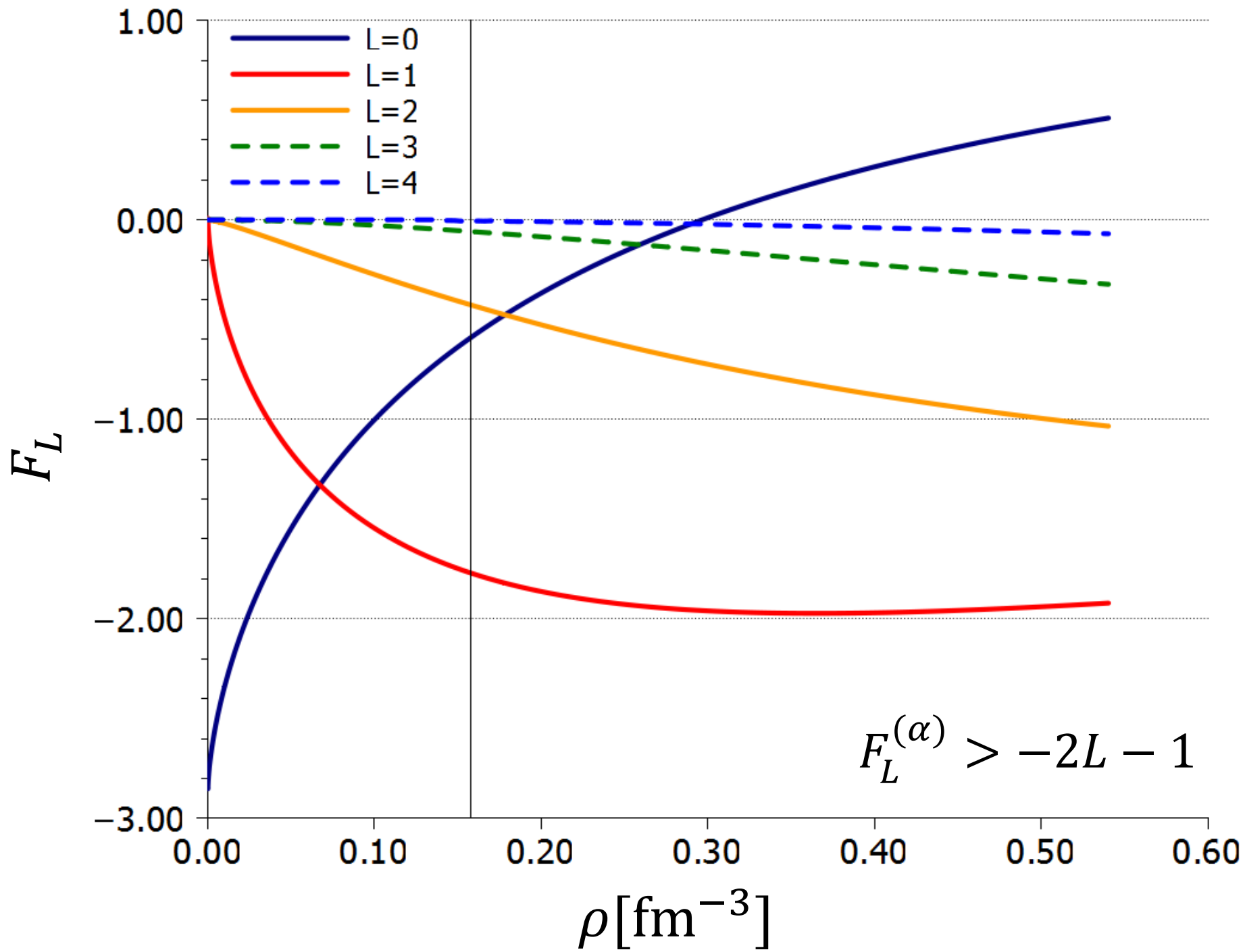
$$\mathcal{F}[v_{\text{Reg}}^{LO}] = D^{(\alpha)}(q) - E^{(\alpha)}(\vec{k} - \vec{k}')$$

