

(Preparations for) Symmetry unrestricted Skyrme mean-field study of heavy nuclei

Wouter Ryssens



26th of June, Warsaw

(Preparations
for)
Symmetry
unrestricted
Skyrme
mean-field
study of heavy
nuclei

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Introduction

MOCCa

Constraints

Tilted Axis
Cranking

Conclusion

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2 MOCCa

3 Constraints

4 Tilted Axis Cranking

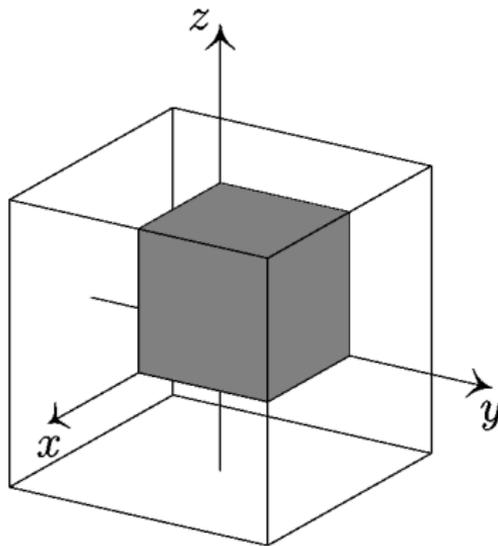
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$$\hat{P}, \hat{R}_z, \check{T}, \check{S}_y^T$$

- Parity
- Signature
- Time Simplex
- Time Reversal



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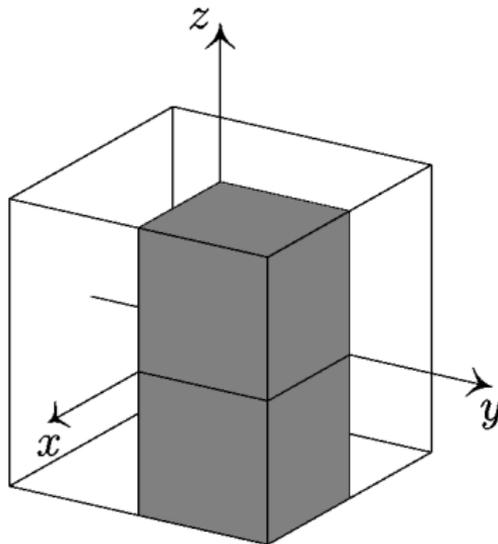
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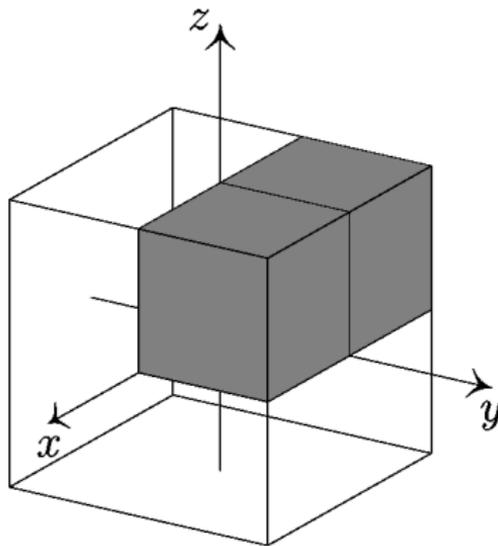
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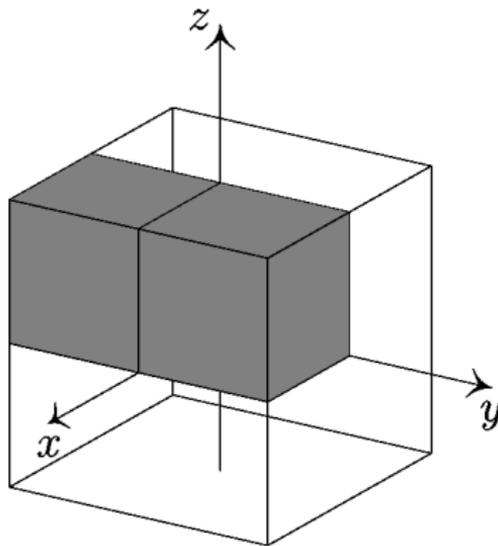
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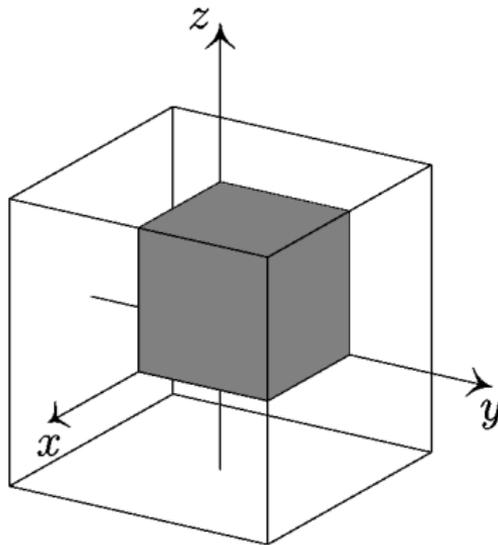
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MOCCa = Modular Cranking Code