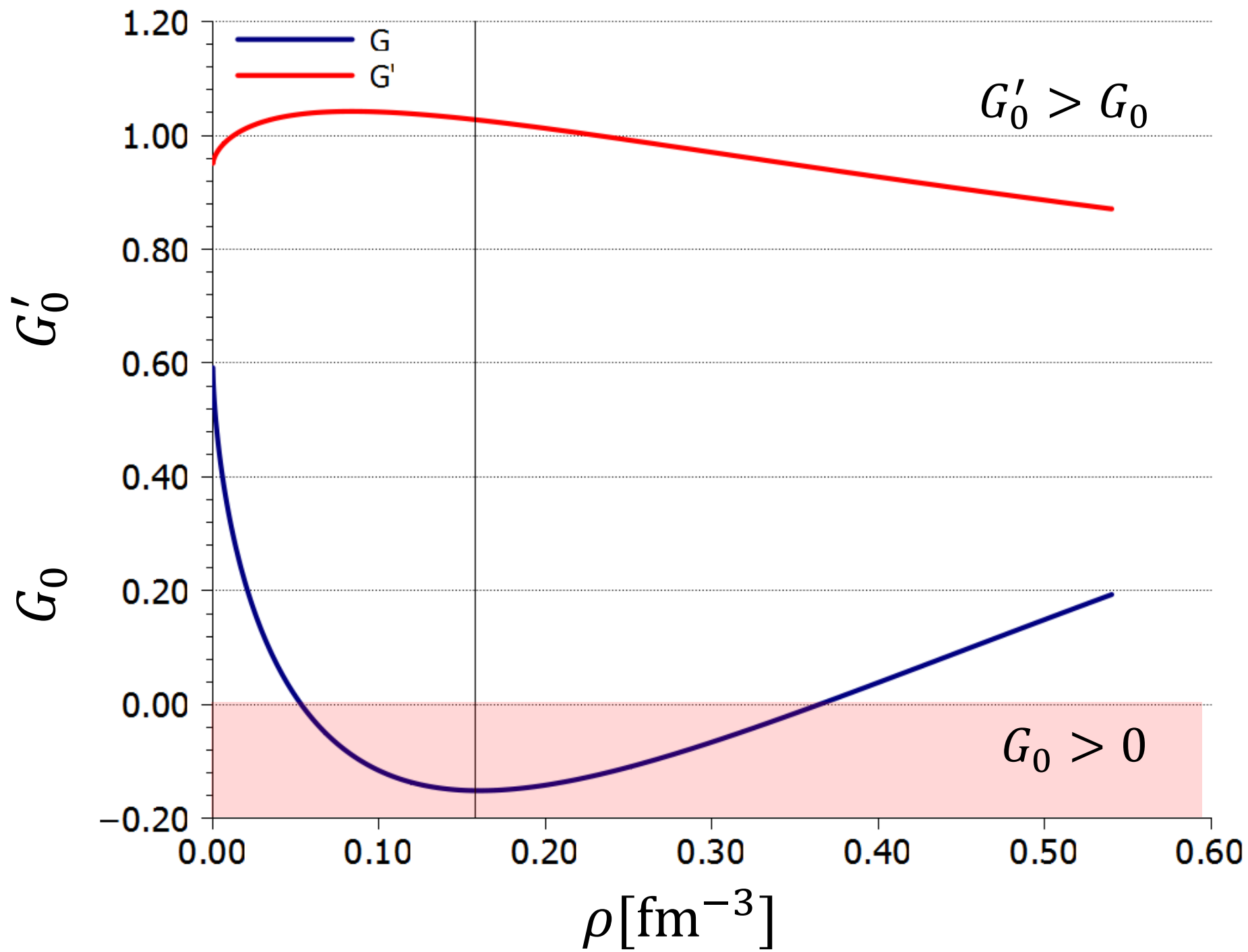
$$f_l = D^{(0,0)} \delta_{l,0} - E^{(0,0)} f_l(a, k_F)$$

$$g_l * \vec{\sigma} \cdot \vec{\sigma} = \left(D^{(1,0)} \delta_{l,0} - E^{(1,0)} f_l(a, k_F) \right)$$