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Done:

- 1 (Almost) all interactions
- 2 Hartree-Fock in any combination of symmetries
- 3 BCS in any spatial symmetry combination you wish
- 4 HFB in any spatial (!) symmetry you wish

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Problem:

- 1 HFB for time-reversal breaking configurations

Breaking symmetries: extra degrees of freedom

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Broken symmetries	Physical	Non-physical
\hat{P}	$\text{Re } \hat{Q}_{30}, \text{Re } \hat{Q}_{32}, \dots$	$\text{Re } \hat{Q}_{10}$
\check{S}_y^T	$\text{Im } \hat{Q}_{20}, \text{Im } \hat{Q}_{42}, \dots$	$\text{Im } \hat{Q}_{22}$
\hat{R}_z	$\text{Re } \hat{Q}_{41}, \text{Re } \hat{Q}_{43}, \dots$	$\text{Re } \hat{Q}_{21}$
\hat{R}_z, \hat{P}	$\text{Re } \hat{Q}_{31}, \text{Re } \hat{Q}_{33}, \dots$	$\text{Re } \hat{Q}_{11}$
\check{S}_y^T, \hat{P}	$\text{Im } \hat{Q}_{32}, \dots$	$\text{Im } \hat{Q}_{11}$
\check{T}	\hat{J}_z	
\check{T}, \hat{R}_z	\hat{J}_x	
\check{T}, \check{S}_y^T	\hat{J}_y	

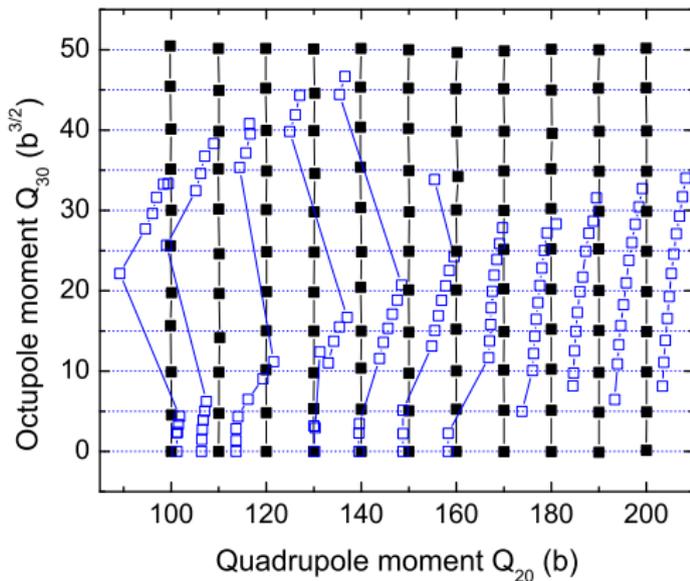
How to constrain quantities in mean-field picture?

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$$\hat{h}^{(i)} \rightarrow \hat{h}^{(i)} + \lambda^{(i)} \hat{O} \text{ with } \langle \hat{O} \rangle = O_0$$

	Form
Constant	$\lambda^{(i)} = C$
Quadratic/Penalty function	$\lambda^{(i)} = C(\langle \hat{O} \rangle^{(i)} - O_0)$
Augmented Lagrangian	$\lambda^{(i)} = \lambda^{(i-1)} + C(\langle \hat{O} \rangle^{(i)} - O_0)$

Augmented Lagrangian is superior



Augmented Lagrangian (black) versus quadratic (blue) constraints for ^{252}Fm using HFBTHO/HFOODD.

Figure from A. Staszczak et al., Eur. Phys. J. A **46**, 85-90 (2010)

However, ...

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$$\lambda^{(i)} = \lambda^{(i-1)} + C_{\hat{O}}(\langle \hat{O} \rangle^{(i)} - O_0)$$

Optimal choice of $C_{\hat{O}}$ depends on:

- Initial mean-field configuration
- O_0
- \hat{O}
- the presence of other constraints!