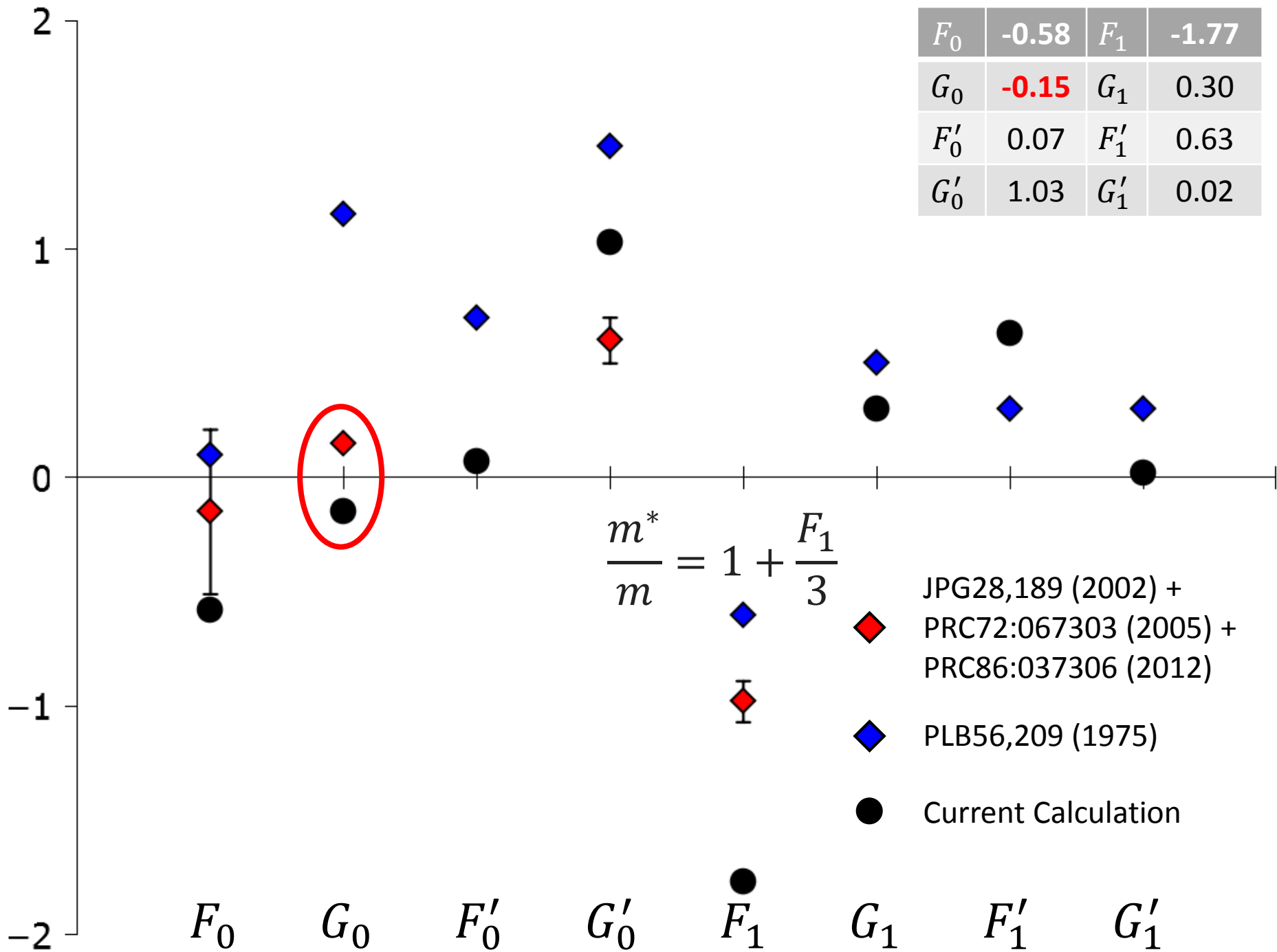
$$f'_l * \vec{\tau} \circ \vec{\tau} = \left(D^{(0,1)} \delta_{l,0} - E^{(0,1)} f_l(a, k_F) \right)$$

$$g'_l * \vec{\sigma} \cdot \vec{\sigma} \vec{\tau} \circ \vec{\tau} = \left(D^{(1,1)} \delta_{l,0} - E^{(1,1)} f_l(a, k_F) \right)$$





F_0	-0.58	F_1	-1.77
G_0	-0.15	G_1	0.30
F'_0	0.07	F'_1	0.63
G'_0	1.03	G'_1	0.02



dziękuję

Multipole Expansion of the Gaussian

$$\begin{aligned} v_{Reg} \propto g_a(\vec{k}_1 - \vec{k}_2) &= \frac{e^{-|\vec{k}_1 - \vec{k}_2|^2/a^2}}{(a\sqrt{\pi})^3} \\ &= \frac{4\pi}{(a\sqrt{\pi})^3} e^{-(k_1^2 + k_2^2)/a^2} \sum_{LM} i_L \left(2 \frac{k_1 k_2}{a^2} \right) Y_M^{L*}(\hat{k}_1) Y_M^L(\hat{k}_2) \\ &\quad \downarrow \\ &\quad (2L + 1) P_L(\hat{k}_1 \cdot \hat{k}_2) / 4\pi \end{aligned}$$

NLO contribution to the Matrix Element

Gradient on wavefunction

$$\vec{\nabla} [\mathcal{R}(r) Y_{m_l}^l(\hat{r})] = \sqrt{\frac{l+1}{2l+1}} \left(\frac{\partial \mathcal{R}(r)}{\partial r} - \frac{l}{r} \mathcal{R}(r) \right) \vec{Y}_{l m_l}^{l+1}(\hat{r}) + \sqrt{\frac{l}{2l+1}} \left(\frac{\partial \mathcal{R}(r)}{\partial r} - \frac{l+1}{r} \mathcal{R}(r) \right) \vec{Y}_{l m_l}^{l-1}(\hat{r})$$

Vector Spherical Harmonics

- Term $\propto \hat{T}_1 = (k'^2 + k^2)/2$

$$\propto g_a(\vec{r}_1 - \vec{r}_2) \vec{\nabla}'_1 \cdot \vec{\nabla}_1 \propto \int \vec{Y}_{l, m_l}^{l+1*}(\hat{r}_1) \cdot \vec{Y}_{l, m_l'}^{l+1}(\hat{r}_1) Y_M^L(\hat{r}_1) dr_1 \int Y_{-m_l}^l(\hat{r}_2) Y_{-m_l'}^l(\hat{r}_2) Y_M^L(\hat{r}_2) dr_2$$

Solved

- Term $\propto \hat{T}_2 \propto \vec{k} \cdot \vec{k}'$

$$\propto g_a(\vec{r}_1 - \vec{r}_2) \vec{\nabla}'_1 \cdot \vec{\nabla}_2 \propto \int \vec{Y}_{l, m_l_1}^{l+1*}(\hat{r}_1) Y_{-m_l_1}^l(\hat{r}_1) Y_M^L(\hat{r}_1) dr_1 \cdot \int \vec{Y}_{l, m_l_2}^{l+1*}(\hat{r}_1) Y_{-m_l_2}^l(\hat{r}_2) Y_M^L(\hat{r}_2) dr_2$$

???