However, ...

$$^{20}\text{Ne}, \text{Re}\langle\hat{Q}_{10}\rangle = 0 \text{ fm}, \text{Re}\langle\hat{Q}_{20}\rangle = 10 \text{ fm}^2, \text{Re}\langle\hat{Q}_{30}\rangle = 10 \text{ fm}^3$$

$$\text{Deviation} = \sum_{l=1,2,3} r_{rms}^{-l} \left(\langle\hat{Q}_{l0}\rangle - Q_{l,desired} \right)^2$$

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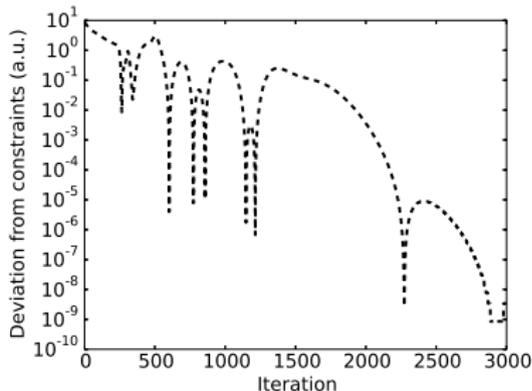
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$$C_{10} = 1, C_{20} = 0.08$$

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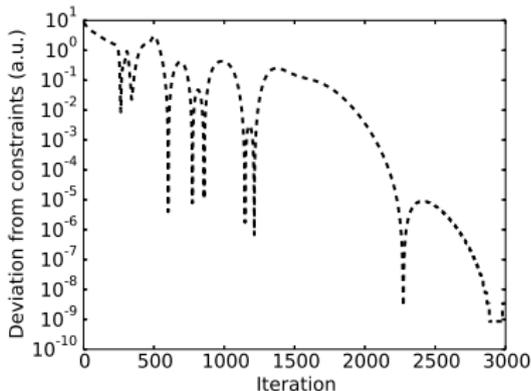
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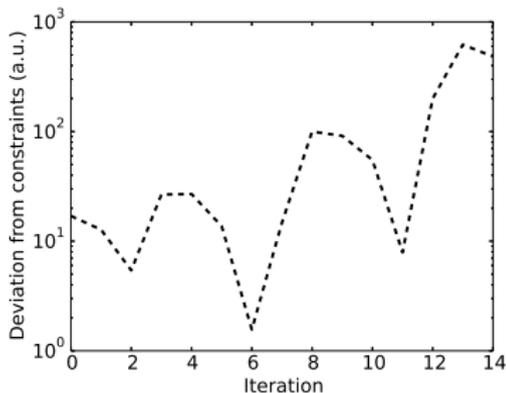
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Predictor-Corrector constraints

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First description:

R.Y. Cusson et al., Z. Phys. A 320, 475-482 (1985)

More complete:

K. Rutz, PhD thesis, Ibidem-Verlag. Frankfurt (1998)

But the procedure described is not exactly the one K. Rutz used, nor the one I use!

Augmented Lagrangian scheme:

- 1 Imaginary time step/gradient descent

$$\Psi^{(i)} = [1 - \epsilon(\hat{h}^{(i-1)} + \lambda^{(i-1)}\hat{O})]\Psi^{(i-1)}$$

- 2 Adapt

$$\lambda^{(i)} = \lambda^{(i-1)} + C(\langle\hat{O}\rangle^{(i)} - O_0)$$

- 3 construct $\hat{h}^{(i)}$

- 4 Repeat

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Predictor-Corrector scheme with two constants $C = 7.0$ and $K = 0.7$:

1 Trial step

$$\chi^{(i)} = \left[1 - \epsilon(\hat{h}^{(i-1)} + \lambda\hat{O}) \right] \Psi^{(i-1)}$$

2 Adapt

$$\lambda^{(i)} = \lambda^{(i-1)} + \frac{C}{A} \left(\langle \hat{O} \rangle^{(trial)} - \langle \hat{O} \rangle^{(i-1)} \right)$$

3 Do a corrective step

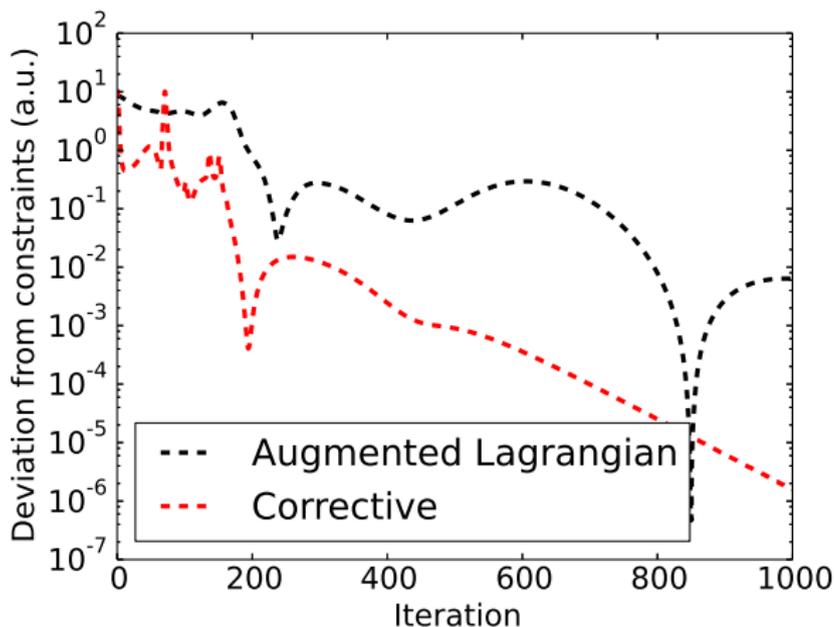
$$\Psi^{(i)} = \left[1 - \frac{K}{A} \left(\langle \hat{O} \rangle^{(trial)} - O_0 \right) \hat{O} \right] \chi^{(i)}$$

4 construct $\hat{h}^{(i)}$

5 repeat

where $A = [\langle \hat{O}^2 \rangle^{(trial)} + 0.00001]^{-1}$

Predictor - Corrector constraints



Predictor-Corrector constraints for ^{20}Ne , constraints on $\hat{Q}_{10}, \hat{Q}_{20}$ and \hat{Q}_{30} .

Tilted axis cranking

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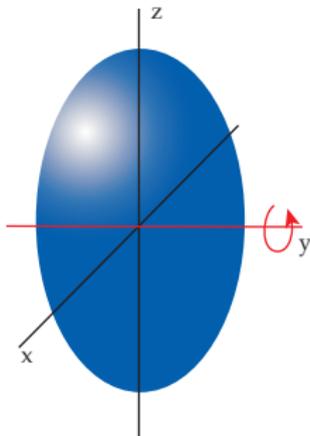
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$$\hat{H}' = \hat{H} - \omega_z \hat{J}_z$$



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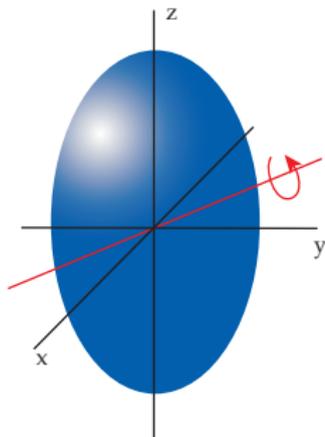
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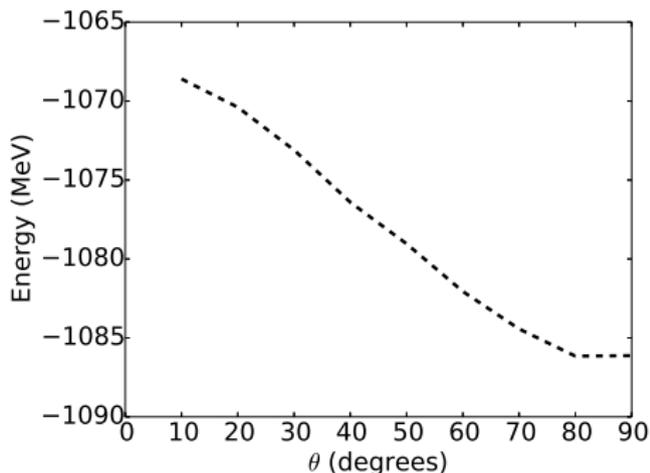
Conclusion

$$\hat{H}' = \hat{H} - \omega_x \hat{J}_x - \omega_y \hat{J}_y - \omega_z \hat{J}_z$$



Y. Shi et al., Phys. Rev. Lett. **108** 092501 (2012)

- $\langle J_x \rangle = 70 \sin(\theta) \hbar$
- $\langle J_z \rangle = 70 \cos(\theta) \hbar$
- $\langle Q_{20} \rangle = 1900 \text{ fm}^2$
- $\langle Q_{22} \rangle = 500 \text{ fm}^2$
- $\langle Q_{21} \rangle = 0 \text{ fm}^2$



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- Future: fix the remaining problems

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- Future: fix the remaining problems
- Future: implement the SLyMRX interactions